

Week 5

Review

- * Operations on quantum states are norm-preserving (Unitary)

$$\rightarrow U^\dagger = U^{-1} \Rightarrow U^\dagger U = U U^\dagger = I$$

- * Unitary operations for a single qubit are rotations on the surface of Bloch sphere

- * Eigenvalues and eigenvectors of unitary matrices (or operations) are special.

$$\rightarrow U|x\rangle = \lambda_x |x\rangle \quad \begin{array}{l} \nearrow \text{eigenvector} \\ \uparrow \text{eigenvalue} \end{array} \quad \begin{array}{l} \checkmark \\ \text{of the form } e^{i\theta} \\ \text{always} \end{array}$$

\rightarrow For two distinct λ_x and λ_y :

$\lambda_x = \lambda_y$: $|x\rangle$ and $|y\rangle$ are orthogonal

$$\langle x|y\rangle = 0$$

* Shor's Algorithm:

$$f(x) = f(y) \text{ for } x \neq y \text{ iff } |x-y| = np \quad \begin{array}{l} \nwarrow \\ \text{period} \end{array}$$

Classically: $O(e^{c \cdot n^{1/3} \cdot (\log n)^{2/3}})$
 \downarrow
 n bits needed to describe the period

Quantum: $O(n^2 \log n \cdot \log \log n)$
 \hookrightarrow A little better than $O(n^3)$

\rightarrow QFT and modular exponentiation are the building blocks of Shor's Alge. \rightarrow Quantum Phase Estimation

\rightarrow Factorisation \longleftrightarrow periodicity

\blacktriangleright QFT: A change of basis from computational basis to Fourier basis.

Comp: $M_0 = \langle 0|0 \rangle$ $M_1 = \langle 1|1 \rangle$

Fourier: $M_+ = \langle +|+ \rangle$ $M_- = \langle -|- \rangle$

$|\tilde{0}\rangle = |+\rangle$; $|\tilde{1}\rangle = |-\rangle$

\rightarrow 1 Qubit QFT is Hadamard gate.

\rightarrow For n qubits, we have $N = 2^n$ basis states

$$|\tilde{x}\rangle = \underset{\substack{\uparrow \\ \text{Fourier basis}}}{\text{QFT}} |x\rangle = \frac{1}{\sqrt{N}} \sum_{y=0}^{N-1} e^{\frac{2\pi i x y}{N}} |y\rangle \rightarrow |y\rangle \text{ comes to binary notation of } y$$

$$\text{ex: } \text{QFT } |1\rangle = \frac{1}{\sqrt{2}} (|0\rangle + e^{i\pi \cdot 1} |1\rangle) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

→ Notation:

$|y\rangle$ is actually $|y_1 y_2 \dots y_n\rangle$

$\therefore \sum_{y=0}^N \rightarrow \sum_{y_0=0}^1 \sum_{y_1=0}^1 \dots \sum_{y_n=0}^1$ Binary version of y

$\therefore |x_1\rangle \otimes |x_2\rangle \otimes |x_3\rangle \otimes \dots \otimes |x_n\rangle = |x\rangle$

\downarrow QFT

$\frac{1}{\sqrt{N}} (|0\rangle + e^{\frac{2\pi i x}{2^1}} |1\rangle) \otimes (|0\rangle + e^{\frac{2\pi i x}{2^2}} |1\rangle) \otimes \dots \otimes (|0\rangle + e^{\frac{2\pi i x}{2^n}} |1\rangle) = |\tilde{x}\rangle$

ex: $n=3$ qubits $\rightarrow N = 2^3 = 8$

$|x\rangle = |5\rangle = |101\rangle = |1\rangle \otimes |0\rangle \otimes |1\rangle$

QFT $|x\rangle = |\tilde{x}\rangle = |\tilde{5}\rangle$

$= \frac{1}{\sqrt{8}} \left(|0\rangle + e^{\frac{2\pi i 5}{2}} |1\rangle \right) \otimes \left(|0\rangle + e^{\frac{2\pi i 5}{4}} |1\rangle \right) \otimes \left(|0\rangle + e^{\frac{2\pi i 5}{8}} |1\rangle \right)$

