



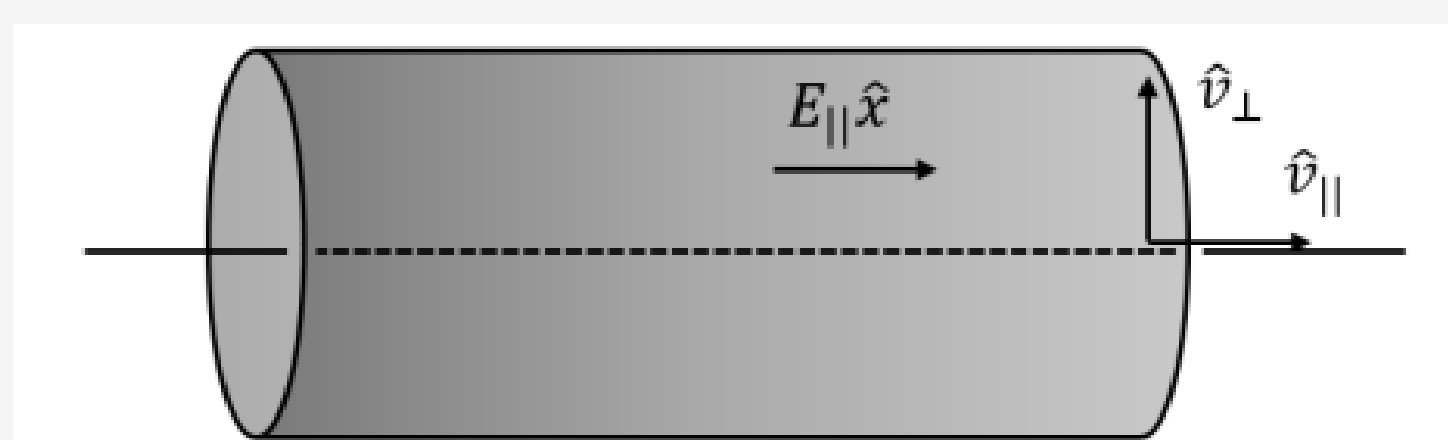
# An Implicit, Low-Rank Solver for 1D2V Vlasov-Fokker-Planck Equation in Cylindrical Coordinates

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## 1D2V Vlasov-Fokker-Planck Model

Laboratory plasma experiments are expensive, necessitating simulations such as the Vlasov-Fokker-Planck model:

$$\frac{\partial f}{\partial t} + v_{\parallel} \frac{\partial f}{\partial x} + E_{\parallel} \frac{\partial f}{\partial v_{\parallel}} = C_{\alpha\alpha}(f) + C_{\alpha e}(f) \quad (1)$$



$f(x, v_{\perp}, v_{\parallel}, t)$  = Ion probability distribution function  
 $E_{\parallel}$  = electric field in parallel ( $x$ ) dimension  
 Coupled with fluid electron pressure equation.

