

Mathematical Methods in Physics

Examination June 2017

Question 1

(a) (Unseen)

(i) It is a vector space. [2 marks]

(ii) It is not a vector space. For instance, scalar multiplication is not a close operation. Also, there are no inverse elements. [2 marks]

(b) (Unseen)

(i) They are linearly dependent. The third is equal to the first minus twice the second. [2 marks]

(ii) They are linearly independent. In fact the expression $c_1(1+x) + c_2(x+x^2) + c_3(1+x^2) = 0$ is only satisfied for $c_1 = c_2 = c_3 = 0$. [2 marks]

(c) (Unseen)

(i) $U_1^{-1} = U_1^\dagger$, $U_2^{-1} = U_2^\dagger$, then

$$(U_1 U_2)^{-1} = U_2^{-1} U_1^{-1} = U_2^\dagger U_1^\dagger = (U_1 U_2)^\dagger.$$

[2 marks]

(ii)

$$(i(A - A^\dagger))^\dagger = -i(A^\dagger - A) = i(A - A^\dagger).$$

[2 marks]

(d) (Unseen)

(i) For instance

$$\underline{r}(t) = x \hat{i} + (3 - x) \hat{j} + (18 - 7x) \hat{k}.$$

[2 marks]

(ii)

$$\underline{r}(t) = (\cos t + 1) \hat{i} + (\sin t - 2) \hat{j} + \hat{k}.$$

[2 marks]

(e) (Unseen)

$$\frac{ds}{du} = \pm \sqrt{\frac{dr}{du} \cdot \frac{dr}{du}} = \pm \sqrt{2},$$

Then

$$\underline{\hat{t}} = \frac{dr}{ds} = \frac{dr}{du} \frac{du}{ds} = \frac{1}{\sqrt{2}}(-\sin u \underline{\hat{i}} + \cos u \underline{\hat{j}} + \underline{\hat{k}}),$$

[2 marks]

and

$$\frac{d\underline{\hat{t}}}{ds} = \frac{d\underline{\hat{t}}}{du} \frac{du}{ds} = -\frac{1}{2}(\cos u \underline{\hat{i}} + \sin u \underline{\hat{j}}).$$

It follows that $\rho = 2$.

[2 marks]

(f) (Unseen)

$$\frac{\partial \underline{r}}{\partial x} = \underline{\hat{i}} + 2x \underline{\hat{j}}, \quad \frac{\partial \underline{r}}{\partial z} = \underline{\hat{k}}.$$

Then

$$d\underline{S} = \left(\frac{\partial \underline{r}}{\partial x} \times \frac{\partial \underline{r}}{\partial z} \right) dx dz = (2x \underline{\hat{i}} - \underline{\hat{j}}) dx dz.$$

[2 marks]

Hence

$$I = \int_0^3 dz \int_0^2 dx (2x^3 - 2) = 12.$$

[2 marks]

(g) (Unseen)

$$c_r = \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{x-irx} dx = \frac{1}{2\pi} \left[\frac{e^{x(1-ir)}}{1-ir} \right]_{-\pi}^{\pi} = \frac{(-1)^r}{2\pi} \left(\frac{e^{\pi} - e^{-\pi}}{1-ir} \right) = \frac{\sinh \pi}{\pi} (-1)^r \frac{1+ir}{1+r^2}.$$

Then

$$e^x = \frac{\sinh \pi}{\pi} \sum_{-\infty}^{\infty} \frac{1+ir}{1+r^2} (-1)^r e^{irx} \quad -\pi \leq x \leq \pi.$$

[4 marks]

(h) (Unseen)

$$\begin{aligned} \bar{f}(s) &= \int_0^{\infty} \cos t H(t-\pi) e^{-ts} dt = \int_{\pi}^{\infty} \cos t e^{-ts} dt \\ &= -\frac{1}{2} \left[\frac{e^{-t(s-i)}}{s-i} + \frac{e^{-t(s+i)}}{s+i} \right]_{\pi}^{\infty} = -e^{-\pi s} \frac{s}{s^2+1}. \end{aligned}$$

[4 marks]

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Question 2

(a) (Unseen)

The eigenvalues are: $\lambda_1 = 1$, $\lambda_2 = -1$ and $\lambda_3 = 3$. [4 marks]

