

1 Kompleks analyse:

1.1 Komplekse tall:

$$z = x + iy, \quad \bar{z} = x - iy, \quad z = r(\cos \theta + i \sin \theta) = re^{i\theta} \quad (1)$$

$r = |z|$ er modulus, θ er argumentet.

$$z + \bar{z} = 2\operatorname{Re}(z) \quad z - \bar{z} = 2i\operatorname{Im}(z), \quad z\bar{z} = |z|^2 = r^2 = x^2 + y^2 \quad (2)$$

$$x = \frac{1}{2}(z + \bar{z}), \quad y = \frac{1}{2i}(z - \bar{z}) \quad (3)$$

$z - z_0| < R$ er alle z innenfor en radius R

1.2 Komplekse røtter:

$$z^{1/n} = \sqrt[n]{r}e^{i\theta/n} = \sqrt[n]{r}(\cos \theta/n + i \sin \theta/n) = \omega_0 \quad (4)$$

Dette gir bare 'the principal root', resten er gitt ved

$$\omega_k = \sqrt[n]{r}e^{i\frac{\theta+2\pi k}{n}} \quad (5)$$

1.3 Analytic functions:

Def: A function is analytic in a region of the complex plane if it has a (unique) derivative at every point in that region.

All analytic functions can be written in terms of $z = x + iy$ alone.

1.3.1 Cauchy-Riemann equation:

$$f(z) = u(x, y) + iv(x, y) \quad (6)$$

$$\frac{\partial u}{\partial x} = -\frac{\partial v}{\partial y}, \quad \frac{\partial u}{\partial y} = \frac{\partial v}{\partial x} \quad (7)$$

If this holds in a region, that f is analytic in this region, and vice versa.

- Regular points: $f(z)$ is analytic
- Singular point/singularities: A point where $f(z)$ is not analytic.
- Isolated singularity: a point where f is not analytic, but is a limit of points where f is analytic.

If f is analytic in some region, it has first order derivatives, then it also has derivatives of all orders in that region.

1.3.2 Harmonic Functions:

If $f = u + iv$ is analytic in a region, then u and v are harmonic:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0 \quad \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = 0 \quad (8)$$

If u is harmonic, one can find a v such that $f = u + iv$ is analytic. v is the harmonic conjugate of u .

1.4 Contour integrals:

$$\int_{\Gamma} f(z)dz = \lim_{z \rightarrow \infty} \sum_{k=1}^{\infty} f(c_k) \Delta z_k \quad (9)$$

For a generalized curve with parametrization $z(t)$:

$$\int_{\gamma} f(t)dz = \int_a^b f(z(t))z'(t)dt \quad (10)$$

1.4.1 An important integral:

$C - r = |z - z_0| = r$:

$$I = \int_{C_r} (z - z_0)^n dz = \begin{cases} 0 & n \neq -1 \\ 2\pi i & n = -1 \end{cases} \quad (11)$$

1.4.2 Upper bound estimate:

Generalized triangle inequality:

$$\left| \sum_k z_k \right| \leq \sum_k |z_k| \quad (12)$$

Applied to Riemann sum we get the upper bound estimate:

$$\left| \int_{\gamma} f(z)dz \right| \leq ML \quad (13)$$

Where $M = \max |f(z)|$ and L is the length of the curve.

1.4.3 Path:

If f is continuous everywhere in D , then contour integrals are independent of paths, and any loop integral is zero. One can also deform a contour without crossing any singularities and get that:

$$\int_{\Gamma_1} f(z)dz = \int_{\Gamma_2} f(z)dz \quad (14)$$

1.4.4 Cauchy Theorem:

If f is analytic in a simply connected domain D with no singularities, and Γ is any closed contour in D , then

$$\int_{\Gamma} f(z)dz = 0 \quad (15)$$

1.4.5 Cauchy's integral formula:

Let Γ be a simple, closed, positively oriented contour. Assume f is analytic in some simply connected domain D containing Γ , and some z_0 is inside Γ . Then:

$$f(z_0) = \frac{1}{2\pi i} \int_{\Gamma} \frac{f(z)}{z - z_0} dz \quad (16)$$

1.4.6 Generalized Cauchy integral formula:

$$f^{(n)}(z) = \frac{n!}{2\pi i} \oint_{\Gamma} \frac{f(w)}{(z - w)^{n+1}} dw \quad (17)$$

