Finite Element Methods

FEM for Poisson problem in 2D - Rectangular element

2016-2 CSE6820

Outline

Introduction

Affine mapping

Triangulation

Basis functions of V_h^k

Mass matrix and Stiffness matrix

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2D Poisson problem

Consider the two dimensional Poisson problem

$$-\Delta u = f \qquad \text{in } \Omega$$
$$u = u_D \qquad \text{on } \partial \Omega$$

Introduction

Weak formulation

Find $u(x, y) \in H_0^1(\Omega)$ such that

$$\int_{\Omega} \nabla u \cdot \nabla v \, d\boldsymbol{x} = \int_{\Omega} f v \, d\boldsymbol{x}, \qquad \forall v \in H^1_0(\Omega).$$

Variational formulation

Find $u_h \in V_h^k$ such that

$$\int_{\Omega} \nabla u_h \cdot \nabla v_h \ d\mathbf{x} = \int_{\Omega} f v_h \ d\mathbf{x}, \qquad \forall v_h \in V_h^k$$

where

Introduction

$$V_h^k = \{ v_h \in H_0^1(\Omega) \mid v_h|_R \in Q_k(R), \ \forall R \in \mathcal{T}_h \},$$

and $Q_k(R)$ is the polynomial function space of degrees $\leq k$ in each variable.

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Affine mapping

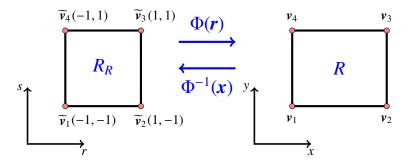


Figure: An affine mapping from the reference rectangle R_Q to a rectangle Q.

Barycentric coordinates

Barycentric coordinates $\{\lambda_i\}_{i=1}^4$ have the properties

$$\begin{cases} 0 \le \lambda_i(\mathbf{x}) \le 1, & i = 1, 2, 3, 4 \\ \sum_{i=1}^4 \lambda_i(\mathbf{x}) = 1. \end{cases}$$

Then we have an affine mapping Φ such that

$$\Phi(\mathbf{r}) = \mathbf{v}_1 + \frac{r+1}{2}(\mathbf{v}_2 - \mathbf{v}_1) + \frac{s+1}{2}(\mathbf{v}_4 - \mathbf{v}_1) = \mathbf{x},$$

where x is a point in R.

Properties

1.
$$\widetilde{\lambda}_k(\mathbf{r}) = \widetilde{\lambda}_i^{1D}(r)\widetilde{\lambda}_i^{1D}(s), \qquad (k = 2(j-1) + i, \quad i, j = 1, 2)$$

2.
$$x_r = \frac{v_2^{(1)} - v_1^{(1)}}{2}$$
, $y_r = \frac{v_2^{(2)} - v_1^{(2)}}{2}$, $x_s = \frac{v_4^{(1)} - v_1^{(1)}}{2}$, $y_s = \frac{v_4^{(2)} - v_1^{(2)}}{2}$

3.
$$r_x = \frac{y_s}{I}$$
, $r_y = -\frac{x_s}{I}$, $s_x = -\frac{y_r}{I}$, $s_y = \frac{x_r}{I}$

4.
$$\lambda_{1}(\mathbf{x}) = \left(\frac{v_{2}^{(1)} - x}{v_{1}^{(1)} - v_{1}^{(1)}}\right) \left(\frac{v_{4}^{(2)} - y}{v_{4}^{(2)} - v_{1}^{(2)}}\right), \qquad \lambda_{2}(\mathbf{x}) = \left(\frac{x - v_{1}^{(1)}}{v_{2}^{(1)} - v_{1}^{(1)}}\right) \left(\frac{v_{4}^{(2)} - y}{v_{4}^{(2)} - v_{1}^{(2)}}\right)$$

$$\lambda_{3}(\mathbf{x}) = \left(\frac{x - v_{1}^{(1)}}{v_{2}^{(1)} - v_{1}^{(1)}}\right) \left(\frac{y - v_{1}^{(2)}}{v_{2}^{(2)} - v_{2}^{(2)}}\right), \qquad \lambda_{4}(\mathbf{x}) = \left(\frac{v_{2}^{(1)} - x}{v_{2}^{(1)} - v_{1}^{(1)}}\right) \left(\frac{y - v_{1}^{(2)}}{v_{2}^{(2)} - v_{2}^{(2)}}\right)$$

