



Norwegian University of
Science and Technology

Department of Mathematical Sciences

Examination paper for **TMA4125/30/35 Mathematics 4N/4D**

Academic contact during examination: X

Phone: Y

Examination date: -

Examination time (from–to): -

Permitted examination support material: Code C: Approved calculator

One yellow, stamped A5 sheet with own handwritten formulas and notes (on both sides)

Other information:

- All answers have to be justified, and they should include enough details in order to see how they have been obtained.
- There are two versions of Problem 4: one for Mathematics 4N and one for Mathematics 4D.

Language: English

Number of pages: 10

Number of pages enclosed: 2

Checked by:

Informasjon om trykking av eksamensoppgave

Originalen er:

1-sidig ☐ 2-sidig ☒

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skal ha flervalgskjema ☐

Date

Signature

Problem 1 [15 points]

Use the Laplace transform to solve the integral equation:

$$y(t) + 2 \int_0^t y(\tau) \sin(t - \tau) \, d\tau = u(t - 1) - u(t - 3),$$

where u is the Heaviside function (unit step function), given by

$$u(t) = \begin{cases} 0 & t \leq 0, \\ 1 & t > 0. \end{cases}$$

Solution: Applying the Laplace transform to the equation and using the convolution theorem (Laplace transform of a (properly defined) convolution is the product of Laplace transforms), we have

$$(\mathcal{L}y)(s) + 2(\mathcal{L}y)(s) \frac{1}{1 + s^2} = \frac{e^{-s} - e^{-3s}}{s}.$$

Here we have also used

$$\mathcal{L}\{\sin(t)\}(s) = \frac{1}{1 + s^2}, \quad \mathcal{L}\{u(\cdot - a)\} = e^{-as}/s \quad (a > 0),$$

and the linearity of the Laplace transform.

Re-arranging the equation, we find

$$(\mathcal{L}y)(s) = \frac{s^2 + 1}{s^2 + 3} \frac{e^{-s} - e^{-3s}}{s} = \frac{e^{-s} - e^{-3s}}{s} - \frac{2}{3} \left(\frac{1}{s} - \frac{s}{s^2 + 3} \right) (e^{-s} - e^{-3s}).$$

The inverse transforms

$$\begin{aligned} \mathcal{L}^{-1} \left\{ \frac{2}{3} \left(\frac{1}{s} - \frac{s}{s^2 + 3} \right) \right\} (t) &= \frac{2}{3} (1 - \cos(\sqrt{3}t)), \\ \mathcal{L}^{-1} \left\{ e^{-s} - e^{-3s} \right\} (t) &= \delta(t - 1) - \delta(t - 3), \end{aligned}$$

hold in the appropriate senses.

Therefore

$$\begin{aligned}
 y(t) &= u(t-1) - u(t-3) - \frac{2}{3} \int_0^t (1 - \cos(\sqrt{2}(t-\tau))) \delta(\tau-1) \, d\tau \\
 &\quad + \frac{2}{3} \int_0^t (1 - \cos(\sqrt{2}(t-\tau))) \delta(\tau-3) \, d\tau \\
 &= u(t-1) - u(t-3) - \begin{cases} 0 & 0 < t < 1 \\ \frac{2}{3} (1 - \cos(\sqrt{3}(t-1))) & 1 < t < 3 \\ \frac{2}{3} (\cos(\sqrt{3}(t-3)) - \cos(\sqrt{3}(t-1))) & t > 3 \end{cases} \\
 &= \begin{cases} 0 & 0 < t < 1 \\ 1 - \frac{2}{3} (1 - \cos(\sqrt{3}(t-1))) & 1 < t < 3 \\ \frac{2}{3} (\cos(\sqrt{3}(t-1)) - \cos(\sqrt{3}(t-3))) & t > 3 \end{cases}.
 \end{aligned}$$

Problem 2 [10 points]

Let f be a 2π -periodic function, defined over $[-\pi, \pi]$ by

$$f(t) = \begin{cases} |x| - \pi/2 & |x| \geq \pi/2, \\ 0 & |x| < \pi/2. \end{cases}$$

Find the real (or complex) Fourier coefficients of f .

