

A Parametric Study of the SASI Comparing General Relativistic and Non-Relativistic Treatments

Samuel J. Dunham

September 11, 2023

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AND KELLY HOLLEY-BOCKELMANN ^{1,5}

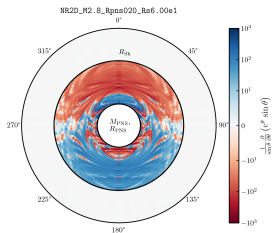
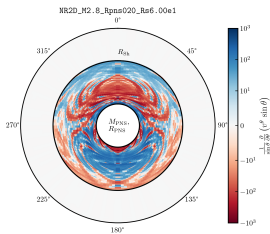
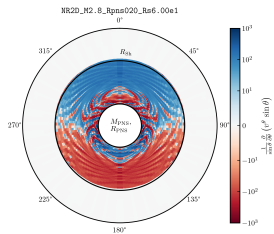
¹*Department of Physics and Astronomy, Vanderbilt University, 6301 Stevenson Center, Nashville, TN 37235*

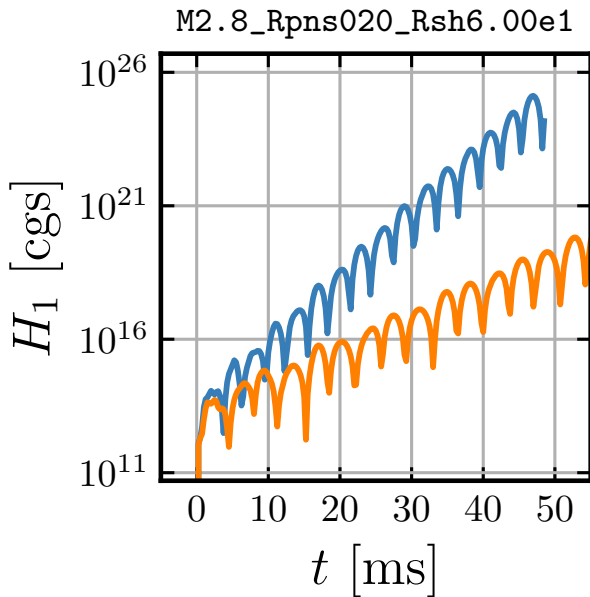
²*Department of Physics and Astronomy, University of Tennessee, Knoxville, Nielsen Physics Building, 401, 1408 Circle Drive, Knoxville, TN 37996*

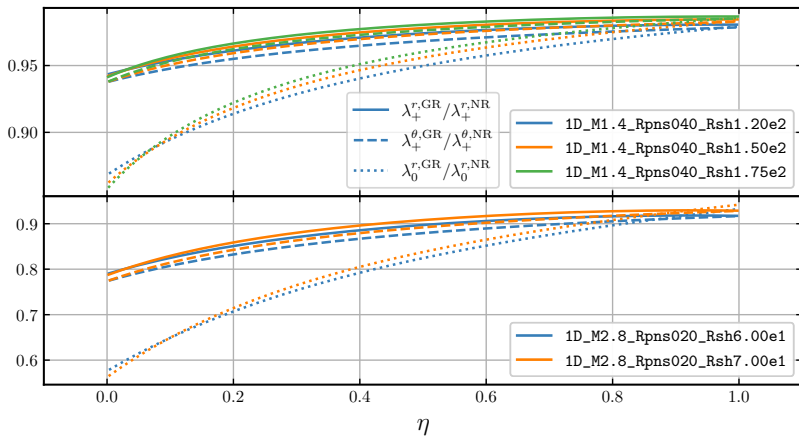
³*Computer Science and Mathematics Division, Oak Ridge National Laboratory, Oak Ridge, TN 37831*

