

Durham
University

Lecture 15: Dark Matter (A)

Cosmology II

Michaelmas Term 2020

Position of this lecture

$$ds^2 = -dt^2 + a^2(t) \left[dr^2 + S_k(r)^2(d\theta^2 + \sin^2\theta d\phi^2) \right] \quad \text{L10}$$

$$d_A(z) = \frac{S_k[r(z)]}{1+z} = \frac{l}{\Theta} \quad \text{L12}$$

$$d_L(z) = (1+z)S_k[r(z)] = \sqrt{\frac{L}{4\pi f}} \quad \text{L11}$$

$$r(z) = c \int_{?}^{?} \frac{dz}{H(z)} = c \int_{?}^{?} \frac{da}{a^2 H(a)}$$

$$r(z) = c \int_{z_{\text{rec}}}^{\infty} \frac{dz}{H(z)} = c \int_0^{a_{\text{rec}}} \frac{da}{a^2 H(a)} \quad \text{L16, 17}$$

$$\begin{aligned} \left(\frac{\dot{a}}{a}\right)^2 &= \frac{8\pi G}{3}\rho - \frac{kc^2}{a^2} + \frac{\Lambda c^2}{3} \\ \ddot{a}/a &= -\frac{4\pi G}{3}\left(\rho + \frac{3P}{c^2}\right) + \frac{\Lambda c^2}{3} \\ \dot{\rho} + 3\frac{\dot{a}}{a}\left(\rho + \frac{P}{c^2}\right) &= 0 \end{aligned}$$

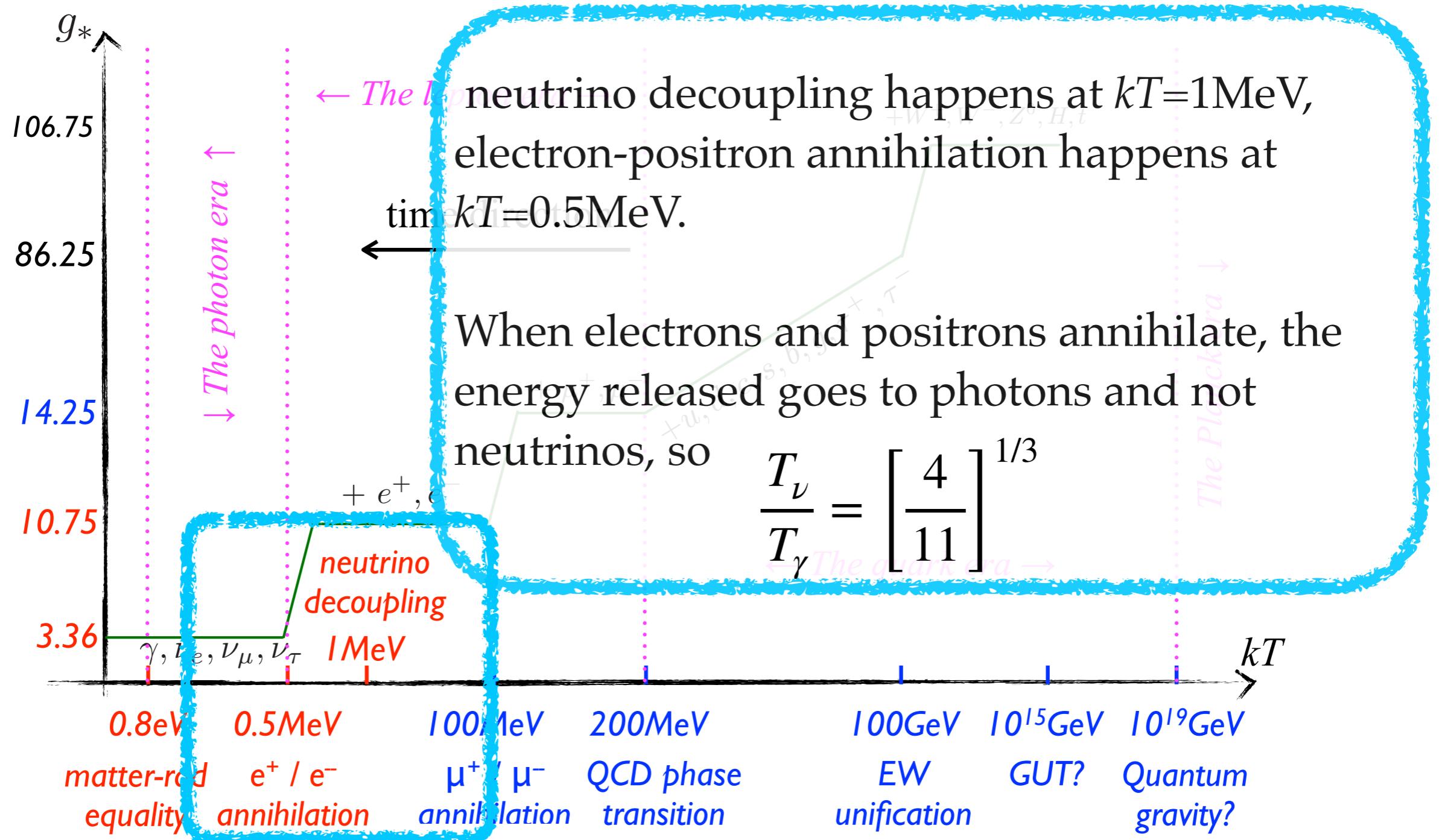
$$\rho_r = \frac{g_*(T)}{2} \frac{4\sigma T^4}{c^3} \quad \text{L13}$$