The forms of the corresponding eigenvectors are: $v_1 = (0, 0, c)^T$, $v_2 = (a, -a, 0)^T$ and $v_3 = (b, b, 0)^T$, with a , b , c constants. The eigenvectors are clearly orthogonal. A possible choice for an orthonormal set of eigenvectors is: $v_1 = (0, 0, 1)^T$, $v_2 = (1/\sqrt{2}, -1/\sqrt{2}, 0)^T$ and $v_3 = (1/\sqrt{2}, 1/\sqrt{2}, 0)^T$. [4 marks]

(b) (Unseen)

For the integral on the left hand side, since the rectangle $ABCD$ lies in the plane $z = 0$, we only need the z -component of the curl

$$(\nabla \times \underline{F})_z = \left(\frac{\partial}{\partial x}((x+y)e^{xy}) - \frac{\partial}{\partial y}(ye^{xy}) \right) \hat{k} = y^2 e^{xy} \hat{k}.$$

[2 marks]

Then

$$d\underline{S} = \hat{k} dx dy,$$

[1 mark]

hence the surface integral is:

$$\begin{aligned} \int_S (\nabla \times \underline{F}) \cdot d\underline{S} &= \int_1^3 dy \int_0^1 dx y^2 e^{xy} = \int_1^3 dy (e^y - 1)y \\ &= y e^y \Big|_1^3 - e^y \Big|_1^3 - \frac{y^2}{2} \Big|_1^3 = 2e^3 - 4. \end{aligned}$$

[3 marks]

For the integral on the right hand side there are four different $d\underline{r}$. Since $d\underline{S}$ has been chosen along the positive k -axis, the integration is anticlockwise along the perimeter of the rectangle $ABCD$:

$$\int_C \underline{F} \cdot d\underline{r} = \int_0^1 dx \underline{F} \cdot \hat{i} \Big|_{y=1} - \int_0^1 dx \underline{F} \cdot \hat{i} \Big|_{y=3} + \int_1^3 dy \underline{F} \cdot \hat{j} \Big|_{x=1} - \int_1^3 dy \underline{F} \cdot \hat{j} \Big|_{x=0},$$

[3 marks]

which is

$$\begin{aligned} &= \int_0^1 e^x dx - \int_0^1 3e^{3x} dx + \int_1^3 (y+1)e^y dy - \int_1^3 y dy \\ &= e - 1 + (1 - e^3) + (4e^3 - 2e - e^3 + e) + \left(\frac{1}{2} - \frac{9}{2}\right) = 2e^3 - 4. \end{aligned}$$

[3 marks]

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Question 3

(a) (Unseen)

(i)

$$\operatorname{div} \underline{v} = 2x + 4y, \quad \operatorname{grad}(\operatorname{div} \underline{v}) = 2\hat{i} + 4\hat{j}, \quad \operatorname{grad}(\operatorname{div} \underline{v}) \cdot \underline{u} = 2x^2 + 4(y - z).$$

[3 marks]

(ii)

$$\nabla \cdot (f \underline{r}) = \nabla f \cdot \underline{r} + f \nabla \cdot \underline{r} = f' \frac{\underline{r} \cdot \underline{r}}{r} + 3f = f'r + 3f.$$

[3 marks]

Set $f'(r)r + 3f(r) = 0$ then

$$f'(r) = -\frac{3f(r)}{r} \longrightarrow \frac{df}{f} = -3\frac{dr}{r} \longrightarrow \ln f = \ln r^{-3} + c' \longrightarrow f = c/r^3.$$

[4 marks]

(b) (Unseen)

(i) The function $f(t)$ is:

$$f(t) = \begin{cases} 1 + t/b & -b \leq t \leq 0 \\ 1 - t/b & 0 \leq t \leq b \end{cases}$$

[2 marks]