$\rightarrow |x_k\rangle \rightarrow |0\rangle + e^{\frac{2\pi i x}{2^k}} |1\rangle \quad [\text{QFT}]$

- Phase is qubit dependent
- Terms with more ^{more} 1's occur more and more \rightarrow ex: $|1011\rangle, |1101\rangle, |1110\rangle$ terms would need 2 phase terms

$$* H|x_k\rangle = \left(|0\rangle + e^{\frac{2\pi i x_k}{2}} |1\rangle \right) / \sqrt{2}$$

if \uparrow 0 or 1

$\xrightarrow{\quad} e^{\frac{2\pi i x_k}{2}} : \begin{matrix} x_k=0 \rightarrow 1 \\ x_k=1 \rightarrow -1 \end{matrix}$

$$* UROT_k |x_j\rangle = e^{\frac{2\pi i}{2^k} x_j} |x_j\rangle$$

Unitary rotation \downarrow

controlled rotation $\xrightarrow{\quad} \begin{matrix} x_j=0 \Rightarrow e^{\frac{2\pi i}{2^k} x_j} |x_j\rangle = |0\rangle \\ x_j=1 \Rightarrow e^{\frac{2\pi i}{2^k} x_j} |x_j\rangle = e^{\frac{2\pi i}{2^k}} |x_j\rangle \end{matrix}$

For 1 qubit $\rightarrow UROT_k = \begin{bmatrix} 1 & 0 \\ 0 & e^{\frac{2\pi i}{2^k}} \end{bmatrix}$

\nwarrow applying a phase if $|1\rangle$ (or)

* Quantum Phase Estimation:

$$U|\psi\rangle = e^{i\theta\psi} |\psi\rangle$$

\uparrow an eigenvector of U

\rightarrow QPE corresponds to measuring the global phase of a state

\rightarrow Assuming that we have the capacity to prepare a given $|\psi\rangle$ and apply U to it:

$$QPE|0\rangle|\psi\rangle = \frac{1}{2} \left[|0\rangle(1+e^{i\theta\psi}) + |1\rangle(1-e^{i\theta\psi}) \right] |\psi\rangle$$

\rightarrow After a lot of measurements, we can estimate $\theta\psi$ based on the $|0\rangle$ and $|1\rangle$ meas. on first qubit

$$\begin{aligned}
\rightarrow U^{2^x} |\psi\rangle &= U^{2^x-1} U |\psi\rangle \\
&= U^{2^x-1} e^{i\theta\psi} |\psi\rangle \\
&= e^{i\theta\psi + i\theta\psi} U^{2^x-2} |\psi\rangle \\
&= e^{2i\theta\psi} |\psi\rangle
\end{aligned}$$

→ If we use n qubits for meas. instead of 1:

$$\begin{aligned}
QPE |0\rangle^{\otimes n} |\psi\rangle &= \left(\frac{1}{\sqrt{2}} \right)^n (|0\rangle + e^{i\theta\psi 2^{n-1}} |1\rangle) \otimes \\
&\quad (|0\rangle + e^{i\theta\psi 2^{n-2}} |1\rangle) \otimes \dots \otimes \\
&\quad (|0\rangle + e^{i\theta\psi 2^0} |1\rangle) |\psi\rangle
\end{aligned}$$

→ QPE is same as QFT except:

$$\begin{array}{ccc}
& \theta_\psi & \rightarrow \frac{\theta_\psi 2\pi}{2^n} \\
\text{QFT} \nearrow \text{phase} & & \nwarrow \text{QPE phase}
\end{array}$$

$$\rightarrow (QFT^{-1})(QPE) |0\rangle^{\otimes n} |\psi\rangle = |2^n \theta_\psi\rangle |\psi\rangle$$

↑
This factor increases the amplitude for phase estimate