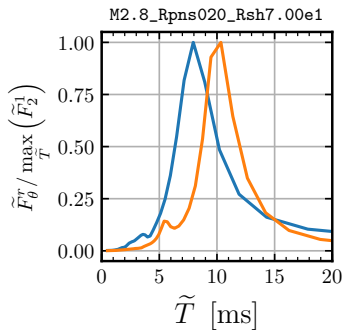
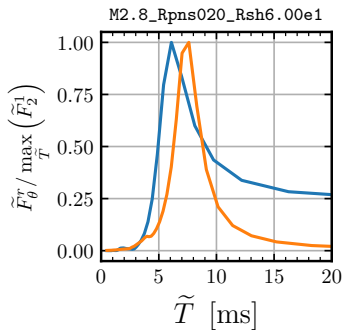
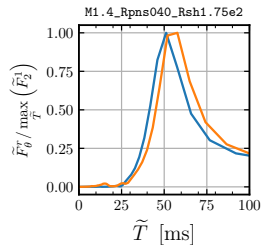
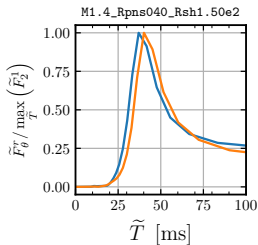
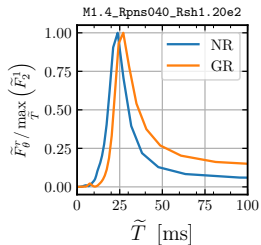
⁴*Department of Physics, North Carolina State University, 2401 Stinson Dr., Raleigh, NC 27607*

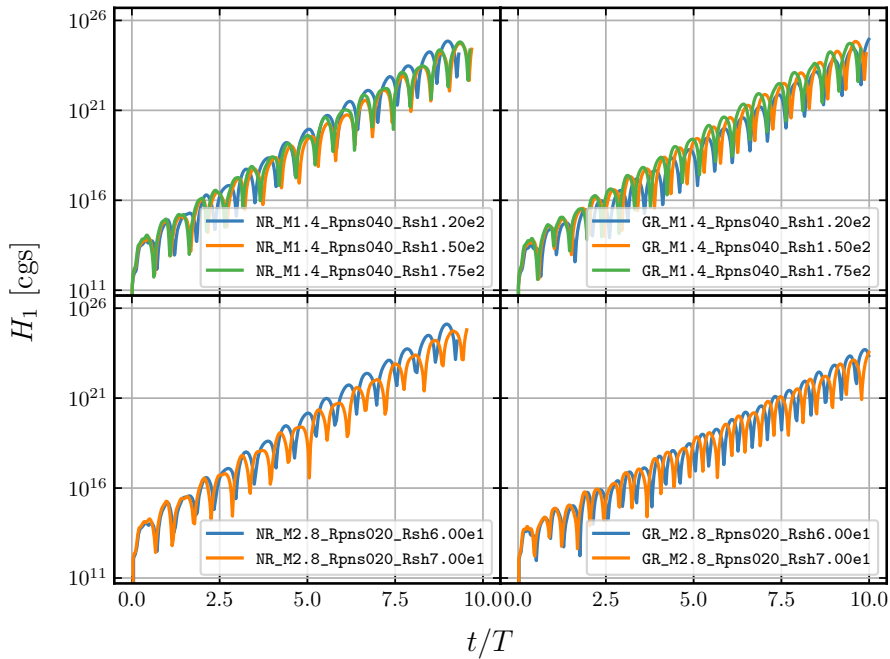
⁵*Department of Life and Physical Sciences, Fisk University, 1000 17th Ave N, Nashville, TN 37208*