Its Fourier transform is:

$$\begin{aligned} \mathcal{F}[f(t)](\omega) &= \frac{1}{\sqrt{2\pi}} \int_{-b}^0 \left(1 + \frac{t}{b}\right) e^{-i\omega t} dt + \frac{1}{\sqrt{2\pi}} \int_0^b \left(1 - \frac{t}{b}\right) e^{-i\omega t} dt \\ &= \sqrt{\frac{2}{\pi}} \int_0^b \left(1 - \frac{t}{b}\right) \cos \omega t dt \\ &= \sqrt{\frac{2}{\pi}} \left(\left[\frac{\sin \omega t}{\omega} \right]_0^b - \left[\frac{t \sin \omega t}{\omega b} \right]_0^b + \frac{1}{b} \left[\frac{-\cos \omega t}{\omega^2} \right]_0^b \right) \\ &= \sqrt{\frac{2}{\pi}} \frac{(1 - \cos \omega b)}{b \omega^2}. \end{aligned}$$

[4 marks]

- (ii) This function, $g(t)$, consists of two segments of the same shape as $f(t)$ in (i), but with each one shifted and scaled. The first segment has width one and it is centered at $t = 1/2$. Hence for this segment $b = 1/2$ and the argument t becomes $(t - 1/2)$. The second segment has width two and it is centered at $t = 2$. Hence it has $b = 1$ and the argument t becomes $(t - 2)$. Hence, the Fourier transform of this function is

$$\mathcal{F}[g(t)](\omega) = \sqrt{\frac{2}{\pi}} \frac{1}{\omega^2} \left(e^{-i\omega/2} 2(1 - \cos \omega/2) + e^{-i2\omega} (1 - \cos \omega) \right).$$

[4 marks]

(a) $\frac{dy}{dx} + \frac{2y}{x} = \frac{4}{x}$ (1)

First solve the homogeneous equation $\frac{dy}{dx} = -\frac{2y}{x}$ [1 mark]

$\Rightarrow y_{\text{hom}} = C/x^2$

To solve the inhomogeneous eqn use the varying const method:
 $C \rightarrow C(x) \Rightarrow y' = \frac{C'}{x^2} - \frac{2C}{x^3}$: plug in the eqn (1)
 $\frac{C'}{x^2} - \frac{2C}{x^3} = \frac{4}{x} \Rightarrow C' = 4x$

$\Rightarrow C = 2x^2 + \text{const.} \Rightarrow y_{\text{inh}} = 2 + \frac{\text{const}}{x^2}$ [2 marks]

[if the answer is correct, will accept any method]

Substitute back into (1)
 $\Rightarrow \frac{-2C}{x^3} + \frac{4}{x} + \frac{2C}{x^3} = \frac{4}{x}$ ✓ [1 mark]

(b) $y'' - 6y' + 5y = 0$ use $y = Ae^{\lambda t}$

$\Rightarrow \lambda^2 - 6\lambda + 5 = 0 \Rightarrow \lambda = 1, 5$ [2 marks]

$y_{\text{hom}} = Ae^{5t} + Be^t$ roots

$y_{\text{part. inhom}} = ax + bx^2 + d$ RHS = $4 + x + 10x^2$

trial f.n.
 plug in eqn: $2b - 6a - 12bx + 5ax + 5bx^2 + 5d = \text{RHS}$

sol-n: $a=5$; $b=2$; $d=6$

$y_{\text{general inho}} = Ae^{5t} + Be^t + 5x + 2x^2 + 6$ ✓ [2 marks]

n^{th} order lin. homo. ODE with const coef-s

(c) $a_n \frac{d^n y}{dx^n} + a_{n-1} \frac{d^{n-1} y}{dx^{n-1}} + \dots + a_0 y = 0$ a_i

One searches for the general sol-n by substituting
 $y(x) = \sum_{i=1}^n c_i e^{\lambda_i x}$ [1 mark]

where each λ_i is a root of the auxiliary eqn:
 $\sum_{j=1}^n a_j \lambda^j = 0$. [clearly there are n roots] [1 mark]

When 1 root e.g. λ_1 has the k -fold degeneracy
 the roots are $\lambda_1, \dots, \lambda_1, \lambda_{k+1}, \lambda_{k+2}, \dots, \lambda_n$
 $\underbrace{\lambda_1, \dots, \lambda_1}_{k \text{ repeated roots } \lambda_1}$

Then the correct gen. sol-n is obtained by:
 $y_{\text{gen}}(x) = (C_1 + C_2 x + \dots + C_k x^{k-1}) e^{\lambda_1 x} + C_{k+1} e^{\lambda_{k+1} x} + \dots + C_n e^{\lambda_n x}$
 $y_{\text{gen}}(x) = (C_1 + C_2 x + \dots + C_k x^{k-1}) e^{\lambda_1 x} + C_{k+1} e^{\lambda_{k+1} x} + \dots + C_n e^{\lambda_n x}$
 [still independent coefficients $c_1 \dots c_n$] [2 marks]

