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Book : p51-62.

Distributions: Classical (or distinguishable) particles we counted the number of microstates such that the order in which particles were arranged mattered ($N!$), e.g.

$$\Omega \rightarrow \frac{N!}{\prod_i n_i!} \rightarrow \text{classical Maxwell-Boltzmann distribution.}$$

Identical/Indistinguishable Particles - this is a purely quantum concept.

Quantum particles - either Fermion or Boson.

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If we have two quantum particles then the wavefunction of the two-particle system is written $\psi(r_1, r_2)$. Let's swap the two particles, so we have system $\psi(r_2, r_1)$. How is $\psi(r_1, r_2)$ described in terms of $\psi(r_2, r_1)$?

Let S be a "swap" operator, i.e.

$$S \psi(r_1, r_2) = x \psi(r_2, r_1)$$

$$S^2 \psi(r_1, r_2) = (\boxed{x^2}) \psi(r_1, r_2)$$

$$\Rightarrow x^2 = 1.$$

$$\Rightarrow x = \pm 1.$$

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We're left with two situations:

$$\psi_F(r_1, r_2) = -\psi_F(r_2, r_1) \leftarrow \text{Fermions}$$

$$\psi_B(r_1, r_2) = \psi_B(r_2, r_1) \leftarrow \text{Bosons.}$$

[Aside: Anyways $x = e^{i\theta}$]

Put 2 Fermions at the same place i.e.

Pauli exclusion.

$$\psi_F(r, r) = -\psi_F(r, r) \Rightarrow \boxed{\psi_F(r, r) = 0}$$

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Fermions Each single particle state will be either occupied ($n_i = 1$) or unoccupied ($n_i = 0$). If state i is g_i -fold degenerate with energy ϵ_i . We want the number of ways g_i slots can be divided into 2 piles - occupied or unoccupied.

One pile (the occupied) will have n_i ($n_i \leq g_i$) and hence the other pile (the unoccupied) will have $g_i - n_i$ slots free.

Using the binomial distribution

$$\Omega(\epsilon_i) = \frac{g_i!}{n_i! (g_i - n_i)!}$$

This is the probability for the i^{th} state so for all states

$$\Omega_{\text{FD}} = \prod_i \frac{g_i!}{n_i! (g_i - n_i)!}$$

As previously done, we maximise entropy under the same constraints
(i.e. constant energy and particle number):

$$\frac{S}{k_B} = \ln \left[\prod_i \frac{g_i!}{n_i! (g_i - n_i)!} \right]$$

$$= \sum_i \left(\ln g_i! - \ln n_i! - \ln (g_i - n_i)! \right) \quad (\text{log rules})$$

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$$= \sum_i \left[(g_i \ln g_i - \underline{g_i}) - (n_i \ln n_i - \underline{n_i}) - ((g_i - n_i) \ln (g_i - n_i) - \underline{(g_i - n_i)}) \right] \quad (\text{Stirling})$$

$$= \sum_i [g_i \ln g_i - n_i \ln n_i - (g_i - n_i) \ln (g_i - n_i)]$$

Now introduce the Lagrange multipliers:

$$\sum_i [g_i \ln g_i - n_i \ln n_i - (g_i - n_i) \ln (g_i - n_i) - \alpha n_i - \beta n_i \epsilon_i]$$

and maximise this with respect to n_i i.e.

$$\partial(\dots) / \partial n_i = 0$$

$$0 = 0 - \ln n_i - 1 + \ln(g_i - n_i) + 1 - \alpha - \beta \epsilon_i \quad (\text{only } i\text{th term in sum survives})$$

