Mathematical Methods in Physics

Examination 2015

Question 1

(a) (Unseen)

(i) Linearly dependent. [1 mark]

(ii) Linearly dependent. [1 mark]

(iii) Linearly dependent. [1 mark]

(iv) Linearly independent. [1 mark]

In order to established whether they are sets of linearly independent or dependent vectors, student can use determinant - if applicable -, inspection - for instance in (i) the third vector is the sum of the first two vectors hence they are not linearly independent - or use the definition of what a set of linearly independent vectors is.

(b) (Unseen)

(i)
$$A = A^{\dagger}, \qquad U^{-1} = U^{\dagger}, \text{ then }$$

$$(U^{-1}AU)^{\dagger} = (U^{\dagger}AU)^{\dagger} = U^{\dagger}A^{\dagger}U = U^{-1}AU.$$

[2 marks]

(ii)
$$U^{-1} = U^{\dagger}$$
 with $U = A + iB$, $A = A^{\dagger}$, $B = B^{\dagger}$ and $AB = BA$ then

$$I = UU^{\dagger} = (A + iB)(A - iB) = A^2 + B^2.$$

[2 marks]

(c) (Unseen)

$$\left(\frac{ds}{dt}\right)^2 = \frac{d\underline{r}}{dt} \cdot \frac{d\underline{r}}{dt} = \left(1 + \frac{9}{4}t\right).$$

[2 marks]

Then

$$L = \int_{0}^{4} \sqrt{\left(1 + \frac{9t}{4}\right)} dt = \frac{8}{27} \left[\left(1 + \frac{9}{4}t\right)^{3/2} \right]_{0}^{4} = \frac{8}{27} \left(10^{3/2} - 1\right).$$

[2 marks]

(d) (Unseen)

(i)
$$\nabla \times \underline{F} = 0$$
. The field is conservative. [1 mark] Then

$$\begin{array}{lcl} \frac{\partial f}{\partial x} & = & 2xy^3z^4 \longrightarrow f = x^2y^3z^4 + g(y,z), \\ \frac{\partial f}{\partial y} & = & 3x^2y^2z^4 \longrightarrow g(y,z) = g(z), \\ \frac{\partial f}{\partial z} & = & 4x^2y^3z^3 \longrightarrow g(z) = c. \end{array}$$

Then $f = x^2y^3z^4 + c$ where c is a constant. [2 marks]

(ii)
$$\nabla \times \underline{F} = (xy\,\hat{\underline{i}} - yz\,\hat{\underline{j}} - 2\,\hat{\underline{k}}) \neq 0$$
. The field is not conservative. [1 mark]

(e) (Unseen)

A suitable parametrisation for the path is:

$$\underline{r}(t) = 2\cos t\,\hat{\underline{i}} + 2\sin t\,\hat{\underline{j}}, \qquad 0 \le t \le \pi.$$

[2 marks]

Then

$$I = \int_{0}^{\pi} \underline{F} \cdot \underline{r}' dt = 6 \int_{0}^{\pi} (-2^{4} \cos^{4} t \sin t + 2^{6} \sin^{6} t \cos t) dt$$
$$= 6 \left(-2^{4} \left[-\frac{\cos^{5} t}{5} \right]_{0}^{\pi} + 2^{6} \left[\frac{\sin^{7} t}{7} \right]_{0}^{\pi} \right) = -\frac{192}{5}.$$

[2 marks]

(f) (Unseen)

$$\nabla \times \underline{F} = -x^2 \, \hat{\underline{k}} = -(a \sin \theta \cos \phi)^2 \, \hat{\underline{k}}.$$
 [1 mark]

The surface element is:

$$d\underline{S} = \left(\frac{\partial \underline{r}}{\partial \theta} \times \frac{\partial \underline{r}}{\partial \phi}\right) = a \sin \theta \, \underline{r} \, d\theta d\phi.$$

[2 mark]

Then

$$I = -a^4 \int_0^{2\pi} \cos^2 \phi \, d\phi \int_0^{\pi/2} \sin^3 \theta \cos \theta \, d\theta = -a^4 \pi \left[\frac{\sin^4 \theta}{4} \right]_0^{\pi/2} = -\frac{a^4 \pi}{4}.$$
 [1 mark]

