#### Lecture slides

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#### Lecture 1

- Introduction to MA42002
- Conservation equations
- Examples of spatially homogeneous models

### Conservation equations

$$\begin{pmatrix} \text{rate of change} \\ \text{in the population density} \end{pmatrix} = \left( \text{spatial movement} \right)$$
 
$$+ \begin{pmatrix} \text{birth, growth, death,} \\ \text{production or degradation} \\ \text{due to chemical reactions} \end{pmatrix}$$

# Spatially homogeneous models (MA32009

revision)

# Example problem - bacteria in a dish

$$N(t + \Delta t) = N(t) + KN(t)\Delta t.$$

## A model for cell growth under nutrient depletion

$$\begin{split} \frac{dN}{dt} &= K(c)N = \kappa cN, \\ \frac{dc}{dt} &= -\alpha \frac{dN}{dt} = -\alpha \kappa cN, \end{split} \tag{1}$$

# Leading to the logistic growth equation

The last equation can be rewritten as

$$\frac{dN}{dt} = \rho N \left(1 - \frac{N}{B}\right) \qquad N(0) = N_0, \tag{2}$$



#### Exercise

Consider a well mixed bio reactor.

A biologist cultures an initial cell population of size  $N_0$  in the bioreactor for 72 h.

Cells undergo division with a period of 14 h.

Each cell produces a non-degradable waste product, W, at rate  $k_1$ .

When total waste levels exceed a threshold,  $W^{st}$ , cell division stops. Otherwise the cell population grows exponentially.

How many cells are there at the end of the experiment?

## Model development

#### Model checklist

- 1. Variables (dependent, indepedent ?)
- 2. Schematic diagram what processes are being modelled?
- 3. Governing equations?
- 4. Define model parameters?
- 5. Initial conditions?

### Exercise solution

#### Recap

- ls course layout clear
- Introduction to conservation equation
- Deriving spatially homogeneous models

#### Lecture 2

- Continue example
- ► Introduce SIR model
- Introduce an activator inhibitor model
- Derive a conservation equation

#### Exercise

Consider a well mixed bio reactor.

A biologist cultures an initial cell population of size  ${\cal N}_0$  in the bioreactor for 72 h.

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Each cell produces a non-degradable waste product,  $\boldsymbol{W}$ , at rate  $k_1$ .

When total waste levels exceed a threshold,  $W^{\ast}$ , cell division stops. Otherwise the cell population grows exponentially.

How many cells are there at the end of the experiment?

# The SIR model (used in Chapter 7)

Consider the SIR model equations:

$$\begin{split} \frac{dS}{dt} &= -rIS, \\ \frac{dI}{dt} &= rIS - aI, \\ \frac{dR}{dt} &= aI. \end{split}$$

What are the variables? What are the parameters?

Identify an expression for the reproduction number,  $R_{\rm 0}.$ 

Hence explain why the condition  $R_0 < 1$  is necessary to avoid an epidemic?

## SIR model Calculations

$$\begin{split} \frac{dS}{dt} &= -rIS, \\ \frac{dI}{dt} &= rIS - aI, \\ \frac{dR}{dt} &= aI. \end{split}$$

# An activator inhibitor model (used in Chapter 6)

Assume that species A is produced at constant rate  $k_1$  and degrades at rate  $k_2. \label{eq:k2}$ 

Assume that B is produced at a constant rate,  $k_4$ .

Consider the reaction schematic

$$2A + B \rightarrow 3A$$
,

with reaction rate  $k_3$ .

Write down governing ODEs.

#### Activator-inhibitor model

Consider the ODEs

$$\begin{aligned} \frac{da}{dt} &= k_1 - k_2 a + k_3 a^2 b, \\ \frac{db}{dt} &= k_4 - k_3 a^2 b, \end{aligned}$$

Identify the steady state of the ODEs. How would you compute linear stability of the steady state?

### Recap

- ▶ Introduced SIR and activator-inhibitor models
- Computed steady states and stability analysis

# Lecture 3 Spatiotemporal models

- Derive conservation PDEs
- Consider different models of fluxes

## Spatiotemporal models - derivation

Consider a spatial domain V. A conservation equation can be written either in terms of the mass or number of particles of a species as follows:

$$\begin{pmatrix} \text{rate of change of} \\ \text{number of particles} \\ \text{per unit time} \end{pmatrix} = \begin{pmatrix} \text{rate of entry of} \\ \text{particles into } V \\ \text{per unit time} \end{pmatrix} - \begin{pmatrix} \text{rate of exit of} \\ \text{particles from } V \\ \text{per unit time} \end{pmatrix} \\ + \begin{pmatrix} \text{rate of degradation} \\ \text{or creation of particles} \\ \text{in } V \text{ per unit time} \end{pmatrix}$$