t  $t + 0.25 T$  $t + 0.5 T$ 

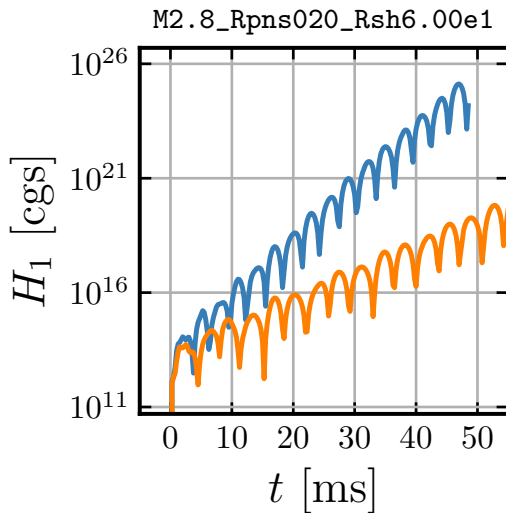




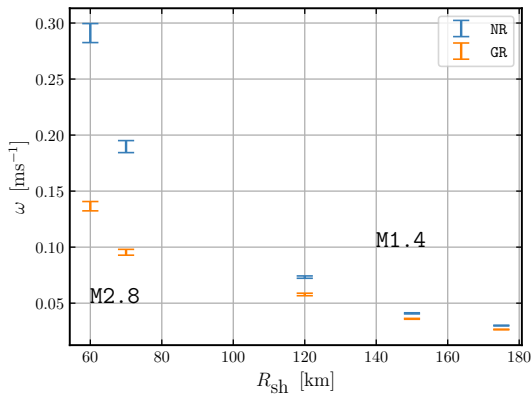




Conclusions

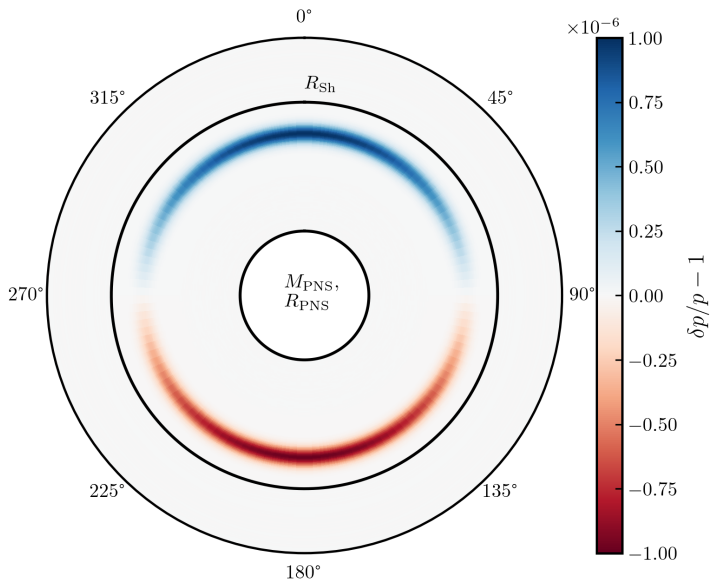


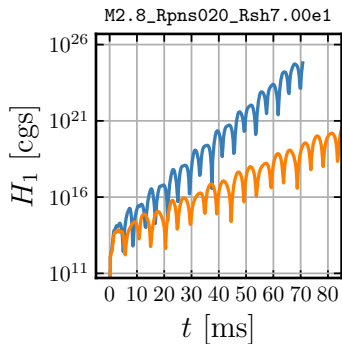
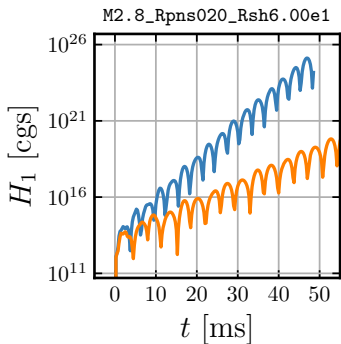
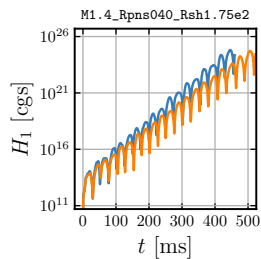
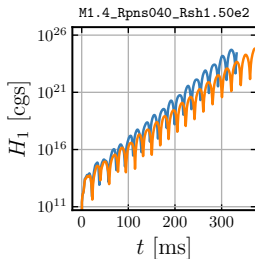
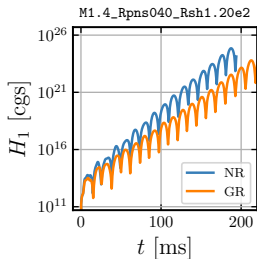
Conclusions



References

- Extended study of ? to include GR
- Showed that GR leads to longer SASI oscillation period than NR
- Showed that GR leads to smaller SASI growth rate than NR
- Found that growth rate is such that ωT is roughly constant for some parameter sets: implications for growth rate mechanism
- Future Work
 - Extend study to 3D
 - Include GR monopole (?)



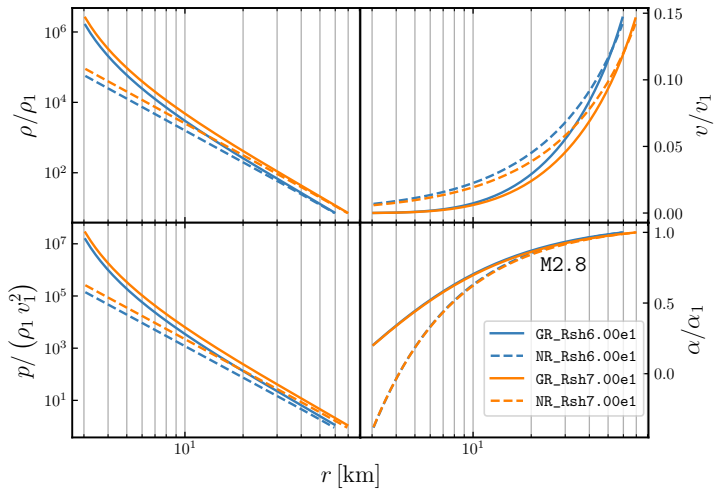


$$A := \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (v^\theta \sin \theta) \quad (?)$$

