



Article

# IBM Quantum Platforms: A Quantum Battery Perspective

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**Abstract:** We characterize for the first time the performances of IBM quantum chips as quantum batteries, specifically addressing the single-qubit Armonk processor. By exploiting the Pulse access enabled to some of the IBM Quantum processors via the Qiskit package, we investigate the advantages and limitations of different profiles for classical drives used to charge these miniaturized batteries, establishing the optimal compromise between charging time and stored energy. Moreover, we consider the role played by various possible initial conditions on the functioning of the quantum batteries. As the main result of our analysis, we observe that unavoidable errors occurring in the initialization phase of the qubit, which can be detrimental for quantum computing applications, only marginally affect energy transfer and storage. This can lead counter-intuitively to improvements of the performances. This is a strong indication of the fact that IBM quantum devices are already in the proper range of parameters to be considered as good and stable quantum batteries comparable to state-of-the-art devices recently discussed in the literature.



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## 1. Introduction

Quantum batteries (QBs) recently emerged as a fast-growing and very active field of research in the domain of quantum technologies [1]. The idea behind these miniaturized devices for energy storage represents a radical change of perspective in the framework of energy manipulation with respect to the electrochemical principles developed in the eighteenth and nineteenth centuries, which are still at the core of modern technology [2,3]. Indeed, here, genuinely non-classical features such as quantum superposition, entanglement, and many-body collective behaviors can be exploited in order to outperform their classical counterparts in terms of stored energy, charging time, average charging power and extractable work [4]. These quantities are the main figures of merits to look at in order to characterize and optimize the performances of these kind of devices. From the practical point of view, quantum batteries could be used in the near future to supply in a fast and controlled way the energy needed for the functioning of more complex quantum devices and sensors, opening the way to new and fascinating technological developments.

Starting from seminal ideas developed in Ref. [5], the theoretical investigations in this domain have been at first characterized by the influence of theorems mediated by quantum information [6–9]. During the years, the studies progressively moved toward more experimentally oriented proposals. They addressed set-ups conveniently designed in such a way to be easily implemented on existing quantum computing platforms such as arrays of artificial atoms [10–16] and systems for cavity and circuit quantum electrodynamics [17,18]. This latter class of set-ups has been also proposed as support for the realization of quantum supercapacitors where the energy is stored by exploiting radiation-induced non-classical arrangement of the electron charge [19].

Very remarkably, the first experimental evidence of a QB has been reported less than one year ago in a system where fluorescent organic molecules play the role of two-level systems embedded in a microcavity [20]. Even more recently, QBs realized with transmon qubits [21] and quantum dots [22] have been reported, further testifying the great ferment around this topic.

The majority of the considered approaches for QBs are based on two-level systems (qubits) promoted from the ground to the excited state by means of the action of another system playing the role of a charger [23,24]. This latter element can be genuinely quantum, such as photons trapped into a cavity, or more simply a classical time-dependent drive directly applied to the qubit [14,25,26]. As stated above, the efficiency of this kind of charging processes is characterized in terms of figures of merits such as the energy stored into the QB, the charging time required to reach its maximum and the average charging power, which is the energy stored in a given time [6].

IBM quantum devices [27] offer the unique opportunity to simulate quantum systems under controlled conditions, leading to an exponentially increasing number of scientific papers covering various branches of research. The following are some examples that go from quantum chemistry and material sciences [28,29] to the analysis of molecular magnetic clusters and spin–spin dynamical correlation functions [30], up to quantum field theories [31,32], high energy physics [33,34] and dark matter [35]. Lots of different publications are defined within real use cases belonging to different industries, including finance, material science and optimization [36–38].

An important step forward in the direction of addressing quantum dynamics has been represented by the implementation of the Pulse tool included in the Qiskit package [39], which opened the way to the possibility to reach an unprecedented level of control over the form and the relevant parameters of a classical drive applied to quantum systems.

The aim of this paper is to realize the first simulation of a classically driven QB applying different controlled pulses on an IBM quantum device. We will focus on the simplest possible machine, the Armonk quantum processor, made by a single transmon qubit. After a proper calibration of the data, we will characterize the charging profile of the QB as a function of the time integral of the envelop function of the pulse. We will determine the constraints on the form of the pulse in order to obtain an universal charging curve, and we will establish the minimum reachable charging time. Without implementing any ad hoc optimization procedure, we observe performances in terms of charging time and stored energy compatible with state-of-the-art experiments in the domain [21]. Moreover, we observe that initialization errors ubiquitously present in the Noisy Intermediate-Scale Quantum devices [40], which can be detrimental in a quantum computation perspective, can lead to an improvement of the performances of these set-ups as a QB.