### Deriving a conservation equation in 1D

$$\frac{\partial}{\partial t} \int_{x}^{x+\Delta x} c(\tilde{x}, t) A d\tilde{x} = J(x, t) A - J(x + \Delta x, t) A + \int_{x}^{x+\Delta x} f(\tilde{x}, t, c(\tilde{x}, t)) A d\tilde{x}.$$
(3)

#### A conservation PDE in 1D

$$\frac{\partial}{\partial t}c(x,t) = -\frac{\partial}{\partial x}J(x,t) + f(x,t,c(x,t)). \tag{4}$$

### Generalising to $\mathbb{R}^n$

$$\frac{\partial}{\partial t} \int_{V} c(x,t) \, dx = - \int_{S} J(x,t) \cdot \mathbf{n} \, d\sigma + \int_{V} f(x,t,c) dx.$$

### Fluxes - Fickian diffusion

$$\mathbf{J} = -D\nabla c,\tag{5}$$

### Fluxes - Nonlinear diffusion

$$D = D(c),$$

D = D(c), e.g.  $D(c) = D_0 c^m,$   $D_0 > 0,$ 

Hence

$$J = -D(c)\nabla c$$

# Fluxes - Convection/advection

$$\mathbf{J} = \mathbf{v}c,\tag{6}$$

#### Fluxes - Taxis

$$\mathbf{J} = \chi(a)c\nabla a,$$

# Domain of definition of the problem

#### Lecture 4

- ▶ Boundary and initial conditions
- Nondimensionalisation
- ▶ Model formulation

# Boundary conditions

- Dirichlet
- Neumann
- Robin

# Initial conditions

# Formulating a model

#### Lecture 5

- Introduce a linear reaction diffusion model
- Diffusion

### Linear reaction diffusion equation

$$\frac{\partial c}{\partial t} = D\nabla^2 c + f(c), \quad c \equiv c(\mathbf{x}, t), \quad \mathbf{x} \in \mathbb{R}^n, \ t > 0.$$

so in 1D Cartesian coordinates

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2} + f(c), \quad x \in \mathbb{R}, \ t > 0.$$

# 1D diffusion equation with delta IC

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2}, \quad x \in \mathbb{R}, \ t > 0. \tag{7}$$

$$c(x_0,0)=\delta_0(x) \qquad x\in\mathbb{R}, \tag{8}$$

where  $\delta_0$  is a  $\it Dirac$  delta distribution (Dirac measure) satisfying

$$\int_{-\infty}^{+\infty} \delta_0(x) = 1 \quad \text{ and } \quad \int_{-\infty}^{+\infty} f(x) \delta_0(x) = f(0), \text{ for continuous } f.$$

#### Numerical solution

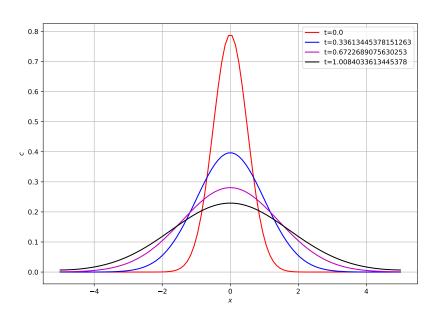


Figure 1: Numerical solution of diffusion equation.

#### An exact solution computed using a similarity variable

Consider the diffusion Equation 7 with initial condition Equation 8. Introduce the similarity variable

$$\eta = \frac{x}{\sqrt{Dt}}$$

and look for solution of the form

$$c(x,t) = \frac{1}{\sqrt{Dt}}F(\eta).$$

Hence it can be shown that the explicit (analytic) solution is given by

$$c(x,t) = \frac{1}{\sqrt{4\pi Dt}} \exp\left(-\frac{x^2}{4Dt}\right). \tag{9}$$

#### The 1D diffusion equation for arbitrary initial condition

For a general initial condition  $c(x,0)=c_0(x)$  for  $x\in\mathbb{R}$ :

$$c(x,t) = \int_{-\infty}^{+\infty} \frac{c_0(y)}{\sqrt{4\pi Dt}} \exp\left(-\frac{(x-y)^2}{4Dt}\right) dy.$$

### Key properties of the (linear) diffusion equation (heat equation)

- ▶ The solution is infinitely smooth.
- The solution c(x,t) stays positive for all t>0 and  $x\in\mathbb{R}$  if c(x,0)>0 for  $x\in\mathbb{R}$ .
- The solution "propagates' with infinite speed i.e. for any t>0, the solution is everywhere in  $\mathbb{R}$ .
- If we change the initial data c(x,0) (continuously) then the solution also changes (continuously).

#### Diffusive transit time

$$D\frac{d^2c}{dx^2} = 0 \quad \text{ in } (0,L), \quad c(0) = C_0, \, c(L) = 0.$$

#### Diffusion as a description of random walk

Suppose that the probability of a particle hopping distance  $\Delta x$  to the right in time  $\Delta t$  is

$$\lambda_R \Delta t$$
.