$$A(r, \theta, t) = \sum_{\ell'=0}^{\infty} G_{\ell'}(r, t) P_{\ell'}(\cos \theta)$$

$$\implies G_\ell(r, t) := \frac{1}{N_\ell} \int_0^\pi A(r, \theta, t) P_\ell(\cos \theta) \sin \theta d\theta$$

$$H_\ell(t) := 4\pi \int_{r_a}^{r_b} [G_\ell(r, t)]^2 [\psi(r)]^6 r^2 dr$$



Parameters we varied:

$\xi =$

$(M_{\text{PNS}}/M_{\odot}) / (R_{\text{PNS}}/20 \text{ km})$

$(?), R_{\text{Sh}} (t = 0)$

Table 1. Model Parameters

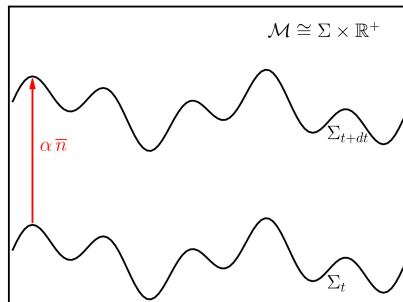
Model	$M_{\text{PNS}} [M_{\odot}]$	$R_{\text{PNS}} [\text{km}]$	$R_{\text{sh}} [\text{km}]$	ξ
M1.4-Rpns040-Rsh1.20e2	1.4	40	120	0.7
M1.4-Rpns040-Rsh1.50e2	1.4	40	150	0.7
M1.4-Rpns040-Rsh1.75e2	1.4	40	175	0.7
M2.8-Rpns020-Rsh6.00e1	2.8	20	60	2.8
M2.8-Rpns020-Rsh7.00e1	2.8	20	70	2.8

NOTE—Model parameters chosen for the 5 models. All models were run with both GR and NR. The first three rows correspond to the low-compactness models and the last two rows correspond to the high-compactness models.

Conformally-Flat Condition

??

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij} dx^i dx^j$$



Conformally-Flat Condition

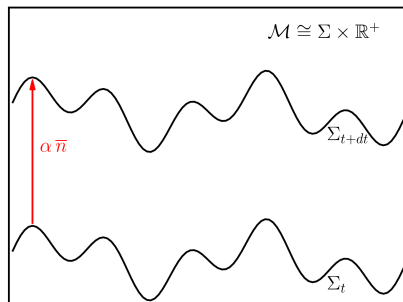
??

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij} dx^i dx^j$$

$$\gamma_{ij} = \psi^4 \bar{\gamma}_{ij}$$

$1 \leq \psi < 2$: Conformal factor

$$\bar{\gamma}_{ij} = \text{diag}(1, r^2, r^2 \sin^2 \theta)$$



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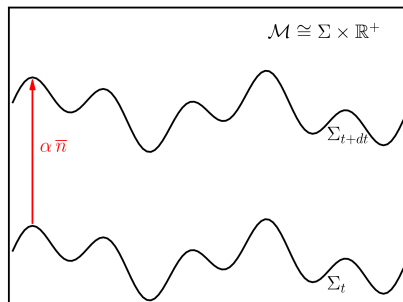
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(Also, maximum slicing condition:

$$K := \text{Tr}_{\gamma_{ij}} (\underline{\underline{K}}) = \partial_t K = 0)$$



Isotropic Coordinates ($c = G = 1$)

?

$$\alpha(r) = \left(1 - \frac{R_{\text{Sc}}}{r}\right) \left(1 + \frac{R_{\text{Sc}}}{r}\right)^{-1}$$

