

Towards Discontinuous Galerkin Methods for General Relativistic Core-Collapse Supernova Simulations



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Outline

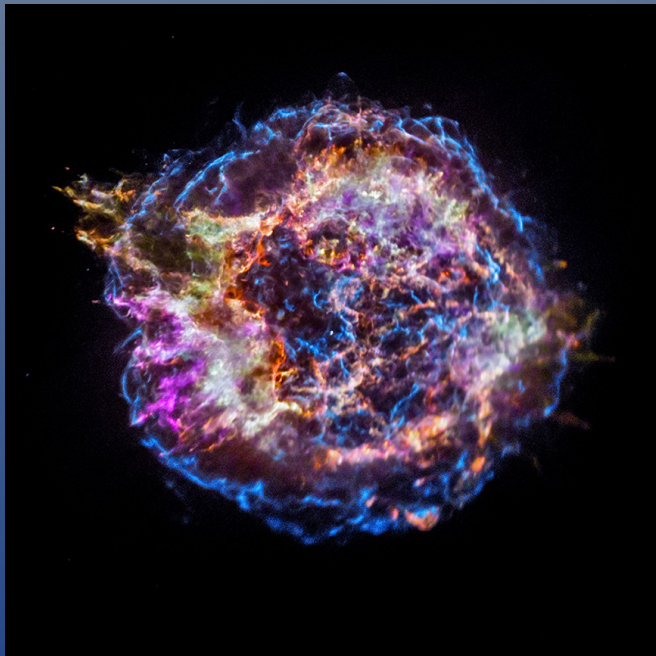
Motivation
Methods
Results
Summary, etc.

Outline

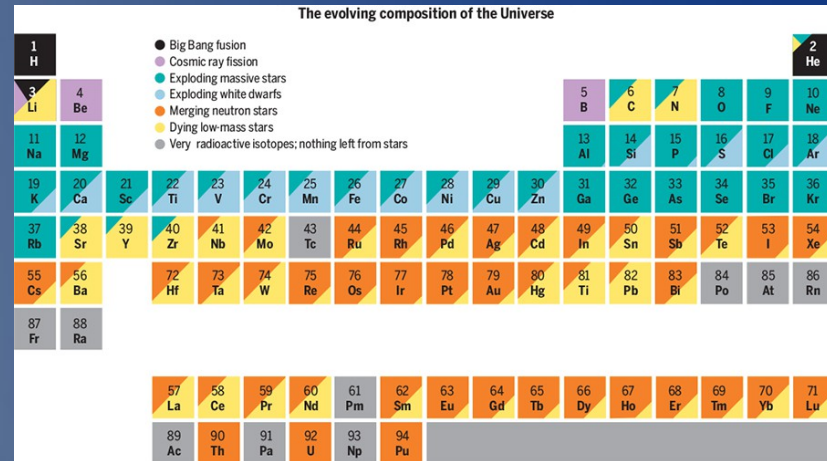
Motivation

Motivation

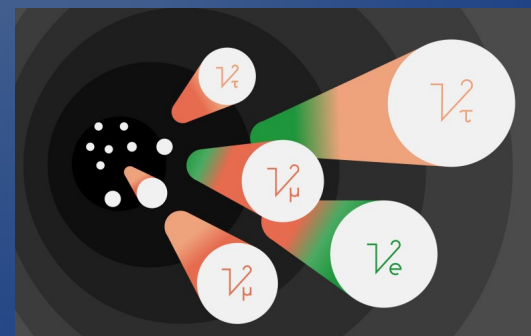
- Why study core-collapse supernovae (CCSNe)?
 - Nucleosynthesis
 - Neutron stars
 - Gravitational waves
 - Neutrinos



APOD



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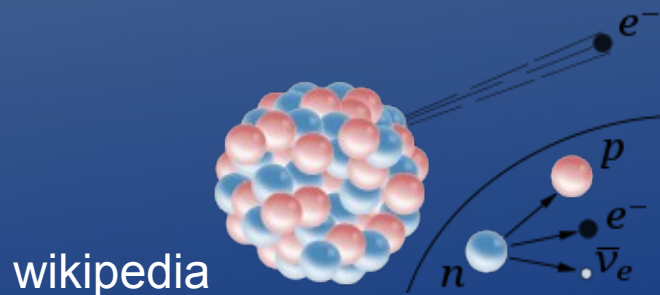
sanfordlab.org

Outline

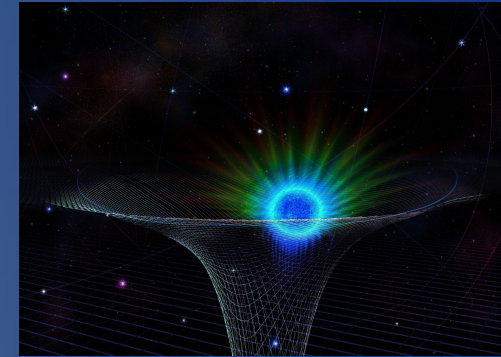
Methods

How to study CCSNe?