Solution: Notice that f is an even function. Therefore the sine coefficients b_n are all zero. And from Chapter 11.2 of Kreyszig 10th ed, the cosine coefficients for $n > 0$ are

$$\begin{aligned}
 a_n &= \frac{2}{\pi} \int_0^\pi f(x) \cos(nx) \, dx \\
 &= \frac{2}{\pi} \int_{\pi/2}^\pi (x - \pi/2) \cos(nx) \, dx \\
 &= \frac{2}{n\pi} (x - \pi/2) \sin(nx) \Big|_{\pi/2}^\pi - \frac{2}{n\pi} \int_{\pi/2}^\pi \sin(nx) \, dx \\
 &= 0 + \frac{2}{n^2\pi} \cos(nx) \Big|_{\pi/2}^\pi \\
 &= \frac{2}{n^2\pi} (\cos(n\pi) - \cos(n\pi/2)),
 \end{aligned}$$

which can be expressed as

$$\begin{aligned}
 &= \frac{2}{n^2\pi} \cdot \begin{cases} 1-1 & n/4 \in \mathbb{Z} \\ 1-(-1) & (n-2)/4 \in \mathbb{Z} \\ -1-0 & (n-1)/2 \in \mathbb{Z} \end{cases} \\
 &= \frac{2}{n^2\pi} \cdot \begin{cases} 0 & n/4 \in \mathbb{Z} \\ 2 & (n-2)/4 \in \mathbb{Z} \\ -1 & (n-1)/2 \in \mathbb{Z} \end{cases}.
 \end{aligned}$$

and the 0th coefficient is the average over the period of length 2π , which is $(\pi/2)^2/(2\pi) = a_0$, by the formula for the area of a triangle.

Therefore the complex coefficients are $c_0 = a_0$, and

$$c_n = (a_{|n|} - \operatorname{sgn}(n)ib_{|n|})/2 = a_{|n|}/2 = \frac{1}{n^2\pi} \cdot \begin{cases} 0 & n/4 \in \mathbb{Z} \\ 2 & (n-2)/4 \in \mathbb{Z} \\ -1 & (n-1)/2 \in \mathbb{Z} \end{cases}, \quad n \neq 0.$$

Problem 3 [10 points]

Define the convolution for functions $f : \mathbb{R} \rightarrow \mathbb{R}$ and $g : \mathbb{R} \rightarrow \mathbb{R}$ as

$$(f * g)(x) = \int_{\mathbb{R}} f(x-y)g(y) \, dy.$$

Use the convolution theorem to compute $f * f$, where

$$f(x) = e^{-x^2/(2a)},$$

and $a > 0$ is a constant.

Solution: The Fourier transform of the Gaussian can be found in a table:

$$\hat{f}(\xi) = \sqrt{a}e^{-\xi^2 a/2}.$$

By the convolution theorem for Fourier transforms,

$$\widehat{f * f} = \hat{f}^2.$$

Therefore

$$\widehat{f * f}(\xi) = a e^{-\xi^2 a}.$$

Again from the formula sheet and the uniqueness of the Fourier transform,

$$(f * f)(x) = \sqrt{\frac{a}{2}} e^{-x^2/(4a)}.$$

Problem 4 TMA4130 Mathematics 4N: [6 points]

Show that the wave equation,

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial x^2},$$

can be put into the form

$$\frac{\partial^2 u}{\partial y \partial z} = 0,$$

via the change-of-variables $y = x + t$, and $z = x - t$.

Solution: A direct computation gives

$$\begin{aligned} \frac{\partial^2 u}{\partial x^2} &= \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial y} \frac{\partial y}{\partial x} + \frac{\partial u}{\partial z} \frac{\partial z}{\partial x} \right) \\ &= \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial y} + \frac{\partial u}{\partial z} \right) \frac{\partial y}{\partial x} + \frac{\partial}{\partial z} \left(\frac{\partial u}{\partial y} + \frac{\partial u}{\partial z} \right) \frac{\partial z}{\partial x} \\ &= \frac{\partial^2 u}{\partial y^2} + 2 \frac{\partial^2 u}{\partial y \partial z} + \frac{\partial^2 u}{\partial z^2}. \end{aligned}$$

Likewise,

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial^2 u}{\partial y^2} - 2 \frac{\partial^2 u}{\partial y \partial z} + \frac{\partial^2 u}{\partial z^2}.$$

Taking a difference and dividing by 4 on both sides yields the desired equation.