$$\Rightarrow \ln\left(\frac{g_i}{n_i} - 1\right) = \alpha + \beta \epsilon_i$$

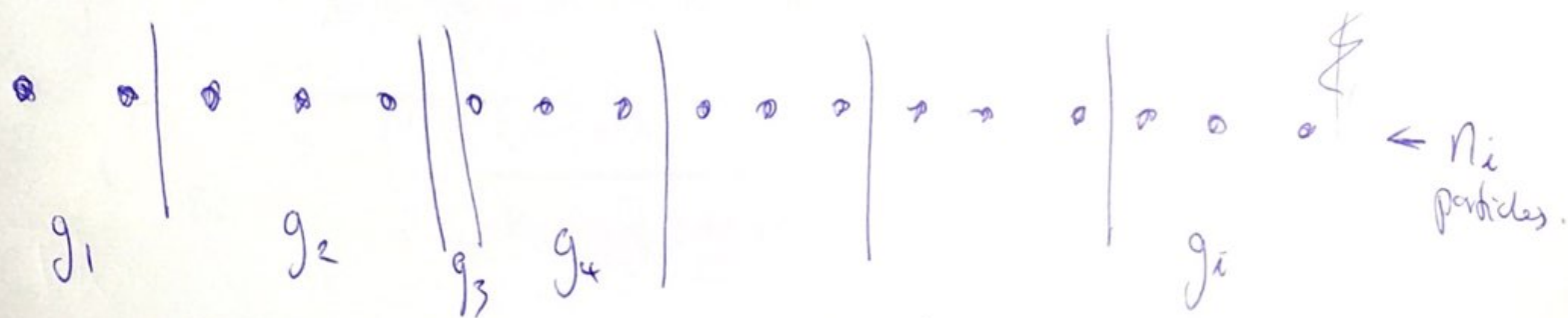
$$\Rightarrow \frac{n_i}{g_i} = f_{FD}(\epsilon_i) = \frac{1}{e^{\alpha} e^{\beta \epsilon_i} + 1}.$$

This is the Fermi-Dirac Distribution function.

Bosons How to count the number of microstates?

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There is no limit on state occupancy, any single particle state can have any number of particles. Let's take a single energy ϵ_i with degeneracy g_i and we have n_i particles to distribute.



We have $g_i - 1$ and n_i particles and so we need the number of arrangements of the $n_i + (g_i - 1)$ objects into 2 piles.

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So n_i particles of energy level ϵ_i in the g_i states

$$\Omega(\epsilon_i) = \frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!} \approx \frac{(n_i + g_i)!}{n_i! g_i!}$$

This is for state i , so for all states we have

$$\Omega_{BE} = \prod_i \frac{(n_i + g_i)!}{n_i! g_i!}$$

So let's play the same game with entropy:

$$S_{k_B} = \ln \left[\prod_i \frac{(n_i + g_i)!}{n_i! g_i!} \right] = \sum_i \left(\ln (n_i + g_i)! - \ln n_i! - \ln g_i! \right) \quad \text{(log rules)} \quad \textcircled{10}$$

$$= \sum_i (n_i + g_i) \ln (n_i + g_i) - (n_i + g_i) - (n_i \ln n_i - n_i) - (g_i \ln g_i - g_i) \quad \text{(Stirling)}$$

$$= \sum_i (n_i + g_i) \ln (n_i + g_i) - n_i \ln n_i - g_i \ln g_i$$

Lagrange to constrain particle number and energy:

$$\sum_i (n_i + g_i) \ln (n_i + g_i) - n_i \ln n_i - g_i \ln g_i - \alpha n_i - \beta n_i \epsilon_i$$

maximise with $\frac{\partial (\dots)}{\partial n_i} = 0$

$$0 = \ln(n_i + g_i) + 1 - \ln n_i - 1 - \alpha - \beta \epsilon_i \quad \text{(no other terms in the sum survives)} \quad (1)$$

$$\Rightarrow \ln \left(\frac{g_i}{n_i} + 1 \right) = \alpha + \beta \epsilon_i$$

$$\text{hence } \frac{n_i}{g_i} = \underbrace{f_{BE}(\epsilon_i)} = \frac{1}{e^{\alpha} e^{\beta \epsilon_i} - 1}$$

This is the Bose-Einstein distribution function.