(g) (Unseen)

$$\hat{f}(w) = \frac{1}{2i} \frac{1}{\sqrt{2\pi}} \int_0^\infty \left(e^{t(-\beta + i\alpha - iw)} - e^{t(-\beta - i\alpha - iw)} \right) dt$$

$$= \frac{1}{2i} \frac{1}{\sqrt{2\pi}} \left[\frac{e^{t(-\beta + i\alpha - iw)}}{-\beta + i\alpha - iw} + \frac{e^{t(-\beta - i\alpha - iw)}}{\beta + i\alpha + iw} \right]_0^\infty = \frac{1}{\sqrt{2\pi}} \frac{\alpha}{\alpha^2 + (\beta + iw)^2}.$$
[4 marks]

(h) (Unseen)

(i)

$$I_1 = 2\pi \left(\int_{-\infty}^{\infty} \frac{\delta(x-\pi)}{|2\pi|} e^{2x} + \int_{-\infty}^{\infty} \frac{\delta(x+\pi)}{|-2\pi|} e^{2x} \right) + 2 = e^{2\pi} + e^{-2\pi} + 2 = (2\cosh\pi)^2.$$

[3 marks]

(ii)
$$I_1 = \frac{1}{\underline{a} \cdot \underline{b} - 2} = \frac{1}{3}.$$

[1 marks]

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Question 2

- (a) (Unseen)
 - (i) The eigenvalues are: $\lambda_1 = 1$, $\lambda_2 = 1 + (\alpha \beta)^{1/2}$ and $\lambda_3 = 1 (\alpha \beta)^{1/2}$. The forms of the corresponding eigenvectors are: $v_1 = (0, 0, z)^T$, $v_2 = (x, x(\beta/\alpha)^{1/2}, 0)^T$ and $v_3 = (x, -x(\beta/\alpha)^{1/2}, 0)^T$. Hence a possible choice for the eigenvectors is: $v_1 = (0, 0, 1)^T$, $v_2 = (1, (\beta/\alpha)^{1/2}, 0)^T$ and $v_3 = (1, -(\beta/\alpha)^{1/2}, 0)^T$. [4 marks]
 - (ii) The eigenvalues are real if the product $\alpha\beta$ is real and positive. The scalar product between pairs of different eigenvectors must be equal to zero in order for the eigenvectors to be orthogonal. This leads to the constraint $|\alpha| = |\beta|$. [3 marks]
 - (iii) The matrix A is Hermitian if $A = A^{\dagger}$. This leads to the constraint $\alpha^* = \beta$, which implies the constraints in (ii) since $|\alpha| = |\beta|$ and $\alpha\beta = \alpha\alpha^* = |\alpha|^2 > 0$. [3 marks]
- (b) (i) (Bookwork)

The divergence theorem states that

$$\int_{V} (\nabla \cdot \underline{F}) dV = \int_{S} \underline{F} \cdot d\underline{S}.$$

The integral on the left is the integral of the divergence of the vector field \underline{F} over the volume enclosed by the surface S. The integral of the right is the integral of the vector field \underline{F} over the surface S. The symbol ∇ is a vector differential operator and dS is a differential vector perpendicular to the surface S. [4 marks]

(ii) (Unseen)

For the integral on the left $\nabla \cdot F = 3x^2 + 2yz + 2z$. Hence the volume integral is:

$$\int_{-2}^{2} dy \int_{0}^{3} dz \int_{-1}^{1} 3x^{2} dx + \int_{-1}^{1} dx \left(\int_{0}^{3} dz \int_{-2}^{2} 2yz dy \right) + \int_{-1}^{1} dx \int_{-2}^{2} dy \int_{0}^{3} 2z dz$$

$$= 12 \left[x^{3} \right]_{-1}^{1} + 2 \left(\int_{0}^{3} dz \left[zy^{2} \right]_{-2}^{2} \right) + 8 \left[z^{2} \right]_{0}^{3} = 24 + 72 = 96.$$

[3 marks]

For the integral on the right there are six different $d\underline{S}$, they are:

$$d\underline{S} = \pm \hat{\underline{i}} \, dy dz, \quad d\underline{S} = \pm \hat{\underline{j}} \, dx dz, \quad d\underline{S} = \pm \hat{\underline{k}} \, dx dy.$$