Problem 4 TMA4135 Mathematics 4D: [6 points]

Compute the Fourier transform of the following function:

$$f(x) = \begin{cases} x + 2 & |x| < 1, \\ 0 & |x| \geq 1. \end{cases}$$

Solution: We have

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-1}^1 e^{-ix\omega} (x+2) dx.$$

Notice that

$$\frac{1}{\sqrt{2\pi}} \int_{-1}^1 e^{-ix\omega} (x+2) dx = \frac{1}{\sqrt{2\pi}} \int_{-1}^1 e^{-ix\omega} x dx + \frac{1}{\sqrt{2\pi}} \int_{-1}^1 2e^{-ix\omega} dx.$$

We have

$$\int_{-1}^1 e^{-ix\omega} x dx = \int_{-1}^1 \left(\frac{e^{-ix\omega}}{-i\omega} \right)' x dx = \frac{e^{-ix\omega}}{-i\omega} x \Big|_{-1}^1 - \int_{-1}^1 \left(\frac{e^{-ix\omega}}{-i\omega} \right) dx,$$

which gives

$$\int_{-1}^1 e^{-ix\omega} x dx = \frac{e^{-i\omega} + e^{i\omega}}{-i\omega} - \frac{e^{-i\omega} - e^{i\omega}}{(-i\omega)^2} = \frac{2\cos(\omega)}{-i\omega} - \frac{2i\sin(\omega)}{\omega^2} = \frac{2i(w\cos(\omega) - \sin(\omega))}{\omega^2}.$$

We also have

$$\int_{-1}^1 2e^{-ix\omega} dx = \frac{2(e^{-i\omega} - e^{i\omega})}{-i\omega} = \frac{4\sin(\omega)}{\omega}.$$

Thus

$$\hat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \left(\frac{2i(w\cos(\omega) - \sin(\omega))}{\omega^2} + \frac{4\sin(\omega)}{\omega} \right).$$

Problem 5 [14 points]

Find the solution to the following initial boundary value problem on $[0, \pi]$ using separation-of-variables:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad u(0, t) = 0, \quad u(\pi, t) = 0, \quad u(x, 0) = \sin(2x) + \sin(4x).$$

Solution: Postulating the ansatz $u(t, x) = F(x)G(t)$, and inserting it into the heat equation we find

$$F(x) \frac{dG}{dt} = \frac{d^2 F}{dx^2} G(t).$$

This gives us

$$\frac{1}{G} \frac{dG}{dt} = \frac{1}{F} \frac{d^2 F}{dx^2},$$

which is an equation with a function of t on the left and a function of an independent variable x on the right. Therefore both sides must be equal to some constant k , i.e.,

$$\frac{1}{G_k} \frac{dG_k}{dt} = \frac{1}{F_k} \frac{d^2 F_k}{dx^2} = k,$$

The boundary conditions applied to the ansatz give:

$$G(t)F(0) = G(t)F(\pi) = 0.$$

The only non-trivial solution is $F(0) = F(\pi) = 0$.

If $k = 0$, we find that

$$\frac{d^2 F_k}{dx^2} = 0 \quad \implies \quad F_0(x) = a_0 x + b_0.$$

This is only compatible with the boundary conditions if $a_0 = b_0 = 0$.

Likewise if $k > 0$, we find

$$\frac{d^2 F_k}{dx^2} = k F_k, \tag{1}$$

which has the general solution

$$F_k(x) = A_k e^{\sqrt{k}x} + B_k e^{-\sqrt{k}x}.$$

Again, there are no non-zero constants A_k , B_k for which F_k can satisfy the boundary conditions.

It remains then to $k = -p^2 < 0$ for existence of non-trivial solutions. We find general solutions to (1) are

$$F_p(x) = A_p \cos(px) + B_p \sin(px),$$

where we have re-labelled the solutions F_k by F_p with $p > 0$.

The boundary conditions impose the following restrictions:

$$F_p(0) = 0 \implies A_p = 0, \quad F_p(\pi) = 0 \implies p \in \mathbb{Z}.$$

But we also assumed (in order not to double-count) that $p > 0$. Hence $p \in \mathbb{N} \setminus \{0\}$. Turning to the equation for G_p , we find

$$G_p(t) = C_p e^{-p^2 t}.$$

Writing $D_p = B_p C_p$, we use the linearity of the PDE to write a solution of the form

$$u(t, x) = \sum_{p=1}^{\infty} D_p \sin(px) e^{-p^2 t}.$$

This is enough for us to fit the initial condition. Assuming continuity in t of the series, we can write

$$f(x) = u(0, x) = \sum_{p=1}^{\infty} D_p \sin(px).$$

By extending f in an odd fashion to $[-\pi, \pi]$, to

$$\tilde{f}(x) = \begin{cases} f(x) & 0 \leq x < \pi \\ -f(-x) & -\pi \leq x < 0 \end{cases},$$