5.
$$\lambda_i(\mathbf{x}) = \widetilde{\lambda}_i(\Phi^{-1}(\mathbf{x}))$$
, $\widetilde{\lambda}_i(\mathbf{r}) = \lambda_i(\Phi(\mathbf{r}))$

5.
$$\lambda_i(\mathbf{x}) = \overline{\lambda}_i(\Phi^{-1}(\mathbf{x})), \qquad \overline{\lambda}_i(\mathbf{r}) = \lambda_i(\Phi(\mathbf{r}))$$
6. $\frac{d}{dx}\lambda_i(\mathbf{x}) = r_x \frac{d}{dx}\widetilde{\lambda}_i(\mathbf{r}), \qquad \frac{d}{dy}\lambda_i(\mathbf{x}) = s_y \frac{d}{ds}\widetilde{\lambda}_i(\mathbf{r})$

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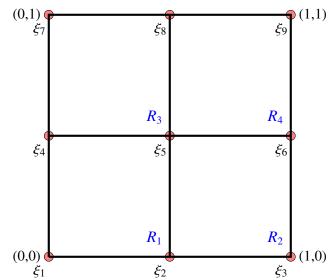
Basis functions of V_h^k

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Triangulation data (P_1)

If $\Omega = [0, 1]^2$, M = 2 and k = 1, data are stored as follows.



Triangulation data (P_1)

$$c4n = \begin{pmatrix} 0 & 1/2 & 1 & 0 & 1/2 & 1 & 0 & 1/2 & 1 \\ 0 & 0 & 0 & 1/2 & 1/2 & 1/2 & 1 & 1 & 1 & 1 \end{pmatrix}$$

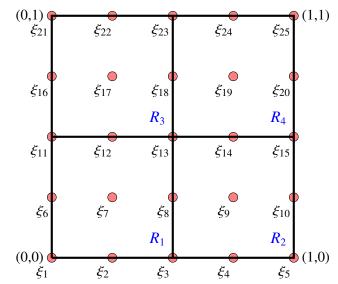
$$n4e = \begin{pmatrix} 1 & 2 & 4 & 5 \\ 2 & 3 & 5 & 6 \\ 5 & 6 & 8 & 9 \\ 4 & 5 & 7 & 8 & 1 \end{pmatrix}$$

$$n4db = (1, 2, 3, 4, 6, 7, 8, 9)$$

$$ind4e = \begin{pmatrix} 1 & 2 & 4 & 5 \\ 2 & 3 & 5 & 6 \\ 4 & 5 & 7 & 8 \\ 5 & 6 & 8 & 9 \end{pmatrix}$$

Triangulation data (P_2)

If $\Omega = [0, 1]^2$, M = 2 and k = 2, data are stored as follows.



Triangulation data (P_2)

$$n4e = \begin{pmatrix} 1 & 3 & 11 & 13 \\ 3 & 5 & 13 & 15 \\ 13 & 15 & 23 & 25 \\ 11 & 13 & 21 & 23 \end{pmatrix}$$

n4db = (1, 2, 3, 4, 5, 6, 10, 11, 15, 16, 20, 21, 22, 23, 24, 25)

$$ind4e = \left(\begin{array}{ccccc} 1 & 3 & 11 & 13 \\ 2 & 4 & 12 & 14 \\ 3 & 5 & 13 & 15 \\ 6 & 8 & 16 & 18 \\ 7 & 9 & 17 & 19 \\ 8 & 10 & 18 & 20 \\ 11 & 13 & 21 & 23 \\ 12 & 14 & 22 & 24 \\ 13 & 15 & 23 & 25 \\ \end{array} \right)$$