Hence the integral becomes

$$\begin{split} &\int_{-2}^{2} dy \int_{0}^{3} dz \, x^{3} \mid_{x=1} - \int_{-2}^{2} dy \int_{0}^{3} dz \, x^{3} \mid_{x=-1} + \int_{-1}^{1} dx \int_{0}^{3} dz \, y^{2} z \mid_{y=2} \\ &- \int_{-1}^{1} dx \int_{0}^{3} dz \, y^{2} z \mid_{y=-2} + \int_{-1}^{1} dx \int_{-2}^{2} dy \, z^{2} \mid_{z=3} - \int_{-1}^{1} dx \int_{-2}^{2} dy \, z^{2} \mid_{z=0} = 96. \end{split}$$
 [3 marks]

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Question 3

- (a) (Unseen)
 - (i) The even extension between 0 and $-\pi$ corresponds to the function $f(x) = -\sin x$. The period is 2π . The Fourier coefficients are:

$$a_0 = \frac{1}{\pi} 2 \int_0^{\pi} \sin x \, dx = \frac{4}{\pi}.$$

[1 mark]

$$a_n = \frac{1}{\pi} 2 \int_0^{\pi} \sin x \cos(nx) dx = \frac{1}{\pi} \int_0^{\pi} (\sin(1+n)x + \sin(1-n)x) dx$$
$$= -\frac{1}{\pi} \frac{((-1)^{1+n} - 1)}{1+n} - \frac{1}{\pi} \frac{((-1)^{1-n} - 1)}{1-n} = -\frac{4}{\pi(n^2 - 1)} \quad \text{for } n \text{ even.}$$

[4 marks]

Hence

$$|\sin x| = \frac{2}{\pi} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{\cos(2nx)}{4n^2 - 1}.$$

[1 mark]

(ii) For
$$x = 0$$

$$0 = \frac{2}{\pi} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{4n^2 - 1},$$

hence

$$\sum_{n=1}^{\infty} \frac{1}{4n^2 - 1} = \frac{1}{2}.$$

For
$$x = \pi/2$$

$$1 = \frac{2}{\pi} - \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{4n^2 - 1}.$$

The sum of this expression and the previous one gives

$$\sum_{n=1}^{\infty} \frac{1}{16n^2 - 1} = \frac{1}{2} - \frac{\pi}{8}.$$

[4 marks]

(b) (Unseen)

(i) $h(u) = \int_0^t ue^{-3(t-u)} du = -\frac{1}{9} + \frac{t}{3} + \frac{e^{-3t}}{9}.$

[3 marks]

(ii) Using the hint

$$\bar{h}(s) = -\frac{1}{9s} + \frac{1}{3s^2} + \frac{1}{9(s+3)} = \frac{1}{s^2(s+3)}.$$

On the other hand

$$\bar{f}(s) = \frac{1}{s^2}, \qquad \bar{g}(s) = \frac{1}{(s+3)},$$

hence $\bar{h}(s) = \bar{f}(s)\bar{g}(s)$.

[3 marks]

(iii) The roots of the polynomial at the denominator are 0, 2, -3. Therefore

$$\bar{m}(s) = \frac{A_1}{s} + \frac{A_2}{(s-2)} + \frac{A_3}{(s+3)},$$

with $A_1 = -1/6$, $A_2 = 3/10$ and $A_3 = -2/15$. Then

$$\mathcal{L}[\bar{m}(s)](t) = -\frac{1}{6} + \frac{3}{10}e^{2t} - \frac{2}{15}e^{-3t}.$$

[4 marks]

(a)
$$x^2 \frac{d^2y}{dx^2} + x \frac{dy}{dx} - 9 y = 5 x^2$$

The solution $(y \propto x^r)$ to the homogeneous equation

$$x^2 \frac{d^2y}{dx^2} + x \frac{dy}{dx} - 9 y = 0$$

is

$$y = Ax^3 + Bx^{-3}$$

[2 marks, bookwork]

Given the RHS, the smart ansatz thus needs to be of the form $y=D\,x^2$ leading to D=-1. [2 marks, unseen]

(b)
$$\frac{d^2y}{dx^2} - 2\frac{dy}{dx} + y = 6 \ x \ e^x$$

Using $y \propto e^{p\ x}$ to solve this equation, we find p=1 (double root), thus leading to

$$y = k_1 \ x \ e^x + k_2 \ e^x$$

[1 mark, bookwork]

The Wronskian is

$$W(x) = -e^{2x}$$

and therefore

$$k_1' = 6 x$$

and

$$k_2' = -6 x^2$$
.

