

AST4320 Oblig 2

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1 Exercise 1

We have the 1D window function

$$W(x) = \begin{cases} 1 & \text{if } |x| < R \\ 0 & \text{else} \end{cases} \quad (1)$$

We now want to find the Fourier transform of this function. We do this straight forward from the definition of the Fourier transform

$$\tilde{W}(k) = \int_{-\infty}^{\infty} W(x) e^{-ikx} dx. \quad (2)$$

Inserting our definition of $W(x)$ we find

$$\tilde{W}(k) = \int_{-\infty}^{-R} 0 \cdot e^{-ikx} dx + \int_{-R}^R 1 \cdot e^{-ikx} dx + \int_R^{\infty} 0 \cdot e^{-ikx} dx = \int_{-R}^R e^{-ikx} dx \quad (3)$$

$$= \left. \frac{i}{k} e^{-ikx} \right|_{-R}^R = \frac{i}{k} (e^{-ikR} - e^{ikR}) = \frac{2 \sin Rk}{k}. \quad (4)$$

Before we plot this function we need to notice that this is a function that we need to be careful with for $k \rightarrow 0$. With the use of L'Hôpital's rule

$$\lim_{k \rightarrow 0} \tilde{W}(k) = \lim_{k \rightarrow 0} \frac{2 \sin Rk}{k} = \lim_{k \rightarrow 0} \frac{2R \cos Rk}{1} = 2R. \quad (5)$$

We now see that $\tilde{W}(k)$ is well defined across the whole real line.

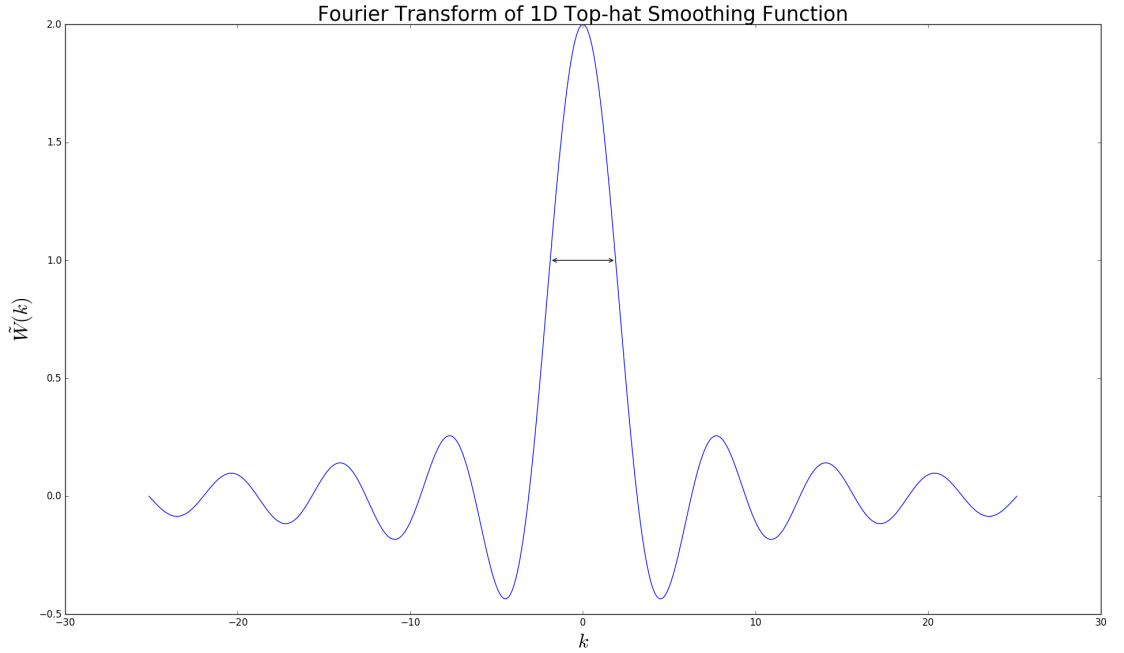


Figure 1: The plot of the Fourier Transform of the top-hat smoothing function $W(x)$. This is plotted with $R = 1$. The FWHM is marked with the arrow.

We now want to find the *full width at half maximum*, FWHM, for this function. We first need to find the half maximum. From fig.1 we see that maximum is at 0. We know from (5) that here $\bar{W}(k=0) = 2R$, so half maximum is R . For the width we just need to find for which k we have $\bar{W} = R$ and multiply it by 2 (this is because \bar{W} is symmetric around 0). So we first need to solve

$$R = \frac{2}{k_{\text{half max}}} \sin R k_{\text{half max}}. \quad (6)$$

This is difficult to solve analytical, but easy numerically. From our program we easily find that

$$\text{FWHM} = 2 \cdot k_{\text{half max}} = 3.77. \quad (7)$$

2 Exercise 2

By doing the random walk, found in *ex2.py*, we get the distribution for δ found in fig. 2. We can see that this follows the Gaussian function $\delta \sim N(0, \sigma(M))$. If we restrict our self to the walks that never cross the critical density δ_c we get the distribution for δ_{nc} seen in fig. 3. To make the hight of the analytical distribution fit that of the simulated data, I had to multiply it with a fudge factor of 2.2. Even with this fudge factor the simulated distribution is skewed to the right, with a dramatic cut off point at $\delta = 1$. This cut off point is expected since we restrict our walks to $\delta < \delta_c = 1$, but the cut off is a bit sharp. With a lower ϵ this becomes smaller, but the code takes to long to run.