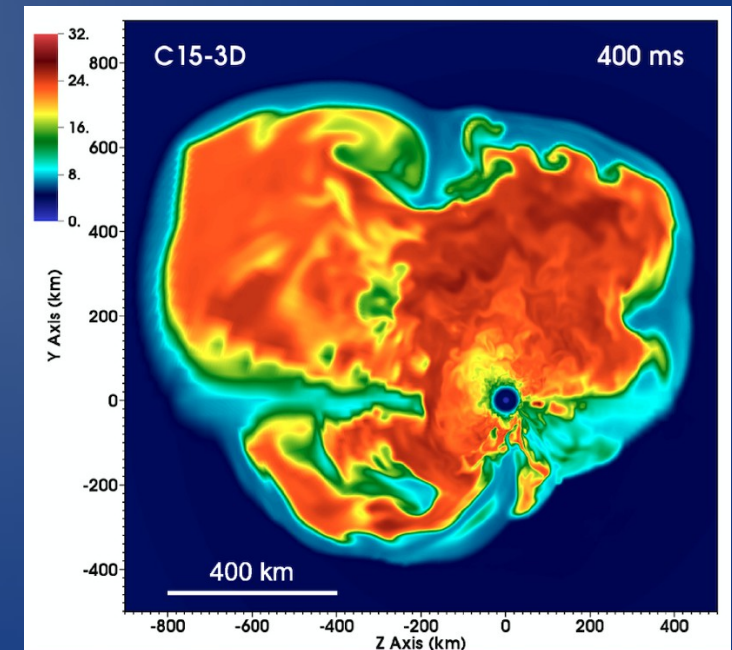
- Simulations!
 - Macrophysics
 - Gravity
 - Hydrodynamics
 - Neutrino transport
 - Microphysics
 - Equation of state
 - $p = p(\rho, T, Y_e)$
 - Weak interactions



phys.org



Lentz et al., 2015



How do **we** study CCSNe?

thornado (<https://github.com/endeve/thornado>)

- toolkit for **high-order neutrino-radiation hydrodynamics**
- Hydrodynamics
 - Discontinuous Galerkin methods
 - Explicit Runge-Kutta methods
- Neutrino Transport
 - Discontinuous Galerkin methods
 - Implicit/Explicit Runge-Kutta methods
- Gravity
 - Continuous Galerkin (finite-element) methods
- Tabulated microphysics from weaklib (<https://github.com/starkiller-astro/weaklib>)
- Parallelism via AMReX (<https://amrex-codes.github.io/amrex/>)

Why GR?

- Results of GR (e.g., Bruenn et al., (2001), Müller et al., (2012))
 - Increased compactness of proto-neutron star
 - Decreased region between gain radius and shock
 - Increased post-shock fluid speed
- Effects
 - Harder neutrino spectra
 - Higher neutrino luminosities
 - Changes in neutrino heating
 - Changes in convection and turbulence

3+1 Hydrodynamics Equations

$$\partial_t (\sqrt{\gamma} \mathbf{U}) + \partial_i (\alpha \sqrt{\gamma} \mathbf{F}^i (\mathbf{U})) = \alpha \sqrt{\gamma} \mathbf{S} (\mathbf{U})$$

$$\begin{pmatrix} \rho \\ v^j \\ e \\ n_e \end{pmatrix} \longrightarrow \begin{pmatrix} D \\ S_j \\ \tau \\ N_e \end{pmatrix} \quad \mathbf{U} = \begin{pmatrix} D \\ S_j \\ \tau \\ N_e \end{pmatrix} = \begin{pmatrix} \rho W \\ \rho h W^2 v_j \\ \rho h W^2 - p - \rho W \\ n_e W \end{pmatrix}$$