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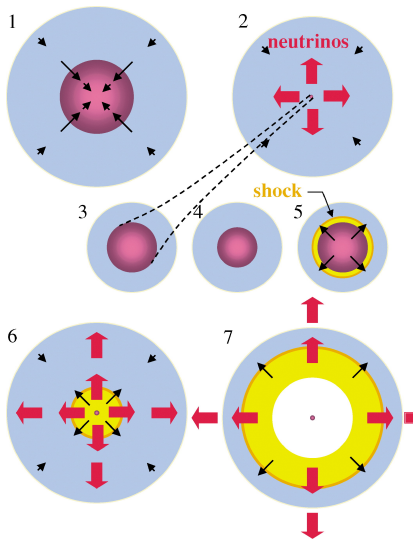
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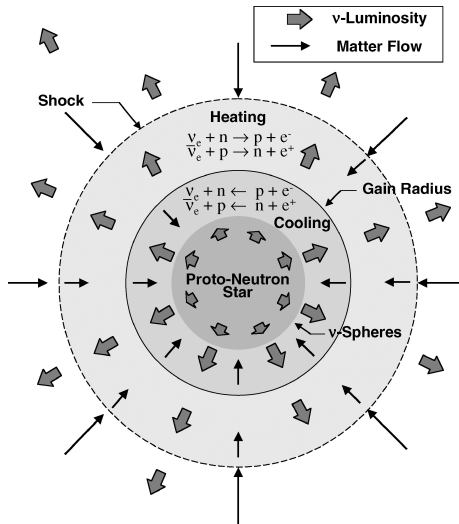
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$$\beta^i = 0$$

$$K_{ij} = 0$$



Mezzacappa A. 2005.
Annu. Rev. Nucl. Part. Sci. 55:467–515



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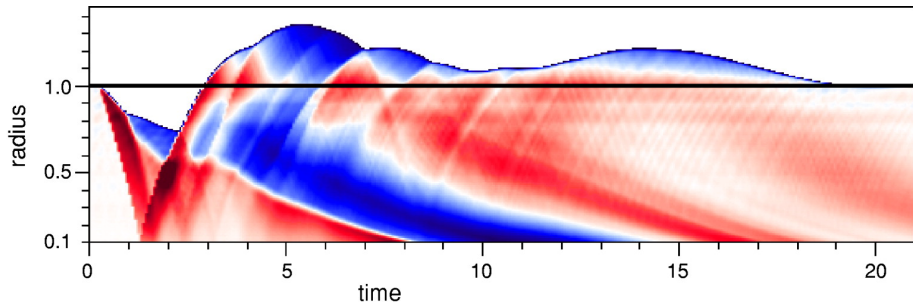


Figure: Equilibrium (white), under- (blue), and over- (red) pressure (?).

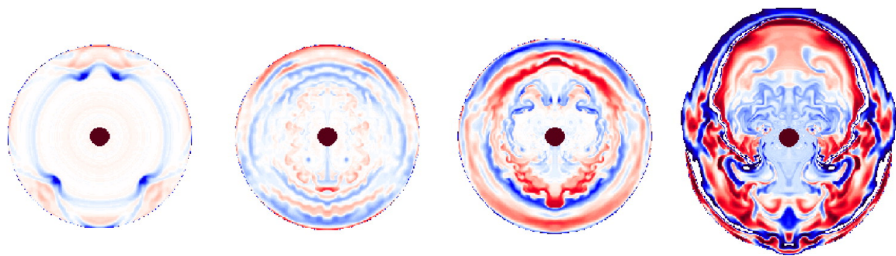
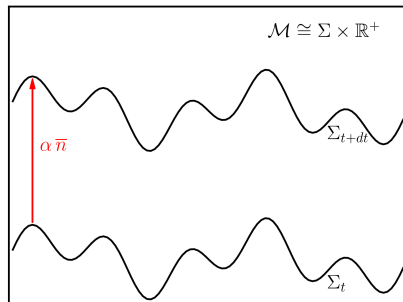


Figure: Equilibrium (white), under- (blue), and over- (red) entropy (?).

$d+1$ Decomposition

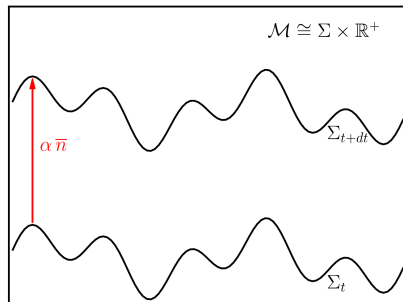
$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$



$d+1$ Decomposition

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\underline{g} : spacetime metric on \mathcal{M}