1.4.7 Cauchy inequality:

Let f be analytic on and inside a circle(C_r) of radius R , centered at z_0 . If $|f(z)| \leq M$ for some z on C_r , then the derivatives satisfy:

$$|f^{(n)}(z_0)| \leq \frac{n!M}{R^n} \quad (18)$$

This gives Liouville's theorem: A function which is analytic and bounded in the entire complex plane, is constant.

1.5 Taylor and Laurent Series:

1.5.1 Taylor:

$$f(z) = f(z_0) + f'(z_0)(z - z_0) + \frac{1}{2!}f''(z_0)(z - z_0)^2 + \dots = \sum_n \frac{f^{(n)}(z_0)}{n!}(z - z_0)^n \quad (19)$$

If f is analytic in the disk $|z - z_0| < R$ then the above Taylor series converges in that disk.[i.e. the disk touching the nearest singularity]

If f is analytic at z_0 , then the Taylor series for df/dz can be obtained by termwise differentiation.

1.5.2 Laurent:

Let f be analytic in the annulus $r < |z - z_0| < R$. Then f can be expanded there as the sum of two series:

$$f(z) = \sum_{k=0}^{\infty} a_k(z - z_0)^k + \sum_{k=1}^{\infty} b_k(z - z_0)^{-k} \quad (20)$$

With

$$a_n = \frac{1}{2\pi i} \oint_{\Gamma} \frac{f(z)}{(z - z_0)^{n+1}} dz \quad b_n = \frac{1}{2\pi i} \oint_{\Gamma} \frac{f(z)}{(z - z_0)^{-n+1}} dz \quad (21)$$

Def: The coefficient b_1 of the $1/(z - z_0)$ term is the residue of $f(z)$ at $z = z_0$.

Laurent series are unique. So to find them we can use

$$\frac{1}{1 - \omega} = \sum_{n=0}^{\infty} \omega^n, \text{ when } |\omega| < 1 \quad (22)$$

1.6 Zeros:

- A zero of a function is a point where f is analytic and $f(z_0) = 0$
- A zero of order m : $f(z_0) = f'(z_0) = \dots = f^{m-1}(z_0) = 0$, $f^m(z_0) \neq 0$
- Can be factorized as: $f(z) = (z - z_0)^m \cdot g(z)$, where $g(z)$ is analytic and $g(z_0) \neq 0$

1.7 Isolated singularities:

Let f have a Laurent series, then we can have:

1.7.1 Removable Singularity/Regular point:

If all $b_n = 0$ at z_0 . $f(z)$ has a limit $z \rightarrow z_0$ and we can be redefined such that f is analytic at z_0 .

1.7.2 Essential Singularity:

Infinitely many b-terms at z_0

1.7.3 Pole of order m

Order m is the highest exponent of the $1/(z - z_0)$ terms.

$$f(z) = \frac{b_m}{(z - z_0)^m} + \dots + \frac{b_1}{z - z_0} + a_0 + a_1(z - z_0) + \dots \quad (23)$$

$f(z)$ can be written as $\frac{g(z)}{(z - z_0)^m}$. A pole of order 1 ($m = 1$) is a Simple pole.

1.8 Residue Theory:

1.8.1 Residue Theorem:

If Γ is a simple, closed, positively oriented contour, and f is analytic on and inside Γ except at the points z_0, z_1, \dots, z_n inside Γ , then

$$\oint_{\Gamma} f(z) dz = 2\pi i \sum_{k=0}^n \text{Res}(f, z_k) \quad (24)$$

1.8.2 Determining the residues:

1: Read off b_1 from the Laurent series.