(d) $1 \cdot y'' + p(x)y' + q(x)y = 0$ \Leftarrow 2nd order linear ODE

Singularities are when $p(x=x_0) = \infty$
 $q(x=x_0) = \infty$

regular-singular pt $x=x_0$ when

$(x-x_0)p(x)$ is analytic at $x \rightarrow x_0$ (finite) ✓
 $(x-x_0)^2 q(x)$ is analytic at $x \rightarrow x_0$ (finite) ✓
 if non-analytic \Rightarrow essential singularity [2 marks]

Legendre eq-n
 $(1-x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + l(l+1)y = 0$
 $x \rightarrow \pm 1$ $\frac{1}{x} \rightarrow$ sing (essential sing at $x \rightarrow \infty$)

regular-sing points are at $x = \pm 1$ [2 marks]

(e) Using the Rodriguez formula we get
 $P_0 = 1$; $P_1(x) = x$; $P_2(x) = \frac{1}{2}(3x^2 - 1)$ [1 mark]
 $P_3(x) = \frac{1}{2}(5x^3 - 3x)$ [1 mark]

$\int_{-1}^1 dx P_l P_k = \begin{cases} \text{even, odd} = 0 \\ l=0, k=2 = \int_{-1}^1 (\frac{3}{2}x^2 dx - \frac{1}{2}) dx = \frac{3}{2} \cdot \frac{2}{3} - \frac{1}{2} \cdot 2 = 1 - 1 = 0 \\ l=1, k=3 = \int_{-1}^1 (5x^4 - 3x^2) dx = \frac{5}{5}x^5 - \frac{3}{3}x^3 \Big|_{-1}^1 = 1 - 1 = 0 \end{cases}$

all \int 's $= 0$
 As expected ✓ [2 marks]

End

(f) $F = -3y$ clearly satisfies $y \frac{\partial F}{\partial x} - x \frac{\partial F}{\partial y} = 3x$
 by substitution [1 mark]

Need now gen. sol. of $y \frac{\partial F}{\partial x} - x \frac{\partial F}{\partial y} = 0$
 $F(x,y) = F(p(x,y))$ $dp = 0 = \frac{\partial F}{\partial x} dx + \frac{\partial F}{\partial y} dy$
 $y \frac{dF}{dp} \frac{\partial p}{\partial x} - x \frac{dF}{dp} \frac{\partial p}{\partial y} = 0 \Rightarrow -y \frac{\partial p}{\partial y} dy = x \frac{\partial p}{\partial x} dx$
 $\Rightarrow y dy + x dx = 0$ $p: x^2 + y^2 = \text{const.}$ [2 marks]

Gen. sol. of the inh. eq-n: $F(x,y) = p(x^2 + y^2) - 3y$
 arbitrary f.n. [1 mark]

(g) 1-dim wave eq-n: $\frac{\partial^2 y}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 y}{\partial t^2}$ [1 mark]

$y(x,t) = y(p); p = x + \lambda t$
 $\frac{\partial^2 y}{\partial x^2} = \frac{d^2 y}{dp^2}$; $\frac{\partial^2 y}{\partial t^2} = \frac{d^2 y}{dp^2} \cdot \lambda^2$
 \Rightarrow eq-n $\Rightarrow 1 = \lambda^2 / c^2 \Rightarrow \lambda = \pm c$

gen. sol-n $y(x,t) = f(x+ct) + g(x-ct)$ [1 mark]

2dim Laplace: $\frac{\partial^2 y}{\partial x^2} + \frac{\partial^2 y}{\partial y^2} = 0$ [1 mark]
 $p = x + \lambda y$
 $\Rightarrow 1 + \lambda^2 = 0 ; \lambda = \pm i$
 $\Rightarrow y(x,t) = f(x+iy) + g(x-iy)$ ✓ [1 mark]

[End]

a) As the Legendre linear eq-n it takes the form:

$$(x-1)^2 \frac{d^2 y}{dx^2} - 2y = 0 \quad \text{i.e. } \alpha x + \beta = x-1 \quad [2 \text{ marks}]$$