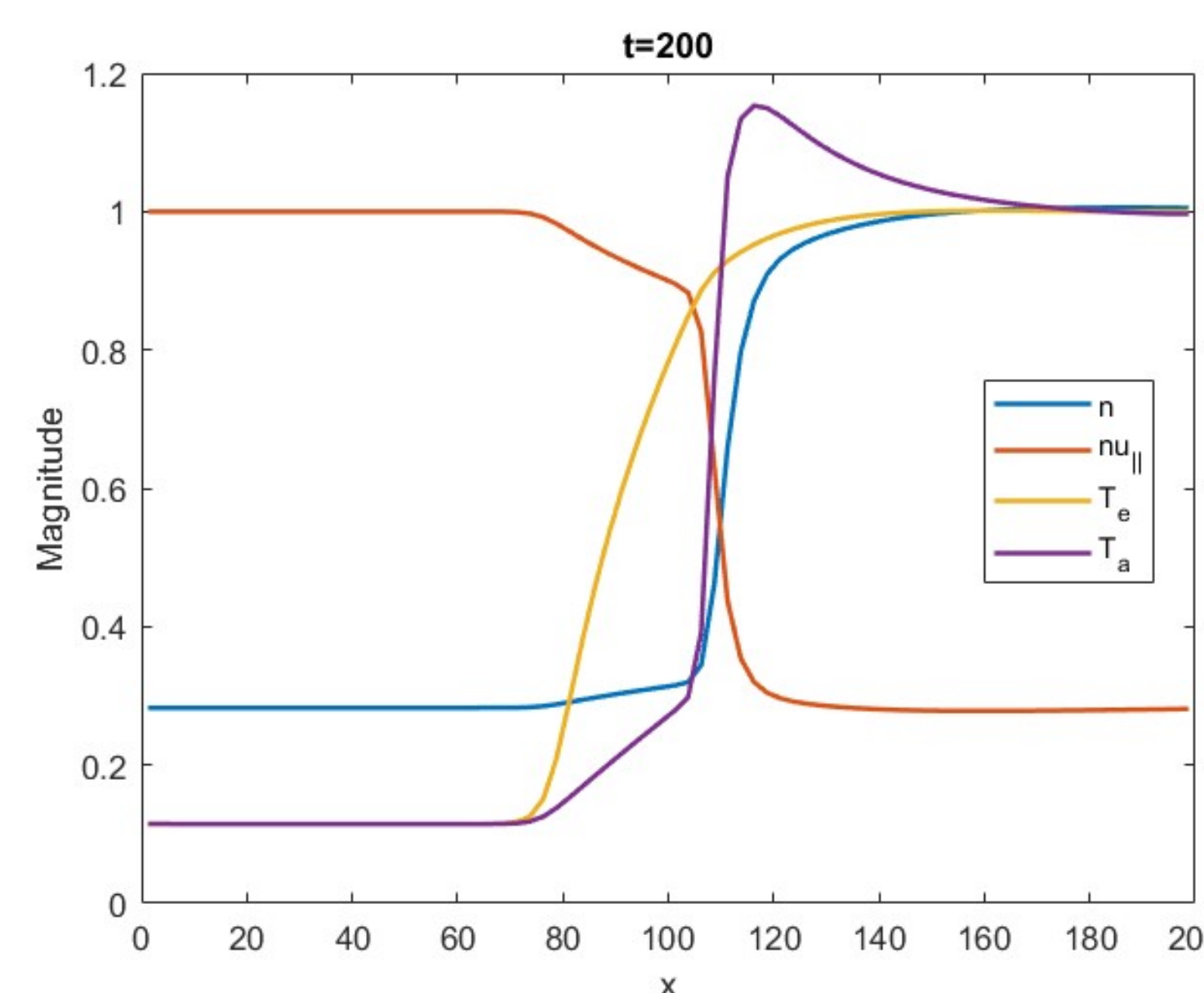