Similarly, the probability of hopping a distance  $\Delta x$  to the left is

$$\lambda_L \Delta t$$
.

#### Numerical simulation

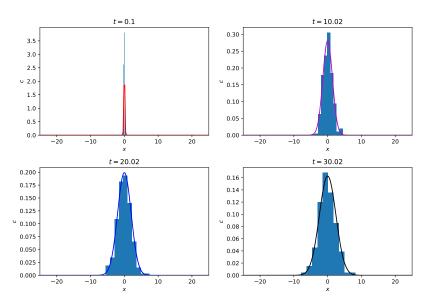


Figure 3: Numerical implementation of random walk

#### Derivation

Let c(x,t) represent the particle density at spatial location x and time t.

A conservation equation for c is given by

$$\begin{split} c(x,t+\Delta t) &= c(x,t) + \lambda_R \Delta t c(x-\Delta x,t) - \lambda_R \Delta t c(x,t) + \\ \lambda_L \Delta t c(x+\Delta x,t) - \lambda_L \Delta t c(x,t). \end{split}$$

#### Lecture 6

- Random walk as a model for the diffusion equation
- Linear reaction diffusion

#### Recap from last week

A conservation equation for c is given by

$$\begin{split} c(x,t+\Delta t) &= c(x,t) + \lambda_R \Delta t c(x-\Delta x,t) - \lambda_R \Delta t c(x,t) + \\ & \lambda_L \Delta t c(x+\Delta x,t) - \lambda_L \Delta t c(x,t). \end{split}$$

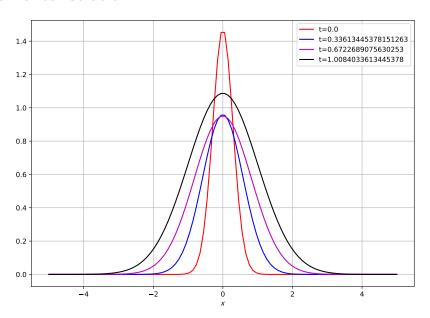
#### Linear reaction term

$$\frac{\partial c}{\partial t} = D \frac{\partial^2 c}{\partial x^2} + \rho c, \quad x \in \mathbb{R}, \ t > 0, \tag{10}$$

where  $\rho \in \mathbb{R}$  is a constant. with initial condition

$$u(x,0) = M\delta_0(x), \quad x \in \mathbb{R}. \tag{11}$$

#### Numerical solution



#### Muskrat invasion dynamics

$$\frac{\partial u}{\partial t} = D\left(\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2}\right) + \rho u, \quad \mathbf{x} = (x_1, x_2) \in \mathbb{R}^2, \ t > 0,$$

with initial condition

$$u(\mathbf{x},0) = M\delta_0(\mathbf{x}), \quad \mathbf{x} \in \mathbb{R}^2.$$
 (12)

$$u_1(\mathbf{x},t) = \frac{M}{4\pi Dt} \exp\left(\rho t - \frac{r_1^2}{4Dt}\right).$$

#### Lecture 7

#### Travelling waves

Travelling wave

A travelling wave is a solution of a PDE that has a constant profile (shape) and a constant propagation speed.

#### Fisher's equation

$$\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} + \rho u (1 - \frac{u}{K}), \qquad x \in \mathbb{R}, \ t > 0$$

with initial condition

$$u(x,0) = u_0(x).$$
 (13)

#### Nondimensional form

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + u(1-u), \qquad x \in \mathbb{R}, \ t > 0$$

with initial condition

$$u(x,0) = u_0(x).$$
 (14)

#### Numerical solution

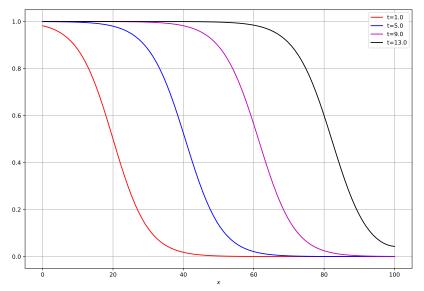


Figure 4: Numerical solution of Fisher's equation.