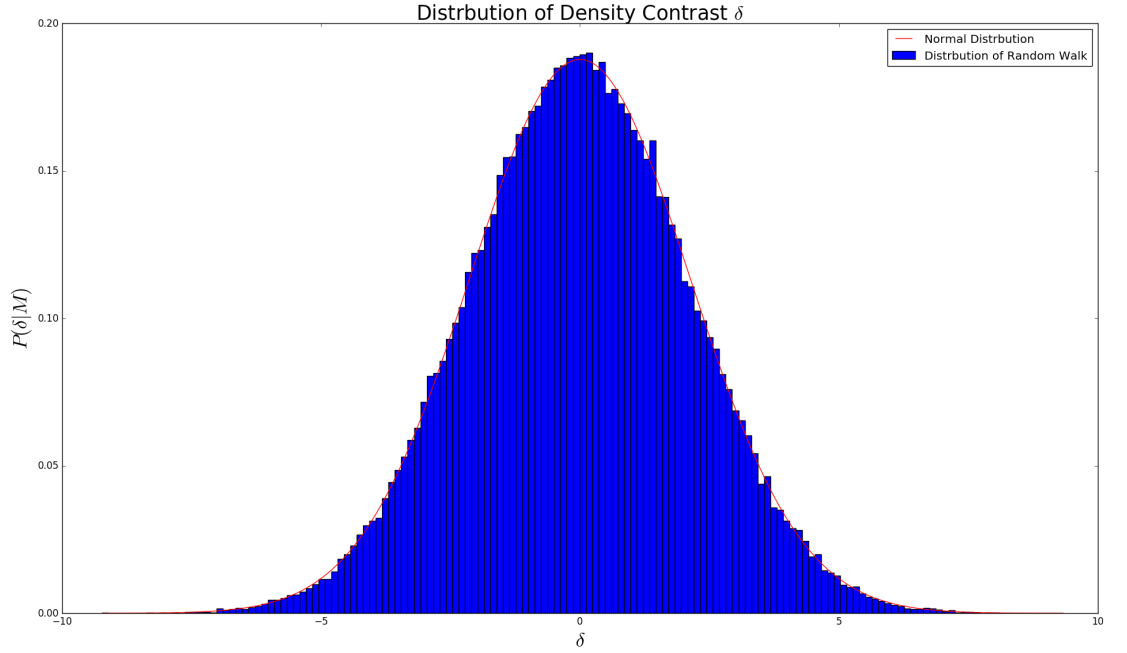


Figure 2: Distribution of the endstep in the random walk, simulated with $\epsilon = 0.1$ and 10^5 walks. We see that this follows a Gaussian distribution with $\sigma^2 = \pi$.

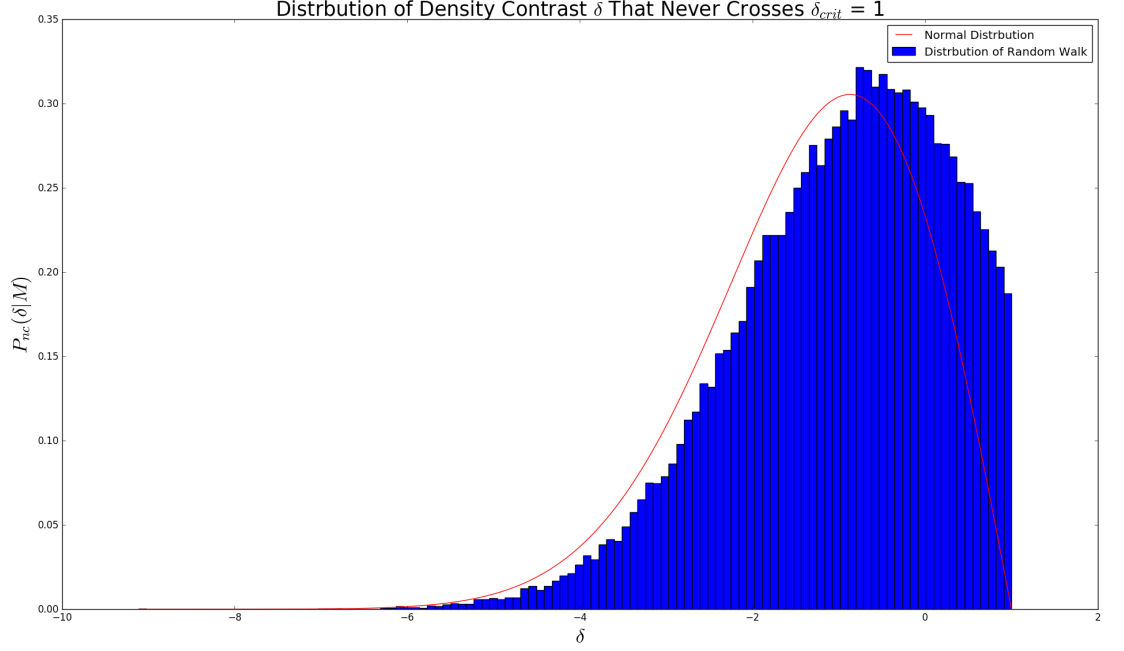


Figure 3: Distribution of the endstep in the random walk, simulated with $\epsilon = 0.1$ and 10^5 walks. We can see that distribution is a bit of the expected distribution, even with a 2.2 fudge factor.

