

# Final Recitation – Information Theory, Application for Quantum Fault Tolerance.

David Ponarovsky

June 25, 2025

# Introduction

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- ▶ Noisy circuit and noisy computation.
- ▶ Classical no-go for reversible noisy computation at logarithmic depth (and polynomial space).
- ▶ Quantum case.
- ▶ Trading (local) entropy for space.

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## Definition

$p$ - Depolarizing Channel. The qubit depolarizing channel with parameter  $p \in [0, 1]$  is the quantum channel  $\mathcal{D}_p$  defined by:

$$\mathcal{D}_p(\rho) = (1 - p)\rho + p \cdot \frac{I}{2}$$

where  $\rho$  is a single-qubit density matrix and  $I$  is the identity matrix.

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$p$ -Noisy Circuit. Given a circuit  $C$  (regardless of the model), its  $p$ -noisy version  $\tilde{C}$  is the circuit obtained by alternately taking layers from  $C$  and then passing each (qu)bit through a  $p$ -Depolarizing channel.



# Classical no-go.

## Informal:

We can't compute when subjected to noise, after a while, the bits become garbage.

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## Claim

Denote by  $X = (X_1, X_2, \dots, X_m)$  and  $Y = (Y_1, Y_2, \dots, Y_m)$  the input and the output distributions of reversible  $p$ -noisy computation at width  $m$  (bits) and depth  $d$ . Then:

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In particular, for  $d = \Omega(\log m)$  we have  $H(Y) \rightarrow m$ .

# Classical no-go.

## Claim

Let  $Y$  be a bit given by moving  $X$  through  $\text{BSC}(p)$ , Then there is  $\gamma_p < 1$  such :

$$1 - H(Y) \leq \gamma(1 - H(X))$$

## Classical no-go.

Denote by  $\delta$  the parameter for which  $X$  distributed as  $\sim \text{Bin}(\frac{1+\delta}{2})$ .  
First notice that:

$$\Pr(Y = 1) = \frac{1+\delta}{2}(1-p) + \frac{1-\delta}{2}p = \frac{1+\delta-2p\delta}{2}$$

So  $Y \sim \text{Bin}(\frac{1+\delta-2p\delta}{2})$ , Or  $\delta \mapsto 1 - 2p\delta$ .

## Classical no-go.

Now expand  $1 - H(X)$  to it's Taylor Series at  $\delta$  gives:

$$1 - H(X) = 1 - \frac{1}{2} \left( (1 + \delta) \log \left( \frac{1 + \delta}{2} \right) + (1 - \delta) \log \left( \frac{1 - \delta}{2} \right) \right)$$

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Denote the above by  $K(\delta)$

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Now, observes that:

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And notice that since  $p < 1$  we have  $\gamma < 1$ , notice also that inequality is symmetric to  $p \mapsto 1-p$ , in particular the entropy is not increase if either  $p = 0$  or  $p = 1$ .



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## Claim

Let  $Y = (Y_1, Y_2, \dots, Y_m)$  be a bit given by moving each of  $X_i \in X = (X_1, X_2, \dots, X_m)$  through  $\text{BSC}(p)$ . Then:

$$m - H(Y) \leq \gamma(m - H(X))$$

# Classical no-go.

Proof.

$$m - H(Y_1, Y_2, \dots, Y_m) = m - \sum_i H(Y_i | Y_1, Y_2, \dots, Y_{i-1})$$

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In particular, for  $d = \Omega(\log m)$  we have  $H(Y) \rightarrow m$ .

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## Claim

Let  $\rho_1$  be a reduce density matrix of  $\rho$  Then:

$$-\text{Tr } \rho \log (\rho_1 \otimes I) = S(\rho_1)$$

## Quantum no-go.

First consider the case in which  $\rho$  is a tensor of  $\rho_1$  namely  $\rho = \rho_1 \otimes \rho_2$ . Then clearly  $\rho$  and  $\log \rho_1 \otimes I$  commute. Denote by  $\lambda_1, \dots, \lambda_n$  and  $\mu_1, \dots, \mu_m$  the eigen values of  $\rho_1$  and  $\rho_2$ . So the trace equals:

$$\sum \lambda_i \mu_j \log(\lambda_i \cdot 1) = \left( \sum \mu_j \right) \left( \sum_i \lambda_i \log \lambda_i \right)$$

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$$\begin{aligned} \sum \lambda_i \mu_j \log(\lambda_i \cdot 1) &= \left( \sum \mu_j \right) \left( \sum_i \lambda_i \log \lambda_i \right) \\ &= (\text{Tr } \rho_2) \sum_i \lambda_i \log \lambda_i = -S(\rho_1) \end{aligned}$$

# Quantum no-go.

Let's use the notation  $\sum_{A_k} \rho|_{A_k}$  to denote the sum over all the reduced matrices over  $k$  qubits.