$$W = (1 - v^i v_i)^{-1/2}$$

$$\rho h = \rho + e + p$$

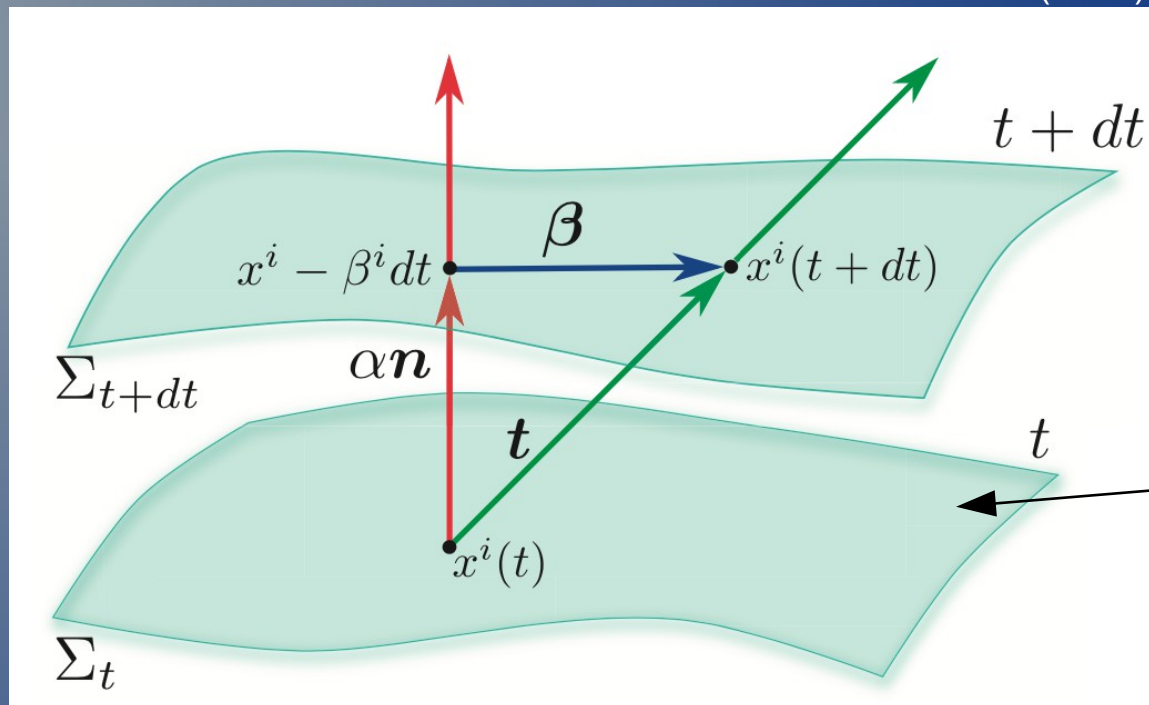
$$P^{ij} = \rho h W^2 v^i v^j + p \gamma^{ij}$$

$$\mathbf{F}^i = \begin{pmatrix} F_D \\ F_{S_j}^i \\ F_\tau \\ F_{N_e} \end{pmatrix} = \begin{pmatrix} D (v^i - \alpha^{-1} \beta^i) \\ P_j^i - \alpha^{-1} \beta^i S_j \\ S^i - D v^i - \alpha^{-1} \beta^i \tau \\ N_e (v^i - \alpha^{-1} \beta^i) \end{pmatrix}$$

$$\mathbf{S} = \begin{pmatrix} 0 \\ \frac{1}{2} P^{ik} \partial_j \gamma_{ik} + \alpha^{-1} S_i \partial_j \beta^i - \alpha^{-1} (\tau + D) \partial_j \alpha \\ P^{ij} K_{ij} - \alpha^{-1} S^j \partial_j \alpha \\ 0 \end{pmatrix}$$

3+1 GR

Rezzolla & Zanotti (2013)



$$\gamma_{ij} = \gamma_{ij}(x)$$

$$\alpha = \alpha(x) \quad \text{Lapse of proper time}$$

$$\beta^i = \beta^i(x) \quad \text{How much coordinates have shifted}$$

$$\alpha, \beta^i \quad \text{Freely specifiable}$$

$$K_{ij} = K_{ij}(\alpha, \beta^i, \gamma_{ij}) \quad \text{How the 3D surface is curved in the 4D manifold}$$

Conformally-Flat Approximation (CFA)

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij} (dx^i + \beta^i dt) (dx^j + \beta^j dt)$$

$$\gamma_{ij} = \psi^4 \text{diag}(\bar{\gamma}_{11}, \bar{\gamma}_{22}, \bar{\gamma}_{33})$$

- Eliminates dynamical degrees of freedom
- Exact in spherical symmetry
- Valid for slowly rotating progenitors (Dimmelmeier et al., (2002))

$$\psi = \psi(x)$$

$$\partial_t \bar{\gamma}_{ij} := 0$$

$$K^i_i = \gamma^{ij} K_{ij} := 0$$

$$\partial_t K^i_i := 0$$

“Maximal Slicing”

Gravity Solver

- Solves CFA equations (coupled system of elliptic PDEs)

$$\bar{\nabla}^2 \psi = - \left(2\pi \psi^5 E + \frac{1}{8} \psi^{-7} \bar{K}_{ij} \bar{K}^{ij} \right)$$

$$\bar{\nabla}^2 (\alpha \psi) = \alpha \psi \left(\frac{7}{8} \psi^{-8} \bar{K}_{ij} \bar{K}^{ij} + 2\pi \psi^4 (E + 2S) \right)$$

