Final Recitation – Information Theory, Application for Quantum Fault Tolerance.

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- Classical no-go for reversible noisy computation at logarithmic depth (and polynomial space).

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- Classical no-go for reversible noisy computation at logarithmic depth (and polynomial space).
- Quantum case.
- ► Trading (local) entropy for space.

Nosiy Circuit.



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Definition

p- Depolarizing Channel. The qubit depolarizing channel with parameter $p \in [0,1]$ is the quantum channel \mathcal{D}_p defined by:

$$\mathcal{D}_{p}(\rho) = (1-p)\rho + p \cdot \frac{l}{2}$$

where ρ is a single-qubit density matrix and I is the identity matrix.

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Definition

p-Noisy Circuit. Given a circuit C (regardless of the model), its p-noisy version \tilde{C} is the circuit obtained by alternately taking layers from C and then passing each (qu)bit through a p-Depolarizing channel.

Informal:

We can't compute when subjected to noise, after a while, the bits become garbage.

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Claim

Denote b $X = (X_1, X_2, ..., X_m)$ and $Y = (Y_1, Y_2, ..., Y_m)$ the input and the output distrubtions of reversible p-noisy computation at widith m (bits) and depth d. Then:

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$$m - H(Y) \le \gamma^d (m - H(X))$$

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$$m - H(Y) \le \gamma^d (m - H(X))$$

In particular, for $d = \Omega(\log m)$ we have $H(Y) \to m$.

Claim

Let Y be a bit given by moving X trough BSC(p), Then there is $\gamma_p < 1$ such :

$$1 - H(Y) \le \gamma \left(1 - H(X)\right)$$

Denote by δ the parameter for which X distributed as $\sim Bin(\frac{1+\delta}{2})$. First notice that:

$$\Pr(Y = 1) = \frac{1+\delta}{2}(1-p) + \frac{1-\delta}{2}p = \frac{1+\delta-2\delta p}{2}$$

So
$$Y \sim \text{Bin}(\frac{1-\delta(1-2p)}{2})$$
, Or $\delta \mapsto 1-2p\delta$.

$$1 - H(X) = 1 - \frac{1}{2} \left((1 + \delta) \log \left(\frac{1 + \delta}{2} \right) + (1 - \delta) \log \left(\frac{1 - \delta}{2} \right) \right)$$

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$$= -\frac{1}{2} \cdot (1+\delta) \sum_{i=1}^{\infty} \frac{(-1)^{n+1} \delta^n}{n} + (1-\delta) \sum_{i=1}^{\infty} \frac{(-1)^{n+1} (-\delta)^n}{n}$$

$$\begin{split} 1 - H(X) &= 1 - \frac{1}{2} \left((1 + \delta) \log \left(\frac{1 + \delta}{2} \right) + (1 - \delta) \log \left(\frac{1 - \delta}{2} \right) \right) \\ &= -\frac{1}{2} \left((1 + \delta) \log \left(\frac{1 + \delta}{2} \right) + (1 - \delta) \log \left(\frac{1 - \delta}{2} \right) \right) \\ &= -\frac{1}{2} \cdot (1 + \delta) \sum_{i=1}^{\infty} \frac{(-1)^{n+1} \delta^n}{n} + (1 - \delta) \sum_{i=1}^{\infty} \frac{(-1)^{n+1} (-\delta)^n}{n} \\ &= -\frac{1}{2} \cdot \sum_{i=1}^{\infty} 2 \frac{\delta^{2n}}{2n} - \sum_{i=1}^{\infty} 2 \frac{\delta^{2n}}{2n - 1} \end{split}$$

$$1 - H(X) = 1 - \frac{1}{2} \left((1 + \delta) \log \left(\frac{1 + \delta}{2} \right) + (1 - \delta) \log \left(\frac{1 - \delta}{2} \right) \right)$$

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$$= -\frac{1}{2} \cdot \sum_{i=1}^{\infty} 2 \frac{\delta^{2n}}{2n} - \sum_{i=1}^{\infty} 2 \frac{\delta^{2n}}{2n-1}$$

$$= \sum_{i=1}^{\infty} \frac{\delta^{2n}}{2n(2n-1)}$$

Now expand 1 - H(X) to it's Taylor Seryias at δ gives:

$$1 - H(X) = 1 - \frac{1}{2} \left((1 + \delta) \log \left(\frac{1 + \delta}{2} \right) + (1 - \delta) \log \left(\frac{1 - \delta}{2} \right) \right)$$

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Denote the above by $K(\delta)$

Now, observes that:

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And notice that since p<1 we have $\gamma<1$, noitce also that inequlity is symmetric to $p\mapsto 1-p$, in paritcular the entropy is not increase if either p=0 or p=1.

Claim

Let $Y = (Y_1, Y_2, ..., Y_m)$ be a bit given by moving each of $X_i \in X = (X_1, X_2, ..., X_m)$ trough BSC(p). Then:

$$m - H(Y) \le \gamma (m - H(X))$$

$$m - H(Y_1, Y_2, ..., Y_m) = m - \sum_{i} H(Y_i | Y_1, Y_2, ..., Y_{i-1})$$

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$$= \gamma (m - H(X))$$

Claim

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Let's continue to the quantum case. From now on we will denote by $I(\rho) = \dim \rho - S(\rho)$.

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Quantum no-go.