## Claim

Let  $\rho$  be a density matrix over  $n$  qubits then:

$$\binom{n}{k}^{-1} \sum_{A_k} I(\rho|_{A_k}) \leq \frac{k}{n} I(\rho)$$

# Quantum no-go.

## Proof

Let  $\rho_1$  be  $\rho^{\otimes k \binom{n}{k}}$  and let  $\rho_2 = \left( \prod_{A_k} \rho|_{A_k} \right)^{\otimes n}$ .

$$0 \geq S(\rho_2|\rho_1) = \mathbf{Tr} \left( \rho_1 (\log \rho_1 - \log \rho_2) \right) = -S(\rho_1) - \mathbf{Tr} \left( \rho_1 \log \rho_2 \right)$$



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Now observes that  $\rho|_{A_K}^{\otimes n}$  is a reduced density matrix of  $\rho_1$ . So we get:

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Now observes that  $\rho|_{A_K}^{\otimes n}$  is a reduced density matrix of  $\rho_1$ . So we get:

$$\begin{aligned} 0 &\leq -k \binom{n}{k} S(\rho) - \sum_{A_k} n S(\rho|_{A_k}) \\ \Rightarrow \sum_{A_k} I(\rho|_{A_k}) &\leq \frac{k}{n} \binom{n}{k} I(\rho) \end{aligned}$$

# Quantum no-go.

## Claim

Let  $\rho$  be a density matrix of  $n$  qubits. Let each qubit be replaced with independent probability  $p$  by a fully mixed qubit denoted by  $v$ , to give the density matrix  $\sigma$ . Then  $I(\sigma) \leq (1 - p) I(\rho)$ .

## Quantum no-go.

Let us write:

$$\sigma = \sum_{k=1}^n \sum_{A_k} p^{n-k} (1-p)^k \rho|_{A_k} \otimes v^{n-k}$$

By the concavity of the entropy (convexity of  $I$ ), We have:

$$I(\sigma)$$

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$$I(\sigma) \leq \sum_{k=1}^n \sum_{A_k} p^{n-k} (1-p)^k [I(\rho|_{A_k}) + (n-k)I(v)]$$

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# Entropy-Space tradeoff.

Zeros Distillation. Consider the unitary majority gate over 3 qubits, which on the computational basis sets the last 3rd bit to be the majority, and acts as follows:

$$M |0, 0, 1\rangle \mapsto |1, 1, 0\rangle$$

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And acts trivially on every other configuration.

# Entropy-Space tradeoff.

## Claim

There is a  $p_0$  such that for any  $p < p_0$ , it follows that if each qubit has a probability less than  $p$  of being at  $|1\rangle$ , then after a cycle of Noise  $\rightarrow$  Majority, the probability of the third qubit being  $|1\rangle$  is still less than  $p$ .

# Entropy-Space tradeoff.

Proof.

After the noise round, the density matrices of each qubit:

$$\leq (1-p)((1-p)|0\rangle\langle 0| + p|1\rangle\langle 1|) + p\frac{1}{2}(|0\rangle\langle 0| + |1\rangle\langle 1|)$$

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After the noise round, the density matrices of each qubit:

$$\begin{aligned} &\leq (1-p) ((1-p) |0\rangle \langle 0| + p |1\rangle \langle 1|) + p \frac{1}{2} (|0\rangle \langle 0| + |1\rangle \langle 1|) \\ &= \left( (1-p)^2 + \frac{1}{2}p \right) |0\rangle \langle 0| + \left( (1-p)p + \frac{1}{2}p \right) |1\rangle \langle 1| \\ &= |0\rangle \langle 0| + (3/2p - p^2) |1\rangle \langle 1| \end{aligned}$$

And for convenience, let's bound by looking at the following noisier state, where each qubit is in:

$$(1-2p) |0\rangle \langle 0| + 2p |1\rangle \langle 1|$$



# Entropy-Space tradeoff.

## Proof.

After the noise round, the density matrices of each qubit:

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The End.

