AST4320 Oblig 2

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1 Exercise 1

We have the 1D window function

$$W(x) = \begin{cases} 1 & \text{if } |x| < R \\ 0 & \text{else} \end{cases} \tag{1}$$

We now want to find the Fourier transform of this function. We do this straight forward from the definition of the Fourier transform

$$\tilde{W}(k) = \int_{-\infty}^{\infty} W(x)e^{-ikx}dx. \tag{2}$$

Inserting our definition of W(x) we find

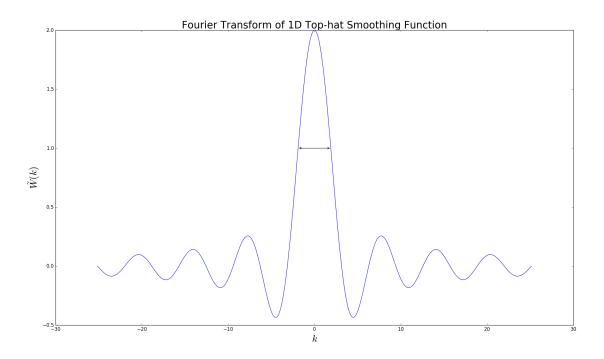
$$\tilde{W}(k) = \int_{-\infty}^{-R} 0 \cdot e^{-ikx} dx + \int_{-R}^{R} 1 \cdot e^{-ikx} dx + \int_{R}^{\infty} 0 \cdot e^{-ikx} dx = \int_{-R}^{R} e^{-ikx} dx$$
 (3)

$$= \frac{i}{k}e^{-ikx}\Big|_{-R}^{R} = \frac{i}{k}\left(e^{-ikR} - e^{ikR}\right) = \frac{2\sin Rk}{k}.$$
 (4)

Before we plot this function we need to notice that this is a function that we need to be careful with for $k \to 0$. With the use of L'Hôpitals rule

$$\lim_{k \to 0} \tilde{W}(k) = \lim_{k \to 0} \frac{2\sin Rk}{k} = \lim_{k \to 0} \frac{2R\cos Rk}{1} = 2R.$$
 (5)

We now see that $\tilde{W}(k)$ is well defined across the whole real line.



Figur 1: The plot of the Fourier Transform of the top-hat smoothing function W(x). This is plotted with R=1. The FWHM is marked with the arrow.

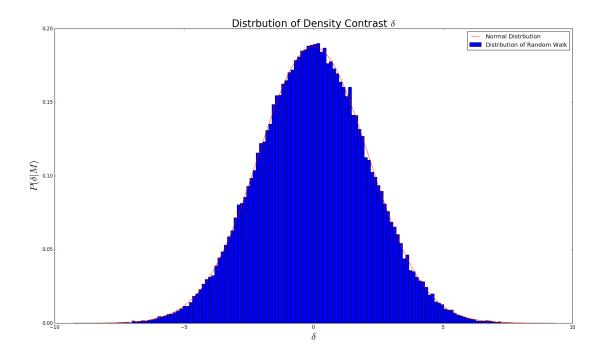
We now want to find the full width at half maximum, FWHM, for this function. We first need to find the half maximum. From fig.2 we see that maximum is at 0. We know from (5) that here $\tilde{W}(k=0)=2R$, so half maximum is R. For the width we just need to find for which k we have $\tilde{W}=R$ and multiply it by 2 (this is because \tilde{W} is symmetric around 0). So we first need to solve

$$R = \frac{2}{k_{\text{half max}}} \sin R k_{\text{half max}}.$$
 (6)

This is difficult to solve analytical, but easy numerically. From our program we easily find that

$$FWHM = 2 \cdot k_{half max} = 3.77. \tag{7}$$

2 Exercise 2



Figur 2: Distribution of the endstep in the random walk. This is plotted with $\epsilon = 0.1$ and $\sigma^2(S_c) = 10^{-4}$. We see that this follows a Gaussian distribution with $\sigma^2 = \pi$.

3 Exercise 3

3.1 a)

We have that the probability of δ provided that δ was never larger than δ_{crit} at some scale M' > M is

$$P_{nc}(\delta|M) = \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) - \exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right]. \tag{8}$$

Since δ never has crossed δ_{crit} it is safe to assume that collapse has never happened here. This means that the probability of a mass at \mathbf{x} is therefore embedded within a *non*-collapsed object of mass > M is therefore

$$\int_{-\infty}^{\delta_{crit}} P_{nc}(\delta|M) d\delta. \tag{9}$$

This means that the probability that this mass is within a collapsed mass is 1 minus this probability. So

$$P(>M) = 1 - \int_{-\infty}^{\delta_{crit}} P_{nc}(\delta|M) d\delta. \tag{10}$$

3.2 b)

Using (11) and (8) we get

$$P(>M) = 1 - \int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) - \exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right] d\delta$$
(11)
$$= 1 - \int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) \right] d\delta + \int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right] d\delta$$
(12)