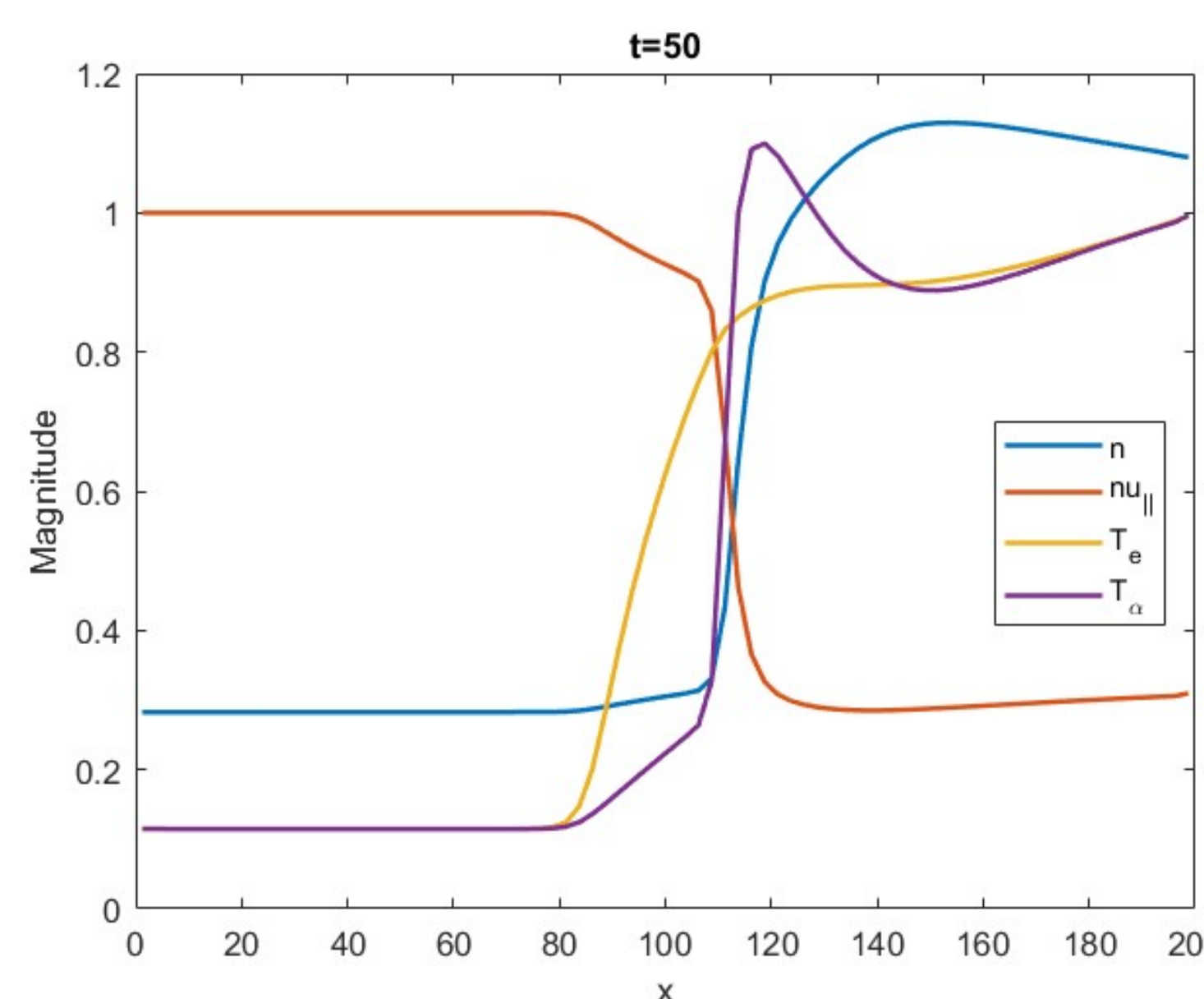
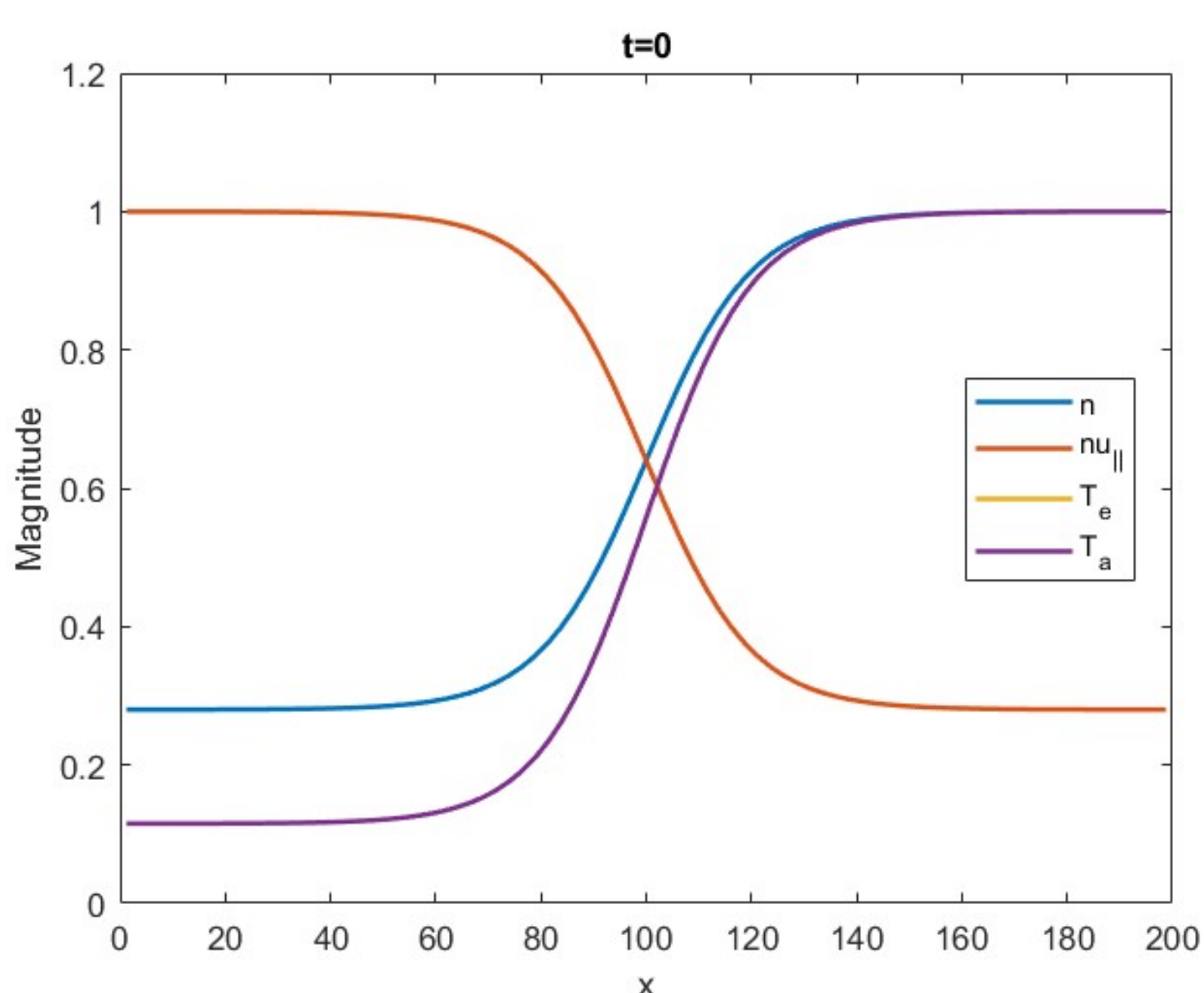
## Discretization: 1st order IMEX RK