The substitution is $x-1 = e^t$

$$\Rightarrow \frac{dy}{dx} = \frac{1}{x-1} \frac{dy}{dt} \quad ; \quad \frac{d^2 y}{dx^2} = \left(\frac{d^2 y}{dt^2} - \frac{dy}{dt} \right) \frac{1}{(x-1)^2} \quad [1 \text{ mark}]$$

\Rightarrow the eq-n becomes: $\frac{(x-1)^2}{(x-1)^2} (\ddot{y} - \dot{y}) - 2y = 0$ [2 marks]

Eq-n w. const coeff.s ✓

b) Use: $y = e^{\lambda t}$ to solve it [1 mark]

\Rightarrow auxiliary eq-n $\lambda^2 - \lambda - 2 = 0 \Rightarrow 2 \text{ sol's } \lambda = 2, -1$

$$y_1 = e^{2t} = e^{2 \log(x-1)} = (x-1)^2 \quad [2 \text{ marks}]$$

$$y_2 = e^{-t} = \frac{1}{1-x}$$

2 particular sol.s [3 marks]

c) Wronskian $W = \det \begin{pmatrix} y_1 & y_2 \\ y_1' & y_2' \end{pmatrix}$ [1 mark]

$$W = |y_1 y_2' - y_2 y_1'| = |1 - 3| = 3 \neq 0$$

hence y_1 and y_2 are linearly independent. [2 marks]

general $y(x) = C_1 (x-1)^2 + C_2 \frac{1}{1-x}$

C_1, C_2 constants, [2 marks]

d) $x=0$ is a regular point of the equation.

Hence search for the power-series in the form

$$y(x) = \sum_{n=0}^{\infty} a_n x^n$$

compute $y'(x)$ and $y''(x)$ [2 marks]

\Rightarrow substitute to the eq-n.

Hence obtain a recursion relation for a_n 's.

It can then be solved to determine 2 particular sol.s

[2 marks]

End

a) $\ddot{x} + \omega_0^2 x = 0 \Rightarrow$ general sol'n : $x = A \cos \omega_0 t + B \sin \omega_0 t$

To remain in equilibrium at rest $\begin{cases} x(0) = 0 \\ \dot{x}(0) = 0 \end{cases} \Rightarrow x(t) = 0$ [1 mark]

b) The inhomogeneous ODE with the force :

$$\ddot{x} + \omega_0^2 x = f_0 \sin(at)$$

Laplace transform : $L[x](s) = \int_0^\infty e^{-st} x(t) dt$ [2 marks]

Laplace tr. of \ddot{x}

$$L[\ddot{x}](s) = s^2 L[x](s) - s x(0) - \dot{x}(0)$$

$= 0$ from initial cond's

$$L[\text{eq'n}] \Rightarrow (s^2 + \omega_0^2) L[x](s) = f_0 L[\sin(at)](s)$$
 [2 marks]

$\frac{as}{(s^2 + a^2)}$ from the [hint.]

⊙ $\Rightarrow L[x](s) = \frac{f_0 a s}{(s^2 + \omega_0^2)(s^2 + a^2)}$ [2 marks]

when $\omega_0^2 \neq a^2 \Rightarrow \frac{f_0 a}{\omega_0^2 - a^2} \left(\frac{s}{s^2 + a^2} - \frac{s}{s^2 + \omega_0^2} \right)$ [2 marks]

use [hint] to do inverse

Laplace transform :

harmonic motion \Leftarrow

$$x(t) = \frac{f_0 a}{\omega_0^2 - a^2} \left(\frac{\sin at}{a} - \frac{\sin \omega_0 t}{\omega_0} \right)$$

[2 marks]

c) for $a^2 = \omega_0^2$ go back to ⊙ : $L[x](s) = \frac{f_0 a s}{(s^2 + a^2)^2}$ [3 marks]

use the 2nd hint to do the inverse Laplace transform.

$$x(t) = \frac{f_0}{2a^2} (\sin at - at \cos at)$$

[2 marks]

↑ this term grows linearly with time.

This trajectory is the resonance motion of the pendulum disturbed by the force w frequency $a^2 = \omega_0^2$

[2 marks]

End.