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I. DIAMETER OF RANDOM CLUSTERS

A. Introduction

1. Potts Model [?]]

a. Generalization of Ising Model to q spin states

b. Applications

Conformal Field Theory

Percolation Theory

Knot Theory

Mathematical Biology

Complex Networks

SLE

- c. $H = -K \sum_{i,j} \delta_{\sigma_i, \sigma_j}$
- d. *Rich phase diagram*
- e. *Mapped onto Random Cluster model for $q \geq 0$*

$q = 1 \rightarrow$ Percolation

$q = 2 \rightarrow$ Ising

- f. *For $q \leq 4$, the model exhibits a second-order phase transition at the critical point*
- g. *For $q > 4$, the transition is first order [?]*

2. Chemical Distance

- a. *Until recently, only studied for Potts $q = 1$*
- b. *Scaling: $\langle l \rangle \propto r^{d_{min}}$*
- c. *We extend study to $q = 1, 2, 3, 4$ 2D Potts Model*
- d. *Use S-W algorithm to generate bonds, clusters*
- e. *Bonds correspond to spin correlations via Random Cluster Model*

3. Diameter

- a. *w , which we define as the longest of all the shortest paths between sites on a cluster*
- b. *Applications / connections*

maximum transport time

correlation lengths

scaling: $\langle w \rangle \propto r^{w_{min}}$

- c. *hypothesis: d_{min} equal to w_{min}*

d. Algorithm

Finding all-pairs shortest paths goes as $O(N^2)$

We suggest a novel, more efficient algorithm

e. Mean Field predictions

At or above critical dim, MFT should apply

underlying graph of connected sites that form the critical cluster should be well approximated by a complete graph of n vertices

complete graph: simple graph in which every pair of vertices is connected by an edge

Shown by Nachmias [?] that diam of complete graph at criticality scales as $w(n) \propto n^{1/3}$

f. We simulate $q = 2, D = 4$ Potts to assess MFT predictions

Since the mapping of the complete (linear) graph to the Potts random graph in 4D is $L^4 = n$, $w(L) \propto L^{4/3}$; thus, we may expect that w_{min} should equal $4/3$ for $q = 2$ in 4D.

B. Methods

1. Swendsen-Wang Algorithm

a. SW algorithm [?] used to generate statistics for models, create the bond-paths studied here

b. Based on work of Fortuin and Kasteleyn [?]

c. Procedure:

Introduce bonds with probability $p(\sigma_i, \sigma_j) = \delta_{\sigma_i, \sigma_j}(1 - e^{-K})$

Create clusters of bonded spins

Choose one of q possible spin states and assign to all sites in the cluster

- d. Reduces critical slowing relative to algorithms that flip individual spins [?], e.g. Metropolis algorithm [?]*
- e. Bonds introduced in SW algorithm correspond to correlations among spins*
- f. We study paths along bonds in these clusters*

2. Determining the Chem Distance and Diameter

- a. Review of Previous methods*

Stanley, Grassberger [?], Leath, Paul [?], etc.

Memory considerations, two seeds, etc.

- b. Leath growth [?]*

using a random number generator, one assigns all the bonds associated with the seed site the status “occupied” or “unoccupied” with probability p

If a bond is assigned “occupied” status, the site to which this bond connects is deemed a “growth site”, and is added to cluster.

All the sites thus added to the cluster in this round form a “chemical shell” of distance l from the seed site.

This process is then continued for subsequent generations of growth trials, each associated with a larger chemical shell; the growth process stops naturally when one of the growth rounds generates no new growth sites.

(Note: sites not added to the cluster in a particular round get another chance to be added to the cluster in subsequent rounds; but, once added, are no longer considered as possible growth sites.)