The present paper is organized as follows. In Section 2, we investigate the model of a qubit coupled with a time-dependent drive as a proper description of the Armonk IBM quantum processor subject to the action of the Pulse tool. The calibration of the system and the analysis of the data provided by the IBM interface are reported in Section 3. The discussions concerning the universality of the charging curve, the possible technical constraints on the suitable forms of the pulses, the achievable minimal charging time and the role played by different initial conditions are reported in Section 4. Section 5 is devoted to the conclusions. Finally, a table at the end of the paper summarizes the relevant quantities involved in our model for the QB.

## 2. Model

We consider a superconducting qubit in the transmon regime [41]. In the working conditions investigated in this paper, it can be seen as an effective two-level system driven in time by a classical pulse. Its Hamiltonian reads ( $\hbar = 1$ )

$$H = H_{QB} + H_C \quad (1)$$

$$= \frac{\Delta}{2}(1 - \sigma_z) + gf(t) \cos(\omega t)\sigma_x \quad (2)$$

where the first term ( $H_{QB}$ ) represents the free Hamiltonian of a QB with a level spacing  $\Delta$  between the ground state  $|0\rangle$  and the excited state  $|1\rangle$ , while the second term ( $H_C$ ) describes the classical charging of the QB itself due to the application of a time-dependent drive compatible with the Qiskit Pulse tool [39]. In the above equation, with  $\sigma_{x,z}$ , we indicate the Pauli matrices along the  $x$  and  $z$  direction, respectively. Notice that one has

$$\sigma_z|0\rangle = |0\rangle \quad (3)$$

$$\sigma_z|1\rangle = -|1\rangle. \quad (4)$$

With  $f(t)$ , we denote a time-dependent adimensional envelop function with maximum amplitude equal to one, whose form will be specified in the following. It is further modulated by a cosine function with frequency  $\omega$ . Moreover,  $g$  represents the coupling between the QB and the classical drive. The aim of this section is to study the dynamics of this system starting from the generic initial wave-function at time  $t = 0$

$$|\Psi(0)\rangle = \alpha|0\rangle + \beta|1\rangle \quad (5)$$

with  $\alpha$  and  $\beta$  complex parameters satisfying  $|\alpha|^2 + |\beta|^2 = 1$ .

Typically, the IBM quantum machines are built in such a way that  $\Delta \gg g$ . In particular, for the Armonk quantum processor used in this work, one has  $\Delta \approx 31.238$  GHz and  $g \approx 0.105$  GHz. Under this condition, for arbitrary values of driving frequency  $\omega$ , it is not possible to achieve a transition from  $|0\rangle$  to  $|1\rangle$ , namely a charging of the QB [14].

This issue can be overcome by properly tuning the systems in the perfectly resonant case  $\Delta = \omega$ . The reason for the peculiarity of this case can be better appreciated moving to a rotating frame, namely by applying the time-dependent rotation

$$S(t) = e^{-i\frac{\Delta}{2}t\sigma_z} \quad (6)$$

to the Hamiltonian in Equation (2) in such a way to obtain

$$H' = SHS^\dagger - iS\frac{dS^\dagger}{dt}, \quad (7)$$

where the time dependence of the operators has been omitted for notational convenience. By further considering the rotating wave approximation [42], which is very well justified under the conditions of resonance and small coupling we are considering, one can write

$$H' \approx \frac{g}{2}f(t)\sigma_x + \frac{\Delta}{2}, \quad (8)$$

where the constant term plays no role in the dynamics and will be neglected in the following.