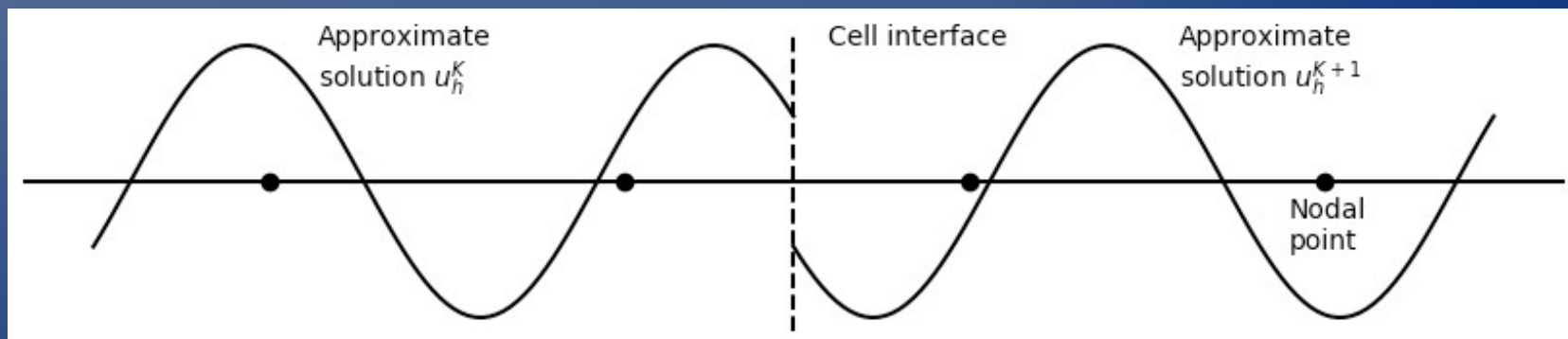
$$\bar{\nabla}^2 \beta^i = 16\pi \alpha \psi^4 S^i - \frac{1}{3} \bar{D}^j \bar{D}_k \beta^k + 2\alpha \psi^{-6} \bar{K}^{ij} \bar{D}_j (\alpha \psi^6)$$

- E, S, Sⁱ are derived from stress-energy tensor of fluid
- Uses hybrid of continuous Galerkin (finite element) and spectral methods

DG Method (1D)

- Discretize computational domain
- Locally approximate solution as polynomial
- Use weak form to evolve in time with SSP-RK methods
- Apply slope- and bound-preserving-limiters at each stage of SSP-RK algorithm

$$u(x, t) \approx u_h(x, t) := \sum_{i=1}^N u_i(t) \ell_i(x)$$



Weak Form (1D)

$$\partial_t \mathbf{U} + \frac{1}{\sqrt{\gamma}} \partial_r (\alpha \sqrt{\gamma} \mathbf{F}) = \mathbf{Q}(\mathbf{U})$$

- Multiply by test function (Lagrange polynomial)
- Integrate over element

Time-dependent metric term incorporated here

$$\int_K (\partial_t \mathbf{U}) \ell_i \sqrt{\gamma} dr = - \left\{ \left[\alpha \sqrt{\gamma} \ell_i(r) \hat{\mathbf{F}} \right]_{r_L}^{r_U} - \int_K \alpha \mathbf{F} \frac{d\ell_i}{dr} \sqrt{\gamma} dr - \int_K \mathbf{Q} \ell_i(r) \sqrt{\gamma} dr \right\}$$

Numerical flux obtained with approximate Riemann solver (e.g., HLL)

Substituting approximation in yields system of ODEs in

$$\mathbf{U}_h(r, t) = \sum_{i=1}^N \mathbf{U}_i(t) \ell_i(r)$$

$$\bar{\mathbf{U}}_h = (\mathbf{U}_1, \dots, \mathbf{U}_N)^T$$

$$\frac{d}{dt} \bar{\mathbf{U}}_h = \mathbf{L}(\bar{\mathbf{U}}_h)$$

Time-Stepping Algorithm

Strong-Stability-Preserving Runge-Kutta (SSP-RK)

Cockburn & Shu, (2001) J. Sci. Comp., Vol. 16, No. 3

$$\frac{d}{dt} \bar{U}_h = L(\bar{U}_h)$$

- Convex combinations of forward-Euler steps

1. $\bar{U}_h^{(0)} = \bar{U}_h^n$

2. For $i = 1, \dots, N_s$ compute the intermediate functions:

$$\bar{U}_h^{(i)} = \Lambda \Pi \left(\sum_{j=0}^{i-1} \alpha_{ij} w_h^{ij} \right), \quad w_h^{ij} = \bar{U}_h^{(j)} + \frac{\beta_{ij}}{\alpha_{ij}} \Delta t^n L(\bar{U}_h^{(j)})$$