$$f_i^k \approx f(x_i, v_{\perp}, v_{\parallel}, t^k)$$

Spatial mesh: Cells  $[x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}}]$

Velocity mesh: Cells  $[v_{\perp, j-\frac{1}{2}}, v_{\perp, j+\frac{1}{2}}] \times [v_{\parallel, l-\frac{1}{2}}, v_{\parallel, l+\frac{1}{2}}]$

$$f_i^{k+1} + \Delta t \left( E_{\parallel}^{k+1} \frac{\partial f_i^{k+1}}{\partial v_{\parallel}} - C_{\alpha\alpha}^{k+1} - C_{\alpha e}^{k+1} \right) = f_i^k - \frac{\Delta t}{\Delta x} v_{\parallel} \left( \hat{f}_{i+\frac{1}{2}}^k - \hat{f}_{i-\frac{1}{2}}^k \right)$$



Evolution of moments  $n, (nu_{\parallel}), T_{\alpha}$ , and  $T_e$  at each spatial node  $x_i$  at times  $t = 0, t = 50, t = 200$ .

## The Algorithm

Naively, storage complexity is  $\mathcal{O}(N_x N_v^2)$

Assume low-rank structure ( $r \ll N_v$ ):  $\mathbf{f}_i^k = \mathbf{V}_{\perp}^k \mathbf{S}^k (\mathbf{V}_{\parallel}^k)^T$