$$g_*(T) = \sum g_{\text{boson}} + \frac{7}{8} \sum g_{\text{fermion}}$$

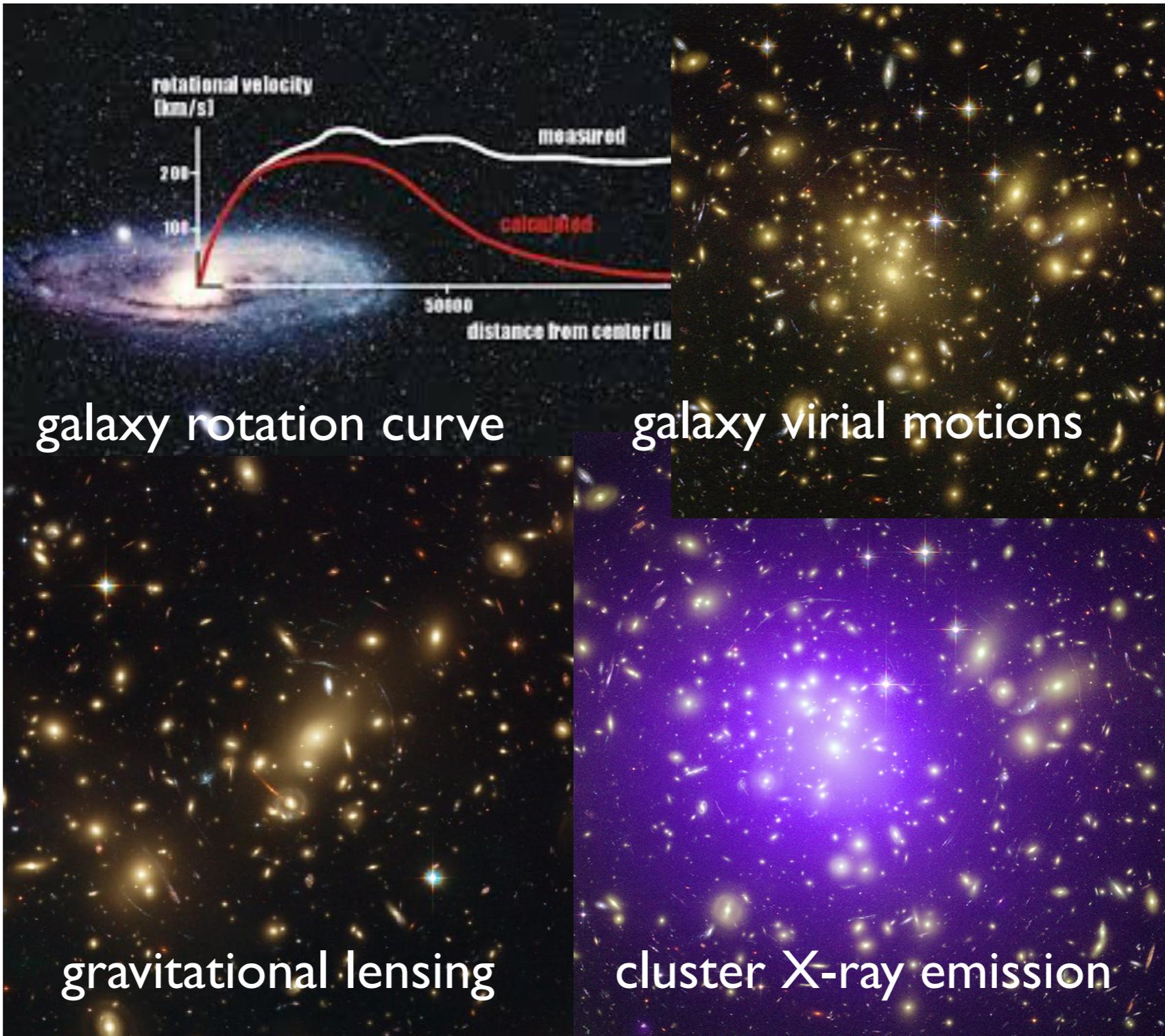
$$\begin{aligned} \frac{N_n}{N_p} &= \exp\left[-\frac{\Delta mc^2}{k_B T}\right] \\ Y_{\text{He}} &= \frac{2N_n}{N_n + N_p} \quad \text{L14} \end{aligned}$$