#### Spatially homogeneous solutions

#### Travelling wave solutions

In travelling wave coordinates

$$\frac{d^2W}{dz^2} + v\frac{dW}{dz} + W(1 - W) = 0.$$

#### A pair of first order ODEs

$$\begin{split} \frac{dW}{dz} &= P = F(W,P), \\ \frac{dP}{dz} &= -vP - W(1-W) = G(W,P). \end{split}$$

#### Numerical solution

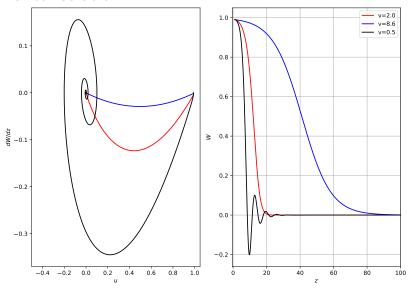


Figure 5: Numerical solution of the travelling wave problem in the phase plane

#### Lecture 8 Recap

- Two steady states (saddle plus stable node)
- Confined set
- no oscillations

#### i PDE

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + u(1 - u), \qquad x \in \mathbb{R}, \ t > 0$$

Travelling wave solution

$$\frac{d^2W}{dz^2} + v\frac{dW}{dz} + W(1 - W) = 0.$$

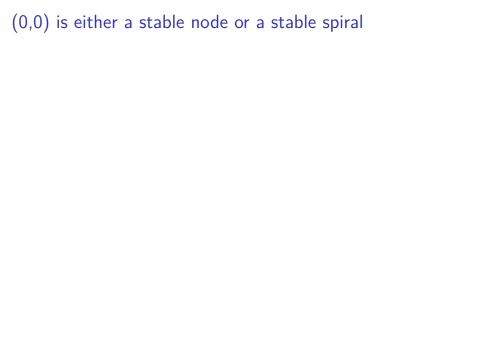
#### Recap

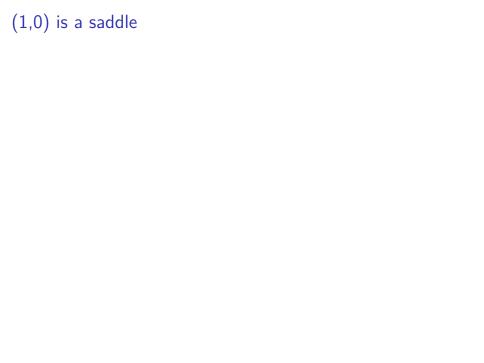
#### i Pair of first order ODEs

$$\begin{split} \frac{dW}{dz} &= P = F(W,P), \\ \frac{dP}{dz} &= -vP - W(1-W) = G(W,P). \end{split}$$

Steady state: (0,0), (1,0)







## A heteroclinic trajectory

## A minimal wave speed

#### Existence of a travelling wave solution

#### Strategy:

- ightharpoonup identify a confined set in  $\Re^2$
- > show no other steady states in confined set
- show no oscillatory solutions

Hence: trajectory that leaves (1,0) via unstable manifold must connect to stable manifold at (0,0)

#### A confined set

Consider

$$T = \{(W, P): 0 \le W \le 1, P \le 0, P \ge \mu W\}$$

for some  $\mu < 0$ .

# A confined set - ctd

#### Lecture 8 Recap

- Two steady states (saddle + stable node)
- Confined set
- No oscillations

#### Finishing off the confined set

#### No oscillations

Bendixson's Negative Criterion, Dulac's Negative Criterion

If there exists a function  $\varphi(W,P)$ , with  $\varphi\in C^1(\mathbb{R}^2)$ , such that

$$\frac{\partial(\varphi F)}{\partial W} + \frac{\partial(\varphi G)}{\partial P},$$

has the same sign  $(\neq 0)$  almost everywhere in a simply connected region (region without holes), then the system

$$\frac{dW}{dz} = F(W, P)$$
$$\frac{dP}{dz} = G(W, P),$$

has no periodic solutions in this region.



For any v>2, there exists a travelling wave solution to Fisher's equation.

# Sign of the wave speed

Consider the travelling wave ODE

$$\frac{d^2W}{dz^2} + vW + W(1-W) = 0$$

# The bistable equation

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + f(u), \qquad x \in \mathbb{R}, \ t > 0, \tag{15}$$

with initial condition

$$u(x,0)=u_0(x), \qquad x\in \mathbb{R}.$$

Let

$$f(0) = f(a) = f(1) = 0$$
, with  $0 < a < 1$ .

There are therefore three spatially uniform steady states  $u_1=0$ ,  $u_2=a,\ u_3=1.$ 

$$f'(0) < 0$$
,  $f'(a) > 0$  and  $f'(1) < 0$ 

$$f = u(u - a)(1 - u),$$

which arises in the study of nerve action potentials along nerve fibres and other problems in *excitable media* 

# Numerical solution

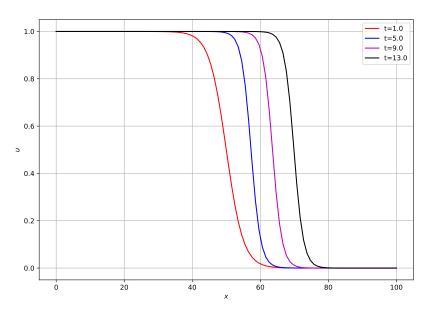


Figure 6: Travelling waves in a numerical solution of bistable PDE.

