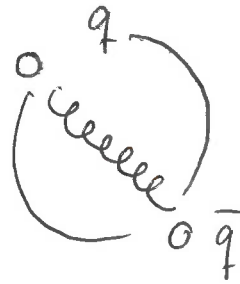


Quarkonia

Positronium



held together by
the e/m force



held together by
the strong force

Compare the spectra

- small distances ~ low radial wavefunctions (small n)
the spectra look similar

For small distances $V(r) \propto \frac{1}{r}$

- large distances ~ spectrum drops like $E_n \propto \frac{1}{n^2}$ for
positronium, not so for charmonium

For large distances $V(r)$ does not scale like $\frac{1}{r}$

- Spin-induced energy splitting is much larger for
quarkonia

The split between these states originates from the
Spin-spin interaction

$$V_{ss}(e^+e^-) = \frac{8\pi}{3} \propto \frac{\vec{S}_1 \cdot \vec{S}_2}{m_e^2} \delta(\vec{r})$$

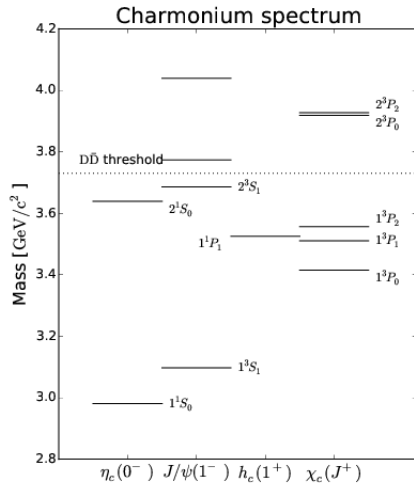
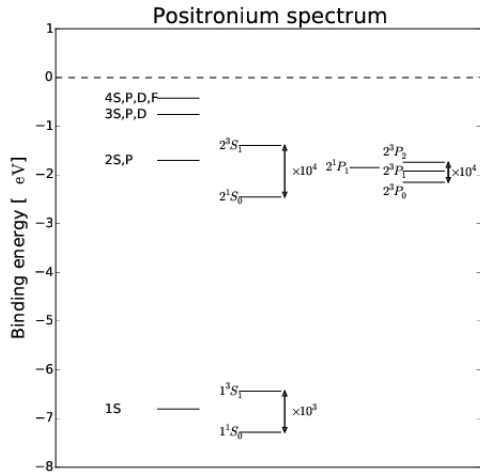
Fine structure constant spin vectors overlap with the origin

$\alpha = \frac{e^2}{4\pi}$

For quarkonium

$$V_{ss}(q\bar{q}) = \frac{32\pi}{8} \propto_s \frac{\vec{S}_q \cdot \vec{S}_{\bar{q}}}{m_q m_{\bar{q}}} \delta(\vec{r})$$

strong coupling constant



Calculate the spin-induced splitting:

we will need $\langle \vec{S}_1 \cdot \vec{S}_2 \rangle$

$$S(S+1) = \langle (\vec{S}_1 + \vec{S}_2)^2 \rangle$$

$$= \langle \vec{S}_1^2 \rangle + 2 \langle \vec{S}_1 \cdot \vec{S}_2 \rangle + \langle \vec{S}_2^2 \rangle$$

$$= \frac{1}{2} \left(\frac{1}{2} + 1 \right) + 2 \langle \vec{S}_1 \cdot \vec{S}_2 \rangle + \frac{1}{2} \left(\frac{1}{2} + 1 \right)$$

q and \bar{q} have spin $\frac{1}{2}$

$$= \frac{3}{2} + 2 \langle \vec{S}_1 \cdot \vec{S}_2 \rangle$$

$$\Rightarrow \langle \vec{S}_1 \cdot \vec{S}_2 \rangle = \frac{2S(S+1) - 3}{4} = \begin{cases} -\frac{3}{4} & \text{for } S=0 \\ +\frac{1}{4} & \text{for } S=1 \end{cases}$$

we get for the contribution E_{ss} to the states was

$$E_{ss} = \langle \psi | V_{ss} | \psi \rangle = \frac{8\pi \alpha_s}{g m_q m_{\bar{q}}} \langle \psi | \delta(\vec{r}) | \psi \rangle \cdot \langle \vec{S}_1 \cdot \vec{S}_2 \rangle$$

$$= \frac{8\pi \alpha_s}{g m_q m_{\bar{q}}} |\psi(0)|^2 \cdot \begin{cases} -3 & \text{for } S=0 \\ 1 & \text{for } S=1 \end{cases}$$

The mass split is therefore

$$\Delta E_{ss} = E_{ss}(S=1) - E_{ss}(S=0) = \frac{32\pi}{g} \frac{\alpha_s}{m_q m_{\bar{q}}} |\psi(0)|^2$$

Here m_q and $m_{\bar{q}}$ are "constituent quark masses," including the binding energy.