$$\begin{aligned} \Omega_\nu &= \frac{\sum_\nu m_\nu c^2}{94 \text{ eV}} \left[\frac{H_0}{100 \text{ km/s/Mpc}} \right]^{-2} \\ \frac{T_\nu}{T_\gamma} &= \left(\frac{4}{11} \right)^{1/3} \quad \text{L15} \end{aligned}$$

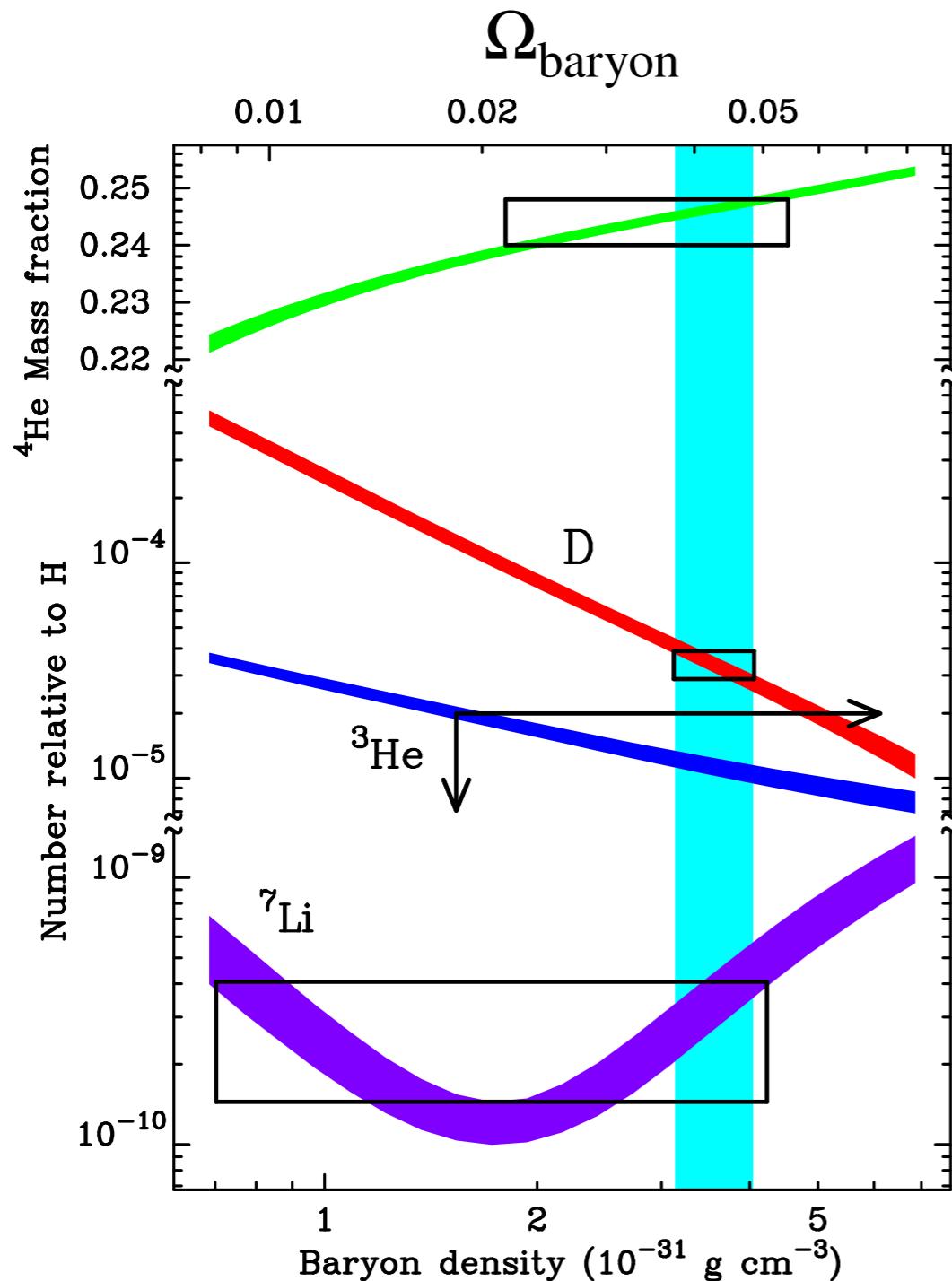
Recap from L13 (relevant materials)