and then periodically over \mathbb{R}

$$\begin{aligned} f_-(x) &= \sum_{n \in \mathbb{Z}} \tilde{f}(x - 2n\pi) \mathbb{1}_{[(2n-1)\pi, (2n+1)\pi)} \\ &= \sum_{n \in \mathbb{Z}} f(x - 2n\pi) \mathbb{1}_{[2n\pi, (2n+1)\pi)} - f(-x + 2n\pi) \mathbb{1}_{[-(2n+1)\pi, -2n\pi)} \end{aligned}$$

we see that $f_-(x)$ is an odd function whose restriction to $[0, \pi)$ is f . Therefore using the inversion formula, we can write

$$D_n = \frac{1}{2\pi} \int_0^\pi f_-(x) \sin(nx) \, dx = \frac{1}{2\pi} \int_0^\pi f(x) \sin(nx) \, dx.$$

This determines a solution $u(t, x)$ completely.

Problem 6 [8 points]

Use the central difference method to discretize the second-order equation:

$$u'' + 4xu = r(x), \quad x \in [2, 5], \quad u(2) = 3, \quad u(5) = 4.$$

That is, with an appropriate discretization of $[2, 5]$ into intervals of length h , using the approximation $U_i \approx u(x_i)$, and writing $R_i = r(x_i)$, write down the discrete approximation to the differential equation involving the second-order central difference of u at x_i .

Solution: The domain can be discretized using

$$h = \frac{5-2}{N}, \quad x_i = 2 + ih, \quad i = 0, \dots, N.$$

Discretizing the equation, we use the central difference approximation

$$u''(x) = \frac{u(x+h) - 2u(x) + u(x-h)}{h^2} + O(h^2),$$

so that

$$u''(x_i) \approx \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2}.$$

Hence the equation can be discretized by the $N + 1$ formulae:

$$U_{i+1} - 2U_i + U_{i-1} + 4x_i h^2 U_i = R_i h^2, \quad n = 1, \dots, N - 1,$$

and

$$U_0 = 3, \quad U_N = 4.$$

Problem 7 [8 points]

a) Given the ordinary differential equation

$$y' = x^2 y, \quad y(0) = 1.$$

Write down the implicit (backward) Euler method for this equation for a given step size h .

b) Choose $h = 0.1$ and compute an approximate value for $y(0.2)$

Solution:

(a) First we discretize the half-line : $\mathbf{x}_n = nh$. The implicit Euler scheme is given by

$$\mathbf{y}_{n+1} = \mathbf{y}_n + h \cdot f(\mathbf{x}_{n+1}, \mathbf{y}_{n+1}).$$

Putting in $f(\mathbf{x}_{n+1}, \mathbf{y}_{n+1}) = \mathbf{x}_{n+1}^2 \mathbf{y}_{n+1} = (n+1)^2 h^2 \mathbf{y}_{n+1}$, we arrive at

$$\mathbf{y}_{n+1} = \mathbf{y}_n + h \mathbf{x}_{n+1}^2 \mathbf{y}_{n+1} = \mathbf{y}_n + (n+1)^2 h^3 \mathbf{y}_{n+1},$$

with initial condition $\mathbf{y}_0 = 1$.

(b) Solving for \mathbf{y}_{n+1} at each step, one finds:

$$(1 - (n+1)^2 h^3) \mathbf{y}_{n+1} = \mathbf{y}_n.$$

First, from the discretization $\mathbf{x}_n = nh$, we can assume that $y(0.2) \approx \mathbf{y}_2$. Therefore calculating as follows:

$$\begin{aligned} \mathbf{y}_1 &= \frac{1}{(1 - h^3)} \mathbf{y}_0 = \frac{1}{0.999}, \\ \mathbf{y}_2 &= \frac{1}{(1 - 2^2 h^3)} \mathbf{y}_1 = \frac{1}{0.996} \cdot \frac{1}{0.999}, \end{aligned}$$

we can conclude that

$$y(0.2) \approx 1.005021085,$$

keeping in mind that the global error is $O(h)$.

Problem 8 [7 points]

Find the interpolation polynomial of lowest degree for the following points:

x_n	-4	0	1	2
$f(x_n)$	-85	-9	0	29

Solution:

The lowest degree polynomial that can be used to interpolate three points is cubic. The solution is $2x^3 + 4x^2 + 3x - 9$.