Quark - antiquark potential

The $q\bar{q}$ potential deviates from the positronium potential for large distances and is similar for small distances. Since there is no suppression for large r the potential appears to be unbounded:

$$V(r) = \underbrace{-\frac{4}{3} \frac{\alpha_s(r)}{r}}_{\text{short range}} + \underbrace{kr}_{\text{long-range}}$$

For

$$\alpha_s \approx 0.15 - 0.25$$

$$k \approx 1 \text{ GeV/fm}$$

$$m_c \approx 1.5 \text{ GeV}$$

we get a good fit to the charmonium spectrum.

Quarkonia decays

There are several ways for quarkonium states to decay:

- E/M decays
- Annihilation
- Decay through $q\bar{q}$ production
- Weak decays

E/M decays

The parity for
a quarkonium state

$n^{2S+1}L_J$ is given by

$$P(n^{2S+1}L_J) = \underbrace{P(q)P(\bar{q})}_{-1} (-1)^L = (-1)^{L+1}$$

(fermions and antifermions have opposite parity)

Electric \uparrow have parity $(-1)^L$

Magnetic transitions have parity $(-1)^{L+1}$

E.g. for $L=1$

$$2^3S_1 \rightarrow 1^3P_3$$

$$J = 0, 1, 2$$

$$\Rightarrow J_i = 1$$

$$J_f :$$

$$|J_i - L| \leq J_f \leq J_i + L$$

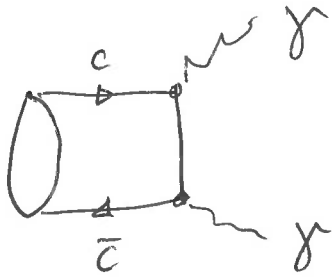
$$\text{Since } P(2^3S_1) = (-1)^{0+1} = -1 \quad \left. \begin{array}{l} \\ \\ \end{array} \right\} \text{EI}$$

$$P(1^3P_3) = (-1)^{1+1} = +1$$

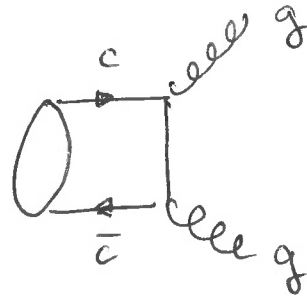
χ - Annihilation

$$J=0$$

$$\eta_c \rightarrow \gamma\gamma$$

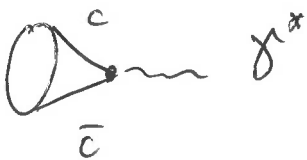


$$\eta_c \rightarrow gg$$

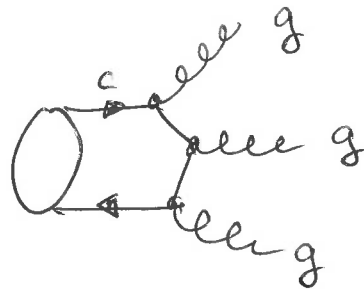


$$J=1$$

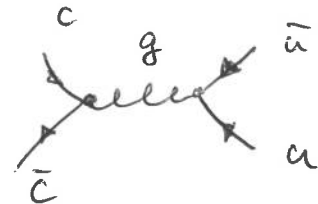
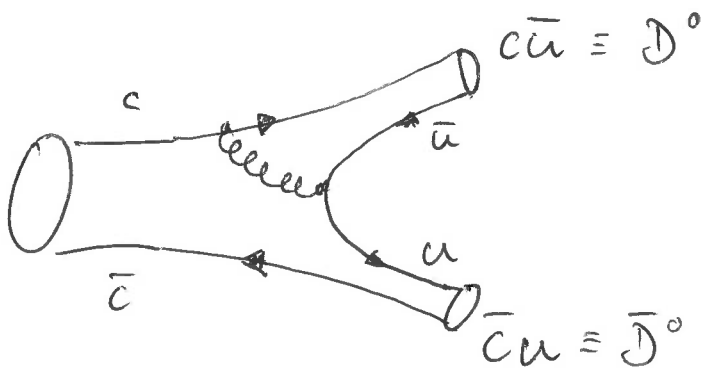
$$J/\psi \rightarrow \gamma^*$$



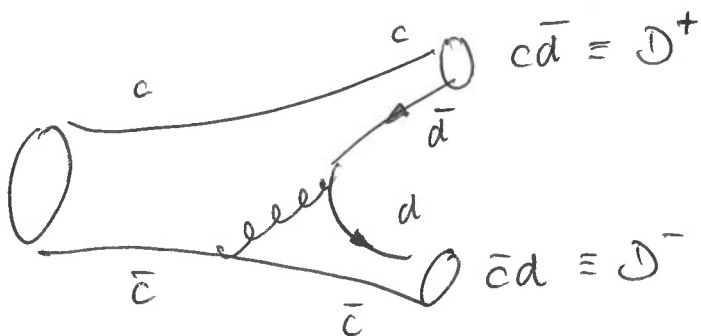
$$J/\psi \rightarrow ggg$$



- $q\bar{q}$ production



$$c\bar{c} \rightarrow D^0 + \bar{D}^0$$



$$c\bar{c} \rightarrow D^+ + D^-$$