Evidence of dark matter (Lecture 6)



Evidence of non-baryon dark matter (15.1. 15.2)



BBN: suggests that baryon density parameter today is $\Omega_{\text{baryon}} \sim 0.04-0.05$ (Lecture 14)

Type Ia supernovae: suggest that the total matter density parameter today is $\Omega_m \sim 0.3$ (Lectures 9 & 11)

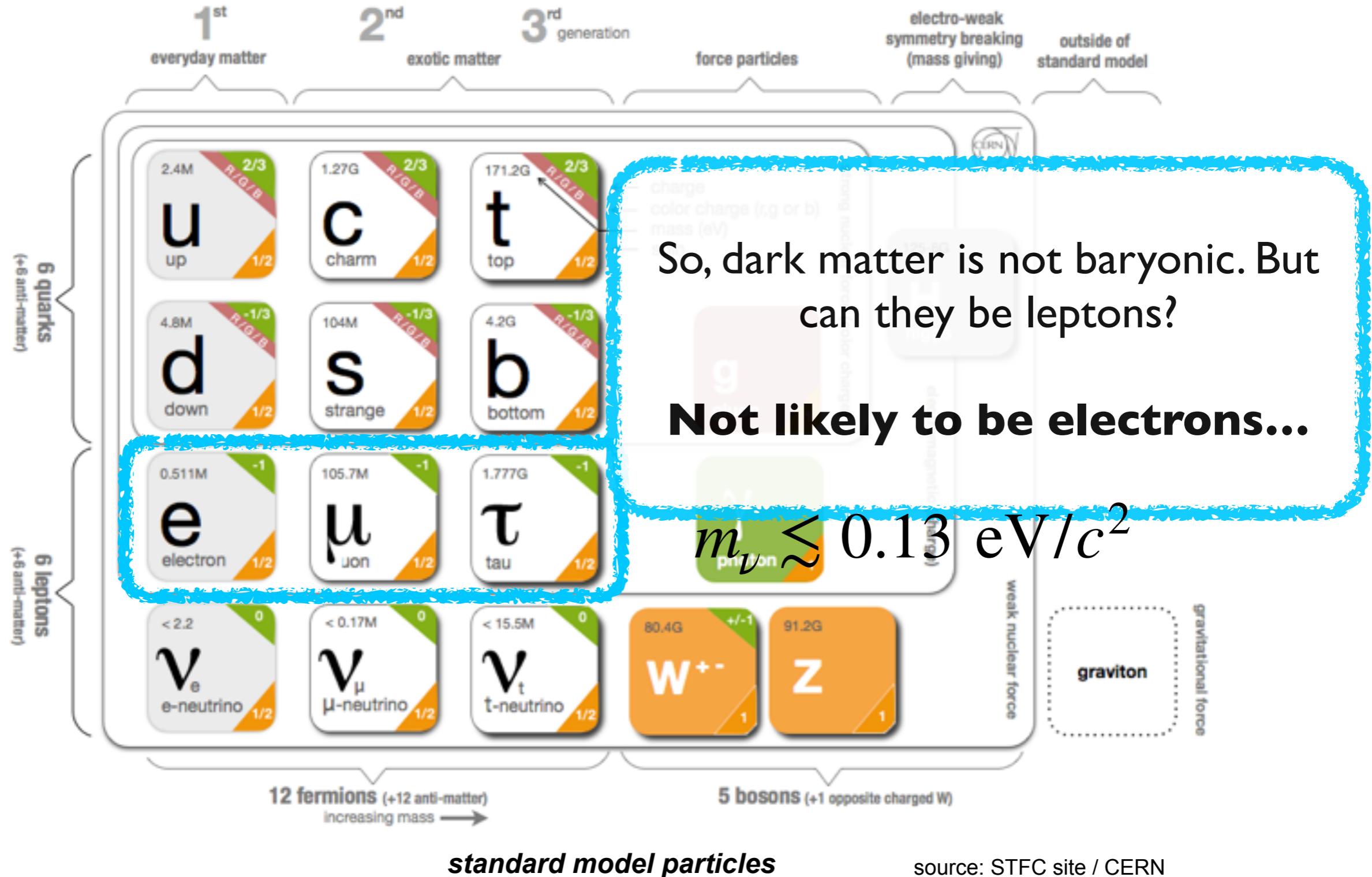
The difference is **dark matter!**

Dark matter (15.3)

