FYS4411 - Project 2

The restricted Boltzmannn machine applied to the quantum many body problem

Daniel Heinesen¹, Gunnar Lange² & Aram Salihi²

¹Department of Theoretical Astrophysics, University of Oslo, N-0316 Oslo, Norway

²Department of Physics, University of Oslo, N-0316 Oslo, Norway

April 19, 2018

 $\begin{aligned} \mathbf{Abstract} \\ \mathbf{NICE} \ \mathbf{ABSTRACT} \end{aligned}$

Contents

1	Introduction	3
2	Theoretical model	3
3	Methods	3
4	Results	3
5	Discussion	3
6	Conclusion 6.1 Outlook	3
Aı	ppendices	4
\mathbf{A}	Finding the derivatives A.1 The local energy	4
	A 2. The derivatives with respect to the parameters	4

1 Introduction

COOL INTRO

2 Theoretical model

MUCH THEORY

3 Methods

MUCHO METHOD

4 Results

GROUNDBREAKING RESULTS

5 Discussion

EXCELLENT DISCUSSION

6 Conclusion

BRILLIANT CONCLUSION

6.1 Outlook

MAGNIFICIENT OUTLOOK

Appendices

A Finding the derivatives

In this section, we derive the expressions for the various derivatives stated in section **SECTION**. We begin with the local energy.

A.1 The local energy

The local energy is given by:

$$E_L = \frac{1}{\Psi} \hat{H} \Psi = \sum_{i=1}^{N} \left(-\frac{1}{2\Psi} \nabla_i^2 \Psi + \frac{1}{2} \omega^2 r_i^2 \right) + \sum_{i < j} \frac{1}{r_{ij}}$$
 (1)

Thus, we must compute:

$$\frac{1}{\Psi}\nabla_i^2\Psi\tag{2}$$

This may be rewritten as:

$$\frac{1}{\Psi}\nabla\left(\Psi\frac{1}{\Psi}\nabla\Psi\right) = \left(\frac{1}{\Psi}\nabla\Psi\right)^2 + \nabla\left(\frac{1}{\Psi}\nabla\Psi\right) = \left[\nabla\log\Psi\right]^2 + \nabla^2\log\Psi \tag{3}$$

The logarithm of our trial wavefunction is given by:

$$\log \Psi = -\log Z - \sum_{i=1}^{M} \left(\frac{(X_i - a_i)^2}{2\sigma^2} \right) + \sum_{j=1}^{N} \log \left(1 + \exp \left(b_j + \sum_{i=1}^{M} \frac{X_i w_{ij}}{\sigma^2} \right) \right)$$
(4)

The derivative with respect to one coordinate is now given by:

$$\frac{\partial \log \Psi}{\partial X_k} = \frac{(a_k - X_k)}{\sigma^2} + \sum_{j=1}^N \frac{w_{kj} \exp\left(b_j + \sum_{i=1}^M \frac{X_i w_{ij}}{\sigma^2}\right)}{\sigma^2 \left(1 + \exp\left(b_j + \sum_{i=1}^M \frac{X_i w_{ij}}{\sigma^2}\right)\right)}$$

$$= \frac{a_k - X_k}{\sigma^2} + \sum_{j=1}^N \frac{w_{kj}}{\sigma^2 \left(1 + \exp\left(-b_j - \sum_{i=1}^M \frac{X_i w_{ij}}{\sigma^2}\right)\right)}$$
(5)

Whereas the second derivative is:

$$\frac{\partial^2 \log \Psi}{\partial X_k^2} = -\frac{1}{\sigma^2} + \sum_{j=1}^N \frac{w_{kj}^2 \exp\left(-b_j - \sum_{i=1}^M \frac{X_i w_{ij}}{\sigma^2}\right)}{\sigma^4 \left(1 + \exp\left(-b_j - \sum_{i=1}^M \frac{X_i w_{ij}}{\sigma^2}\right)\right)^2}$$
(6)

The local energy can now be found by using equation 3, and inserting equation 5 and equation 6.

A.2 The derivatives with respect to the parameters

For our optimization method, we require:

$$\frac{1}{\Psi} \frac{\partial \Psi}{\partial \alpha_k} = \frac{\partial}{\partial \alpha_k} \log \Psi \tag{7}$$

Where α_k is any of the variational parameters a, b or w. These derivatives are given by:

$$\frac{\partial \log \Psi}{\partial a_k} = \frac{a_k - X_k}{\sigma^2} \tag{8}$$

$$\frac{\partial \log \Psi}{\partial b_k} = \frac{\exp\left(b_k + \sum_{i=1}^M \frac{X_i w_{ik}}{\sigma^2}\right)}{1 + \exp\left(b_k + \sum_{i=1}^M \frac{X_i w_{ik}}{\sigma^2}\right)} = \frac{1}{1 + \exp\left(-b_k - \sum_{i=1}^M \frac{X_i w_{ik}}{\sigma^2}\right)}$$
(9)

Finally, the derivative with respect to the weights, w_{kl} is given by:

$$\frac{\partial \log \Psi}{\partial w_{kl}} = \frac{X_k \exp\left(b_l + \sum_{i=1}^M \frac{X_i w_{il}}{\sigma^2}\right)}{\sigma^2 \left(1 + \exp\left(b_l + \sum_{i=1}^M \frac{X_i w_{il}}{\sigma^2}\right)\right)} = \frac{X_k}{\sigma^2 \left(1 + \exp\left(-b_l - \sum_{i=1}^M \frac{X_i w_{il}}{\sigma^2}\right)\right)}$$
(10)