# Lecture 9

▶ Bistable equation

# Travelling wave ansatz

$$\frac{d^2W}{dz^2} + v\frac{dW}{dz} + f(W) = 0,$$

In the phase plane

$$\frac{dW}{dz} = P = F(W,P),$$
 
$$\frac{dP}{dz} = -vP - f(W) = G(W,P),$$



# The sign of $\boldsymbol{v}$

### Lecture 10

Recap: travelling wave analysis of bistable equation

$$\frac{dW}{dz} = P = F(W,P),$$
 
$$\frac{dP}{dz} = -vP - f(W) = G(W,P),$$

# Numerical shooting

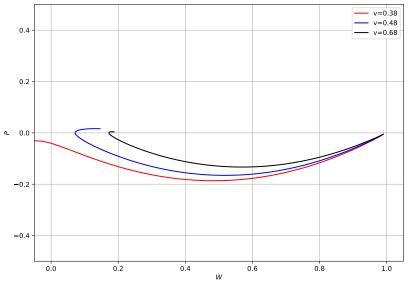


Figure 7: Using a shooting method to investigate travelling wave solutions. Continuity arguments suggest that there exists a travelling

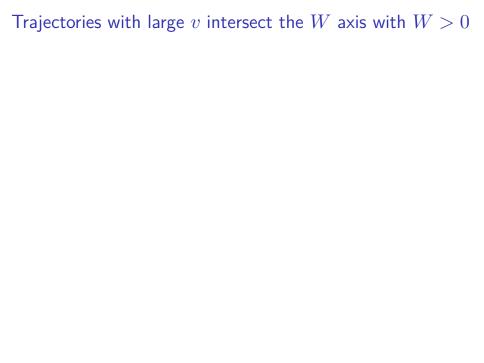
A shooting method to prove the existence of a traveling wave

Outline

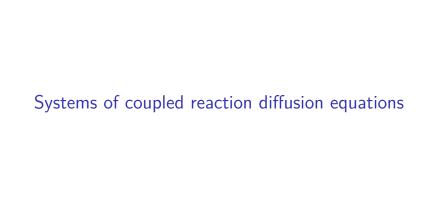
Trajectories with small  $\boldsymbol{v}$  intersect the P axis with P<0

Assume that

$$\int_{0}^{1} f(u) du > 0.$$



Continuity there exists a trajectory with intermediate  $\boldsymbol{v}$  that passes through the origin



# Lotka Volterra with diffusion

$$\begin{split} \frac{\partial u}{\partial t} &= \rho \, u \left( 1 - \frac{u}{K} \right) - \alpha \, u \, n + D_u \Delta u, \\ \frac{\partial n}{\partial t} &= \beta \, u \, n - \gamma \, n + D_n \Delta n, \end{split} \tag{16}$$

### Nondimensional form

$$\frac{\partial u}{\partial t} = u(1 - u - n) + D\frac{\partial^2 u}{\partial x^2} = f(u, n) + D\frac{\partial^2 u}{\partial x^2}, \quad x \in \mathbb{R}, t > 0,$$

$$\frac{\partial n}{\partial t} = a n(u - b) + \frac{\partial^2 n}{\partial x^2} = g(u, n) + \frac{\partial^2 n}{\partial x^2}, \quad x \in \mathbb{R}, t > 0,$$
(17)

# Spatially homogeneous steady states

### Lecture 11

Recap

$$\begin{split} \frac{\partial u}{\partial t} &= u(1-u-n) + D\frac{\partial^2 u}{\partial x^2} \ = \ f(u,n) + D\frac{\partial^2 u}{\partial x^2}, \quad x \in \mathbb{R}, t > 0 \,, \\ \frac{\partial n}{\partial t} &= a \, n(u-b) + \frac{\partial^2 n}{\partial x^2} \ = g(u,n) + \frac{\partial^2 n}{\partial x^2}, \qquad x \in \mathbb{R}, t > 0, \end{split}$$

Spatially hom. steady states:

- $\triangleright$  (0,0) extinction lin. unstable
- $\blacktriangleright$  (1,0) no predator, lin unstable
- $\blacktriangleright$  (b, 1-b) coexistence lin. stable

Question: do travelling wave solutions exist that connect the spatially homogeneous stable steady state to either of the unstable steady states?