- c. Leath growth most appropriate for what we’re measuring*

Can't use two-seed method; we must find all possible paths

3. Procedure for $q > 1$

- a. Generate a new cluster configuration using the Swendsen-Wang algorithm (see above) with periodic boundary conditions. The identification of connected clusters in this steps allows us to determine the largest cluster in the system.
- b. Choose a random site s on this cluster as the seed site.
- c. Beginning with the seed site s , determine all sites in the largest cluster by “growing” along satisfied cluster bonds (this process does not change the bonds that were determined in step 1).
- d. The chemical shell reached in the final step of this growth process, $shell_{final}$, is considered to be the randomly-chosen chemical distance on the largest critical cluster, and is added to our statistics for the chemical distance.
- e. All the i sites at the end of this growth process whose nearest neighbors are all occupied are deemed to be perimeter sites, p_i . This set includes all of the external perimeter sites of the cluster.
- f. A similar Leath growth process is preformed using each of the perimeter sites as seeds, and $shell_{final_i}$ from each of these growth processes is stored.
- g. The diameter for the largest cluster is then $\max\{shell_{final_i}\}$
- h. This method for finding the diameter is an improvement over the naive N^2 algorithm for solving the all-pairs maximum shortest path problem on the paths formed along cluster bonds. It is expected to scale as $O(pN)$, where p is the number of perimeter sites on the largest critical cluster.

4. Procedure for $q > 1$

- a. For $q = 1$, it is possible to grow a cluster from a seed site.
- b. Diameter must have its endpoints on perimeter sites
- c. Any “pins”, or singly-connected paths on the external perimeter of the cluster, contain sites that can be eliminated as possible diameter endpoints

d. Straightforward to show that the existence of such a “pin” also allows us to eliminate as candidate diameter endpoints that lie within the “body” of the cluster as well

e. ‘Proof’ of / argument for the algorithm:

P : the set of all sites on the pin P

let p_{tip} : the site that is the outermost tip of a given pin (i.e., the site with only one nearest neighbor) and p_{attach} the site that attaches this pin to the body of the cluster (i.e., a site with more than 2 nearest neighbors)

Imagine that we were to include as a candidate site in S some site from P that was not p_{tip} , resulting in a candidate diameter D' ; it would be immediately clear that rejecting this site in favor of p_{tip} would result in a new candidate diameter $D'' > D'$. We can therefore exclude all sites in P that are closer than p_{tip} to S .

(?) Similar considerations (PROVE THIS?) allow us to additionally exclude from S all sites in N that have a chemical distance from p_{attach} less than or equal to the chemical distance between p_{tip} and p_{attach} (i.e., the length of the pin).

Initiate, for every site s in S , a “Leath growth” search that examines the chemical distance between s and every other site on the cluster, terminating when all cluster sites have been examined.

The maximum chemical distance found across all such searches is then D .

We thus need only consider a relatively small proportion [quantify this proportion, on average] of cluster sites as possible diameter endpoints, greatly reducing the number of “Leath scans” required in order to determine the diameter exactly

Note that this method does not work for periodic boundary conditions, however; we must therefore grow clusters from a seed site, retaining only those clusters that do not grow to touch the boundaries of the lattice.

f. Procedure

Choose a growth seed site in the center of the lattice

Perform a Leath growth from this site until the cluster dies, or reaches the boundaries of the maximum lattice size of L_{max} . If any cluster site borders L_{max} , begin again at step 1.

Identify all the perimeter sites in the cluster by choosing all sites in the final growth step that are perimeter sites (i.e., those that have less than the maximum number of allowed nearest neighbors). In this geometry, all the sites in the final chemical shell will be external perimeter sites.

Identify all the “pins” among these perimeter sites by performing a Leath growth from each pin site until one finds a site that is not singly-connected. All of the sites in the “neck” of the pin are eliminated from consideration as diameter endpoints.

Beginning from the point of attachment of the pin to the body of the cluster, continue the Leath scan until one has achieved a chemical shell equal to the distance (along sites) between the point of attachment and the end of the pin. All of sites thus scanned are also eliminated from consideration as diameter endpoints.