[2 marks, unseen]

Integrating, we obtain $k_1 = 3 x^2 + A$ and $k_2 = -2 x^3 + B$ and therefore

$$y = (3 x^2 + A) x e^x + (-2 x^3 + B) e^x$$

[1 mark, unseen]

$$(1 - 2x)\frac{dy}{dx} + 2y = 0$$

immediately leads to

$$y = k \exp\left(-\int^x \frac{dx'}{1 - 2x'}\right)$$

with k a constant so we obtain

$$y = k (1 - 2x).$$

[2 marks, unseen]

The equation

$$(1 - 2x)\frac{dy}{dx} + 2y + y^3 = 0$$

is a Bernoulli equation (non linear) and we can use $v=y^{3-1}$ to solve it. [2 marks, bookwork]

(d)

$$\frac{d^4x}{du^4} - 16x = 0$$

 $x \propto e^{r y}$ leads to $r = \pm 2$ and $r = \pm 2i$ and therefore

$$x = \alpha e^{2y} + \beta e^{-2y} + \gamma e^{2iy} + \delta e^{-2iy}$$

[2 marks, unseen]

A possible RHS is of the exponential form, i.e.

$$\frac{d^4x}{dy^4} - 16x = Ce^{By}$$

in which case a solution is $y = D e^{By}$ with $D = C/(B^4 - 16)$ but we could have also

$$\frac{d^4x}{dy^4} - 16x = f_x(y)$$

with $f_x(y) = -16 \ A \ y^n$ and n < 4 (with $x = A \ y^n$ a solution of this equation). [2 marks, unseen]

(e) The Cartesian surface and lack of dissipation imply

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)\psi = C\frac{\partial^2\psi}{\partial t^2}$$

with C a coefficient. [2 marks, bookwork] Therefore the solution is of the form

$$y = A\cos(k(x+y) \pm wt) + B\sin(k(x+y) \pm wt)$$

[2 marks, bookwork]

(f) The generic expression of a Legendre polynomial of order n, where n is an odd number is

$$P_{l=n} = \sum_{i=odd}^{n} \alpha_i \ x^i$$

Therefore the polynomial

$$P_3(x) = \frac{1}{2^3 3!} \frac{d^3}{dx^3} (x^2 - 1)^3 = \frac{1}{2} (5x^3 - 3x)$$

which contains only odd powers of x and ends at x^3 is as expected. [4 marks,bookwork]

(g) The convention for spherical harmonics is Y_{lm} and the first (l=0), second (l=1), third (l=2) harmonics are called monopole (or s-wave), dipole (or p-wave), quadrupole (or d-wave). The quadrupole is a pair dipole and can therefore be represented as 2 orthogonal '8' shapes. [4 marks,bookwork]

(a) For the oscillations to be sustainable with time, must contain a second derivatives in t. Hence

$$k^2 \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2}$$

with k a constant. [2 marks,bookwork]

(b) $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} - k \frac{\partial u}{\partial t}$

contains a first order term $\frac{\partial u}{\partial t}$ which implies that there is dissipation with time. [2 marks,bookwork]

(c) We can solve the above equation using the separation of variables:

$$u = X(x) T(t)$$

[2 marks,bookwork]

(d) Using

$$u = X(x) T(t)$$

we find:

$$T''X = c^2 T X'' - k T' X$$

which gives

$$\frac{T''}{T} + k \frac{T'}{T} = c^2 \frac{X''}{X}$$

once we divide by u. [1 mark,bookwork]

From the above equation, we conclude that

$$\frac{T''}{T} + k \frac{T'}{T} = sign\lambda^2 = c^2 \frac{X''}{X}$$

[1 mark,unseen] meaning that

$$c^2 \, \frac{X^{\prime \prime}}{X} = \lambda^2$$

and therefore

$$X = A \sin(\frac{\lambda}{c}x) + B \cos(\frac{\lambda}{c}x)$$
 if sign = -1

and

$$X = A \sinh(\frac{\lambda}{c}x) + B \cosh(\frac{\lambda}{c}x)$$
 if sign = +1

[2 marks,unseen]

The boundary condition u(x=0,t)=u(x=L,t)=0 means that B=0 and the solution must be

$$X = A \sin(\frac{\lambda}{c}x)$$

with $\frac{\lambda}{c} = \frac{2 \pi}{L}$ which also selects the negative sign. [2 marks,bookwork]