3. $\bar{U}_h^{(n+1)} = \bar{U}_h^{N_s}$

$$\alpha_{ij}, \beta_{ij} > 0, \quad \sum_{j=0}^{i-1} \alpha_{ij} = 1$$

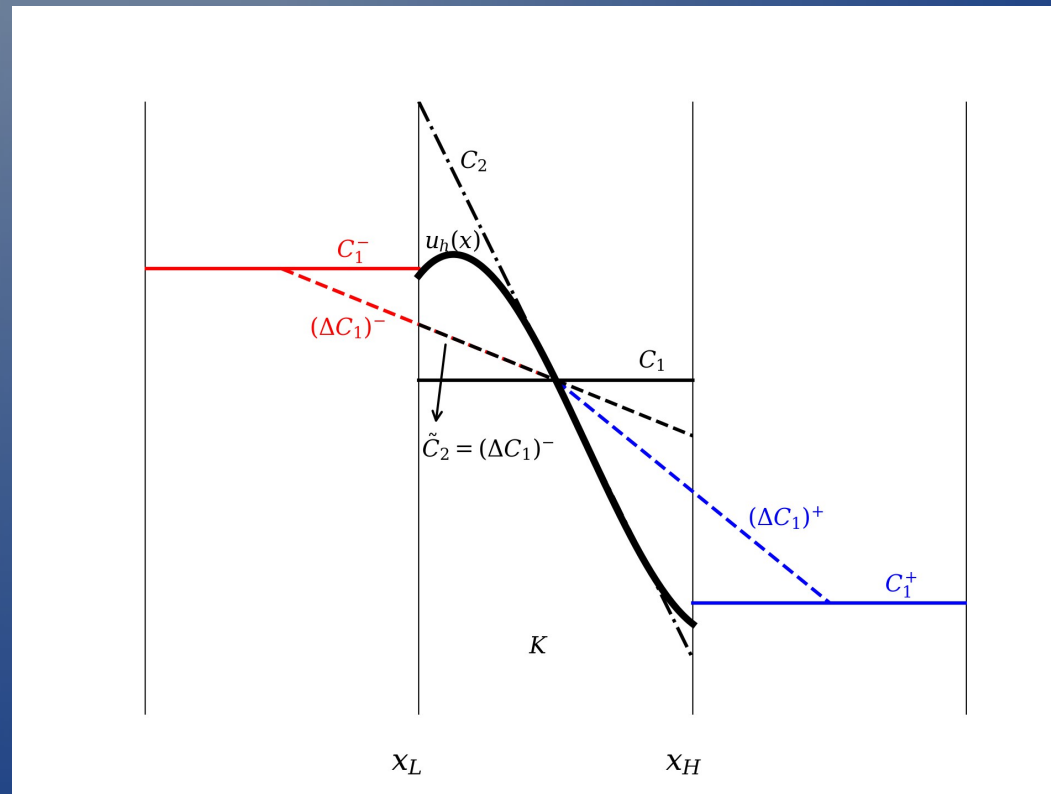
Limiters

Slope-Limiter

- Higher-order polynomials can develop unphysical oscillations
- MinMod limiter
 - Map to spectral representation
 - Filter out high-order terms
 - Map back

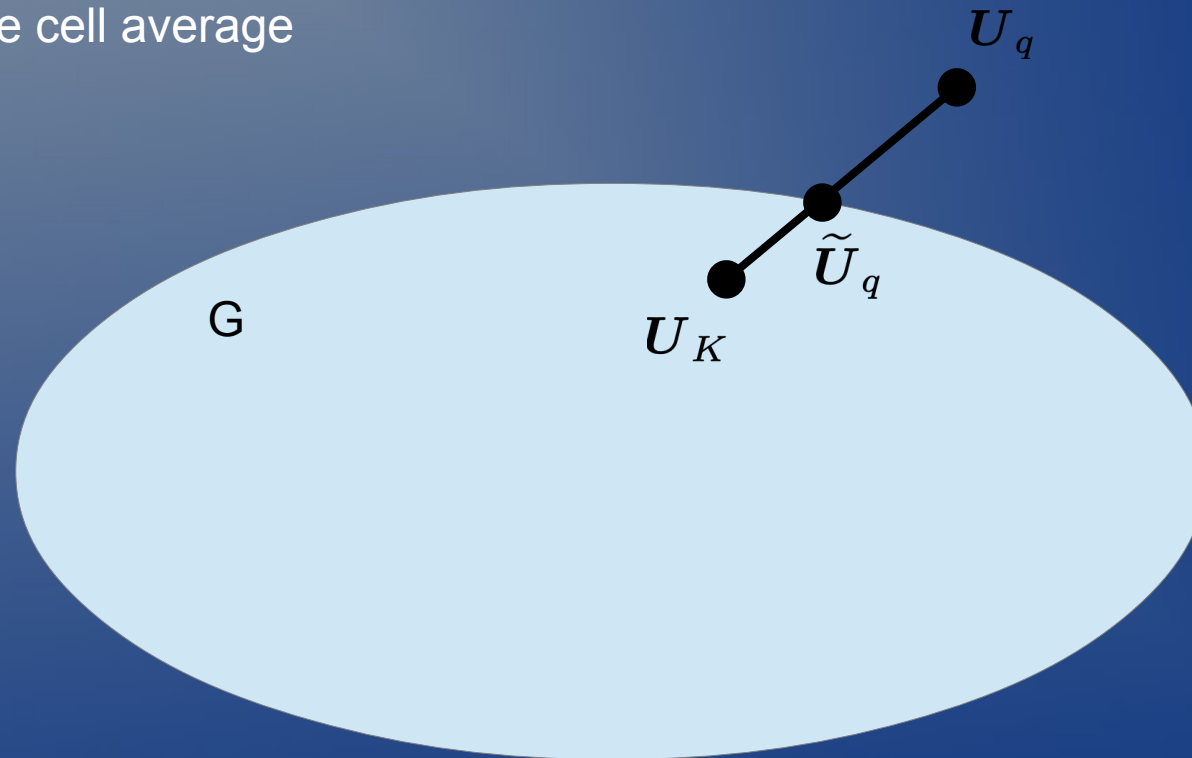
$$u_h(x, t) = \sum_{i=1}^N u_i(t) \ell_i(x) = \sum_{n=1}^N c_n(t) P_n(x)$$