mesh FEM2D Rec

The following Matlab code generates an uniform rectangular mesh on the domain $[xl, xr] \times [yl, yr]$ in 2D with M_x elements along x-direction and M_y elements along y-direction. Also this code returns an index matrix for continuous k-th order polynomial approximations.

```
function [c4n,n4e,ind4e,inddb]=mesh_FEM2D_Rec(x1,xr,y1,yr,Mx,My,k)
ind4e = zeros((k+1)^2,Mx*My);
tmp = (1:k:k*Mx)' * ones(1,My) ...
    + ones(Mx,1) * (0:k*(k*Mx+1):((k*Mx+1)*((My-1)*k+1)-1));
tmp = tmp(:)':
for i=1:k+1
    ind4e((j-1)*(k+1)+(1:(k+1)), :) = repmat(tmp+(j-1)*(k*Mx+1),k+1,1) ...
        +repmat(0:k.Mx*Mv.1)':
end
n4e = ind4e([1 k+1 (k+1)^2 (k*(k+1)+1)],:);
inddb = unique([1:(k*Mx+1), (k*Mx+1):(k*Mx+1):(k*Mx+1)*(k*My+1), ...
(k*Mx+1)*(k*My+1):-1:(k*My*(k*Mx+1)+1), (k*My*(k*Mx+1)+1):-(k*Mx+1):1]);
x=linspace(xl,xr,k*Mx+1);
v=linspace(vl.vr.k*Mv+1):
y=repmat(y,k*Mx+1,1);
x=repmat(x,k*My+1,1)';
c4n = [x(:), y(:)]';
end
```

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Basis functions

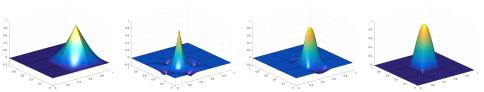
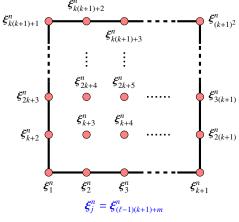


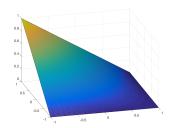
Figure: Global basis functions of V_h^1 (left) and V_h^2 (others) on an interval.

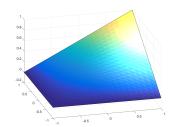
$$\psi_i(\boldsymbol{\xi}_j) = \delta_{ij} = \begin{cases} 1 & \text{if } i = j, \\ 0 & \text{if } i \neq j, \end{cases}$$
$$\sum_{i=1}^{N} \psi_i(\boldsymbol{x}) = 1 \qquad \forall x \in \Omega$$

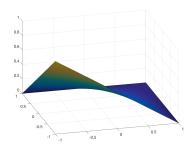


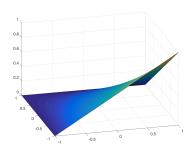
 $(\ell : \text{row number}, m : \text{column number})$

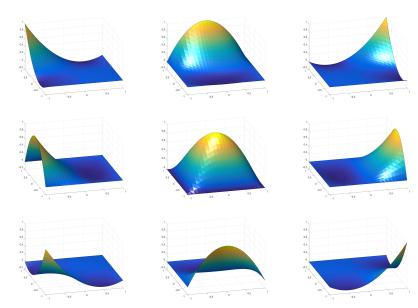
Figure: Node numbering in the *n*-th rectangular element











$$\psi_i^n(\mathbf{x}) = \phi_\ell(x)\phi_m(y)$$

$$\psi_i^n(\mathbf{x}) = 0 \quad \text{if } x \in \Omega \setminus R_n$$

$$\sum_{i=1}^{(k+1)^2} \psi_i^n(\mathbf{x}) = 1 \quad \forall x \in R_n.$$

Interpolation

Using these basis functions, the interpolate function of a function f(x) can be written as follows

$$If(\mathbf{x}) = \sum_{i=1}^{N} f_i \psi_i(\mathbf{x})$$

where $f_i = f(\boldsymbol{\xi}_i)$.