The above expressions for the rotated Hamiltonian together with the fact that the considered rotation does not affect the initial state of the system, namely

$$|\Psi'(0)\rangle = S(0)|\Psi(0)\rangle = |\Psi(0)\rangle, \quad (9)$$

allows to analytically solve the dynamics of the considered qubit. Indeed, one can introduce the eigenstates of the  $\sigma_x$  operator

$$|\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle) \quad (10)$$

in such a way that the initial wave-function can be written as

$$|\Psi'(0)\rangle = \left(\frac{\alpha + \beta}{\sqrt{2}}\right)|+\rangle + \left(\frac{\alpha - \beta}{\sqrt{2}}\right)|-\rangle \quad (11)$$

$$= C_+(0)|+\rangle + C_-(0)|-\rangle. \quad (12)$$

On this basis, the time evolution of the coefficients  $C_{\pm}$  satisfies

$$\frac{dC_{\pm}}{dt} = \mp i\frac{g}{2}f(t)C_{\pm}(t) \quad (13)$$

and consequently

$$|\Psi'(t)\rangle = C_+(t)|+\rangle + C_-(t)|-\rangle, \quad (14)$$

with

$$C_{\pm}(t) = C_{\pm}(0)e^{\mp i\frac{g}{2}\int_0^t f(\tau)d\tau}. \quad (15)$$

In the following, we will focus on a situation, relevant for the considered quantum simulations, where the state of the system is measured at a given time  $t = t_m$  such that

$$C_{\pm}(t_m) = C_{\pm}(0)e^{\mp i\frac{\theta(t_m)}{2}} \quad (16)$$

with

$$\theta(t_m) = g \int_0^{t_m} f(\tau)d\tau. \quad (17)$$

Going back to the original basis, one obtains directly

$$|\Psi(t_m)\rangle = S^{\dagger}(t_m)|\Psi'(t_m)\rangle \quad (18)$$

$$= e^{i\varphi(t_m)} \left[ \alpha \cos \frac{\theta(t_m)}{2} - i\beta \sin \frac{\theta(t_m)}{2} \right] |0\rangle$$

$$+ e^{-i\varphi(t_m)} \left[ \beta \cos \frac{\theta(t_m)}{2} - i\alpha \sin \frac{\theta(t_m)}{2} \right] |1\rangle \quad (19)$$

with

$$\varphi(t_m) = \frac{\Delta}{2}t_m. \quad (20)$$

According to this analysis, the energy stored into the QB at the measurement time ( $t_m$ ) is given by [23]

$$E(t_m) = \frac{\Delta}{2}\langle\Psi(t_m)|(1 - \sigma_z)|\Psi(t_m)\rangle \quad (21)$$

$$= \Delta\langle\Psi(t_m)|1\rangle\langle1|\Psi(t_m)\rangle \quad (22)$$

$$= \Delta|\langle\Psi(t_m)|1\rangle|^2 \quad (23)$$

$$= \Delta P_1(\theta(t_m)) \quad (24)$$

where the last line indicates the probability to find the QB in the excited state ( $|1\rangle$ ) when the measurement is carried out. In the following, we will omit the dependence on  $t_m$  for notation convenience.

Taking into account the reparametrization

$$\alpha = \sqrt{a} \quad (25)$$

$$\beta = \sqrt{1-a}e^{-i\phi}, \quad (26)$$

with  $a$  and  $\phi$  real numbers, one can finally write

$$E(a, \phi, \theta) = \Delta \left[ a \sin^2 \frac{\theta}{2} + 2\sqrt{a}\sqrt{1-a} \sin \phi \sin \frac{\theta}{2} \cos \frac{\theta}{2} + (1-a) \cos^2 \frac{\theta}{2} \right], \quad (27)$$

where we have made explicit the dependence of the stored energy on the relevant parameters characterizing the initial wave-function of the qubit and the phase associated to the envelop function of the applied pulse. In the following, the results of the simulations carried out as a function of  $\theta$  will be fitted by means of Equation (27) in order to extract the values of the parameters  $a$  and  $\phi$ . This will allow us to reconstruct the initial state of the QB and characterize its performances in terms of stored energy. Notice that a similar analysis has been carried out very recently in an experimental work devoted to the characterization of a semiconducting quantum dot as a QB [22].

