

## Theoretical Physics 2019/20 — Solution of Problem QT2.3

- (a) As mentioned near the beginning of the question, any spin state can be represented by a non-zero two-component column vector. However, any two-component column vector can be represented as a linear combination of  $\chi_a$  and  $\chi_b$  since these two vectors form a basis for the Hilbert space of two-component column vectors. Since a non-zero vector orthogonal to both  $\chi_a$  and  $\chi_b$  cannot be written as a linear combination of these two basis vectors, such a state does not exist.

[1 mark]

- (b) Since  $c_a$ ,  $c_b$ ,  $\chi_a$  and  $\chi_b$  are constant,  $\chi(t)$  depends on  $t$  only through the two exponentials. Therefore

$$\begin{aligned} i\hbar \frac{\partial \chi}{\partial t} &= i\hbar [c_a \chi_a (-iE_a/\hbar) \exp(-iE_a t/\hbar) + c_b \chi_b (-iE_b/\hbar) \exp(-iE_b t/\hbar)] \\ &= c_a \chi_a E_a \exp(-iE_a t/\hbar) + c_b \chi_b E_b \exp(-iE_b t/\hbar). \end{aligned} \quad (1)$$

Moreover,

$$\begin{aligned} H \chi(t) &= c_a H \chi_a \exp(-iE_a t/\hbar) + c_b H \chi_b \exp(-iE_b t/\hbar) \\ &= c_a E_a \chi_a \exp(-iE_a t/\hbar) + c_b E_b \chi_b \exp(-iE_b t/\hbar). \end{aligned} \quad (2)$$

Comparing Eqs. (1) and (2) shows that indeed

$$i\hbar \frac{\partial \chi}{\partial t} = H \chi(t).$$

[1 mark]

- (c) We need to show that the norm of  $\chi(t)$  is 1 when  $|c_a|^2 + |c_b|^2 = 1$ . The square of the norm of  $\chi(t)$  is  $(\chi(t), \chi(t))$ . Now,

$$\begin{aligned} (\chi(t), \chi(t)) &= c_a^* c_a (\chi_a, \chi_a) \exp[i(E_a - E_a)t/\hbar] + c_a^* c_b (\chi_a, \chi_b) \exp[i(E_a - E_b)t/\hbar] \\ &\quad + c_b^* c_a (\chi_b, \chi_a) \exp[i(E_b - E_a)t/\hbar] + c_b^* c_b (\chi_b, \chi_b) \exp[i(E_b - E_b)t/\hbar]. \end{aligned} \quad (3)$$

Since  $(\chi_a, \chi_a) = 1$  and  $(\chi_a, \chi_b) = 0$  (remember that the question says that these two column vectors are orthonormal),

$$(\chi(t), \chi(t)) = c_a^* c_a + c_b^* c_b = |c_a|^2 + |c_b|^2.$$

Thus  $\chi(t)$  is normalized if  $|c_a|^2 + |c_b|^2 = 1$ .

[1 mark]

- (d) As defined by Eq. (3) of the question,  $\chi(t)$  is a linear combination of the two vectors  $\chi_a$  and  $\chi_b$ . Since  $\chi_a$  and  $\chi_b$  are orthonormal, the coefficients of this superposition can be found simply by calculating the inner product of  $\chi(t)$  with each of these two basis vectors: At  $t = 0$ ,  $\exp(-iE_a t/\hbar) = \exp(-iE_b t/\hbar) = 1$ , and therefore  $\chi(t=0) = c_a \chi_a + c_b \chi_b$ . Thus  $(\chi_a, \chi(t=0)) = c_a (\chi_a, \chi_a) + c_b (\chi_a, \chi_b)$ .

Since  $(\chi_a, \chi_a) = 1$  and  $(\chi_a, \chi_b) = 0$ ,  $(\chi_a, \chi(t=0)) = c_a$ . Likewise,  $(\chi_b, \chi(t=0)) = c_b$ . The question says that  $\chi(t) = \chi_+$  at  $t = 0$ . Thus

$$c_a = (\chi_a, \chi_+) = \frac{1}{\sqrt{k^2 + 1}} \begin{pmatrix} -k & 1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{-k}{\sqrt{k^2 + 1}},$$

$$c_b = (\chi_b, \chi_+) = \frac{1}{\sqrt{k^2 + 1}} \begin{pmatrix} 1 & k \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{k^2 + 1}}.$$

(No complex conjugation is indicated in the row vectors because  $k$  is real since  $B_x$ ,  $B_z$  and  $|\mathbf{B}|$  are real.)

Check:

$$|c_a|^2 + |c_b|^2 = \frac{k^2}{k^2 + 1} + \frac{1}{k^2 + 1} = 1.$$

[3 marks for this part of the problem.]

- (e) The probability of finding the electron in the spin state represented by  $\chi_+$  at time  $t$  is  $|(\chi_+, \chi(t))|^2$ ,

Calculating  $|(\chi_+, \chi(t))|^2$  is not particularly difficult: Using the fact that  $(\chi_+, \chi_a) = (\chi_a, \chi_+)$  since  $(\chi_+, \chi_a)$  is real, and similarly for  $(\chi_+, \chi_b)$ ,

$$\begin{aligned} (\chi_+, \chi(t)) &= c_a(\chi_+, \chi_a) \exp(-iE_a t/\hbar) + c_b(\chi_+, \chi_b) \exp(-iE_b t/\hbar) \\ &= \frac{-k}{\sqrt{k^2 + 1}} \frac{-k}{\sqrt{k^2 + 1}} \exp(-iE_a t/\hbar) + \frac{1}{\sqrt{k^2 + 1}} \frac{1}{\sqrt{k^2 + 1}} \exp(-iE_b t/\hbar) \\ &= \frac{k^2}{k^2 + 1} \exp(-i\mu|\mathbf{B}|t/\hbar) + \frac{1}{k^2 + 1} \exp(i\mu|\mathbf{B}|t/\hbar). \end{aligned}$$

Recall that  $k = B_x/(B_z + |\mathbf{B}|)$ . Thus  $k = \pm 1$  when  $B_y = B_z = 0$  since  $B_x = \pm|\mathbf{B}|$  in that case. Accordingly, the probability of finding the electron in the state  $\chi_+$  is

$$[(1/2) \exp(-i\mu|\mathbf{B}|t/\hbar) + (1/2) \exp(i\mu|\mathbf{B}|t/\hbar)]^2 = \cos^2(\mu|\mathbf{B}|t/\hbar).$$

[3 marks]

- (f) Since  $\chi(t)$  is normalized and  $\chi_a$  and  $\chi_b$  are orthonormal, these two probabilities are, respectively,  $|c_a \exp(-iE_a t/\hbar)|^2$  and  $|c_b \exp(-iE_b t/\hbar)|^2$  (i.e., the square of the modulus of the coefficients of  $\chi_a$  and  $\chi_b$  in  $\chi(t)$ ). That is,  $|c_a|^2$  and  $|c_b|^2$  (recall that  $|\exp(ix)| = 1$  if  $x$  is a real number). In order for  $|c_a|^2$  to be equal to  $|c_b|^2$ ,  $k^2$  must be equal to 1, which means that the field must be oriented in the positive or the negative  $x$ -direction.

[1 mark for (f)]