• Numerical solution

$$u_h = \sum_{i=1}^N u_i \psi_i$$

Local solution

$$u_h\Big|_{R_n} = \sum_{i=1}^{(k+1)^2} u_i^n \psi_i^n$$

Derivatives

$$\nabla u_h = \sum_{i=1}^{N} u_i \nabla \psi_i,$$

$$\nabla u_h \Big|_{R_n} = \sum_{i=1}^{(k+1)^2} u_i^n \nabla \psi_i^n.$$

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Variational formulation

For a test function $\psi_i(x) \in V_h^k$, the variational formulation can be rewritten as

$$\sum_{i=1}^{N} u_{i} \int_{\Omega} \nabla \psi_{i} \cdot \nabla \psi_{j} \, d\mathbf{x} = \int_{\Omega} f \psi_{i} \, d\mathbf{x}$$

Finite element system

$$Au = b$$

where

$$(A)_{ij} = \int_{\Omega} \nabla \psi_i \cdot \nabla \psi_j \, d\mathbf{x}$$
$$(\mathbf{b})_i = \int_{\Omega} f \psi_i \, d\mathbf{x}$$
$$(\mathbf{u})_j = u_j.$$

Finite element system

$$Au = Mf$$

where

$$(A)_{ij} = \int_{\Omega} \nabla \psi_i \cdot \nabla \psi_j \, d\mathbf{x}$$
$$(\mathbf{u})_j = u_j$$
$$(M)_{ij} = \int_{\Omega} \psi_i \psi_j \, d\mathbf{x}$$
$$(\mathbf{f})_j = f_j.$$

Matrices

$$A = \sum_{\ell=1}^{M} A_{R_n}$$

$$M = \sum_{\ell=1}^{M} M_{R_n}$$

where $(k + 1)^2$ -by- $(k + 1)^2$ matrices A_{R_n} and M_{R_n} are defined as

$$(A_{R_n})_{ij} = \int_{R_n} \nabla \psi_i^n(x) \cdot \nabla \psi_j^n \, dx$$
$$(M_{R_n})_{ij} = \int_{R_n} \psi_i^n \psi_j^n \, dx$$

where $1 \le i, j \le (k+1)^2$.

Differentiation matrix

$$(Dx)_{ij} = \frac{\partial \psi_j}{\partial x}(\boldsymbol{\xi}_i), \qquad (Dy)_{ij} = \frac{\partial \psi_j}{\partial y}(\boldsymbol{\xi}_i).$$

$$\frac{\partial}{\partial x}\psi_i(\mathbf{x}) = \sum_{j=1}^N \frac{\partial \psi_i}{\partial x}(\boldsymbol{\xi}_j)\psi_j(\mathbf{x}) = (Dx^t)_i \boldsymbol{\psi}$$

$$\frac{\partial}{\partial y}\psi_i(x) = \sum_{i=1}^N \frac{\partial \psi_i}{\partial y}(\boldsymbol{\xi}_j)\psi_j(x) = (Dy^t)_i \boldsymbol{\psi}$$

where

$$(Dx^{t})_{i} = \left[\frac{\partial \psi_{i}}{\partial x}(\xi_{1}) \cdots \frac{\partial \psi_{i}}{\partial x}(\xi_{N})\right]$$
$$(Dy^{t})_{i} = \left[\frac{\partial \psi_{i}}{\partial y}(\xi_{1}) \cdots \frac{\partial \psi_{i}}{\partial y}(\xi_{N})\right]$$
$$\psi = \left[\psi_{1}(x) \cdots \psi_{N}(x)\right]^{t}$$

Properties of D

•
$$Dx = \sum_{n=1}^{M^2} Dx_{R_n}, \qquad (Dx_{R_n})_{ij} = \frac{\partial \psi_j^n}{\partial x} (\boldsymbol{\xi}_i^n)$$

•
$$Dy = \sum_{n=1}^{M^2} Dx_{R_n}, \qquad (Dy_{R_n})_{ij} = \frac{\partial \psi_j^n}{\partial y} (\boldsymbol{\xi}_i^n)$$