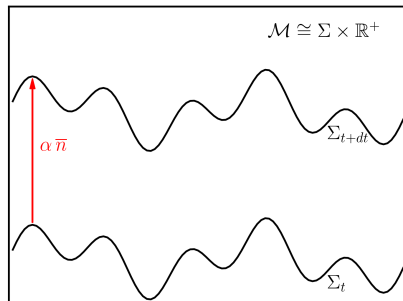


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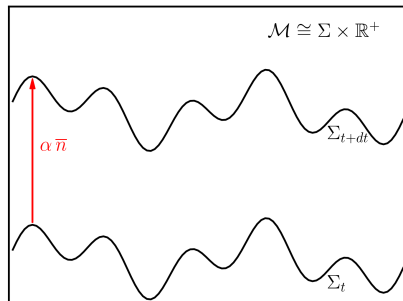
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$\underline{\underline{\gamma}}$: spatial three-metric on Σ_t



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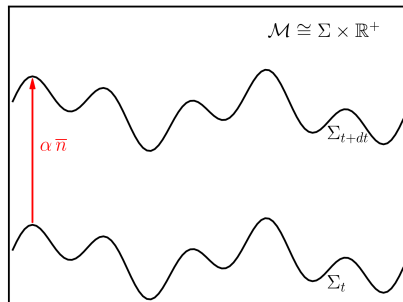
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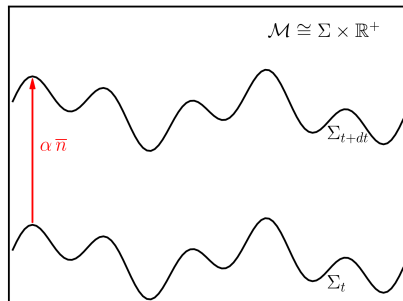
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\bar{n} : Eulerian four-velocity

$$(\underline{\underline{g}}(\bar{n}, \bar{n}) = \bar{n} \cdot \bar{n} = -1)$$



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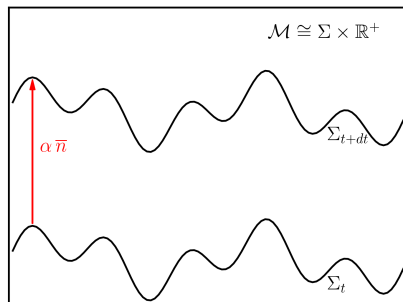
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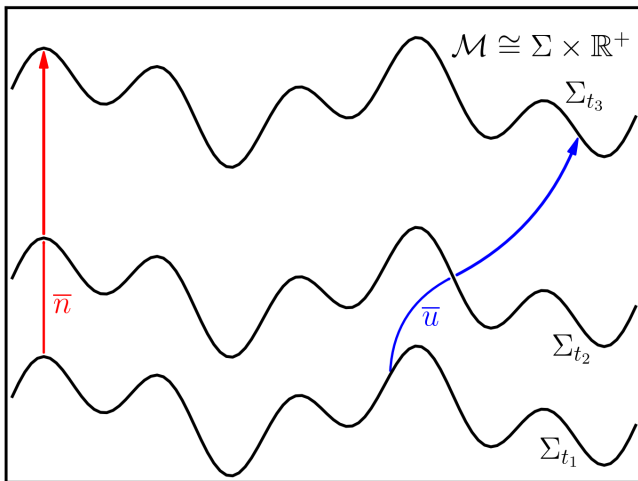
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(Also, K : Extrinsic curvature)





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Units defined such that $c = G = 1$

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$\bar{\nabla} \cdot \bar{J} = 0$ (\bar{J} : baryon mass density current four-vector)

$\bar{\nabla} \cdot \bar{\bar{T}} = \bar{0}$ ($\bar{\bar{T}}$: Rank (2, 0) stress-energy tensor)

$\bar{J} := \rho \bar{u}$ (ρ : **comoving** baryon mass density)

$\bar{\bar{T}} := \rho h \bar{u} \otimes \bar{u} + p \bar{\bar{g}}$ (p : comoving thermal pressure,

$h := 1 + (e + p) / \rho$: specific enthalpy, e : comoving internal energy density)

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Five equations with six unknowns ☹