2: Simple poles:

$$\text{Res}(z_0) = b_1 = \lim_{z \rightarrow z_0} (z - z_0) f(z) \quad (25)$$

Finite answer only if the pole is of first order.

$$f(z) = \frac{P(z)}{Q(z)} \Rightarrow \text{Res}(z_0) = \frac{P(z_0)}{Q'(z_0)} \quad (26)$$

3: Multiple poles: If f has a pole of order m at z_0 , then

$$\text{Res}(z_0) = \lim_{z \rightarrow z_0} \left[\frac{1}{(m-1)!} \frac{d^{m-1}}{dz^{m-1}} ((z - z_0)^m f(z)) \right] \quad (27)$$

Ok to overshoot with value if m .

1.9 Applications to Real Integrals:

1.9.1 Type 1:

Rational and finite functions of $\sin \theta$ and $\cos \theta$ over the interval $[0, 2\pi]$. Use:

- $z = e^{i\theta}$, $d\theta = dz/iz$
- $\cos \theta = \frac{1}{2}(e^{i\theta} + e^{-i\theta}) = \frac{1}{2}(z + 1/z)$, $\sin \theta = \frac{1}{2i}(e^{i\theta} - e^{-i\theta}) = \frac{1}{2i}(z - 1/z)$
- then use residue theorem.

1.9.2 Type 2a

Integrals of rational functions from $-\infty$ to ∞

$$I = \int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} dx \quad (28)$$

- Make a contour γ_ρ from $-\rho$ to ρ
- Add a second contour from ρ via a the complex plane (half circle in the upper part of the complex plane) back to $-\rho$, C_ρ .
- use the residue theorem. Remember that the singularities have to be in the upper part
- Show that the contribution from C_ρ vanishes as $\rho \rightarrow \infty$

1.9.3 Type 2b:

$$I = \int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} \cos(mx) dx, \quad I = \int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} \sin(mx) dx \quad (29)$$

Alt 1 (Always safe):

use $\cos(mx) = \frac{1}{2}(e^{imx} + e^{-imx})$

$$I = \frac{1}{2} \int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} e^{imx} dx + \int_{-\infty}^{\infty} \frac{P(x)}{Q(x)} e^{-imx} dx \quad (30)$$

For the first term, use a closed contour in the upper part of the complex plane, for the second term use one in the lower half.

Alt 2: safe as long as P/Q is real. Note that $\cos(mx) = \text{Re}(e^{imx})$ (and $\sin(mx) = \text{Im}(e^{imx})$). We can then use $\cos(mx) \rightarrow e^{imx}$ and then take the real part at the end (or the imaginary if we have $\sin(mx)$)

1.9.4 Jordan's lemma:

If $m > 0$ (real) and P/Q is the quotient of two polynomials such that $\deg(Q) \geq \deg(P) + 1$, then

$$\lim_{\rho \rightarrow \infty} \int_{C_\rho^+} \frac{P}{Q} e^{imz} dz = 0 \quad (31)$$

Where C_ρ^+ is the contour in the upper plane. Same holds for $m < 0$ but with C_ρ^- in the lower plane.

1.9.5 Type 3:

Singularities on the real plane. We get Principal Values.

$$PV \int_a^b f(x) dx = \lim_{r \rightarrow 0} \int_a^{c-r} f(x) dx + \int_{c+r}^b f(x) dx \quad (32)$$

Where c is a singularity. If the singularities are simple poles the integral can be calculated with the residue theorem.

$$PV \int_{-\infty}^{\infty} f(x) dx = 2\pi i \sum_k \text{Res}(z_k) + \pi i \sum_j \text{Res}(z_j) \quad (33)$$

Where z_k are singularities on the upper half plane, and z_j are singularities on the real line.