(e) Using

$$\frac{T''}{T} + k \frac{T'}{T} = -\lambda^2 = c^2 \frac{X''}{X}$$

we obtain the following equation for the time dependence:

$$T'' + kT' + \lambda^2 T = 0$$

[1 mark,unseen]

which admits a solution in $T \propto e^{rt}$ [1 mark,bookwork] with

$$r_{\pm} = \frac{-k \pm \sqrt{k^2 - 4\lambda^2}}{2}$$

[2 marks,unseen]

and therefore $T \propto A \ e^{r_+ \ t} + B \ e^{r_- \ t}$

(f) The vibrations are thus described by the function:

$$u \propto \sin(\frac{\lambda}{c}x) \left(A e^{r_+ t} + B e^{r_- t}\right)$$

with

$$r_{\pm} = \frac{-k \pm \sqrt{k^2 - 4\lambda^2}}{2}$$

When $k=2\,\lambda,\,r_{\pm}<0$ and the vibration is exponentially damped with time.

When $k>2\,\lambda$, the term in $e^{-(\frac{k+\sqrt{k^2-4\lambda^2}}{2})\,t}$ decreases with time while the other exponential vanishes so the solution is damped.

When $k < 2 \, \lambda$, the exponentials contain an imaginary part but the oscillations are damped.

At last when k = 0, there is no damping, just oscillations.

[4 marks,unseen]

(a) To solve

$$x^{2}y'' + (2p+1)xy' + (\alpha^{2}x^{2r} + \beta^{2})y = 0$$

we need to use the Frobenius method where

$$y = x^{\rho} \sum_{n} a_{n} x^{n}$$

Inserting y and its derivatives in the equation we obtain for n=0 (and lowest order x^{ρ}) [2 marks, bookwork]:

$$a_0[\rho(\rho-1) + (2p+1)\rho + \beta^2] = 0$$

which gives

$$\rho^2 + 2 p \rho + \beta^2$$

[1 mark, unseen] and therefore

$$\rho_{\pm} = -p \pm q$$

with $q = \sqrt{p^2 - \beta^2}$ [2 marks, unseen]

Hence one solution of this equation has the form

$$y = x^{-p} [x^{-q} \sum_{n} a_n x^n].$$

[1 mark, unseen]

(b) Writing the above solution as $y = x^{-p} v$ and inserting the derivatives

$$y' = -p x^{-p-1} v + x^{-p} v'$$

[1 mark,bookwork]

$$y'' = p(p+1)x^{-p-2}v - 2px^{-p-1}v' + x^{-p}v''$$

[2 marks, bookwork]

we find

$$v'' x^{-p+2} + v'x^{-p+1} + v(-p^2 + \beta^2 + \alpha^2 x^{2r}) x^{-p}$$

[1 mark, unseen]

that is equivalent to

$$x^{2}v'' + xv' + v(\alpha^{2}x^{2r} - (p^{2} - \beta^{2})) = 0$$

(c) When r = 1 the equation reads

$$x^{2}y'' + (2p+1)xy' + (\alpha^{2}x^{2} + \beta^{2})y = 0$$

which has for solution

$$y = x^{-p}v$$

with v solution to the Bessel equation [2 marks, bookwork]

$$x^{2}v'' + xv' + v(\alpha^{2}x^{2} - (p^{2} - \beta^{2})) = 0$$

that is

$$v = A J_q(\alpha x) + B Y_q(\alpha x)$$

[2 marks,bookwork]

implying that

$$y = x^{-p} \left(A J_q(\alpha x) + B Y_q(\alpha x) \right)$$

(d) Multiplying by t^2 the equation for the background in presence of axions, we obtain:

$$t^2\ddot{\phi} + 3ht\dot{\phi} + m^2t^2\phi = 0$$

[2 marks, bookwork] which can be written as

$$x^{2}y'' + (2p+1)xy' + (\alpha^{2}x^{2r} + \beta^{2})y = 0$$

providing that $3 h = (2p+1), \alpha = m, r=1$ and $\beta = 0$. [2 marks,unseen] Therefore the solution is

$$\phi = t^{-p} (A J_n(m t) + B Y_n(m t))$$

with $p = \frac{3h-1}{2}$. [2 marks,unseen]