Problem 9 [10 points]

Recall the following difference formula for a four times continuously differentiable function $f: \mathbb{R} \rightarrow \mathbb{R}$:

$$f''(a) \text{ can be approximated by } \frac{f(a+h) - 2f(a) + f(a-h)}{h^2}.$$

Assume that $|f^{(4)}(x)| \leq 1$ for all $x \in \mathbb{R}$. Use fourth order Taylor expansion to show the following error estimate

$$\left| f''(a) - \frac{f(a+h) - 2f(a) + f(a-h)}{h^2} \right| \leq \frac{h^2}{12}.$$

Solution: Use Taylor's expansion.

Problem 10 [12 points]

a) Show that the equation

$$e^{-x^2} = x$$

has a unique solution on the real line.

b) Write down a bisection method for this equation, using $[0, 1]$ as the initial interval, and compute the next iteration.

c) Write down the Newton method for this equation. Compute the next iteration x_1 , using $x_0 = 0.5$ as the initial point.

Solution: (a) Since the left-hand side of the equation is always positive, we cannot have any solution for x non-positive. Consider the function

$$f: [0, \infty) \rightarrow \mathbb{R}, \quad f(x) = e^{-x^2} - x$$

We have

$$f'(x) = -2xe^{-x^2} - 1 = -(2xe^{-x^2} + 1) < 0$$

which shows that f is strictly decreasing. Since $f(0) = 1 > 0$ and $f(1) = e^{-1} - 1 < 0$ the function f has a unique zero on the interval $[0, 1]$ and no zeros outside this interval.

(b) [Method p. 2 in notes] Let $a_0 = 0$ and $b_0 = 1$ and $c_0 = (a_0 + b_0)/2 = 0.5$. Since $f(c_0) = e^{-0.5^2} - 0.5 = 0.28 > 0$, we let $a_1 = c_0 = 0.28$ and $b_1 = b_0 = 1$. We find then $c_1 = (a_1 + b_1)/2 = 0.64$.

(c) The Newton iteration equals

$$x_{k+1} = x_k - \frac{f(x_k)}{f'(x_k)}, \quad k = 0, 1, 2, \dots$$

We have

$$f'(x) = -2xe^{-x^2} - 1 = -(2xe^{-x^2} + 1),$$

and thus

$$x_0 = 0.5, \quad x_1 = x_0 - \frac{f(x_0)}{f'(x_0)} = 0.5 - \frac{f(0.5)}{f'(0.5)} \approx 0.656735.$$

Fourier Transform

$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(w) e^{iwx} \, dw$	$\hat{f}(w) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-iwx} \, dx$
e^{-ax^2}	$\frac{1}{\sqrt{2a}} e^{-w^2/4a}$
$e^{-a x }$	$\sqrt{\frac{2}{\pi}} \frac{a}{w^2 + a^2}$
$\frac{1}{x^2 + a^2}$	$\sqrt{\frac{\pi}{2}} \frac{e^{-a w }}{a}$
$\begin{cases} 1 & \text{for } x < a \\ 0 & \text{otherwise} \end{cases}$	$\sqrt{\frac{2}{\pi}} \frac{\sin wa}{w}$

Laplace Transform

$f(t)$	$F(s) = \int_0^{\infty} e^{-st} f(t) \, dt$
$\cos(\omega t)$	$\frac{s}{s^2 + \omega^2}$
$\sin(\omega t)$	$\frac{\omega}{s^2 + \omega^2}$
$\cosh(\omega t)$	$\frac{s}{s^2 - \omega^2}$
$\sinh(\omega t)$	$\frac{\omega}{s^2 - \omega^2}$
t^n	$\frac{\Gamma(n+1)}{s^{n+1}},$ for $n = 0, 1, 2, \dots$, $\Gamma(n+1) = n!$
e^{at}	$\frac{1}{s - a}$
$\delta(t - a)$	e^{-as}

$$\int x^n \cos ax \, dx = \frac{1}{a} x^n \sin ax - \frac{n}{a} \int x^{n-1} \sin ax \, dx$$

$$\int x^n \sin ax \, dx = -\frac{1}{a} x^n \cos ax + \frac{n}{a} \int x^{n-1} \cos ax \, dx$$