2 Tensor:

2.1 Cartesian Tensor

Transform properly under rotation of Cartesian coordinate system.

$$e'_i \cdot e_j = \cos \theta_{ij} \equiv A_{ij} \quad (34)$$

2.1.1 Transformation of a position vector

$$\vec{r} = x_i e_i = x'_j e'_j \quad (35)$$

$$x'_i = \vec{r} \cdot e'_i = e'_i \cdot e_j x_j = A_{ij} x_j \quad (36)$$

Reverse:

$$x_i = e_i \cdot \vec{r} = e_i \cdot e'_j x'_j = A'_{ji} x'_j = A_{ij}^T x'_j \quad (37)$$

$$A^{-1} = A^T \quad (38)$$

2.1.2 Cartesian vector/tensor

$$v' = Av, \quad T'_{kl} = A_{ki}A_{lj}T_{ij} \quad (39)$$

$$\vec{j} = \sigma \vec{E} \text{ Ohm's law} \quad (40)$$

2.1.3 Inertia Tensor

Rigid body rotation around fixed axis:

$$\vec{L} = I\vec{\omega} \quad (41)$$

Rotation around a point. I is a rank 2 tensor

$$L_i = I_{ij}\omega_j \quad (42)$$

Determined from

$$\vec{L} = \sum_k m_k r_k \times (\omega \times r_k) \quad (43)$$

Uniform mass density: sum goes to integral. I is symmetric: can find coord. system in which I is diagonal. Eigenvectors (axes of new coord. system): Principal axes of inertia.

2.2 Levi-Civita and Kronicker Delta

$$\delta_{ij} = \begin{cases} 1 & i = j \\ 0 & \end{cases} \quad (44)$$

$$\epsilon_{ijk} = \begin{cases} 1 & \text{even permutation} \\ -1 & \text{odd permutation} \\ 0 & \end{cases} \quad (45)$$

$$\epsilon_{ijk}\epsilon_{imn} = \delta_{jm}\delta_{kn} - \delta_{jn}\delta_{km} \quad (46)$$

Inner product: $\vec{u} \cdot \vec{c} = u_i v_j \delta_{ij} = u_i u_i$

Cross Product: $\vec{A} \times \vec{B} = \epsilon_{ijk} A_j B_k$

Curl: $(\nabla \times \vec{V})_i = \epsilon_{ijk} \frac{\partial}{\partial x_j} V_k$

3 Calculus of Variation:

Minimize:

$$I = \int_{x_1}^{x_2} F(x, y, y') dx \quad (47)$$

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0 \quad (48)$$

Can simplify Euler-Lagrange by change of variables:

$$x' = \frac{1}{y'}, \quad dx = x' dy \text{ or } dy = y' dx \quad (49)$$

3.1 Optics – Fermat's Principle:

$$P = \int n ds, \quad ds = \sqrt{dx^2 + dy^2} \quad (50)$$

4 ODE:

4.1 Linear First order DE:

$$y' + P(x)y = Q(x) \quad dy + [Py - Q]dx = 0 \quad (51)$$

4.1.1 Exact DE's

$$du = \frac{\partial u}{\partial x}dx + \frac{\partial u}{\partial y}dy, \quad M(x, y)dx + N(x, y)dy = 0 \quad (52)$$

The last DE is called exact if the LHS is a total differential,

$$M = \frac{\partial u}{\partial x}; \quad N = \frac{\partial u}{\partial y} \quad (53)$$

This gives:

$$du = 0 \Leftrightarrow u = \text{const} \quad (54)$$

4.1.2 Integrating factor:

$$y(x) = \frac{1}{\mu(x)} \left[\int \mu(x)Q(x)dx + C \right], \quad \mu(x) = e^{\int P(x)dx} \quad (55)$$

If $Q(x) = 0$, the homogeneous case:

$$y(x) = \frac{C}{\mu(x)} = Ce^{-\int P(x)dx} \quad (56)$$

4.2 Ordinary 2nd order DE:

$$y'' + P(x)y' + Q(x)y = R(x), \quad y(x) = y_h + y_p = c_1y_1 + c_2y_2 + y_p \quad (57)$$

If y_1 and y_2 are linearly independent, it must hold that the Wronskian determinant

$$W(x_0) = \begin{vmatrix} y_1(x_0) & y_2(x_0) \\ y_1'(x_0) & y_2'(x_0) \end{vmatrix} \neq 0 \quad \forall x_0 \quad (58)$$

4.3 Homogeneous equations ($R(x) = 0$)

4.3.1 Variations of Constants:

if y_1 is a solution, an other linearly independent solution can be found as

$$y_2(x) = c(x)y_1(x) \quad (59)$$

Where $c(x)$ can be found from the DE.