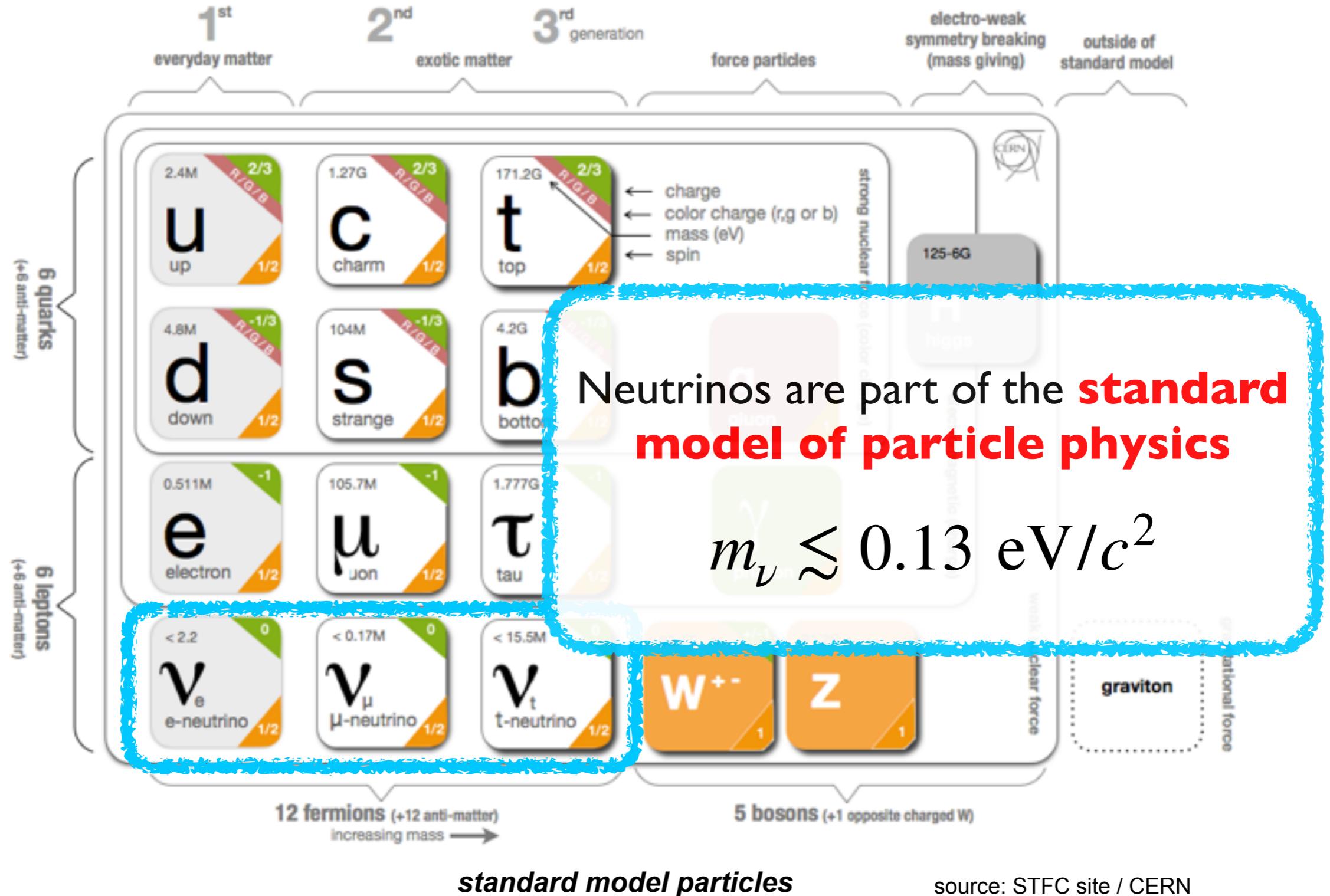
What do we know about dark matter?

- $\Omega_{\text{DM}0} \sim 0.25$
- it is **non-relativistic** matter today: $\rho_{\text{DM}} \propto a^{-3}$
- weakly interacting with other particles