Numerics

- Newton's method: $x_{k+1} = x_k - \frac{f(x_k)}{f'(x_k)}$.
- Newton's method for system of equations: $\vec{x}_{k+1} = \vec{x}_k - JF(\vec{x}_k)^{-1}F(\vec{x}_k)$, with $JF = (\partial_j f_i)$.
- Lagrange interpolation: $p_n(x) = \sum_{k=0}^n \frac{l_k(x)}{l_k(x_k)} f_k$, with $l_k(x) = \prod_{j \neq k} (x - x_j)$.
- Interpolation error: $\epsilon_n(x) = \prod_{k=0}^n (x - x_k) \frac{f^{(n+1)}(t)}{(n+1)!}$.
- Chebyshev points: $x_k = \cos\left(\frac{2k+1}{2n+2}\pi\right)$, $0 \leq k \leq n$.
- Newton's divided difference: $f(x) \approx f_0 + (x - x_0)f[x_0, x_1] + (x - x_0)(x - x_1)f[x_0, x_1, x_2] + \dots + (x - x_0)(x - x_1) \dots (x - x_{n-1})f[x_0, \dots, x_n]$, with $f[x_0, \dots, x_k] = \frac{f[x_1, \dots, x_k] - f[x_0, \dots, x_{k-1}]}{x_k - x_0}$.
- Trapezoid rule: $\int_a^b f(x) \, dx \approx h \left[\frac{1}{2}f(a) + f_1 + f_2 + \dots + f_{n-1} + \frac{1}{2}f(b) \right]$.
Error of the trapezoid rule: $|\epsilon| \leq \frac{b-a}{12} h^2 \max_{x \in [a,b]} |f''(x)|$.
- Simpson rule: $\int_a^b f(x) \, dx \approx \frac{h}{3} [f_0 + 4f_1 + 2f_2 + 4f_3 + \dots + 2f_{n-2} + 4f_{n-1} + f_n]$.
Error of the Simpson rule: $|\epsilon| \leq \frac{b-a}{180} h^4 \max_{x \in [a,b]} |f^{(4)}(x)|$.
- Gauss–Seidel iteration: $\mathbf{x}^{(m+1)} = \mathbf{b} - \mathbf{L}\mathbf{x}^{(m+1)} - \mathbf{U}\mathbf{x}^{(m)}$, with $\mathbf{A} = \mathbf{I} + \mathbf{L} + \mathbf{U}$.
- Jacobi iteration: $\mathbf{x}^{(m+1)} = \mathbf{b} + (\mathbf{I} - \mathbf{A})\mathbf{x}^{(m)}$.
- Euler method: $\mathbf{y}_{n+1} = \mathbf{y}_n + h\mathbf{f}(x_n, \mathbf{y}_n)$.
- Improved Euler method: $\mathbf{y}_{n+1} = \mathbf{y}_n + \frac{1}{2}h[\mathbf{f}(x_n, \mathbf{y}_n) + \mathbf{f}(x_n + h, \mathbf{y}_{n+1}^*)]$, where $\mathbf{y}_{n+1}^* = \mathbf{y}_n + h\mathbf{f}(x_n, \mathbf{y}_n)$.
- Classical Runge–Kutta method: $\mathbf{k}_1 = h\mathbf{f}(x_n, \mathbf{y}_n)$,
 $\mathbf{k}_2 = h\mathbf{f}(x_n + h/2, \mathbf{y}_n + \mathbf{k}_1/2)$, $\mathbf{k}_3 = h\mathbf{f}(x_n + h/2, \mathbf{y}_n + \mathbf{k}_2/2)$,
 $\mathbf{k}_4 = h\mathbf{f}(x_n + h, \mathbf{y}_n + \mathbf{k}_3)$, $\mathbf{y}_{n+1} = \mathbf{y}_n + \frac{1}{6}\mathbf{k}_1 + \frac{1}{3}\mathbf{k}_2 + \frac{1}{3}\mathbf{k}_3 + \frac{1}{6}\mathbf{k}_4$.
- Backward Euler method: $\mathbf{y}_{n+1} = \mathbf{y}_n + h\mathbf{f}(x_{n+1}, \mathbf{y}_{n+1})$.
- Finite differences: $\frac{\partial u}{\partial x}(x, y) \approx \frac{u(x+h, y) - u(x-h, y)}{2h}$, $\frac{\partial^2 u}{\partial x^2}(x, y) \approx \frac{u(x+h, y) - 2u(x, y) + u(x-h, y)}{h^2}$.
- Crank–Nicolson method for the heat equation: $r = \frac{k}{h^2}$,
 $(2 + 2r)u_{i,j+1} - r(u_{i+1,j+1} + u_{i-1,j+1}) = (2 - 2r)u_{ij} + r(u_{i+1,j} + u_{i-1,j})$.