4.3.2 Homogeneous DE's with constant coefficients

$$y'' + ay' + by = 0, \quad \lambda^2 + a\lambda + b = 0 \quad \lambda_{\pm} = \frac{1}{2}[-a \pm \sqrt{a^2 - 4b}] \quad (60)$$

Case 1: $\lambda_- \neq \lambda_+$, both real ($a^2 - 4b > 0$), gives two linearly independent solutions:

$$y(x) = C_1e^{\lambda_+x} + C_2e^{\lambda_-x} \quad (61)$$

Case 2: Double root $\lambda_+ = \lambda_- \equiv \lambda = \frac{-a}{2}$

$$y(x) = (Ax + B)e^{\lambda x} \quad (62)$$

Case 3: Complex root ($a^2 - 4b < 0$)

$$\lambda_{\pm} = -\frac{a}{2} \pm i\sqrt{4b - a^2} = -\frac{a}{2} \pm i\omega \quad (63)$$

$$y(x) = Ae^{-a/2+i\omega} + Be^{-a/2-i\omega} = e^{-ax/2} (A \cos(\omega x) + B \sin(\omega x)) \quad (64)$$

$$= e^{-ax/2} (\tilde{A}e^{i\omega x} + \tilde{B}e^{-i\omega x}) = ke^{-ax/2} \sin(\omega x + \phi) \quad (65)$$

4.3.3 Euler-Cauchy:

$$x^2 y'' + a_1 x y' + a_0 y = 0, \text{ or } y'' + \frac{a_1}{x} y' + \frac{a_0}{x^2} y = 0 \quad (66)$$

Use:

$$x = e^z, \quad z = \ln x, \quad dx = e^z dz = x dz \quad (67)$$

For $x > 0$

$$x = -|x| = -e^z, \quad z = \ln |x|, \quad dx = -e^z dz = x dz \quad (68)$$

For $x < 0$. We then get

$$\frac{d^2 y(z)}{dz^2} + (a_1 - 1) \frac{dy(z)}{dz} + a_0 y(z) = 0 \quad (69)$$

Solve then transform back to $y(x)$. No solution at $x = 0$, and different coefficients for cases $x < 0$ and $x > 0$!

4.4 Inhomogeneous DE

$$y(x) = y_h(x) + y_p(x) \quad (70)$$

4.4.1 Methods of Undetermined Coefficients:

$$y'' + ay' + by = R(x) \quad (71)$$

Where $R(x)$ is simple. If $R(x)$ is a sum of simple functions, then y_p is also a sum.

Case 1: $R(x) = Ae^{kx}$. Assume α, β are the roots of the characteristic equation of the homogeneous solution.

- a) If $k \neq \alpha, \beta$. Try $y_p = Be^{kx}$ (Find B from the DE)
- b) $k = \alpha$ or $k = \beta$. Try $y_p = Cxe^{kx}$
- c) $k = \alpha = \beta$, so $y_h = (Ax + B)e^{kx}$, try: $y_p = Dx^2e^{kx}$

Case 2: $R(x) = A \sin kx$ or $\cos kx$. Has the form $y_p = B \cos kx + C \sin kx$. Efficient: Solve or $\tilde{R}(x) = e^{ikx}$ and take Re or Im at the end.

Case 3: $R(x) = e^{kx} P_n(x)$. P_n and Q_n are here polynomials of order n .