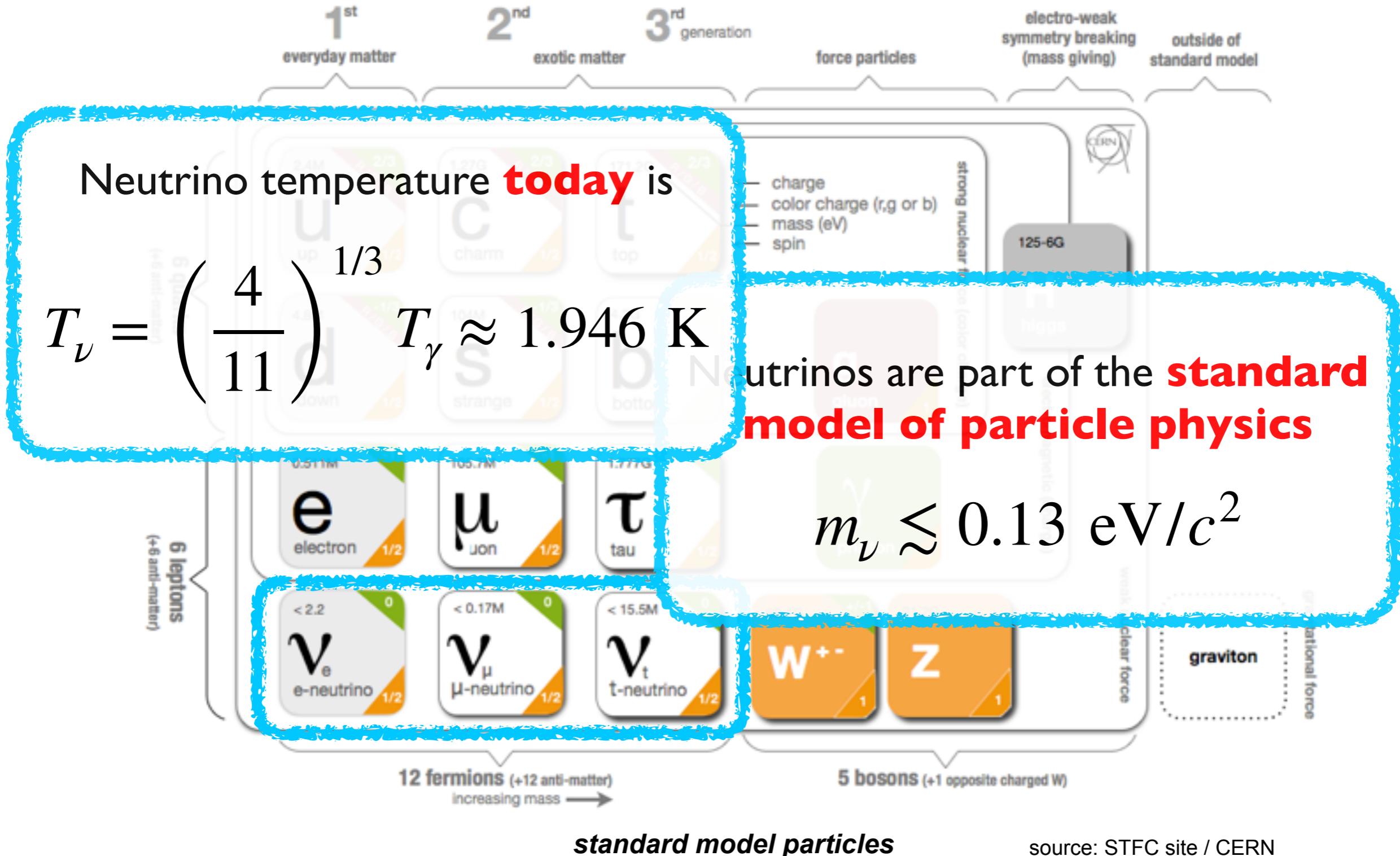
Hot dark matter: neutrinos



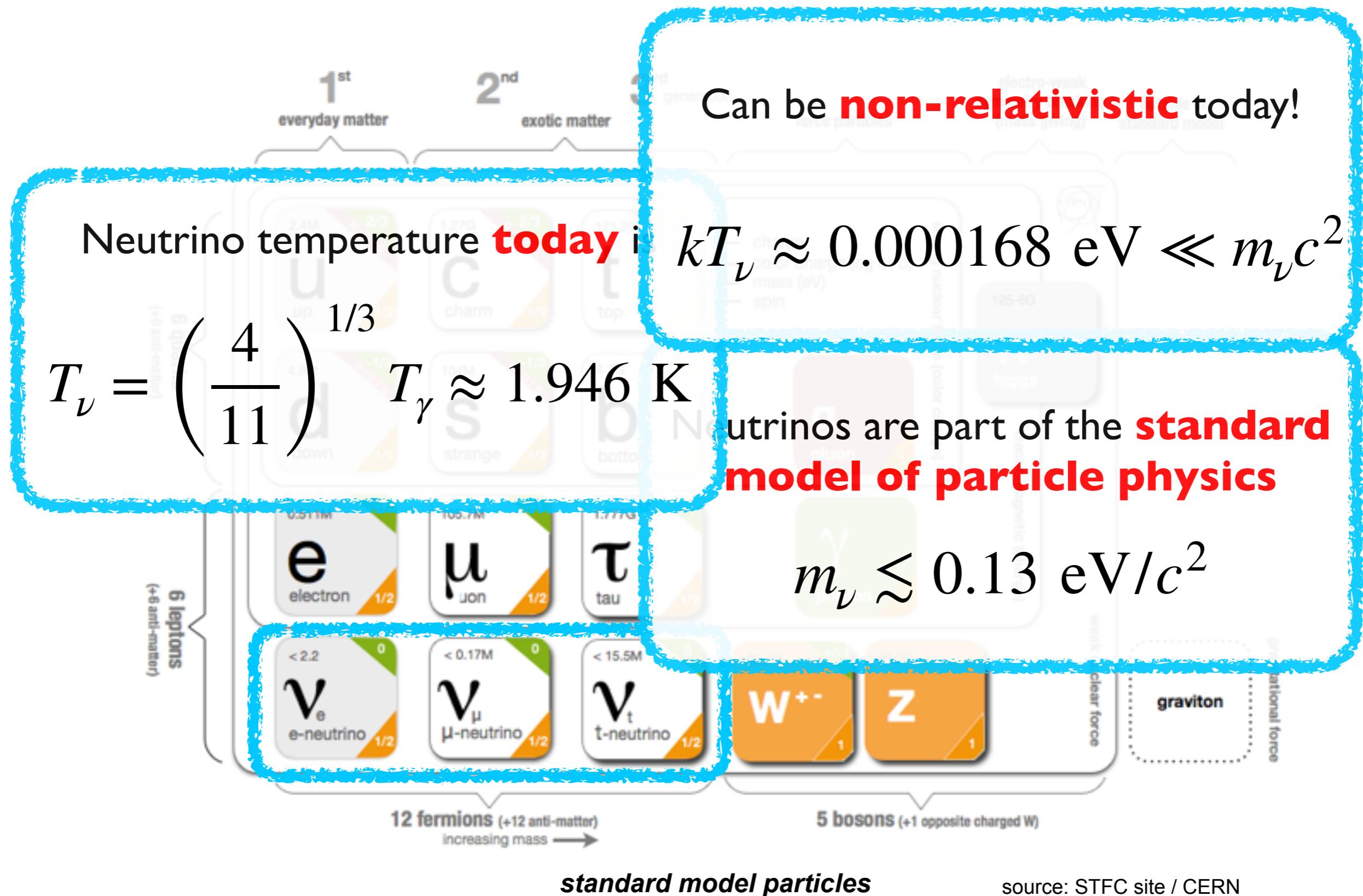
Hot dark matter: neutrinos



Hot dark matter: neutrinos



Hot dark matter: neutrinos



source: STFC site / CERN

Hot dark matter: neutrinos

Neutrinos are called **hot dark matter** because they are highly relativistic at decoupling (**since when the total neutrino number in the universe stays nearly constant**).

But can neutrinos account for all dark matter, $\Omega_{\text{DM}0} \sim 0.25$?

Let's derive its abundance.

Hot dark matter: neutrinos

$$\Omega_\nu = \frac{8\pi G \rho_\nu}{3H^2}$$

$$\rho_\nu = n_\nu m_\nu$$

$$\Omega_\nu = \frac{\sum m_\nu c^2}{94 \text{ eV}} \left[\frac{H_0}{100 \text{ km/s/Mpc}} \right]^{-2}$$

key eqn of L15

neutrino
number density

neutrino
particle mass

$$n_\nu(a) \propto a^{-3}$$

$$n_\nu(a) = n_\nu(a_{\text{dec}}) \frac{a_{\text{dec}}^3}{a^3}$$

