

3.2.3 WeakLib Opacities

Opacity tables with all electron-type neutrino and antineutrino interaction rates based on (Bruenn, 1985) had been created using **WeakLib**. (And bremsstrahlung) Following Bruenn using a multigroup flux-limited diffusion approximation (MGFLDA), the Boltzmann transport equations of the 0th-order moment and the 1st-order moment are (Equation (A12) in Bruenn (1985))

$$\begin{aligned} & \frac{1}{c} \frac{\partial}{\partial t} \psi^{(0)} + \frac{1}{3r^2} \frac{\partial}{\partial r} \left(r^2 \psi^{(1)} \right) + \frac{1}{3c} \frac{\partial \ln \rho}{\partial t} \left(\omega \frac{\partial}{\partial \omega} \psi^{(0)} \right) \\ & = j(\omega) \left(1 - \psi^{(0)} \right) - \Psi^{(0)} / \lambda^{(a)}(\omega) + A^{(0)}(\omega) \psi^{(0)} + B^{(0)}(\omega) \psi^{(1)} + C^{(0)}(\omega), \end{aligned} \quad (3.36)$$

and (A13)

$$\begin{aligned} & \frac{\partial}{\partial r} \psi^{(0)} + \left[-\frac{2}{c} \frac{\partial \ln \rho}{\partial t} \psi^{(1)} - \frac{2v}{cr} \psi^{(1)} + \frac{3}{5c} \frac{\partial \ln \rho}{\partial t} \frac{1}{\omega^3} \frac{\partial}{\partial \omega} \left(\omega^4 \psi^{(1)} \right) + \frac{4}{5} \frac{v}{cr} \frac{1}{\omega^3} \frac{\partial}{\partial \omega} \left(\omega^4 \psi^{(1)} \right) \right] \\ & = - \left[j(\omega) + 1/\lambda^{(a)}(\omega) \right] \psi^{(1)} + A^{(1)}(\omega) \psi^{(0)} + B^{(1)}(\omega) \psi^{(1)} + C^{(1)}(\omega), \end{aligned} \quad (3.37)$$

respectively.

Neutrino Absorption on Nucleons and Nuclei (AbEm)

Neutrino absorptivity on nucleons and nuclei, which is $[j(\omega) + 1/\lambda^{(a)}(\omega)]$ on the right-hand-side of Equation (3.37), is tabulated in **WeakLib** AbEm table in unit of cm^{-1} . It's also known as inverse mean free path on free nucleons and nuclei. AbEm tables are four-dimensional table with (ω, ρ, T, Y_e) . To interpolate the AbEm opacity with a given ω array for (ρ_i, T_i, Y_{ei}) state, one can

1 **CALL** LogInterpolateSingleVariable_1D3D_Custom

The interpolated value, $[j(\omega) + 1/\lambda^{(a)}(\omega)]$, is plotted directly in example plot Figure 3.1.

Neutrino Isoenergetic Scattering on Nucleons and Nuclei (Iso)

The 0-th and 1-st order of Legendre coefficients of neutrino isoenergetic scattering kernel on nucleons and nuclei, which are $\Phi_{0,\text{IS}}$ and $\Phi_{1,\text{IS}}$ with $\frac{4\pi}{c(2\pi\hbar c)^3} \omega^2$ and the coefficients 1/2 and 3/2, respectively, are tabulated in **WeakLib** Iso table in unit of cm^{-1} . So that (Equation (A41) in Bruenn (1985))

$$B_{\text{IS}}^{(1)}(\omega) = -\frac{2\pi}{c(2\pi\hbar c)^3} \omega^2 [\Phi_{1,\text{IS}}(\omega) - \Phi_{0,\text{IS}}(\omega)], \quad (3.38)$$

which is the only non-zero term of $A_{\text{IS}}^{(\alpha)}(\omega)$, $B_{\text{IS}}^{(\alpha)}(\omega)$, and $C_{\text{IS}}^{(\alpha)}(\omega)$, is given by

$$\text{TableValue}_0 - \frac{\text{TableValue}_1}{3}.$$

Iso tables are five-dimensional table with $(\omega, l, \rho, T, Y_e)$, where l is the order index of the Legendre coefficients: $l = 1$ for 0-th order (TableValue_0) and $l = 2$ for 1-st order (TableValue_1). To interpolate the Iso opacity with a given ω array for (ρ_i, T_i, Y_{ei}) state, one can

1 **CALL** LogInterpolateSingleVariable_1D3D_Custom

once if only needs 0-th order, and twice if needs both 0-th and 1-st order Legendre coefficients of neutrino isoenergetic scattering kernel. $-B_{\text{IS}}^{(1)}(\omega)$ is plotted in example plot Figure 3.1.

