

3 Counting Microstates (microcanonical ensemble).

①

Book = p13-17.

Which distribution of particles is the most likely?

From thermodynamics, equilibrium is when entropy is maximised.

Therefore the macrostate is an entropy maximum.

Boltzmann that entropy S is some function of the number of microstates Ω , i.e. $S(\Omega)$.

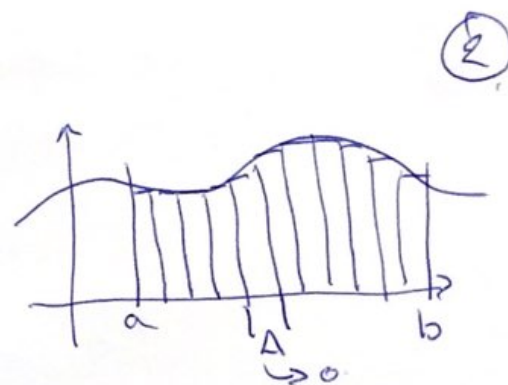
What form? We know that for independent events Ω is

multiplicative, i.e. $\Omega_{AB} = \Omega_A \Omega_B$, ~~not $\Omega_{AB} = \Omega_A + \Omega_B$~~

Boltzmann thought that

$$S = k_B \ln \Omega$$

Boltzmann constant.



Stirling's Approx

For $N \rightarrow \infty$ we have $\ln N! \approx N \ln N - N \{ + o(\ln N) \}$

Proof $\ln N! = \sum_{k=1}^N \ln k \cdot \underbrace{1}_{\substack{\uparrow \\ \Delta \rightarrow 0}} \approx \int_1^N \underbrace{\ln k}_{\substack{\uparrow \\ \Delta \rightarrow 0}} \underbrace{dk}_{\substack{\uparrow \\ \Delta \rightarrow 0}} = [k \ln k - k]_1^N = \underline{N \ln N - N}$

Sometimes written $N! \approx \left(\frac{N}{e}\right)^N$

With large numbers, is it a good enough approximation to

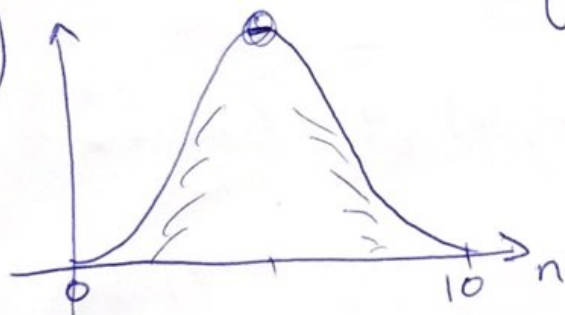
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say that $\Omega \approx \Omega(\max\{n_i\})$?

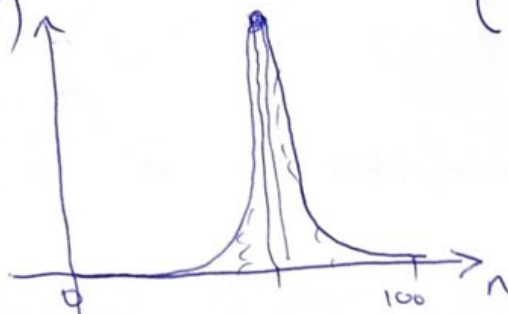
so $S \approx S(\max\{n_i\})$.

Not intuitive at first sight.

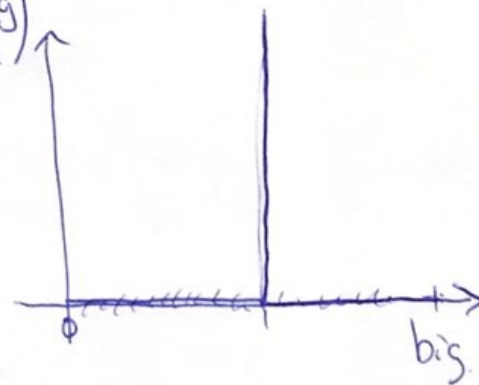
$\binom{10}{n}$



$\binom{100}{n}$



$\binom{\text{big}}{n}$



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Let's have states i with energy ϵ_i and occupation n_i

$$\sum_i n_i = N \quad \text{and} \quad \sum_i n_i \epsilon_i = U.$$

We want the distribution $\{n_1, n_2, \dots\}$

First calculate the entropy: $\Omega(\{n_i\}) = \frac{N!}{\prod_i n_i!}$

$$S(\{n_i\}) = k_B \ln \left(\frac{N!}{\prod_i n_i!} \right) \leftarrow \text{a maximum of equilibrium.}$$

$$\begin{aligned} \frac{S(\{n_i\})}{k_B} &= \ln N! - \ln \left(\prod_i n_i! \right) = \ln N! - \sum_i \ln n_i! \quad (\text{log rules}) \\ &= (\cancel{N \ln N} - \cancel{N}) - \sum_i (\cancel{n_i \ln n_i} - \cancel{n_i}) \end{aligned}$$

$$\frac{S(\{n_i\})}{k_B} = \sum_i n_i \ln N - \sum_i n_i \ln n_i$$

$$= - \sum_i n_i (\ln n_i - \ln N) = - \sum_i n_i \ln \left(\frac{n_i}{N} \right)$$

$$\frac{1}{N} \frac{S(\{n_i\})}{k_B} = - \sum_i \frac{n_i}{N} \ln \frac{n_i}{N}$$

Let's define $p_i = \frac{n_i}{N}$ probability particle is in state i (with energy ϵ_i)

$$\Rightarrow S(\{p_i\}) = - N k_B \sum_i p_i \ln p_i$$

Gibbs expression for entropy.

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Statistical entropy (Boltzmann)

$$S = k_B \ln \Omega \geq 0$$

Gibbs entropy

$$S(\{p_i\}) = -N k_B \sum_i p_i \ln p_i \geq 0$$

Bridges thermodynamic and probabilistic concepts of entropy.

N being large means $N!$ is huge (really huge!!) - fluctuations are extremely rare - essentially the 2nd law of thermodynamics.

Find the most probable distribution.

[look at video on Lagrangian multipliers]

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$$S(\{n_i\}) = N \ln N - \sum_i n_i \ln n_i \leftarrow \text{maximize}$$

with constraints $\sum_i n_i = N$ and $\sum_i n_i \epsilon_i = U$.

For these two constraints we introduce 2 Lagrange multipliers (α , β).
and maximise:

$$N \ln N - \sum_i n_i \ln n_i - \alpha \sum_i n_i - \beta \sum_i n_i \epsilon_i$$

with respect to the occupation numbers of the states, i.e. n_i

We have to take the derivatives $\partial/\partial n_i$ for each i .

⑧

$$\partial/\partial n_i \left[N \ln N - \sum_i n_i \ln n_i - \alpha \sum_i n_i - \beta \sum_i n_i \epsilon_i \right] = 0$$

$$\Rightarrow 0 + \underbrace{(-\ln n_i - 1)}_{\substack{\text{only the } i^{\text{th}} \\ \text{term survives}}} - \alpha \cdot 1 - \beta \epsilon_i \cdot 1 = 0$$

$$\Rightarrow \ln n_i = \overbrace{-1 - \alpha}^A - \beta \epsilon_i$$

$$\Rightarrow n_i = e^A e^{-\beta \epsilon_i}$$