Storage reduction:  $\mathcal{O}(N_x(r^2 + 2N_v r))$

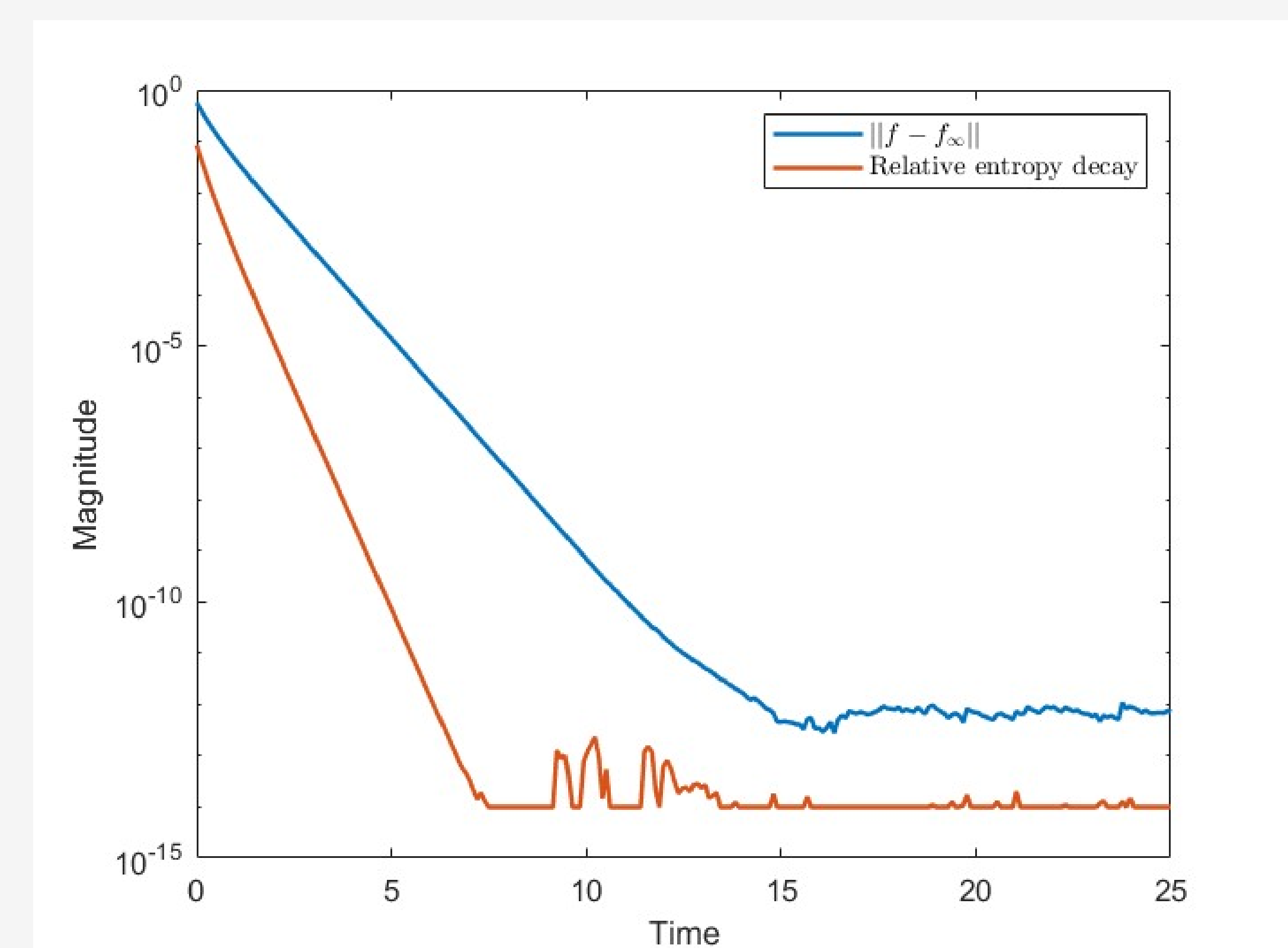
**At each timestep  $t^k$ :**

- Compute moments at next timestep  $n^{k+1}, (nu_{\parallel})^{k+1}, T_{\alpha}^{k+1}, T_e^{k+1}$  using Newton method
- **At each spatial node  $x_i$ :**
  - Compute numerical fluxes  $\hat{f}_{i\pm\frac{1}{2}}^k$
  - Compute electric field: 
$$E_{\parallel}^{k+1} = \frac{1}{q_e n_e^{k+1}} \frac{(n_e T_e)_{i+1}^{k+1} - (n_e T_e)_{i-1}^{k+1}}{2\Delta x}$$
  - Compute collision operators  $C_{\alpha\alpha}^{k+1}, C_{\alpha e}^{k+1}$
  - Solve for  $f_i^{k+1}$  using low-rank projection [1].
 
$$\left. \begin{aligned} \mathbf{K}^k &= \mathbf{V}_{\perp}^k \mathbf{S}^k (\mathbf{V}_{\parallel}^k)^T \mathbf{V}_{\parallel, \star}^{k+1} \\ \mathbf{L}^k &= \mathbf{V}_{\parallel}^k (\mathbf{S}^k)^T (\mathbf{V}_{\perp}^k)^T \mathbf{V}_{\perp, \star}^{k+1} \\ \mathbf{S}^k &= (\mathbf{V}_{\perp, \star}^{k+1})^T \mathbf{V}_{\perp}^k \mathbf{S}^k (\mathbf{V}_{\parallel}^k)^T \mathbf{V}_{\parallel, \star}^{k+1} \end{aligned} \right\} \text{Basis update}$$
  - Galerkin projection
  - Solve for  $\mathbf{K}^{k+1}, \mathbf{L}^{k+1}, \mathbf{S}^{k+1} \mapsto \mathbf{V}_{\perp}^{k+1}, \mathbf{V}_{\parallel}^{k+1}, \mathbf{S}^{k+1}$
  - Apply conservative truncation procedure [2].

## Dougherty-Fokker-Planck Equation 0D2V

$$f_t = C_{\alpha\alpha}(f) = \frac{n_{\alpha}}{T_{\alpha}^{3/2}} \nabla_{\mathbf{v}} \cdot ((\mathbf{v} - \mathbf{u}_{\alpha}) f + \frac{T_{\alpha}}{m_{\alpha}} \nabla_{\mathbf{v}} f)$$