### 3. Calibration

Before entering into the details of the QB charging, we need to discuss the nature of the data returned by the IBM platform and the way to analyze them. The measurement of the state of a given qubit after the application of an external drive is done through a readout in the so-called dispersive regime [43]. Here, a harmonic oscillator (resonator) is weakly coupled to the QB. The system is then described by the Jaynes–Cummings Hamiltonian [42]

$$H_{RO} = H_{QB} + \omega_r b^\dagger b + \lambda (b^\dagger \sigma_- + b \sigma_+) \quad (28)$$

where  $\lambda$  indicates the strength of the matter–radiation dipolar coupling and with  $b^\dagger/b$  and  $\sigma_\pm$  ladder operators for the harmonic oscillator and the effective spin associated to the QB, respectively. When the two parts of the system are kept far from resonance, namely under the condition  $\lambda \ll |\Delta - \omega_r|$ , it is possible to perform a Schrieffer–Wolff transformation and to expand Equation (28) up to the second order in the coupling, obtaining the effective Hamiltonian [42]

$$H_{eff} = H_{QB} + \omega_r b^\dagger b - \frac{2\lambda^2}{|\Delta - \omega_r|} \sigma_z \left( b^\dagger b + \frac{1}{2} \right). \quad (29)$$

This leads to a shift in the frequency of the oscillator, which depends on the state of the QB.

Under these conditions, a monochromatic microwave with frequency  $\Omega$  applied to the resonator is modified in such a way that

$$\cos \Omega t \rightarrow \mathcal{A} \cos(\Omega t + \chi). \quad (30)$$

Namely, it acquires a different amplitude  $\mathcal{A}$  and phase  $\chi$ . These quantities depend on the properties of the resonator and consequently are unambiguously related to the state of the considered qubit. Getting rid of the time dependence of the signal and taking into account the complex representation of the transmitted wave, one can write

$$\mathcal{A} e^{i\chi} = I + iQ, \quad (31)$$

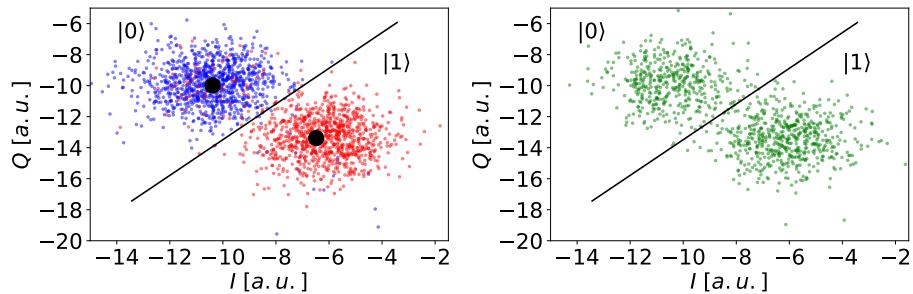
with  $I$  and  $Q$  real numbers.

Every measurement of the qubit state is then reported as a point in the  $(I, Q)$  plane [44]. In order to accumulate proper statistics, the machine performs multiple runs (1024 in default settings). They are typically very scattered, requiring a calibration to extract meaningful information from them. Focusing now on the specific case of the Armonk single-qubit device, one has that at the beginning of every run, it is initialized to the ground state, so as a first thing, we have characterized the ground state  $|0\rangle$  of the system measuring it just after the initialization of the machine. Then, the excited state  $|1\rangle$  has been obtained initializing

the Armonk qubit in  $|0\rangle$  and applying a built-in pulse with  $\theta = \pi$  (see Equation (17)) of the form

$$f(t) = \sqrt{\frac{\pi}{2}} \frac{1}{g\sigma} e^{-\frac{(t-t_m/2)^2}{2\sigma^2}} \quad (32)$$

with standard deviation  $\sigma = t_m/8$ , being  $t_m = 600$  ns the measuring time. A typical distribution of the points in the  $(I, Q)$  plane associated to the ground and excited state is reported in the left panel of Figure 1.



**Figure 1.** (Color online). Left panel. Example of data distribution associated to the measurements of the state  $|0\rangle$  (blue dots) and  $|1\rangle$  (red dots) in the  $(I, Q)$  plane (in arbitrary units) of the Armonk single-qubit device. Big black dots indicate the centers of the two distributions, while the black line separates them. Notice that despite the distortion due to the aspect ratio of the plot, this line is perpendicular to the segment connecting the two centers. The efficiency of the considered separation is roughly 97.4% for the ground state and 92.7% for the excited state. Right panel. Example of a distribution associated to the measurement of a state with  $\theta = \pi/2$  (green dots). The distribution presents now two separated lobes, one in each sector. The energy stored in the system is related to the fraction of the dots composing the lobe in the  $|1\rangle$  sector with respect to the total runs. In the present case, this is given by  $P_1 \approx 0.567$ . For each state shown in the plot, we have considered 1024 runs.