•
$$\nabla \psi_i^n(\mathbf{x}) = \left((Dx_{R_n}^t)_i \boldsymbol{\psi}^n, \ (Dy_{R_n}^t)_i \boldsymbol{\psi}^n \right)$$

•
$$\nabla u_h(\boldsymbol{\xi}_m) = ((Dx)_m \boldsymbol{u}, (Dy)_m \boldsymbol{u})$$

•
$$\nabla u_h(\boldsymbol{\xi}_m^n) = ((Dx_{R_n})_m \boldsymbol{u}, (Dy_{R_n})_m \boldsymbol{u})$$

Mass matrix

$$(M)_{ij} = \int_{R_n} \psi_i(\mathbf{x}) \psi_j(\mathbf{x}) \ d\mathbf{x} = J \int_{R_R} \widetilde{\psi}_i(\mathbf{r}) \widetilde{\psi}_j(\mathbf{r}) \ d\mathbf{r} = J(M_R)_{ij}$$

$$\Rightarrow M = JM_R$$

Stiffness matrix

$$(S)_{ij} = J\Big(r_x^2(S_R^{rr})_{ij} + s_y^2(S_R^{ss})_{ij}\Big)$$

where

$$(S_R^{rr})_{ij} = (Dr_R^t M_R Dr_R)_{ij}$$

$$(S_R^{ss})_{ij} = (Ds_R^t M_R Ds_R)_{ij}$$

$$\Rightarrow \qquad S = J(r_x^2 S_R^{rr} + s_y^2 S_R^{ss})$$

$$\begin{split} \widetilde{\psi}_{1}(\boldsymbol{r}) &= \widetilde{\phi}_{1}(r)\widetilde{\phi}_{1}(s) = \widetilde{\lambda}_{1}^{1D}(r)\widetilde{\lambda}_{1}^{1D}(s), & \nabla \widetilde{\psi}_{1}(\boldsymbol{r}) = \left(-\frac{1}{2}\widetilde{\lambda}_{1}^{1D}(s), -\frac{1}{2}\widetilde{\lambda}_{1}^{1D}(r)\right), \\ \widetilde{\psi}_{2}(\boldsymbol{r}) &= \widetilde{\phi}_{2}(r)\widetilde{\phi}_{1}(s) = \widetilde{\lambda}_{2}^{1D}(r)\widetilde{\lambda}_{1}^{1D}(s), & \nabla \widetilde{\psi}_{2}(\boldsymbol{r}) = \left(\frac{1}{2}\widetilde{\lambda}_{1}^{1D}(s), -\frac{1}{2}\widetilde{\lambda}_{2}^{1D}(r)\right), \\ \widetilde{\psi}_{3}(\boldsymbol{r}) &= \widetilde{\phi}_{1}(r)\widetilde{\phi}_{2}(s) = \widetilde{\lambda}_{1}^{1D}(r)\widetilde{\lambda}_{2}^{1D}(s), & \nabla \widetilde{\psi}_{3}(\boldsymbol{r}) = \left(-\frac{1}{2}\widetilde{\lambda}_{2}^{1D}(s), \frac{1}{2}\widetilde{\lambda}_{1}^{1D}(r)\right), \\ \widetilde{\psi}_{4}(\boldsymbol{r}) &= \widetilde{\phi}_{2}(r)\widetilde{\phi}_{2}(s) = \widetilde{\lambda}_{2}^{1D}(r)\widetilde{\lambda}_{2}^{1D}(s), & \nabla \widetilde{\psi}_{4}(\boldsymbol{r}) = \left(\frac{1}{2}\widetilde{\lambda}_{2}^{1D}(s), \frac{1}{2}\widetilde{\lambda}_{2}^{1D}(r)\right), \end{split}$$

$$M_R = \frac{1}{9} \begin{pmatrix} 4 & 2 & 2 & 1 \\ 2 & 4 & 1 & 2 \\ 2 & 1 & 4 & 2 \\ 1 & 2 & 2 & 4 \end{pmatrix}.$$