$$n_\nu(a_{\text{dec}}) = \left[\frac{\text{neutrino energy density}}{\text{mean energy per neutrino}} \right]_{\text{dec}} \rightarrow 3kT_\nu$$

$$= \left[\frac{\frac{6}{2} \cdot \frac{7}{8} \cdot \frac{40 T_\nu^3}{c}}{3 k T_\nu} \right]_{\text{dec}}$$

$$= \frac{70}{2k_c} T_{\nu, \text{dec}}^3$$

// total number density of all three flavours of neutrinos
and their anti-neutrinos.

$$n_{\nu_e + \bar{\nu}_e}^{\text{dec}} = n_{\nu_\mu + \bar{\nu}_\mu}^{\text{dec}} = n_{\nu_\tau + \bar{\nu}_\tau}^{\text{dec}} = \frac{1}{3} \cdot \frac{70}{2k_c} \cdot T_{\nu, \text{dec}}^3 = \frac{70}{6k_c} T_{\nu, \text{dec}}^3$$

$$\begin{aligned} n_{\nu_e + \bar{\nu}_e}(a) &= n_{\nu_\mu + \bar{\nu}_\mu}(a) = n_{\nu_\tau + \bar{\nu}_\tau}(a) = \frac{70}{6k_c} T_{\nu, \text{dec}}^3 \left(\frac{a_{\text{dec}}}{a} \right)^3 \\ &= \frac{70}{6k_c} T_\nu^3(a) \end{aligned}$$