In order to discriminate between the two states of the system, we have written a Python routine able to determine the center of the relative distributions (big black dots) and identify the line perpendicular to the segment connecting them and passing through its middle point. This approach implicitly assumes that both distributions have the same spread. This line divides the  $(I, Q)$  plane into two parts that we can identify as the  $|0\rangle$  and the  $|1\rangle$  sector, respectively. It is worth noting that due to the intrinsic errors affecting the machine, a small fraction (few percent) of the blue dots corresponding to ground-state measurement fall into the part of the plane associated to the excited state and vice versa. The evaluation of these errors, as well as a reasonable hypothesis of their origin, will be relevant points of the analysis of the results.

According to the picture discussed above, the energy stored in the Armonk qubit seen as a QB, proportional to the probability for a state to be measured in  $|1\rangle$  after the application of a generical pulse with  $0 \leq \theta \leq \pi$  (see Equation (24)), can be evaluated directly as the ratio of the number of points falling in the  $|1\rangle$  sector with respect to total runs performed by the machine (see the right panel of Figure 1 for an example).

## 4. Results

### 4.1. Universal Charging Behavior and Technical Constraints on the Pulses

In this section, we consider the charging curve of the Armonk QB for different profiles of the time-dependent envelop function  $f(t)$ . On a very general ground, we observe that Equation (27) does not explicitly depend on the functional form of this envelop but only on the parameter  $\theta$ , which is proportional to its integral up to the measurement time (see Equation (17)). From the mathematical point of view, there are infinite possible functions  $f(t)$  showing the same value of  $\theta$  for a fixed measurement time  $t_m$ . Therefore, we expect

this universal behavior to emerge in the measurements. Moreover, we attribute deviations with respect to the ideal curve

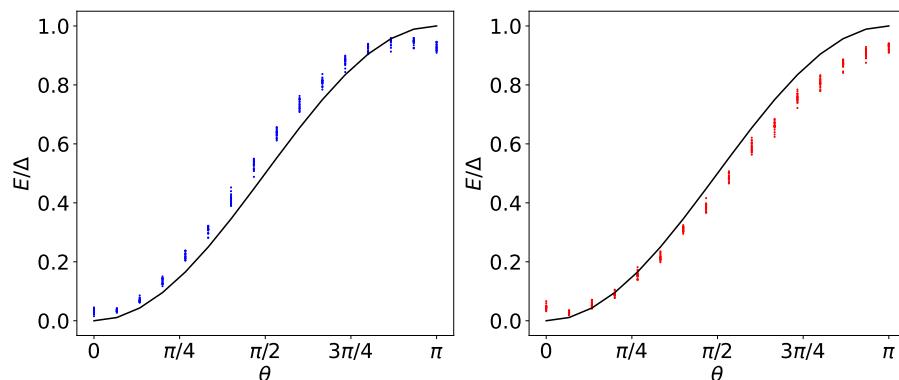
$$E(1, \phi, \theta) = \Delta \sin^2 \frac{\theta}{2}, \quad (33)$$

corresponding to the perfect initialization in the ground state, to errors occurring in this phase.