June 25, 2025



# Why Tracing-out the first and second qubit doesn't work?.

June 25, 2025

# Entropy-Space tradeoff.

## Claim

Consider the noisy zero  $\rho = (1 - \frac{p}{2}) |0\rangle \langle 0| + \frac{p}{2} |1\rangle \langle 1|$ . Then:

$$\mathbf{Tr}_{[1,2]} M \rho^3 = \left(1 - \frac{3}{4} p^2 + \frac{1}{4} p^3\right) |0\rangle \langle 0| + , ,$$



## Entropy-Space tradeoff.

$$\rho^{\otimes 3} = \left(1 - \frac{p}{2}\right)^3 |000\rangle \langle 000|$$

$$M_{\rho^{\otimes 3}}$$

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$$\begin{aligned}\rho^{\otimes 3} = & \left(1 - \frac{p}{2}\right)^3 |000\rangle \langle 000| \\ & + \left(1 - \frac{p}{2}\right)^2 \frac{p}{2} (|001\rangle \langle 001| + |010\rangle \langle 010| + |100\rangle \langle 100|)\end{aligned}$$

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## Entropy-Space tradeoff.

$$\mathbf{Tr}_{[1,2]} M_{\rho^{\otimes 3}} = \left( \left(1 - \frac{p}{2}\right)^3 + 3 \left(1 - \frac{p}{2}\right)^2 \frac{p}{2} \right) |0\rangle \langle 0| +$$

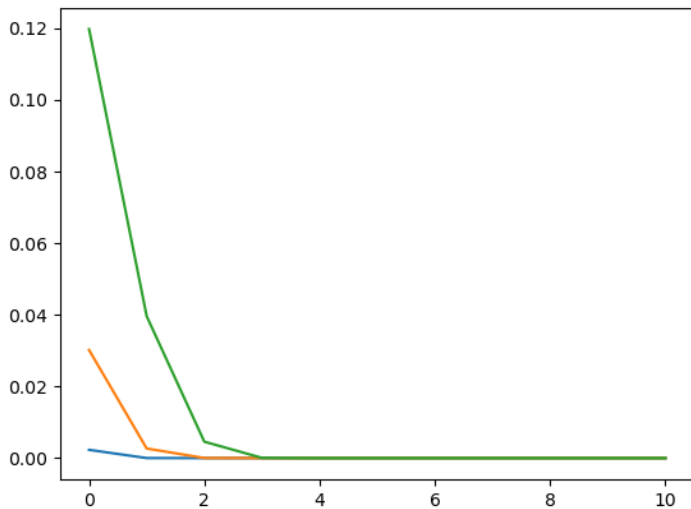
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$$\left(1 - \frac{3}{4}p^2 + \frac{1}{4}p^3\right) |0\rangle \langle 0| + , ,$$

$\mapsto$

$=$

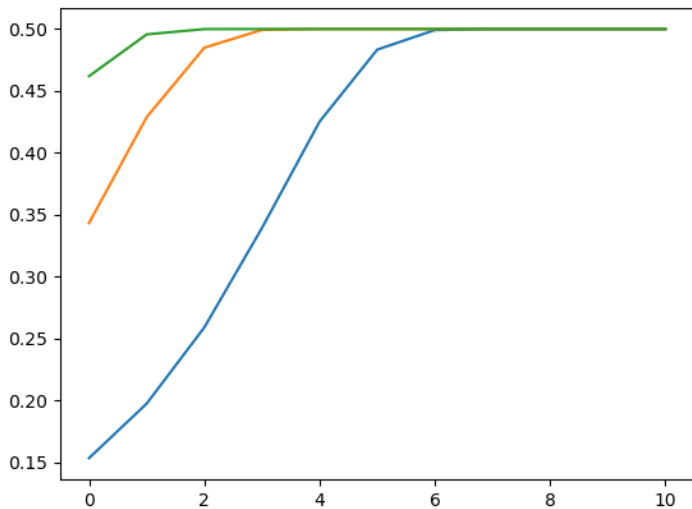
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# Your Title Here





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