$$Dr_R = \frac{1}{2} \begin{pmatrix} -1 & 1 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & 0 & -1 & 1 \\ 0 & 0 & -1 & 1 \end{pmatrix}$$

$$Ds_R = \frac{1}{2} \begin{pmatrix} -1 & 0 & 1 & 0 \\ 0 & -1 & 0 & 1 \\ -1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \end{pmatrix}.$$

$$S_R^{rr} = Dr_R^t M_R Dr_R = \frac{1}{6} \begin{pmatrix} 2 & -2 & 1 & -1 \\ -2 & 2 & -1 & 1 \\ 1 & -1 & 2 & -2 \\ -1 & 1 & -2 & 2 \end{pmatrix}$$
$$S_R^{ss} = Ds_R^t M_R Ds_R = \frac{1}{6} \begin{pmatrix} 2 & 1 & -2 & -1 \\ 1 & 2 & -1 & -2 \\ -2 & -1 & 2 & 1 \\ -1 & -2 & 1 & 2 \end{pmatrix}.$$

P_2 matrices

$$\begin{split} \widetilde{\psi}_{1}(\boldsymbol{r}) &= \widetilde{\phi}_{1}(r)\widetilde{\phi}_{1}(s), & \nabla \widetilde{\psi}_{1}(\boldsymbol{r}) = \left((-2\widetilde{\lambda}_{1}^{1D}(r) + \frac{1}{2})\widetilde{\phi}_{1}(s), \ (-2\widetilde{\lambda}_{1}^{1D}(s) + \frac{1}{2})\widetilde{\phi}_{1}(r) \right), \\ \widetilde{\psi}_{2}(\boldsymbol{r}) &= \widetilde{\phi}_{2}(r)\widetilde{\phi}_{1}(s), & \nabla \widetilde{\psi}_{2}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{1}^{1D}(r) - 2\widetilde{\lambda}_{2}^{1D}(r))\widetilde{\phi}_{1}(s), \ (-2\widetilde{\lambda}_{1}^{1D}(s) + \frac{1}{2})\widetilde{\phi}_{2}(r) \right), \\ \widetilde{\psi}_{3}(\boldsymbol{r}) &= \widetilde{\phi}_{3}(r)\widetilde{\phi}_{1}(s), & \nabla \widetilde{\psi}_{3}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{2}^{1D}(r) - \frac{1}{2})\widetilde{\phi}_{1}(s), \ (-2\widetilde{\lambda}_{1}^{1D}(s) + \frac{1}{2})\widetilde{\phi}_{3}(r) \right), \\ \widetilde{\psi}_{4}(\boldsymbol{r}) &= \widetilde{\phi}_{1}(r)\widetilde{\phi}_{2}(s), & \nabla \widetilde{\psi}_{4}(\boldsymbol{r}) = \left((-2\widetilde{\lambda}_{1}^{1D}(r) + \frac{1}{2})\widetilde{\phi}_{2}(s), \ (2\widetilde{\lambda}_{1}^{1D}(s) - 2\widetilde{\lambda}_{2}^{1D}(s))\widetilde{\phi}_{1}(r) \right), \\ \widetilde{\psi}_{5}(\boldsymbol{r}) &= \widetilde{\phi}_{2}(r)\widetilde{\phi}_{2}(s), & \nabla \widetilde{\psi}_{5}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{1}^{1D}(r) - 2\widetilde{\lambda}_{2}^{1D}(r))\widetilde{\phi}_{2}(s), \ (2\widetilde{\lambda}_{1}^{1D}(s) - 2\widetilde{\lambda}_{2}^{1D}(s))\widetilde{\phi}_{3}(r) \right), \\ \widetilde{\psi}_{6}(\boldsymbol{r}) &= \widetilde{\phi}_{3}(r)\widetilde{\phi}_{2}(s), & \nabla \widetilde{\psi}_{6}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{2}^{1D}(r) - \frac{1}{2})\widetilde{\phi}_{2}(s), \ (2\widetilde{\lambda}_{1}^{1D}(s) - 2\widetilde{\lambda}_{2}^{1D}(s))\widetilde{\phi}_{3}(r) \right), \\ \widetilde{\psi}_{7}(\boldsymbol{r}) &= \widetilde{\phi}_{1}(r)\widetilde{\phi}_{3}(s), & \nabla \widetilde{\psi}_{7}(\boldsymbol{r}) = \left((-2\widetilde{\lambda}_{1}^{1D}(r) + \frac{1}{2})\widetilde{\phi}_{3}(s), \ (2\widetilde{\lambda}_{2}^{1D}(s) - \frac{1}{2})\widetilde{\phi}_{1}(r) \right), \\ \widetilde{\psi}_{8}(\boldsymbol{r}) &= \widetilde{\phi}_{2}(r)\widetilde{\phi}_{3}(s), & \nabla \widetilde{\psi}_{8}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{1}^{1D}(r) - 2\widetilde{\lambda}_{2}^{1D}(r))\widetilde{\phi}_{3}(s), \ (2\widetilde{\lambda}_{2}^{1D}(s) - \frac{1}{2})\widetilde{\phi}_{2}(r) \right), \\ \widetilde{\psi}_{9}(\boldsymbol{r}) &= \widetilde{\phi}_{3}(r)\widetilde{\phi}_{3}(s), & \nabla \widetilde{\psi}_{9}(\boldsymbol{r}) = \left((2\widetilde{\lambda}_{1}^{1D}(r) - 2\widetilde{\lambda}_{2}^{1D}(r))\widetilde{\phi}_{3}(s), \ (2\widetilde{\lambda}_{2}^{1D}(s) - \frac{1}{2})\widetilde{\phi}_{3}(r) \right), \end{aligned}$$