The results concerning Gaussian profiles of the form

$$f(t) = \mathcal{N} e^{-\frac{(t-t_m/2)^2}{2\sigma^2}} \quad (34)$$

are reported in Figure 2. Due to the overall constraint given by Equation (17), the amplitude  $\mathcal{N}$  and the standard deviation  $\sigma$  cannot be tuned independently. Therefore, we have compared two possible complementary approaches. In the left panel, generalizing what was calculated in Equation (32), the standard deviation  $\sigma$  has been kept fixed to a convenient fraction of the measurement time  $t_m$ , while the amplitude  $\mathcal{N} = \theta / (\sqrt{2\pi} g \sigma)$  (always satisfying  $\mathcal{N} < 1$ ) has been varied in order to return the proper value of  $\theta$ . In the right panel, we proceeded in the opposite way, keeping fixed the amplitude to the maximum possible value  $\mathcal{N} = 1$  and varying the standard deviation according to  $\sigma = \theta / (\sqrt{2\pi} g \mathcal{N})$ . It is evident that in both cases, the data depart from the ideal behavior of Equation (33). Moreover, despite being equivalent from the mathematical point of view, these two approaches lead to different results. This is due to the fact that the second approach is affected by a technical problem. For the considered values of  $\theta$ , the support of  $f(t)$  is reasonably different from zero only in a narrow time window. Due to the fact that the Pulse tool discretizes the signals with a minimum time step of  $\delta = 0.222$  ns, this peaked envelop function is strongly affected by this and is characterized by an effective area smaller with respect to what was expected. Even if this error can be mitigated by considering lower amplitudes  $\mathcal{N}$ , the first approach was more suitable for our purposes with different choices of standard deviation satisfying  $\sigma < t_m / 5$  (not shown), leading to almost superposing distributions of the points. This latter condition is justified by the need to minimize the error made by cutting the envelop function at a finite measurement time  $t_m$ . While this error is typically under control for fast decaying envelop functions such as the Gaussian considered here, broader profiles such as the Lorentzian one require to be properly adjusted in order to compensate for this missing-tail effect and recover the predicted universal behavior.



**Figure 2.** (Color online). Energy stored in the QB (in units of  $\Delta$ ) as a function of the area  $\theta$ . The black curves correspond to the ideal charging profile with  $a = 1$  (see Equation (33)). The left panel shows the charging achieved using a Gaussian pulse with fixed standard deviation  $\sigma = t_m/8$ , being  $t_m = 600$  ns the measurement time, and amplitude  $\mathcal{N} = \theta / (\sqrt{2\pi} g \sigma)$ . The right panel shows the charging due to a Gaussian pulse with fixed amplitude equal to  $\mathcal{N} = 1$  and adjustable standard deviation  $\sigma = \theta / (\sqrt{2\pi} g \mathcal{N})$ . The vertical spread of the points is due to the fact that we have considered 20 measurements for each value of  $\theta$ .

The above discussion indicates that a not too narrow and fast-decreasing envelop function represents the better choice for a pulse leading to a suitable charging of the considered QB. Once these requirements are properly fulfilled, it is also possible to reduce the charging time with respect to the value reported in Figure 2. This point, together with a fit of the data in order to determine the actual initial state of the system, will be investigated in the following.

#### 4.2. Best Fit of the Data and Characterization of the QB Performances

In Figure 3, we consider curves obtained averaging over  $N = 20$  measurements carried out with the same form of the envelop function  $f(t)$ . The first one (left panel) is achieved with a measurement time  $t_m^{(1)} = 600$  ns (default value for the Armonk machine), while the second with  $t_m^{(2)} = 135$  ns, which is a value still consistent with the pulse constraints discussed above. Through best fit of the reported data according to Equation (27), one obtains

$$a^{(1)} = 0.981 \pm 0.003 \quad (35)$$

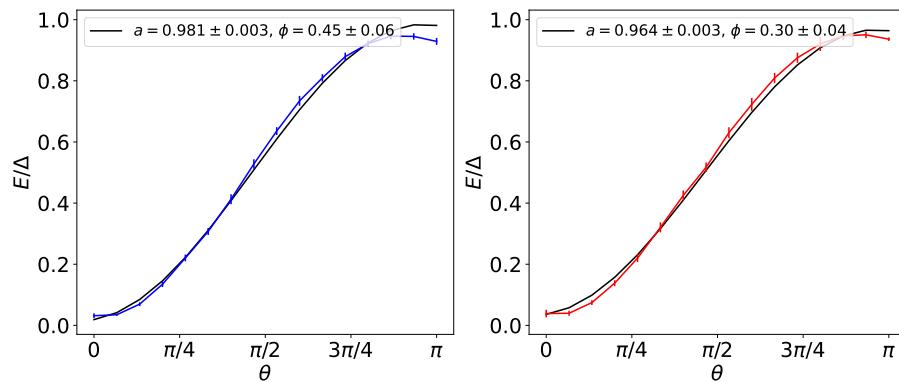
$$\phi^{(1)} = 0.45 \pm 0.06 \quad (36)$$

in the first case and

$$a^{(2)} = 0.964 \pm 0.003 \quad (37)$$

$$\phi^{(2)} = 0.30 \pm 0.06 \quad (38)$$

in the second. Notice that measurements carried out at an intermediate time  $t_m^{(2)} < t_m < t_m^{(1)}$  (not shown) are compatible to what was observed in the two extreme cases. Conversely, for  $t_m < t_m^{(2)}$ , the effects associated to the previously described discretization of the pulse lead to a less efficient charging. Moreover, the results are not strongly affected by the numbers of measurements as long as  $N \geq 10$ . Here, we have considered  $N = 20$  as a good compromise between the quality of the reported results and the time required to realize every point in the plots.