Pochik et al., (2020)



Bound-Preserving Limiter

- The cell-averages are guaranteed to be physical, but not the nodal points (Qin et al. (2016))
- “Set of admissible states”, G , forms a convex set
- Unphysical nodal points are damped toward the cell average



AMReX

<https://amrex-codes.github.io/amrex/>

- Parallel/AMR framework developed primarily from Lawrence Berkeley National Laboratory
- Block-Structured AMR

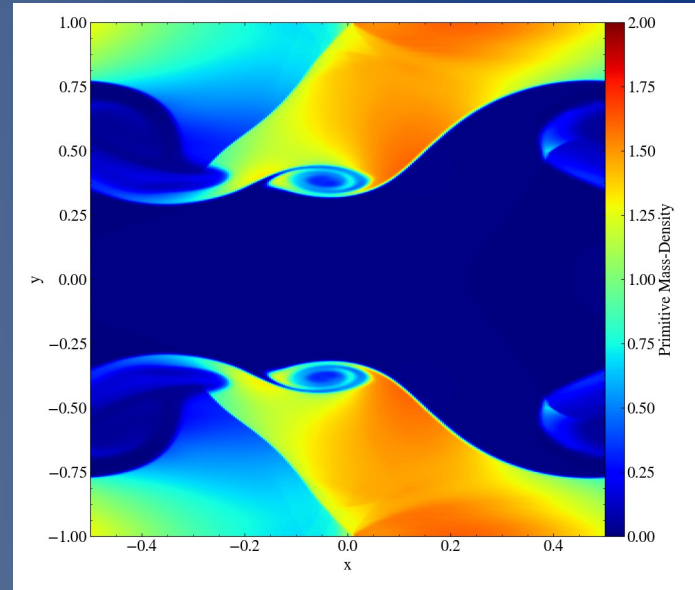


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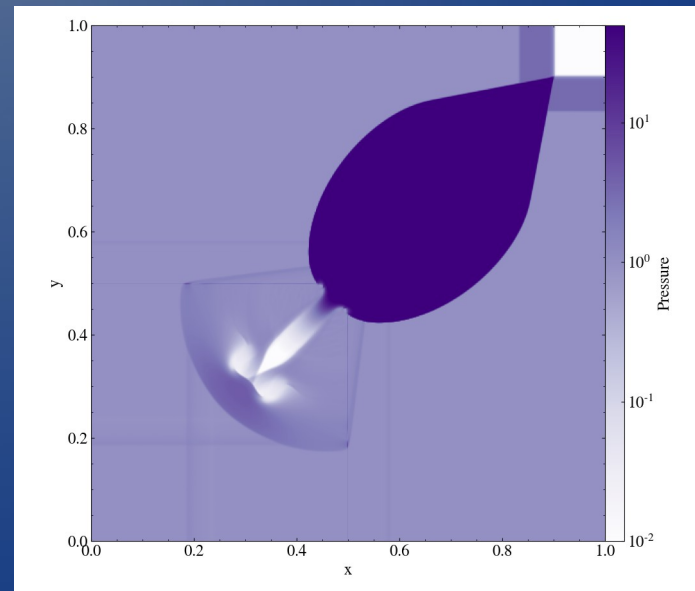
Results

Results

- Kelvin—Helmholtz instability
- Initial conditions from Radice & Rezzolla (2012)
- 256 x 512 elements
- 3rd order methods
- Run with AMReX