# A travelling wave that connects (1,0) and (b,1-b)

$$u(x,t) = W(x+vt) = W(z), \quad v > 0,$$
  
 $n(x,t) = N(x+vt) = N(z), \quad v > 0.$ 

# The limit of fast diffusing predator

## Three first order ODEs

$$\frac{dW}{dz} = \frac{1}{v}W(1 - W - N) = F(W, N, P),$$

$$\frac{dN}{dz} = P = G(W, N, P),$$

$$\frac{dP}{dz} = vP - aN(W - b) = R(W, N, P).$$
(18)





### Lecture 12

- $\blacktriangleright$  steady states: (0,0,0), (1,0,0), (b,1-b,0)
- $\blacktriangleright$  Heteroclinic trajectory from (1,0,0) to (b,1-b,0)
- (1,0,0) has a 2 dim unstable manifold
- ightharpoonup Eigenvalues at (b,1-b,0) satisfy

$$\lambda^3 - \lambda^2(v - \frac{b}{v}) - \lambda b - \frac{1}{v}ab(1 - b) = p(\lambda) = 0.$$

# TPs are independent of $\boldsymbol{a}$

 $\ensuremath{\mathsf{p}}(\ )$  has a real positive root and two roots with negative real part

# Plotting the cubic

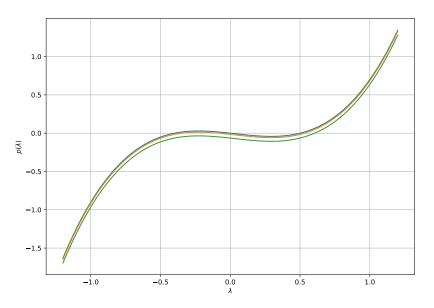


Figure 8: Plot of cubic.

# Aggregation via chemotaxis

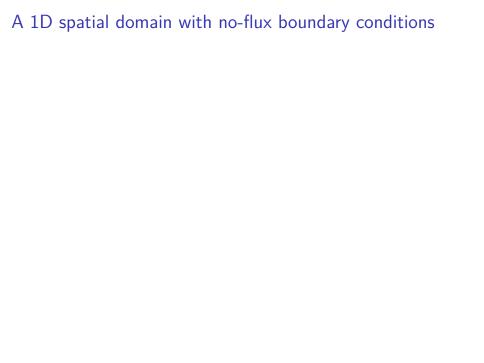
- Dictyostelium discoideum (Dicty) is a slime-mold that is widely studied experimentally as a model organism.
- under nutrient starvation, it exhibits complex collective behaviour
- individual amoebae that constitute a slime-mold exhibit a range of phenomena also observed in mammalian cells e.g. differentiation, proliferation, migration.



Figure 9: Spiral wave patterns underly Dictystelium aggregation.

How do simple rules give rise to complex behaviours?

# A chemotactic model



### Lecture 13

Domain:

$$x \in [0, L], t > 0$$

PDE:

$$\frac{\partial n}{\partial t} = D_n \frac{\partial^2 n}{\partial x^2} - \chi_0 \frac{\partial}{\partial x} \left( n \frac{\partial a}{\partial x} \right),$$

$$\frac{\partial a}{\partial t} = D_a \frac{\partial^2 a}{\partial x^2} + \mu n - \delta a,$$

Boundary conditions:

$$\frac{\partial a}{\partial x} = \frac{\partial n}{\partial x} = 0, \quad x = 0, L.$$

(19)

ICs:

### numerical solution

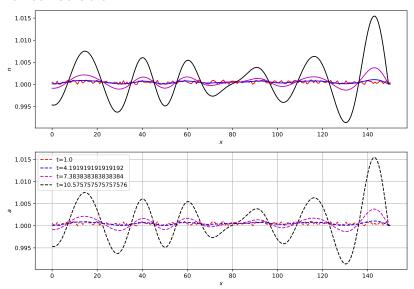


Figure 10: Numerical solution of bacterial chemtaxis model.

# Conservation of cell number

# Spatially homogeneous solutions

# Linearisation about the spatially homogeneous steady state

$$n(x,t) = n^* + \tilde{n}(x,t), \quad a(x,t) = a^* + \tilde{a}(x,t)$$

## Separable solution

$$\tilde{n}(t,x) = u(t)\phi_1(x), \quad \tilde{a}(t,x) = v(t)\phi_2(x)$$

#### The elliptic problem

$$\begin{split} \frac{d^2\phi}{dx^2} &= -k^2\phi & \text{ in } (0,L),\\ \frac{d\phi}{dx} &= 0 & \text{ for } x=0, \; x=L. \end{split}$$

$$\phi_1 = \phi 2 = \phi$$

## Linear system solution

$$u(t) = C_1 e^{\lambda t}$$
 and  $v(t) = C_2 e^{\lambda t}$ 

## Lecture 13 - Recap

Domain:

$$x \in [0, L], t > 0$$

PDE:

$$\frac{\partial n}{\partial t} = D_n \frac{\partial^2 n}{\partial x^2} - \chi_0 \frac{\partial}{\partial x} \left( n \frac{\partial a}{\partial x} \right),$$

$$\frac{\partial a}{\partial t} = D_a \frac{\partial^2 a}{\partial x^2} + \mu n - \delta a,$$