**Figure 3.** (Color online). Best fit of the energy stored into the QB (in units of  $\Delta$ ) as a function of  $\theta$  (black curves). Data correspond to Gaussian pulses with  $\sigma = t_m/8$ , being  $t_m = 600$  ns (blue curve in the left panel) and  $t_m = 135$  ns (red curve in the right panels). Every point is obtained averaging over 20 different measurements (see Figure 2), and the bars take into account the standard error associated to this average.

These data indicate that the Armonk qubit behaves as a QB with non-zero, but limited, initial energy which can be almost completely charged with maximum stored energy  $E_{max}^{(1/2)}$  exceeding 95% in both cases. This corresponds to a maximal stored energy of  $E_{max}^{(1/2)} \approx 20 \mu\text{eV}$  ( $\approx 200 \text{ mK}$  in temperature units). Moreover, the second case is characterized by a greater average charging power, namely the energy stored with respect to the charging

time (which is necessary lower with respect to the measured time  $t_m$ ) [23]. In addition, the observed charging time and stored energy is comparable with the recently reported state-of-the-art measurements on the first charging step of a three-level transmon QB [21]. Considering even lower values of the measurement time, and consequently of the standard deviation of the Gaussian pulse, the data show a progressively stronger deviation from the predicted universal charging curve as a consequence of the emergence of the previously described technical limitations.

It is worth noticing that the reported measurement (and consequently charging) times are substantially shorter with respect to both the decay ( $T_1 = 165 \mu\text{s}$ ) and dephasing ( $T_2 = 214 \mu\text{s}$ ) time reported for the machine. According to this, once a pulse with a given value of  $\theta$  is applied to the device, the corresponding energy stays confined into the QB for a time of the order of  $T_1$  as long as no other pulse is applied. This represents a very strong indication of the fact that IBM quantum machines are already, without any ad hoc optimizations, in the proper range of parameters to be considered as quite efficient and stable QBs.

To conclude this section, we observe that the fluctuating phase acquired by the state at the level of the initialization, which is a major problem in a quantum computation perspective, can lead counter-intuitively to an improvement of the device as a QB. Indeed, the measured value of phase parameter  $\phi$  is positive in the majority of the investigated cases. According to this, the value of  $\theta$  at which the maximum charge is reached can occur for a value  $\theta_{max} < \pi$ , requiring a lower effective area of the envelop function to be achieved with respect to the case of a thermal state with the same initial probability of being in the ground state, which is equivalent to the results of Equation (27) at  $\phi = 0$ . This leads to an improvement of the performance of the QB. The effects of an imperfect initialization are even more evident by considering more general initial conditions, which are discussed in the following.

#### 4.3. More General Initial Conditions

The initialization error discussed above can have an impact on the functioning of the QB also when more general quantum superposition states are considered as initial conditions. These states can be obtained starting from the ground state and applying a proper built-in unitary operator. As an example, in the following, we will consider the two operators

$$U = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \quad (39)$$

and

$$V = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & +i \\ -i & -1 \end{pmatrix}. \quad (40)$$

For an ideal QB, the energy stored should be

$$E\left(\frac{1}{2}, 0, \theta\right) = \frac{\Delta}{2} \quad (41)$$

in the first case, showing no charging/discharging dynamics due to the fact that the action of the drive only induces the rotation of this state on the equator of the Bloch sphere, and

$$E\left(\frac{1}{2}, \frac{\pi}{2}, \theta\right) = \frac{\Delta}{2}(1 + \sin \theta) \quad (42)$$

in the second case, requiring half of the energy ( $\Delta/2$ ) and half of the envelop area ( $\theta/2$ ) to realize a complete charging.