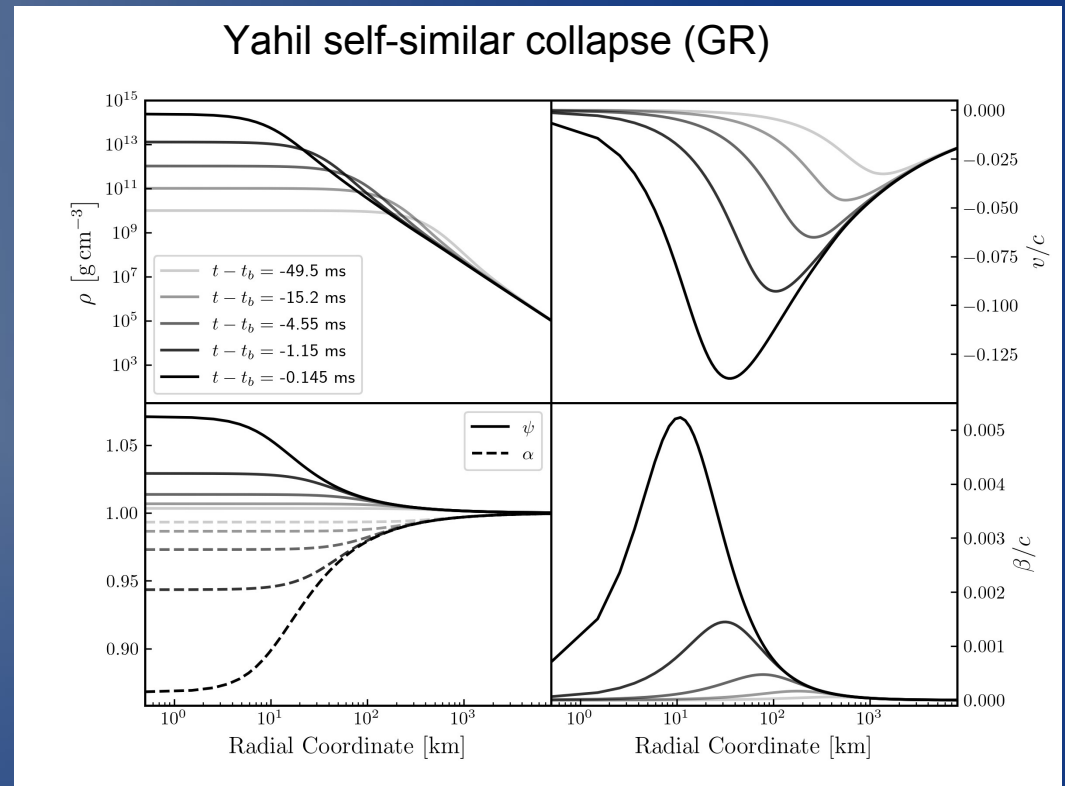


- 2D Riemann Problem
- Initial conditions from Del Zanna & Bucciantini (2002)
- 512 x 512 elements
- 3rd order methods
- Lorentz factor ~ 7
- Run with AMReX



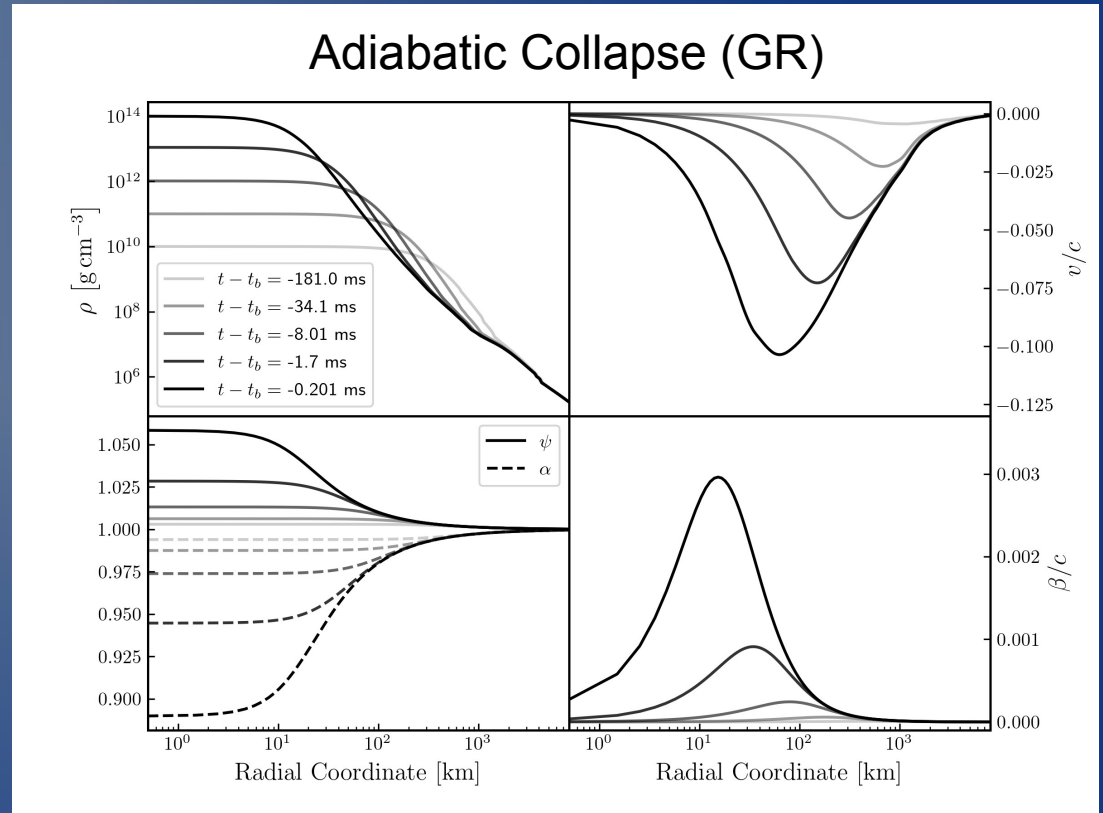
Results

- Polytropic self-similar collapse (Yahil, A. (1983))
 - $p = K \rho^\Gamma$
 - $\Gamma = 1.3$
- Simulates homologous collapse of stellar core
- Evolved from $t - t_b = -150$ ms until central density exceeds $1e15 \text{ g/cm}^3$
- 3rd order methods
- Geometric grid
 - 256 elements
 - $1e5$ km
 - 1 km inner cell width