3 Exercise 3

3.1 a)

We have that the probability of δ provided that δ was never larger than δ_{crit} at some scale $M' > M$ is

$$P_{nc}(\delta|M) = \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) - \exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right]. \quad (8)$$

Since δ never has crossed δ_{crit} it is safe to assume that collapse has never happened here. This means that the probability of a mass at \mathbf{x} is therefore embedded within a *non*-collapsed object of mass $> M$ is therefore

$$\int_{-\infty}^{\delta_{crit}} P_{nc}(\delta|M) d\delta. \quad (9)$$

This means that the probability that this mass is within a collapsed mass is 1 minus this probability. So

$$P(> M) = 1 - \int_{-\infty}^{\delta_{crit}} P_{nc}(\delta|M) d\delta. \quad (10)$$

3.2 b)

Using (11) and (8) we get

$$P(> M) = 1 - \int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) - \exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right] d\delta \quad (11)$$

$$= 1 - \underbrace{\int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{\delta^2}{2\sigma^2(M)}\right) \right] d\delta}_{I_1} + \underbrace{\int_{-\infty}^{\delta_{crit}} \frac{1}{\sqrt{2\pi}\sigma(M)} \left[\exp\left(-\frac{(2\delta_{crit} - \delta)^2}{2\sigma^2(M)}\right) \right] d\delta}_{I_2}. \quad (12)$$

We then look at this term by term:

I_1 :

We know that we want something with an error function. To get this we see that we need to introduce the variable

$$u = \frac{\delta}{\sqrt{2}\sigma}, \quad (13)$$

which gives

$$\Rightarrow \frac{du}{d\delta} = \frac{1}{\sqrt{2}\sigma} \Rightarrow d\delta = \sqrt{2}\sigma du. \quad (14)$$

We need to take care of the upper limit u_{\uparrow} . We see that this becomes

$$\delta_{\uparrow} = \delta_{crit} = \sqrt{2}\sigma u_{\uparrow} \Rightarrow u_{\uparrow} = \frac{\delta_{crit}}{\sqrt{2}\sigma} \equiv \frac{\nu}{\sqrt{2}}. \quad (15)$$

We can then solve I_1 as

$$I_1 = \frac{1}{2} \int_{-\infty}^{\nu/\sqrt{2}} \frac{2}{\sqrt{\pi}} e^{-u^2} du. \quad (16)$$

Looking this up on integration tables we find that this is

$$I_1 = \frac{1}{2} \left[1 + \operatorname{erf}\left(\frac{\nu}{\sqrt{2}}\right) \right] \quad (17)$$

I_2 :

This integral is quite similar to I_1 , so we use a similar variable change

$$u = \frac{2\delta_{crit} - \delta}{\sqrt{2}\sigma}, \quad (18)$$

which gives us

$$\Rightarrow \frac{du}{d\delta} = -\frac{1}{\sqrt{2}\sigma} \Rightarrow d\delta = -\sqrt{2}\sigma du. \quad (19)$$

We now need to look at the limits. While the upper limit stays the same, we need to take a look at the lower limit

$$\delta_{\downarrow} = x = 2\delta_{crit} - \sqrt{2}u_{\downarrow}\sigma \Rightarrow u_{\downarrow} = \frac{2\delta_{crit} - x}{\sqrt{2}\sigma}, \quad (20)$$

then taking the limit of x we get

$$u_{\downarrow} = \lim_{x \rightarrow -\infty} \frac{2\delta_{crit} - x}{\sqrt{2}\sigma} = \infty. \quad (21)$$

We can then use the minus sign from (19) to turn the limits around and get

$$I_2 = \frac{1}{2} \int_{\nu/\sqrt{2}}^{\infty} \frac{2}{\sqrt{\pi}} e^{-u^2} du = \frac{1}{2} \operatorname{erfc}\left(\frac{\nu}{\sqrt{2}}\right) = \frac{1}{2} \left[1 - \operatorname{erf}\left(\frac{\nu}{\sqrt{2}}\right) \right] \quad (22)$$

We thus get the probability

$$P(> M) = 1 - I_1 + I_2 = 1 - \frac{1}{2} \left[1 + \operatorname{erf}\left(\frac{\nu}{\sqrt{2}}\right) \right] + \frac{1}{2} \left[1 - \operatorname{erf}\left(\frac{\nu}{\sqrt{2}}\right) \right] = 1 - \operatorname{erf}\left(\frac{\nu}{\sqrt{2}}\right) = 2P(\delta > \delta_{crit} | M). \quad (23)$$

And there is no need for the fudge factor!