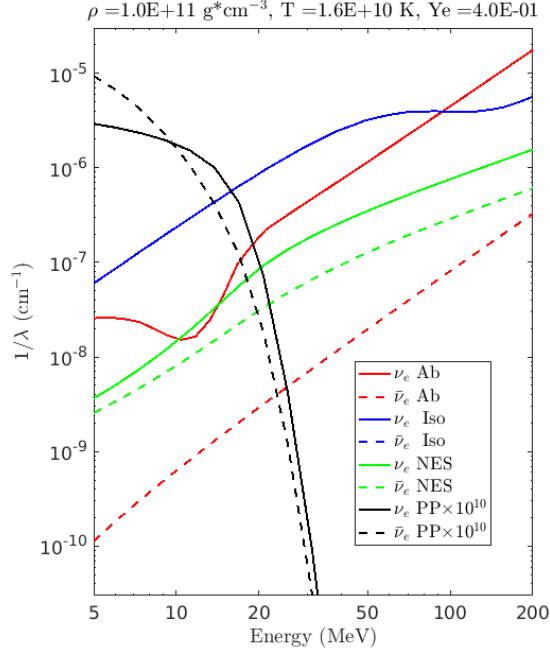


Figure 3.1: Example of WeakLib opacities. Adapted from Bruenn (1985) Figure 36(a).

Neutrino-Electron Scattering (NES)

The coefficient function H_l^I and H_l^{II} with $l = 0, 1$ of the first two Legendre coefficients of neutrino-electron scattering kernel with $\frac{2\pi}{c(2\pi\hbar c)^3}$ and $\int dE_e F_e(E_e) [1 - F_e(E_e + \omega - \omega')]$ are tabulated in WeakLib NES table. Following Bruenn (1985) Equation (C50),

$$\Phi_{l,\text{NES}}^{\left\{ \begin{smallmatrix} \text{in} \\ \text{out} \end{smallmatrix} \right\}} = \frac{G^2}{\pi\omega^2\omega'^2} \int dE_e F_e(E_e) [1 - F_e(E_e + \omega - \omega')] \left\{ \begin{array}{c} \exp[-\beta(\omega - \omega')] \\ 1 \end{array} \right\} \times \left[(C_V + C_A)^2 H_l^I(\omega, \omega', E_e) + (C_V - C_A)^2 H_l^{II}(\omega, \omega', E_e) \right], \quad (3.39)$$

and Equation (A38)

$$B_{\text{NES}}^{(1)}(\omega) = -\frac{2\pi}{c(2\pi\hbar c)^3} \int_0^\infty \omega'^2 d\omega' \left\{ \Phi_{0,\text{NES}}^{\text{in}}(\omega, \omega') \psi^{(0)}(\omega') + \Phi_{0,\text{NES}}^{\text{out}}(\omega, \omega') [1 - \psi^{(0)}(\omega')] \right\}. \quad (3.40)$$

NES tabs are five-dimensional table with $(\omega', \omega, l, T, \eta)$, where $\eta = \mu_e/k_B T$. To interpolate the NES opacity with a same given energy array for ω and ω' and state (ρ_i, T_i, Y_{ei}) , one needs to interpolate for η first. Then interpolate for $H_l^{I/II}$ by

1 **CALL** LogInterpolateSingleVariable_2D2D_Custom

Next, he/she needs to assemble $\Phi_{l,\text{NES}}^{\text{In/Out}}$ using the corresponding coefficients C_V and C_A and integrate over electron energy E_e . $-B_{\text{NES}}^{(1)}(\omega)$ is plotted with $\psi^{(0)}(\omega') = 0$ in example plot Figure 3.1.

Neutrino Production from Pair Process (Pair)

The coefficient function J_l^I and J_l^{II} with $l = 0, 1$ of the first two Legendre coefficients of pair process kernel with $\frac{2\pi}{c(2\pi\hbar c)^3}$ and $\int_0^{\omega+\omega^+} dE_e F_e(E_e) [1 - F_e(E_e)] [1 - F_{e^+}(\omega + \omega' - E_e)]$ (annihilation coefficient) are tabulated in **WeakLib** Pair table. Following [Bruenn \(1985\)](#) (A62), the Legendre coefficients of the pair process kernels are

$$\Phi_{l,\text{TP}}^{\left\{ \begin{smallmatrix} p \\ a \end{smallmatrix} \right\}} = \frac{2G^2}{2\pi} \int_0^{\omega+\omega^+} dE_e \left\{ \begin{array}{c} F_e(E_e) F_{e^+}(\omega + \omega' - E_e) \\ [1 - F_e(E_e)] [1 - F_{e^+}(\omega + \omega' - E_e)] \end{array} \right\} \times (C_V + C_A)^2 J_l^I(\omega, \omega', E_e) + (C_V - C_A)^2 J_l^{II}(\omega, \omega', E_e) \quad (3.41)$$

and (A47)

$$B_{\text{TP}}^{(1)}(\omega) = -\frac{2\pi}{c(2\pi\hbar c)^3} \int_0^\infty \omega'^2 d\omega' \left\{ \Phi_{0,\text{TP}}^p(\omega, \omega') [1 - \bar{\psi}^{(0)}(\omega')] + \Phi_{0,\text{TP}}^a(\omega, \omega') \bar{\psi}^{(0)}(\omega') \right\}. \quad (3.42)$$

Same as NES table, Pair tabs are five-dimensional table with $(\omega', \omega, l, T, \eta)$. And the interpolation is also the same as NES: find the η first,

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then assemble and integral. $-B_{\text{TP}}^{(1)}(\omega)$ is plotted with $\bar{\psi}^{(0)}(\omega') = 0$ in example plot [Figure 3.1](#).

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