Boundary conditions:

$$\frac{\partial a}{\partial x} = \frac{\partial n}{\partial x} = 0, \quad x = 0, L.$$

(20)

ICs:

#### Method

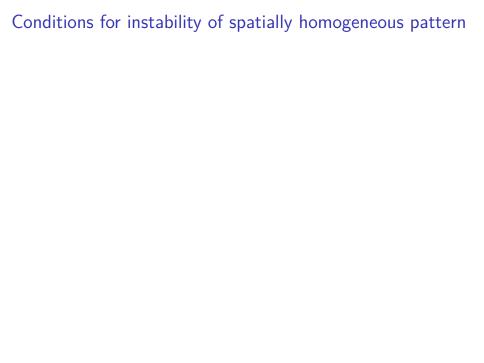
- Linearise about the steady state
- ▶ Separation of variables  $\tilde{n}(x,t) = u(t)\phi(x)$   $\tilde{a}(x,t) = v(t)\phi(x)$
- Eigenvalues of the Laplacian operator

$$\frac{d^2\phi}{dx^2} = -k^2\phi, \quad k = \frac{\bar{n}\pi}{L}, \bar{n} \in Z$$

- Linear ODEs in u(t) and v(t)
- For instability of spatially homogeneous steady state, require  $\Re\{\lambda>0\}$ .

## Eigenvalue equation

$$\lambda^2 + \left(D_n k^2 + D_a k^2 + \delta\right) \lambda + D_n k^2 \left(D_a k^2 + \delta\right) - \mu \chi_0 n^* k^2 = 0. \label{eq:lambda}$$



# Lecture 16 - Diffusion driven instability

$$\begin{split} \frac{\partial A}{\partial t} &= F(A,B) + D_A \nabla^2 A, \\ \frac{\partial B}{\partial t} &= G(A,B) + D_B \nabla^2 B, \end{split}$$

#### Reaction kinetics

Schnackenberg

$$F(A,B) = k_1 - k_2 A + k_3 A^2 B, \quad G(A,B) = k_4 - k_3 A^2 B$$

Gierer Meinhardt:

$$F(A,B) = k_1 - k_2 A + \frac{k_3 A^2}{B}, \quad G(A,B) = k_4 A^2 - k_5 B$$

Thomas:

$$\begin{split} F(A,B) &= k_1 - k_2 A - H(A,B), \\ G(A,B) &= k_4 A^2 - k_4 B - H(A,B), \\ H(A,B) &= \frac{k_5 A B}{k_6 + k_7 + k_8 A^2}. \end{split}$$

# Nondimensionalisation of Schnakenberg model

Using the scaling

$$u = A \left(\frac{k_3}{k_2}\right)^{1/2}, \quad v = B \left(\frac{k_3}{k_2}\right)^{1/2}, \quad t^* = \frac{D_A t}{L^2}, \quad x^* = \frac{x}{L},$$

$$\frac{\partial u}{\partial t} = \gamma (a - u + u^2 v) + \nabla^2 u = \gamma f(u, v) + \nabla^2 u,$$

$$\frac{\partial v}{\partial t} = \gamma (b - u^2 v) + d\nabla^2 v = \gamma g(u, w) + d\nabla^2 v,$$
(21)

# Interpretation of Schnackenberg model: short range activation/long range inhibition

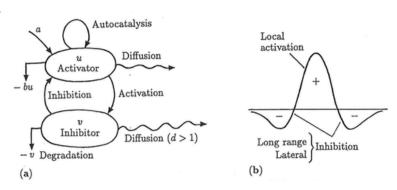
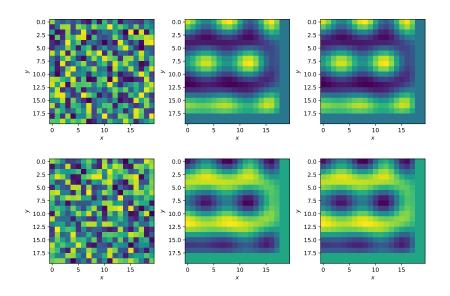


Figure 11

#### General form for nondimensionalised RD model

$$\begin{split} \frac{\partial u}{\partial t} &= \gamma f(u,v) + \nabla^2 u, \\ \frac{\partial v}{\partial t} &= \gamma g(u,w) + d\nabla^2 v, \end{split}$$

#### Numerical solution



# Deriving general conditions for diffusion-driven instability

Let  $\Omega \subset \mathbb{R}^n$  be a domain with smooth (sufficiently regular) boundary  $\partial \Omega$ , with outward unit normal  $\mathbf{n}$ .

$$\begin{split} \frac{\partial u}{\partial t} &= \gamma \, f(u,v) + \nabla^2 u, & x \in \Omega, \quad t > 0, \\ \frac{\partial v}{\partial t} &= \gamma \, g(u,v) + d \nabla^2 v, & x \in \Omega, \quad t > 0, \end{split} \tag{22}$$