Perform Leath growths from all of the remaining perimeter sites p_i , collecting the maximum chemical shells reached in each instance; the largest of these chemical shells is then the diameter of the cluster.

g. Comparison with 'regular' Leath growth method

We compared this method to the method described for $q > 1$, and found that the fraction of perimeter sites eliminated as candidates for diameter endpoints was approximately $X\%$ in our runs with $L_{max} = XX$.

h. Label update procedure In order to determine which sites have been visited in the above-described Leath growth, we must assign each site a label. Because resetting all N

labels is costly, we instead update the value of the label at each time step. In order to increase the efficiency of the algorithm

5. *Simulation Details*

a. Overview

We used the Swendsen-Wang algorithm to simulate Potts Models 2D at criticality for values of L between 8 and L_{max} for our measurements of l , and 4 and L_{max} for our measurements of w . For $q = 2$ in 4D, L ranged between 4 and L_{max} . All simulations began in a random configuration.

b. Values of p_{add} used

For $q = 1$ in 2D, p_{add} is known exactly (REF). For $q = 2, 3, 4$ in 2D, $p_{add} = X$ (REF), X (REF), and X (REF), respectively. For $q = 2$ in 4D, $p_{add} = X$ (REF).

c. Thermalization

For $q > 1$, the simulations require some time to achieve an equilibrium state, and should therefore be thermalized. Accordingly, each simulation for system size L was run for at least $X\tau_{int,m}$ before measurements were taken, where $\tau_{int,m}$ was the estimated integrated autocorrelation time for the mass of the largest cluster for that value of L .

A table of integrated autocorrelation times for the largest system sizes measured is provided (Table)

d. Run times

In 2D, our simulations were run for a length of $X\tau_{int,m}$; for measurements of w , and for $X\tau_{int,m}$ for measurements of l .

For our 4D, $q = 2$ measurements, simulations were run for a length of $X\tau_{int,m}$ for our measurements of l .

Some of our simulations consisted of a single, long run; others were the result of combining data from several runs begun from different initial random number generator seeds.

e. Random Number Generator

Random numbers for the simulations were generated using the Mersenne Twister method (REF: Matsumoto + Nishimura 1998), with parameters chosen to provide a period of at least X (determine this)

f. Tests of the algorithm

As a check on our simulation methods, we also measured the mass of the largest cluster for each lattice size L in order to determine the fractal dimension. The agreement between our values and the latest from the literature was good

g. CPU Time

The CPU time for simulations measuring the diameter w was approximately $XL^2\mu s$ / iteration; for l it was approximately $XL^2\mu s$ / iteration, when run on the

h. Blocking Method

We used the 'blocking' method [?] to extract the proper standard deviation for chemical distance and diameter from our measurements.

This method works by clustering the measurements of the quantity O into blocks of size s ; the average of O is then found for each block independently; the standard deviation in O is then taken to be the standard deviation in these block averages

$$\sigma = \sqrt{\frac{\langle m^2 \rangle - \langle m \rangle^2}{n-1}}, \text{ where } n \text{ is the number of blocks}$$

i. Fitting Methods

For $q = 1, 2, 3$, we attempted fits using the Ansätze $y = aL^b$ and $y = aL^b + L/c$, including in the fit data points down to L value of L_{min} , where L_{min} was the smallest value of L that still yielded a reasonable goodness-of-fit value, Q

The fitting form $y = aL^b$ provided the best fits for all values of q .

For $q = 4$, we also attempted a fit of the form $y = A + B \log L$; the fit was not as good as the Ansatz $y = aL^b$.

j. Comparison, chem dist and diameter

k. Comparison of results with those of Deng et. al

Our numerical results appear to match the conjecture of Deng et al. [?] within error for $q = 1$ and $q = 2$; for $q = 3$, we find [wait until results of new blocking analysis are in]. For $q = 4$, we were unable to find a fit of high quality; but our results seem to support Deng et. al's conjecture

l. Discussion of systematic errors

C. Data Analysis

D. Results and Discussion