However, both the initialization and the application of a unitary operator are intrinsically affected by errors, as can be seen from the corresponding measurements reported in

Figure 4. Again, we observe that these errors can affect the functioning of the Armonk qubit as a QB. Indeed, even if limited, the first state (left panel) shows an unexpected charging dynamics, while for the second state, we observe a charging curve quite close to the ideal one, even if the maximum charging is not reached. Both these behaviors can be explained with the fact that the actual initial states are different with respect to what expected. In particular, we have

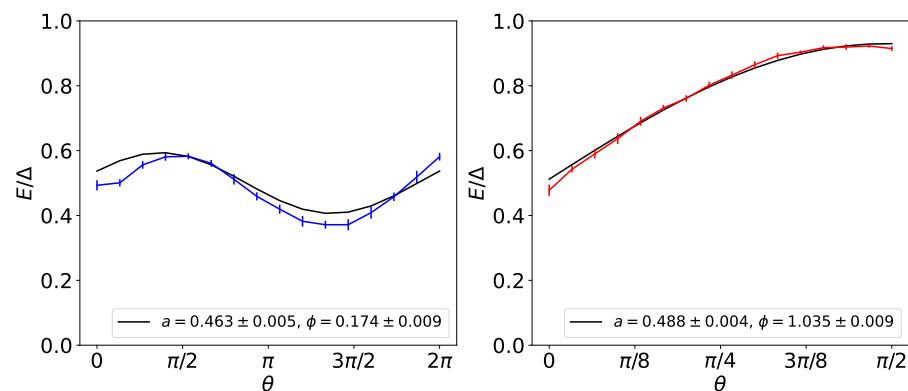
$$a^{(U)} = 0.463 \pm 0.005 \quad (43)$$

$$\phi^{(U)} = 0.174 \pm 0.009 \quad (44)$$

and

$$a^{(V)} = 0.488 \pm 0.004 \quad (45)$$

$$\phi^{(V)} = 1.035 \pm 0.009. \quad (46)$$



**Figure 4.** (Color online). Best fit of the energy stored into the QB (in units of  $\Delta$ ) as a function of  $\theta$  (black curves). Data correspond to Gaussian pulses applied after the action of a unitary matrix  $U$  (blue curve in the left panel) or  $V$  (red curve in the right panels) over the ground state of the system. These transformations are defined in Equations (39) and (40) of the main text, respectively. Every point is obtained averaging over 10 different measurements, and the bars take into account the standard error associated to this average. In both cases, the measurement time has been fixed at  $t_m = 600$  ns.

According to this, while the values of the amplitude are not very far from the ideal condition (few percent mismatch), the phase is again strongly affected by fluctuations.

## 5. Conclusions

In the present paper, we have investigated the performances of the IBM Armonk single qubit in terms of energy storage and charging time by exploiting the Pulse tool included in the Qiskit package. This represents, as far as we know, the first actual simulation of a quantum battery using these kinds of quantum devices.

With our analysis, we have demonstrated that by reasonably choosing a not very narrow and fast decaying classical drive, it is possible to achieve very good energy storage (exceeding 95%) in a very short time (less than 135 ns) with respect to the typical relaxation and dephasing time of the considered device. These performances are comparable with what was observed in very recent state-of-the-art experiments realized using superconducting qubits or semiconducting quantum dots. This indicates that the Armonk qubit, together with an analogous machine in the same range of parameters, is already well designed to be seen as good and stable quantum batteries.

Remarkably enough, their performances can be further improved by errors in the initialization state, such as phase fluctuations, which conversely have a negative impact on the functioning of the device as qubit for quantum computation proposals.

We think that this timely analysis could open new and fascinating perspectives in the fast developing field of quantum batteries and in the more general context of energy transfer devices, addressing for example the controlled application of sequences of pulses in multi-qubit geometries.

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## Abbreviation

The following abbreviations is used in this manuscript:

QB Quantum Battery

## Nomenclature

List of the Relevant Parameters in the Model:

$\Delta$	Level spacing of the qubit
$g$	Qubit–radiation coupling
$f(t)$	Envelop of the applied pulse
$\omega$	Frequency of the modulating cosine
$\theta$	Effective integral of the pulse
$t_m$	Measurement time
$a$	Ground-state probability at $t = 0$
$\phi$	Relative phase at $t = 0$

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