P_2 matrices

$$M_R = \frac{1}{225} \begin{pmatrix} 16 & 8 & -4 & 8 & 4 & -2 & -4 & -2 & 1 \\ 8 & 64 & 8 & 4 & 32 & 4 & -2 & -16 & -2 \\ -4 & 8 & 16 & -2 & 4 & 8 & 1 & -2 & -4 \\ 8 & 4 & -2 & 64 & 32 & -16 & 8 & 4 & -2 \\ 4 & 32 & 4 & 32 & 256 & 32 & 4 & 32 & 4 \\ -2 & 4 & 8 & -16 & 32 & 64 & -2 & 4 & 8 \\ -4 & -2 & 1 & 8 & 4 & -2 & 16 & 8 & -4 \\ -2 & -16 & -2 & 4 & 32 & 4 & 8 & 64 & 8 \\ 1 & -2 & -4 & -2 & 4 & 8 & -4 & 8 & 16 \end{pmatrix}$$

P_2 matrices

$$Dr_R = \frac{1}{2} \begin{pmatrix} -3 & 4 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & -4 & 3 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -3 & 4 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & -4 & 3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & 4 & -1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & -4 & 3 \end{pmatrix}$$

$$Ds_R = \frac{1}{2} \begin{pmatrix} -3 & 0 & 0 & 4 & 0 & 0 & -1 & 0 & 0 \\ 0 & -3 & 0 & 0 & 4 & 0 & 0 & -1 & 0 \\ 0 & 0 & -3 & 0 & 0 & 4 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & -4 & 0 & 0 & 3 & 0 & 0 \\ 0 & 1 & 0 & 0 & -4 & 0 & 0 & 3 & 0 \\ 0 & 0 & 1 & 0 & 0 & -4 & 0 & 0 & 3 \end{pmatrix}$$

P₂ matrices

P₂ matrices

MatrixforPoisson_2D_Rec

This Matlab code generates the mass matrix M_R , the stiffness matrices Srr_R , Sss_R and the differentiation matrices Dr_R , Ds_R for continuous k-th order polynomial approximations on the reference rectangle R_R .