- a) If $k \neq \alpha, \beta$. Try $y_p = Q_n(x)e^{kx}$ (Find B from the DE)
- b) $k = \alpha$ or $k = \beta$. Try $y_p = xQ_n(x)e^{kx}$
- c) $k = \alpha = \beta$, so $y_h = (Ax + B)e^{kx}$, try: $y_p = x^2Q_n(x)e^{kx}$

4.4.2 More General: y_p from factorization:

$$y'' + P(x)y' + Q(x)y = R(x) \quad (72)$$

Assume $u(x)$ is a known solution of the homogeneous DE. Make the ansatz $y_p = u(x)v(x)$. $v(x)$ is found from inserting into the DE. If we define $v' = w$ we'll get

$$w' + \left(\frac{2u}{u'} + P \right) w = \frac{R}{u} \quad (73)$$

Solve by integrating factor. And then get

$$v = \int w dx \quad (74)$$

4.4.3 Variation of Parameters:

NB! Write DE on the standard form $y'' + Py' + Q = R$. This works if the homogeneous solutions are fully known!

$$y_p(x) = -y_1 \int \frac{y_2 R}{W} dx + y_2 \int \frac{y_1 R}{W} dx \quad (75)$$

Where W is the Wronskian

$$W(x_0) = \begin{vmatrix} y_1(x_0) & y_2(x_0) \\ y_1'(x_0) & y_2'(x_0) \end{vmatrix} \quad (76)$$

4.5 Fröbenius method:

Any DE of the form

$$y'' + \frac{B(x)}{x}y' + \frac{C(x)}{x^2}y = 0 \quad (77)$$

Where B and C are analytic in the singular point $x = 0$ (if the whole equation is analytic here, we can use power series).

$$y(x) = x^s \sum_n a_n x^n \quad (78)$$

4.5.1 Indicial equation:

Write $x^2 y'' + x b(x) y' + c(x) y = 0$, with $b(x) = b_0 + b_1 x^1 + \dots$, $c(x) = c_0 + c_1 x^1 + \dots$ and $y(x) = a_0 x^s + a_1 x^{s+1} + \dots$. If we insert these series in to the DE, and equate the coefficients for the lowest possible power x^s we get

$$s(s-1) + b_0 s_0 + c_0 = 0 \quad (79)$$

a) Two distinct roots s_1, s_2 and $s_1 - s_2 \neq \text{integer}$: Two linearly independent solutions.

$$y_i = x^{s_i} \sum_n a_n x^n \quad (80)$$

b) Two distinct roots s_1, s_2 but $s_1 - s_2 = \text{integer}$. Choose $s_1 > s_2$

- Often s_2 gives a complete solution. Always try the smallest root first.
- Sometimes s_2 does not give a solution, s_1 always does: Find y_1 , then $y_2 = c(x)y_1$ from variations of constants.
- Double root $s_1 = s_2 = s$. Then find y_1 from s , then $y_2 = c(x)y_1$

4.5.2 Hermite Eq:

From harmonic oscillator: $-\psi'' + x^2 \psi = (2n+1)\psi$. Use factorization $\psi(x) = e^{-x^2/2} y(x)$. We then get the Hermite DE:

$$y'' - 2xy' + ny = 0 \quad (81)$$

Solution

$$a_{k+2} = -\frac{2(n-k)}{(k+1)(k+2)} a_k, \quad H_n(x) = (-1)^n e^{x^2} \frac{d^2}{dx^2} e^{-x^2} \quad (82)$$

4.6 Greens functions:

$$Dy(x) = R(x), \quad D = \frac{d^2}{dx^2} + P(x)\frac{d}{dx} + Q(x) \quad (83)$$

$$DG(x, z) = \delta(x - z), \quad y(a) = y(b) = 0 \quad (84)$$

$$y(x) = \int_a^b G(x, z)R(z)dz \quad (85)$$

G is continuous at $x = z$, but its derivatives has a discontinuity of 1

$$G(z + \epsilon, z) - G(x - \epsilon, z) = 0, \quad G'(z + \epsilon, z) - G'(x - \epsilon, z) = 1 \quad (86)$$

5 Fourier Series:

5.1 Some orthogonality Relations:

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} \cos(nx)dx = \frac{1}{2\pi} \int_{-\pi}^{\pi} \sin(nx)dx = 0 \quad (87)$$

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} \cos(mx) \sin(nx)dx = 0 \quad (88)$$

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} \cos(nx) \cos(mx)dx = \begin{cases} 0 & n \neq m \\ 1/2 & n = m \neq 0 \\ 1 & n = m = 0 \end{cases} \quad (89)$$