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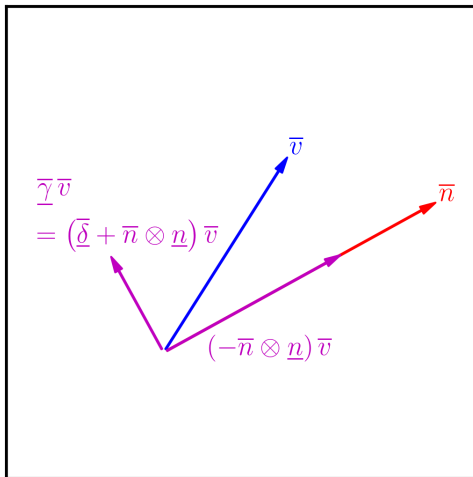
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Five equations with six unknowns ☹

Close with an equation of state: $p = p(e) := (\Gamma - 1) e$, $\Gamma = 4/3$

Valencia Decomposition



Extensible to higher-rank tensors!

$$E := n_{\mu'} n_{\nu'} T^{\mu' \nu'}$$

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$$S^{\mu} := -\gamma^{\mu}_{\mu'} n_{\nu'} T^{\mu' \nu'}$$

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Math...

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

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$$\sqrt{\gamma} = \psi^6 \sqrt{\bar{\gamma}}$$

NR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho$$

$$S_j := \rho v_j$$

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho W$$

$$S_j := \rho h W^2 v_j$$

$$\tau := E - D = \rho h W^2 - p - \rho W$$

$$\alpha = (1 - R_{\text{Sc}}/r) / (1 + R_{\text{Sc}}/r)$$

$$\sqrt{\gamma} = \psi^6 \sqrt{\bar{\gamma}}$$

NR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho$$

$$S_j := \rho v_j$$

$$\tau = e + \frac{1}{2} \rho v^i v_i$$

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho W$$

$$S_j := \rho h W^2 v_j$$

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$$\sqrt{\gamma} = \psi^6 \sqrt{\bar{\gamma}}$$

NR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho$$

$$S_j := \rho v_j$$

$$\tau = e + \frac{1}{2} \rho v^i v_i$$

$$\alpha = 1$$

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho W$$

$$S_j := \rho h W^2 v_j$$

$$\tau := E - D = \rho h W^2 - p - \rho W$$

$$\alpha = (1 - R_{\text{Sc}}/r) / (1 + R_{\text{Sc}}/r)$$

$$\sqrt{\gamma} = \psi^6 \sqrt{\bar{\gamma}}$$

NR

$$\mathbf{U} := (D, S_j, \tau)^\top$$

$$D := \rho$$

$$S_j := \rho v_j$$

$$\tau = e + \frac{1}{2} \rho v^i v_i$$

$$\alpha = 1$$

$$\sqrt{\gamma} = \sqrt{\bar{\gamma}}$$

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\mathbf{F}^i(\mathbf{U}) = \begin{pmatrix} \rho W v^i \\ \rho h W^2 v^i v_j + p \delta^i_j \\ (\rho h W^2 - \rho W) v^i \end{pmatrix}$$

NR

$$\mathbf{F}^i(\mathbf{U}) = \begin{pmatrix} \rho v^i \\ \rho v^i v_j + p \delta^i_j \\ (\rho h_{\text{NR}} + \frac{1}{2} v^j v_j) v^i \end{pmatrix}$$

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_i [\alpha \sqrt{\gamma} \mathbf{F}^i(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\mathbf{S}(\mathbf{U}) = \begin{pmatrix} 0 \\ \frac{1}{2} \alpha P^{ik} \partial_j \gamma_{ik} - E \partial_j \alpha \\ -S^j \partial_j \alpha \end{pmatrix}$$

NR

$$\mathbf{S}(\mathbf{U}) = \begin{pmatrix} 0 \\ \frac{1}{2} P^{ik} \partial_j \bar{\gamma}_{ik} - \rho \partial_j \Phi \\ -S^j \partial_j \Phi \end{pmatrix}$$

$\Phi(r) := -M/r$

Steady-State Solutions (pre-shock)

$$\cancel{\partial_t} \mathbf{U}^0 + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

NR

Steady-State Solutions (pre-shock)

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

NR

$$4\pi r^2 \rho v = -\dot{M}$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

NR

$$4\pi r^2 \rho v = -\dot{M}$$

$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

Steady-State Solutions (pre-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_1 \rho^\Gamma$$