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FEMforPoisson 2D Rec

The following Matlab code solves the Poisson problem. In order to use this code, mesh information (c4n, n4e, n4db, ind4e), matrices $(M_R, S_P^{rr}, S_P^{ss})$, the source f, and the boundary condition u_D. Then the results of this code are the numerical solution u, the global stiffness matrix A, the global load vector b and the freenodes.

```
function [u, A, b, fns] = ...
    FEMforPoisson_2D_Rec(c4n,n4e,n4db,ind4e,M_R,Srr_R,Sss_R,f,u_D)
A = sparse(length(c4n).length(c4n)):
b = zeros(length(c4n),1);
u = b:
for j=1:length(n4e)
    xr = (c4n(1,n4e(2,j))-c4n(1,n4e(1,j)))/2;
    ys = (c4n(2,n4e(4,j))-c4n(2,n4e(1,j)))/2;
    J = xr*vs:
    rx=ys/J; sy=xr/J;
    A(ind4e(:,j),ind4e(:,j)) = A(ind4e(:,j),ind4e(:,j)) ...
        + J*(rx^2*Srr_R + sy^2*Sss_R);
    b(ind4e(:,j)) = b(ind4e(:,j)) + J*M_R*f(c4n(:,ind4e(:,j))');
end
fns = setdiff(1:length(c4n), n4db);
u(fns) = A(fns.fns) \setminus b(fns):
end
```

ComputeErrorFEM_2D_Rec

The following Matlab code computes the semi H1 error between the exact solution and the numerical solution.

```
function error = ...
    ComputeErrorFEM_2D_Rec(c4n,n4e,ind4e,M_R,Dr_R,Ds_R,u,ux,uy)
error = 0:
for j=1:size(ind4e,2)
   xr = (c4n(1.n4e(2.i))-c4n(1.n4e(1.i)))/2:
    vs = (c4n(2,n4e(4,i))-c4n(2,n4e(1,i)))/2;
   J = xr*vs:
   rx=vs/J; sy=xr/J;
    Dex=ux(c4n(:,ind4e(:,j))') - rx*Dr_R*u(ind4e(:,j));
    Dev=uv(c4n(:,ind4e(:,j))') - sy*Ds_R*u(ind4e(:,j));
    error=error+J*(Dex'*M_R*Dex+Dey'*M_R*Dey);
end
error=sqrt(error);
end
```

main_FEMforPoisson_2D_Rec

The following Matlab code solves the Poisson problem by using several matlab codes such as mesh_FEM2D_Rec_rectangle, MatrixforPoisson_2D_Rec, FEMforPoisson_2D_Rec and ComputeErrorFEM_2D_Rec.

```
iter = 10:
xl = 0; xr = 1; yl = 0; yr = 1; k = 2; M = 2.^(1:iter);
f=@(x) 2*pi^2*sin(pi*x(:,1)).*sin(pi*x(:,2));
u D=0(x) x(:.1)*0:
ux=@(x) pi*cos(pi*x(:.1)).*sin(pi*x(:.2)):
uy=@(x) pi*sin(pi*x(:,1)).*cos(pi*x(:,2));
error=zeros(1.iter):
h=1./M;
for j=1:iter
    [c4n, n4e, ind4e, n4db] = ...
        mesh_FEM2D_Rec_rectangle(xl, xr, yl, yr, M(j), M(j), k);
    [M_R, Srr_R, Sss_R, Dr_R, Ds_R] = MatrixforPoisson_2D_Rec(k);
    u = FEMforPoisson_2D_Rec(c4n,n4e,n4db,ind4e,M_R,Srr_R,Sss_R,f,u_D);
    error(j) = ComputeErrorFEM_2D_Rec(c4n,n4e,ind4e,M_R,Dr_R,Ds_R,u,ux,uy);
end
```

Example

Consider the domain $\Omega = [0, 1]^2$. The source term f is chosen such that

$$u = \sin(\pi x)\sin(\pi y)$$

is the analytical solution to the Poisson problem.

Convergence history

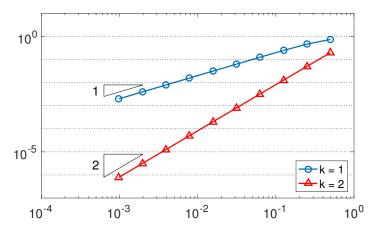


Figure: Convergence history for Example