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} \sin(nx) \sin(mx)dx = \begin{cases} 0 & n \neq m \\ 1/2 & n = m \neq 0 \\ 0 & n = m = 0 \end{cases} \quad (90)$$

$$\int_{-\pi}^{\pi} e^{inx} e^{-imx} dx = 2\pi \delta_{nm} \quad (91)$$

5.2 Dirichlet Conditions:

Sufficient conditions for the Fourier series to converge:

- A finite number of maxima and minima in the basic interval
- A finite number of finite discontinuities [bonded]
- At point y_0 where $f(x)$ has a discontinuity, the Fourier series converges at the midpoint.
- The Fourier Series may be integrated term by term
- If $f'(x)$ satisfies the Dirichlet conditions, the Fourier series may be differentiated term by term

5.3 Interval $[0, 2\pi]$:

$$f(x) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx) \quad (92)$$

$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(nx)dx, \quad b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx)dx \quad (93)$$

5.4 Complex Series:

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}, \quad c_k = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-ikx} dx \quad (94)$$

Can find a sin-cos series from the complex, and vice versa.

5.5 Other intervals:

We can simply change the interval to another basic interval of length $2L$:

$$x \rightarrow \frac{\pi x}{L}, \quad \frac{1}{\pi} \int_{-\pi}^{\pi} \rightarrow \frac{1}{L} \int_{-L}^L \quad (95)$$

5.6 Even and Odd Functions:

Even: $f(-x) = f(x)$, ex: $(\cos x)$

Odd: $f(-x) = -f(x)$, ex: $(\sin x)$

$even \cdot even = odd \cdot odd = even$, $even \cdot odd = odd$

$$\int_{-L}^L f(x) dx = \begin{cases} 0 & \text{f is odd} \\ 2 \int_0^L f(x) dx & \text{f is even} \end{cases} \quad (96)$$

5.6.1 Fourier of odd Function, sine-series:

$$f(x) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right), \quad b_n = \frac{2}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx \quad (97)$$

5.6.2 Fourier of even Function, cos-series:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right), \quad a_n = \frac{2}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx \quad (98)$$

A given function may be represented by several different Fourier Series depending on the physics/context. Given a function defined only on half the interval $[0, L]$. We may either define an **even** extension of $f(x)$ to the period $2L$ by only using the cosine-series; or we may define an **odd** extension by using the sine-series.

5.7 Parseval's Theorem:

$$\frac{1}{2L} \int_{-L}^L |f(x)|^2 dx = \sum_{n=-\infty}^{\infty} |c_n|^2 = \left(\frac{1}{2}a_0\right)^2 + \frac{1}{2} \sum_{n=1}^{\infty} a_n^2 + \frac{1}{2} \sum_{n=1}^{\infty} b_n^2 \quad (99)$$

$$\int_{-\infty}^{\infty} |f(x)|^2 = \int_{-\infty}^{\infty} |F(k)|^2 \quad (100)$$

6 Fourier Transforms:

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx = \mathcal{F}[f(x)] \quad (101)$$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k) e^{ikx} dk \quad (102)$$

6.1 Fourier Integral Theorem:

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{ikx} \underbrace{\int_{-\infty}^{\infty} f(\tilde{x}) e^{-ik\tilde{x}} d\tilde{x}}_{\sqrt{2\pi}F(k)} dk \quad (103)$$

Is OK if $f(x)$ satisfied the Dirichlet conditions on every finite interval, and $\int_{-\infty}^{\infty} |f(x)| dx$ is finite.

6.2 Fourier Transforms of Derivatives:

$$\mathcal{F}[f^{(n)}(x)] = (ik)^n \mathcal{F}[f(x)] \quad (104)$$

6.2.1 With DE:

$$y'' + ay' + b = f(x), \quad Y(k) = \mathcal{F}[y(x)] = \frac{F(k)}{-k^2 + aik + b} \quad (105)$$

$$y(x) = \mathcal{F}^{-1}[Y(k)] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{ikx} Y(k) dk \quad (106)$$