NR

$$4\pi r^2 \rho v = -\dot{M}$$

$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

$$p = K_1 \rho^\Gamma$$

Jump Conditions

$$U_1 \neq U_2$$

$$F^r(U_1) = F^r(U_2)$$

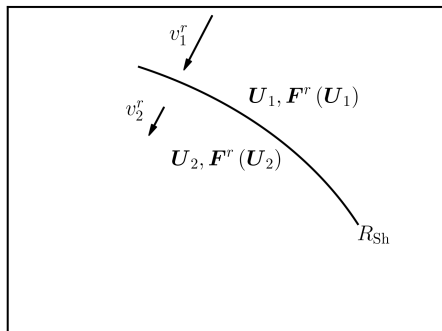
Yields:

$$\rho_2$$

$$e_2(p_2)$$

$$K_2(> K_1)$$

$$v_2^r$$



Steady-State Solutions (post-shock)

$$\partial_t \vec{U} + \frac{1}{\sqrt{\gamma}} \partial_r [\alpha \sqrt{\gamma} \mathbf{F}^r(\mathbf{U})] = \mathbf{S}(\mathbf{U})$$

GR

$$\alpha \psi^6 W \times 4\pi r^2 \rho v = -\dot{M}$$

$$\alpha h W = \mathcal{B}$$

$$p = K_2 \rho^\Gamma$$

NR

$$4\pi r^2 \rho v = -\dot{M}$$

$$\frac{1}{2} v^2 + h_{\text{NR}} + \Phi = \mathcal{B}_{\text{NR}}$$

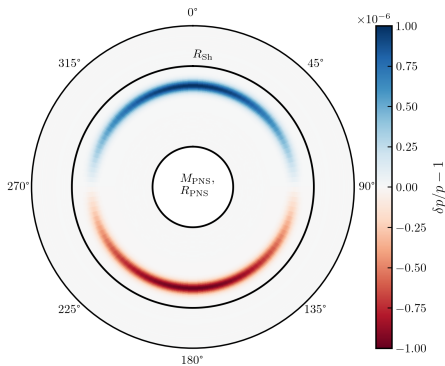
$$p = K_2 \rho^\Gamma$$

$$\eta(r) := \frac{r - R_{\text{PNS}}}{R_{\text{Sh}} - R_{\text{PNS}}}$$

$$\frac{\delta p(\eta, \theta)}{p(\eta_c)} = 10^{-6} \times \exp \left[\frac{-(\eta - \eta_c)^2}{2\sigma^2} \right] \cos \theta$$

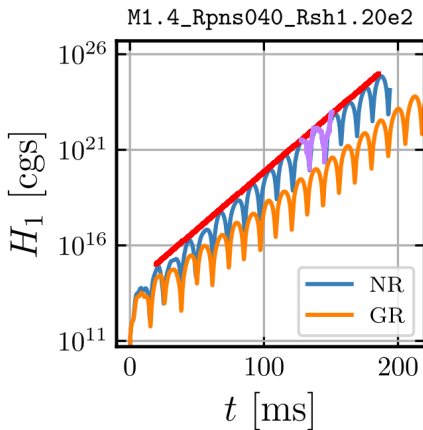
$$\eta_c = 0.75$$

$$\sigma = 0.05$$



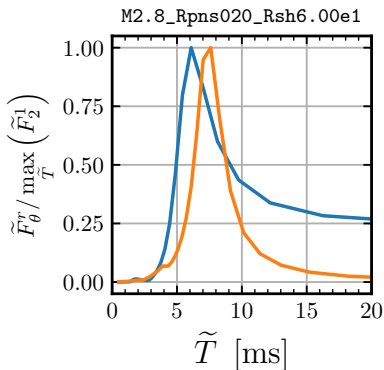
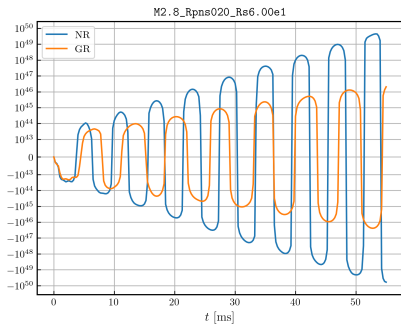
$$F(t) = F(0) e^{2\omega t} \sin^2\left(\frac{2\pi t}{T} + \delta\right)$$

(?)



$$F_{\theta}^r := \alpha \psi^6 h W^2 \times \sqrt{\gamma} \rho v^r v_{\theta}$$

$$\tilde{F}_{\theta}^r := \text{FFT} \{F_{\theta}^r\}$$



T defined as the unique \tilde{T} such that $\tilde{F}_{\theta}^r (\tilde{T}) = 1$