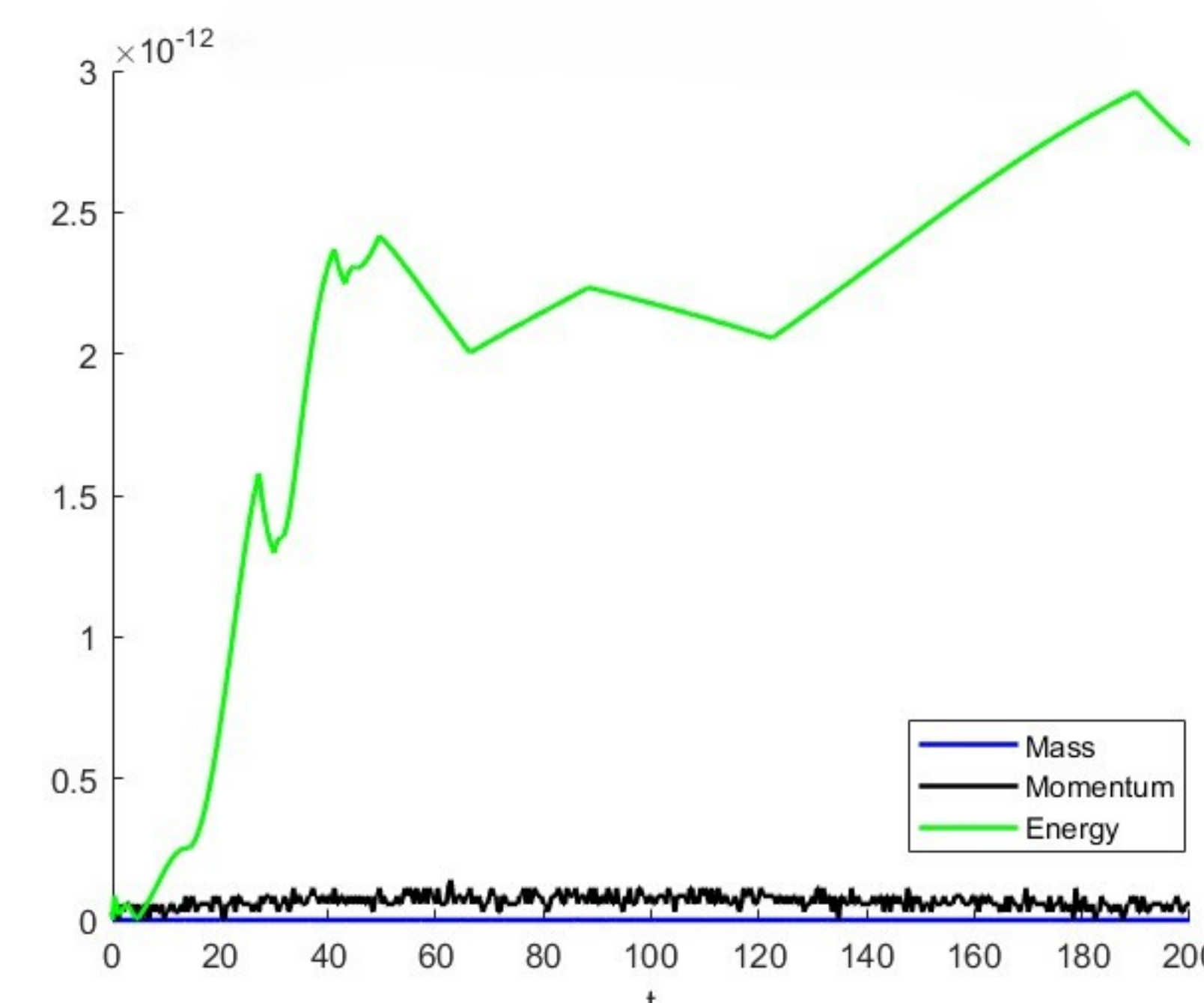
$$f_M(n, u_{\perp}, u_{\parallel}, T) = \frac{n}{(2\pi RT)^{3/2}} \exp\left(-\frac{(v_{\perp} - u_{\perp})^2 + (v_{\parallel} - u_{\parallel})^2}{2RT}\right)$$