$T_\nu \propto \frac{1}{a}$

$$\Omega_{\nu 0} = \frac{8\pi G}{3H_0^2} \rho_{\nu 0}$$

$$= \frac{8\pi G}{3H_0^2} \left[n_{e+\bar{e}} m_{\nu e} + n_{\nu_\mu + \bar{\nu}_\mu} m_{\nu_\mu} + n_{\nu_\tau + \bar{\nu}_\tau} m_{\nu_\tau} \right]_{a=1}$$

$$= \frac{8\pi G}{3H_0^2} \frac{1}{6} \frac{\sigma}{k_C} T_{\nu 0}^3 \underbrace{(m_{\nu e} + m_{\nu_\mu} + m_{\nu_\tau})}_{\sum m_\nu}$$

$$= \frac{8\pi G}{c^2 3H_0^2} \frac{1}{6} \frac{\sigma}{k_C} T_{\nu 0}^3 \underbrace{\sum m_\nu c^2}_{\text{unit of energy: eV}}$$

$$= \sum m_\nu c^2 \times \left[\frac{H_0}{100 \text{ km/s/Mpc}} \right]^{-2} \times \boxed{\frac{8\pi G}{3 \times 100^2 \times c^2} \cdot \frac{1}{6} \frac{\sigma}{k_C} T_{\nu 0}^3}$$

$$\approx \frac{\sum m_\nu c^2}{q_0 \text{ eV}} \left[\frac{H_0}{100 \text{ km/s/Mpc}} \right]^{-2}$$

//

$$T_{\nu 0} \approx 1.946 \text{ K}$$

$$100 \text{ km/s/Mpc} \approx 3.24 \times 10^{-18} \text{ s}^{-1}$$

$$\rightarrow \approx 7 \times 10^{16} \text{ J}^{-1}$$

$$\approx \frac{1}{q_0 \text{ eV}}$$

For $\sum m_\nu c^2 \approx 0.1 \text{ eV}$, $H_0 = 70 \text{ km/s/Mpc}$,

we have: $\Omega_{\nu 0} \approx \frac{0.1}{q_0} \times 0.7^{-2} \approx 0.002 < 0.25!$

Cold dark matter (15.3.2)

$$m_\nu \lesssim 0.13 \text{ eV}/c^2$$

experiment bound
on neutrino mass

$$\Omega_\nu = \frac{\sum m_\nu c^2}{94 \text{ eV}} \left[\frac{H_0}{100 \text{ km/s/Mpc}} \right]^{-2} \ll 0.25$$

Neutrinos cannot be the whole picture of dark matter.
And unfortunately no other fundamental particles in the standard model can be a viable candidate for dark matter either.

Cold dark matter (15.3.2)

Beyond standard models offer a plethora of new particles. One example is supersymmetry (SUSY), a Grand Unified theory.

	SUSY
SM fermions	bosonic partners

Leptons → Sleptons
neutrino → sneutrino
Quarks → Squarks

	SUSY
SM bosons	fermionic partners

Photon → Photino
Higgs → Higgsino
Graviton → Gravitino
W → Wino

Cold dark matter (15.3.2)

- very **heavy** masses: $100 \text{ GeV}/c^2$
- interact with other particles species **weakly** (typical strength of weak interactions)
- **WIMPs** (Weakly Interacting Massive Particles)

SUSY
SM bosons
fermionic partners

Leptons \rightarrow Sleptons
neutrino \rightarrow sneutrino
Quarks \rightarrow Squarks

Beyond standard models offer a plethora of new particles. One example is supersymmetry (SUSY), a Grand Unified theory

SUSY fermionic partners

SM bosons

Photon \rightarrow Photino
Higgs \rightarrow Higgsino
Graviton \rightarrow Gravitino
 $W \rightarrow$ Wino

Cold dark matter (15.3.2)

Like neutrinos, these SUSY particles were also in thermal equilibrium with the rest of the plasma in the very early universe. When decoupling, they are non-relativistic, so that they are called **cold dark matter**.

A similar logic as for neutrinos can be used to calculate the abundance of WIMPs, but the calculation is beyond the scope of this course. We just note that WIMPs can naturally have an abundance of $\Omega_{WIMP0} \sim 0.25$.

A summary of dark matter

So we have two classes of candidates

- **hot dark matter** (e.g., neutrinos): **already detected** but only a small contribution to all dark matter
- **cold dark matter** (e.g., WIMPs): **undetected** in **particle colliders** or **terrestrial experiments** so far, but the preferred and widely accepted model, whose predictions agree with cosmological observations very well

Key topics and learning outcomes

Key topics

- hot dark matter
- cold dark matter

Learning outcome

- Be aware of evidences of dark matter and that dark matter cannot be just dark baryons
- Know that neutrinos can be a (hot) dark matter candidate
- Be able to calculate the density parameter of neutrinos given their masses