Claim

Let ρ_1 be a reduce density matrix of ρ Then:

$$-\mathsf{Tr}\ \rho\log\left(\rho_1\otimes I\right)=S(\rho_1)$$

First consider the case in which ρ is a tensor of ρ_1 namely $\rho=\rho_1\otimes\rho_2$, Then clearly ρ and $\log\rho_1\otimes I$ commute. Denote by $\lambda_1,...\lambda_n$ and $\mu_1,...\mu_m$ the eigen values of ρ_1 and ρ_2 . So the trace equals:

$$\sum \lambda_i \mu_j \log(\lambda_i \cdot 1) = \left(\sum \mu_j\right) \left(\sum_i \lambda_i \log \lambda_i\right)$$

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Then consider the case where ρ is entangle. Let's span ρ over base at the form $\{\rho_1 \otimes \rho_2^i\}$ We get:

$$\mathsf{Tr} \; \rho \log \left(\rho_1 \otimes I \right) = \mathsf{Tr} \; \sum_i \mathsf{a}_i \rho_1 \otimes \rho_2^i \log \left(\rho_1 \otimes I \right)$$
$$= - \sum_i \mathsf{a}_i S(\rho_1) = - S(\rho_1)$$

Let's use the notation $\sum_{A_k} \rho|_{A_k}$ to denote the sum over all the reduced matrices over k qubits.

Claim

Let ρ be a density matrix over n qubits then:

$$\binom{n}{k}^{-1} \sum_{A_k} I(\rho|_{A_k}) \le \frac{k}{n} I(\rho)$$

Proof

Let
$$\rho_1$$
 be $\rho^{\otimes k\binom{n}{k}}$ and let $\rho_2 = \left(\prod_{A_k} \rho|_{A_k}\right)^{\otimes n}$.

$$0 \leq S(\rho_2||\rho_1) = \mathsf{Tr} \; \left(\rho_1 \left(\log \rho_1 - \log \rho_2\right)\right) = -S(\rho_1) - \mathsf{Tr} \; \left(\rho_1 \log \rho_2\right)$$

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Now observes that $\rho|_{A_k}^{\otimes n}$ is a reduced density matrix of ρ_1 . So we get:

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Now observes that $\rho|_{A_k}^{\otimes n}$ is a reduced density matrix of ρ_1 . So we get:

$$0 \le -k \binom{n}{k} S(\rho) + \sum_{A_k} n S(\rho|_{A_k})$$

$$\Rightarrow \sum_{A_k} I(\rho|_{A_k}) \le \frac{k}{n} \binom{n}{k} I(\rho)$$

Claim

Let ρ be a density matrix of n qubits. Let each qubit be replaced with independent probability ρ by a fully mixed qubit denoted by v, to give the density matrix σ . Then $I(\sigma) \leq (1-\rho) \, I(\rho)$.

Let us write:

$$\sigma = \sum_{k=1}^{n} \sum_{A_k} p^{n-k} (1-p)^k \rho|_{A_k} \otimes v^{n-k}$$

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$$= (1-p) I(\rho)$$

Zeros Distillation. Consider the unitary majority gate over 3 qubits, which on the computational basis sets the last 3rd bit to be the majority, and acts as follows:

$$M |0,0,1\rangle \mapsto |1,1,0\rangle$$

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And acts trivially on every other configuration.

Claim

There is a p_0 such that for any $p < p_0$, it follows that if each qubit has a probability less than p of being at $|1\rangle$, then after a cycle of Noise \rightarrow Majority, the probability of the third qubit being $|1\rangle$ is still less than p.

Proof.

After the noise round, the density matrices of each qubit¹:

$$\leq (1-
ho)\left((1-
ho)\ket{0}ra{0}+
ho\ket{1}ra{1}
ight)+
horac{1}{2}\left(\ket{0}ra{0}+\ket{1}ra{1}
ight)$$

 $^{^{1}}$ The inequality refers to the $|1\rangle\langle 1|$ coefficient.

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After the noise round, the density matrices of each qubit¹:

$$\leq (1-p)\left((1-p)\left|0\right\rangle\left\langle 0\right|+p\left|1\right\rangle\left\langle 1\right|\right)+p\frac{1}{2}\left(\left|0\right\rangle\left\langle 0\right|+\left|1\right\rangle\left\langle 1\right|\right)$$

$$=\left((1-p)^{2}+\frac{1}{2}p\right)\left|0\right\rangle\left\langle 0\right|+\left((1-p)p+\frac{1}{2}p\right)\left|1\right\rangle\left\langle 1\right|$$

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After the noise round, the density matrices of each qubit¹:

$$\leq (1 - p) ((1 - p) |0\rangle \langle 0| + p |1\rangle \langle 1|) + p \frac{1}{2} (|0\rangle \langle 0| + |1\rangle \langle 1|)$$

$$= \left((1 - p)^2 + \frac{1}{2}p \right) |0\rangle \langle 0| + \left((1 - p) p + \frac{1}{2}p \right) |1\rangle \langle 1|$$

$$= " |0\rangle \langle 0| + (3/2p - p^2) |1\rangle \langle 1|$$

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$$= " |0\rangle \langle 0| + (3/2p - p^2) |1\rangle \langle 1|$$

And for convenience, let's bound it by looking at the following noisier state, where each qubit is in:

$$(1-2p)|0\rangle\langle 0|+2p|1\rangle\langle 1|$$

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$$(2p)^3 + 3(1-2p)(2p)^2$$

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And for $p_0 = \frac{1}{6}$ we have that this probability is less than p.

The End.

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