Boundary and initial conditions

$$\nabla u \cdot \mathbf{n} = 0, \quad \nabla v \cdot \mathbf{n} = 0, \quad x \in \partial \Omega, \quad t > 0,$$
  
$$u(x,0) = u_0(x), \quad v(x,0) = v_0(x), \quad x \in \Omega.$$
 (23)

# Conditions for diffusion driven instability

$$\begin{split} f_u + g_v &< 0, \\ f_u g_v - f_v g_u &> 0, \\ df_u + g_v &> 0, \\ (df_u + g_v)^2 - 4d(f_u g_v - f_v g_u)^2 &< 0, \end{split}$$

# A spatially homogeneous steady-state

A *spatially homogeneous steady-state* of Equation 22 and Equation 23 satisfies

$$f(u_0,v_0) = g(u_0,v_0) = 0. \\$$

For linear stability

$$\begin{split} f_u + g_v &< 0, \\ f_u g_v - f_v g_u &> 0 \end{split}$$

Here

$$f_u = \frac{\partial f}{\partial u_{(u_0, v_0)}}$$

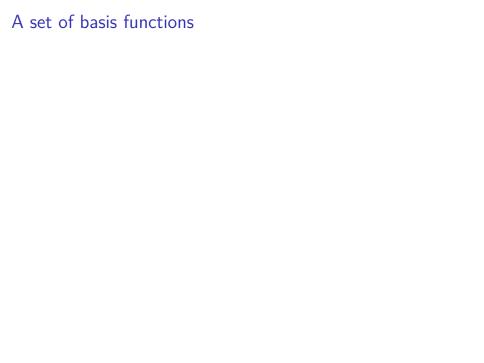
etc.

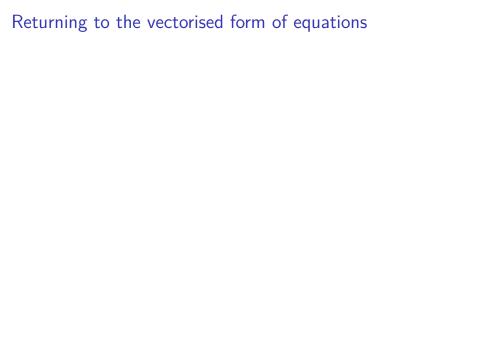
# Spatially dependent perturbations

 $u(x,t) = u_0 + \tilde{u}(x,t), \quad v(x,t) = v_0 + \tilde{v}(x,t), \qquad \|\tilde{u}(x,t)\| \ll 1, \quad \|\tilde{v}(x,t)\| \ll 1,$ 

## Separation of variables

$$V(x,t) = \begin{pmatrix} \bar{u}(t)\varphi_1(x) \\ \bar{v}(t)\varphi_2(x) \end{pmatrix},$$



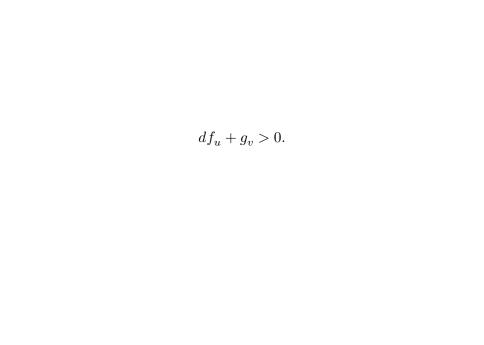


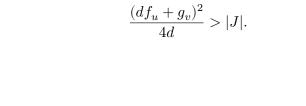
# A modified Jacobian for spatially heterogeneous perturbations

### Lecture 18 Recap

For linear instability of spatially homogeneous steady state:

$$\det(\tilde{J})=h(k^2)=dk^4-\gamma(d\,f_u+g_v)k^2+\gamma^2\det(J)<0.$$





## Instability for a limited range of wavenumbers

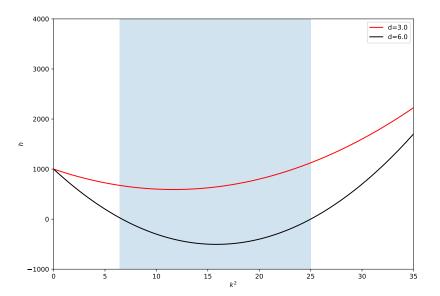


Figure 12: A plot of  $h(k^2)$  plotted against  $k^2.$  Shaded region denotes

## Summary - DDI conditions

$$\begin{split} f_u + g_v &< 0, \\ f_u g_v - f_v g_u &> 0, \\ df_u + g_v &> 0, \\ (df_u + g_v)^2 - 4d(f_u g_v - f_v g_u)^2 &< 0, \end{split}$$