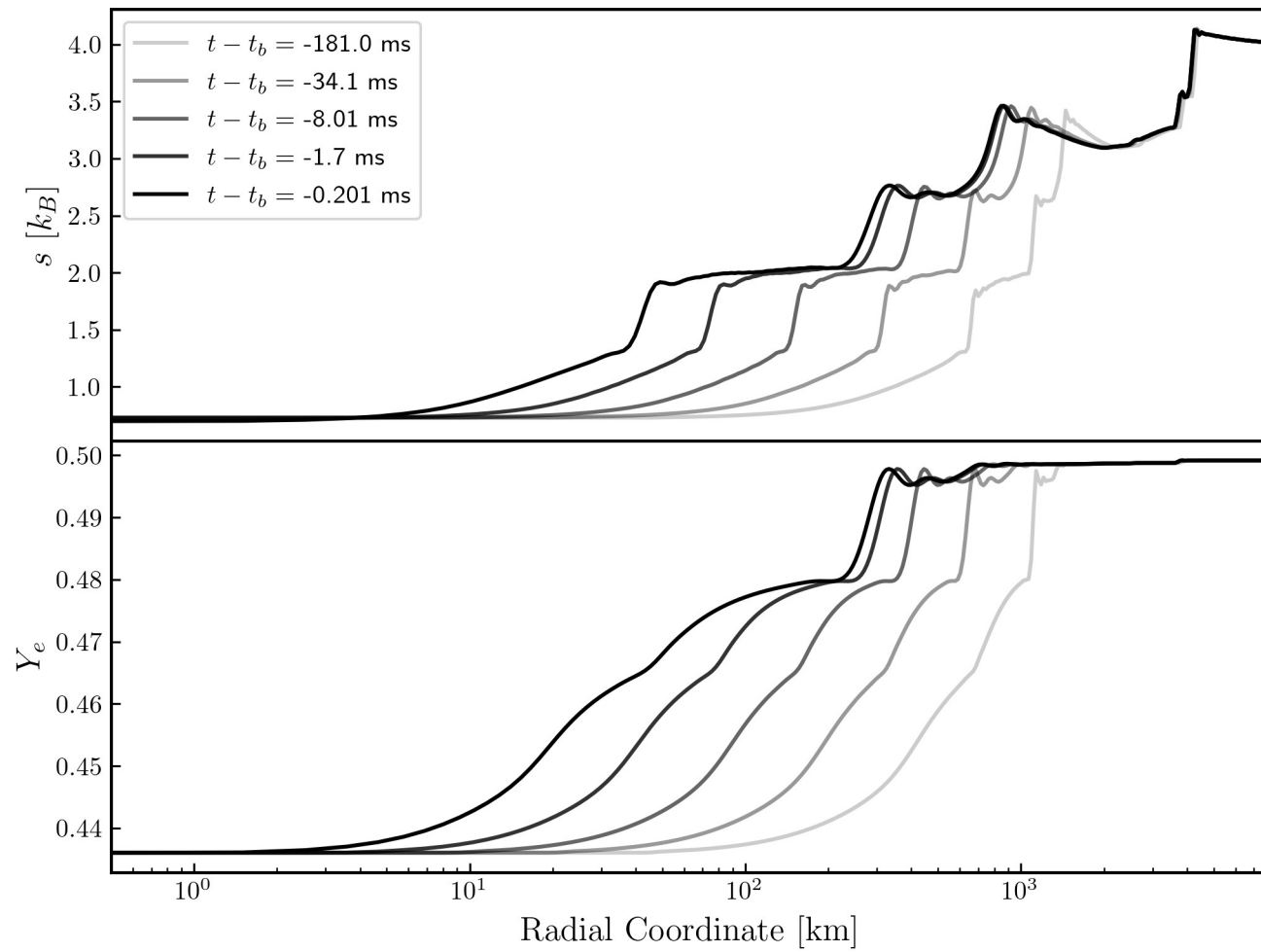
Results

- Adiabatic collapse of 15 Solar mass progenitor (Woosley & Heger (2007))
- SFHo EOS (Steiner et al., (2013a))
- 3rd order SSPRK method
- 2nd order DG method
- Geometric grid
 - 256 elements
 - 8e3 km
 - 1 km inner cell width



Compare to Newtonian results
in Pochik et al., (2020)

Results



Outline

Summary, etc.

Summary

- Developing thornado for CCSN simulations
 - Relativistic hydrodynamics
 - Relativistic gravity
 - Tabulated, nuclear EOS
- Evolved polytropic self-similar collapse
- Evolved adiabatic collapse of realistic 15 Solar mass progenitor with nuclear EOS

Ongoing/Future Work

- Evolve adiabatic collapse further
- Multi-D test simulations
- Incorporate neutrino transport solver
- Port/optimize hydro solver on GPUs
- Implements mesh refinement via AMReX
- CCSN simulations!

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