## Relativistic Electrodynamics, Workshop 3

## Relativistic Addition of 3-Velocities

In this exercise you will derive the relativistic rule for combining velocities. In Galilean mechanics the rule is simple: The velocity measured by observer A,  $\underline{v}_A$ , compared to that measured by observer B,  $\underline{v}_B$ , is  $\underline{v}_A = \underline{v}_B + \underline{v}_{B,A}$ , where  $\underline{v}_{B,A}$  is the velocity of observer B measured by observer A. This obviously cannot hold in special relativity, since faster-than-light velocities could be obtained this way.

Consider two frames of inertia S and S', with S' moving along the direction of x-axis with velocity u, as measured by an observer in S. An observer at rest in S' measures the 3-velocity of a trajectory to be v'.

1. Use the relation between the 0-component of the 4-velocities  $v^{\mu}, v'^{\mu}$  to derive the following relation

$$\gamma(v) = \gamma(u) \ \gamma(v') \left( 1 + \frac{v_x' u}{c^2} \right) \tag{1}$$

2. Find the 3-velocity, v, measured by an observer in S, by explicitly constructing the 4-velocity measured in S' and transform this into S. The result is

$$v_x = \frac{u + v_x'}{1 + uv_x'/c^2} \quad v_y = \frac{v_y'}{\gamma(u)(1 + uv_x'/c^2)} \quad v_z = \frac{v_z'}{\gamma(u)(1 + uv_x'/c^2)}.$$
 (2)

3. Show that for small velocities  $|u|, |\underline{v}'| \ll c$  this reduces to the simple Galilean result.

Note that for collinear velocities, i.e.  $\underline{v}$  aligned with  $\underline{u}$  (the x-axis), the result is symmetric under exchanging u and  $v_x'$ , just as the Galilean result is symmetric. However, when  $\underline{v}$  is not along the boost-axis, then the result is not symmetric, i.e. in general the velocity of a trajectory is different whether it is measured as  $\underline{v}_1$  by an observer moving relative to you with velocity  $\underline{v}_2$ , as opposed to a measurement of  $\underline{v}_2$  by an observer moving relative to you with velocity  $\underline{v}_1$ .

## Solution

The Lorentz transformation going from S' to S is given by

$$\Lambda_{S' \to S^{\mu}_{\nu}} = \begin{pmatrix} \gamma(u) & \gamma(u)u/c & 0 & 0\\ \gamma(u)u/c & \gamma(u) & 0 & 0\\ 0 & 0 & 1 & 0\\ 0 & 0 & 0 & 1 \end{pmatrix}.$$
 (3)

The four-velocities are given by  $v^{\mu}=\gamma(v)(c,\underline{v})$  etc. Therefore, the relationship between  $v^0$  and  $v'^0$  is

$$\gamma(v)c = \gamma(u) \ \gamma(v') \ (c + v_x'u/c) = \gamma(u) \ \gamma(v') \ (1 + v_x'u/c^2) \ c, \tag{4}$$

and therefore

$$\gamma(v) = \gamma(u) \ \gamma(v') \left( 1 + v_x' u/c^2 \right). \tag{5}$$

The transformation of the x-component of the 3-velocities we find by studying the Lorentz-transformation of the 1-component of the 4-vector:

$$\gamma(v)c_x = \gamma(v')\gamma(u)(uc/c + v_x') = \gamma(v')\gamma(u)(v_x' + u) \tag{6}$$

$$\therefore v_x = \frac{v_x' + u}{1 + v_x' u/c^2},\tag{7}$$

using Eq. (5).

For the y and z components we find

$$\gamma(v)v_y = \gamma(v') \ v_y' \tag{8}$$

$$\gamma(v)v_y = \gamma(v') \ v'_y$$

$$v_y = \frac{v'_y}{\gamma(u)(1 + v'_x u/c^2)},$$
(8)

and similarly for  $v_z$ .

Expanding for  $u/c \ll 1, v_x'/c \ll 1$  and keeping just the first order in u and  $v_x'$  we

$$\underline{v} = \underline{u} + \underline{v}'$$
 non-relativistic limit!! (10)