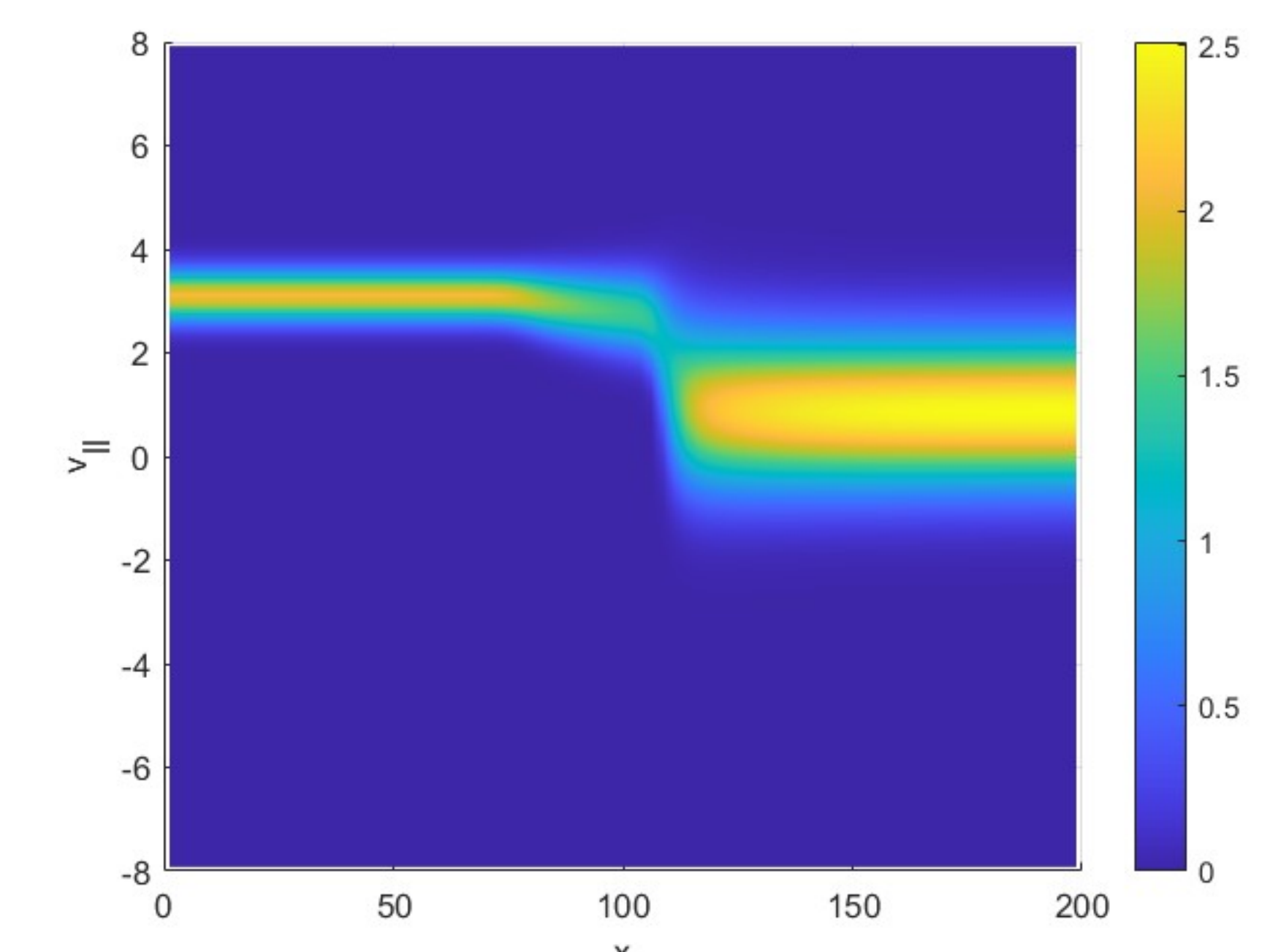
$L_1$  drive to equilibrium solution and relative entropy decay up to machine precision

## References

- [1] Nakao, J., Qiu, J.M. and Einkemmer, L., 2025. Reduced Augmentation Implicit Low-rank (RAIL) integrators for advection-diffusion and Fokker-Planck models. SIAM Journal on Scientific Computing, 47(2), pp.A1145-A1169.
- [2] Guo, W. and Qiu, J.M., 2024. A Local Macroscopic Conservative (LoMaC) low rank tensor method for the Vlasov dynamics. Journal of Scientific Computing, 101(3), p.61.



Mass, momentum, and energy conservation.



Numerical solution at  $t = 200$  integrated along  $v_{\perp}$ .