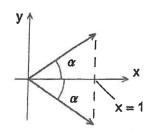
Problem Set 1 - Solutions

1.6 \star b·c = 1 - s², which is zero if and only if $s = \pm 1$. The vectors b and c make equal angles, α , above and below (or below and above) the x axis. The angle between them can be 90° only if $\alpha = 45^{\circ}$.



 $1.8 \star$ (a) Starting from the definition (1.7), we see that

$$\mathbf{r} \cdot (\mathbf{u} + \mathbf{v}) = \sum_{i} r_i (u_i + v_i) = \sum_{i} (r_i u_i + r_i v_i) = \sum_{i} r_i u_i + \sum_{i} r_i v_i = \mathbf{r} \cdot \mathbf{u} + \mathbf{r} \cdot \mathbf{v}$$

where the second equality follows from the distributive property of ordinary numbers. third equality is just a rearrangement of the six terms of the sum, and the last is just definition (1.7) of the two scalar products.

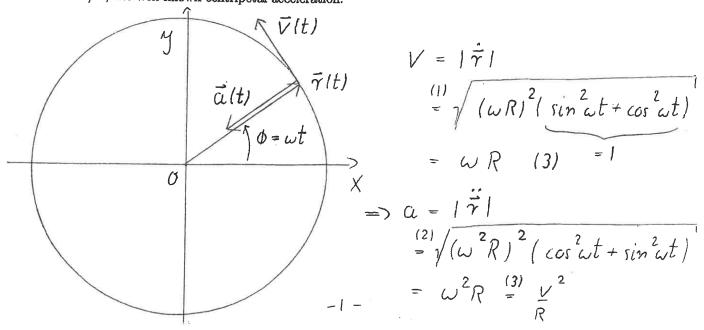
(b) Starting again from (1.7), we find

$$\frac{d}{dt}(\mathbf{r}\cdot\mathbf{s}) = \frac{d}{dt}\sum r_i s_i = \sum \left(r_i \frac{ds_i}{dt} + \frac{dr_i}{dt} s_i\right) = \mathbf{r}\cdot\frac{d\mathbf{s}}{dt} + \frac{d\mathbf{r}}{dt}\cdot\mathbf{s}.$$

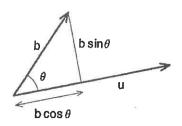
1.10 * The particle's polar angle is $\phi = \omega t$, so $x = R\cos(\omega t)$ and $y = R\sin(\omega t)$ or

$$\mathbf{r} = \hat{\mathbf{x}}R\cos(\omega t) + \hat{\mathbf{y}}R\sin(\omega t).$$
 Differentiating, we find that $\dot{\mathbf{r}} = -\hat{\mathbf{x}}\omega R\sin(\omega t) + \hat{\mathbf{y}}\omega R\cos(\omega t)$ and then (1)
$$\ddot{\mathbf{r}} = -\hat{\mathbf{x}}\omega^2 R\cos(\omega t) - \hat{\mathbf{y}}\omega^2 R\sin(\omega t) = -\omega^2 \mathbf{r} = -\omega^2 R\hat{\mathbf{r}}.$$
 (2)

That is, the acceleration is antiparallel to the radius vector and has magnitude $a = \omega^2 R = v^2/R$, the well known centripetal acceleration.



1.13 \star If θ is the angle between b and the unit vector u, then $\mathbf{u} \cdot \mathbf{b} = b \cos \theta$ is the component of b in the direction of u. Similarly, $|\mathbf{u} \times \mathbf{b}| = b \sin \theta$ is the component of b perpendicular to u. Thus the result $b^2 = (\mathbf{u} \cdot \mathbf{b})^2 + (\mathbf{u} \times \mathbf{b})^2$ is simply a statement of Pythagoras' theorem.



or, following hint:

$$(\hat{u} \cdot \vec{b})^2 + (\hat{u} \times \vec{b})^2 = |\hat{u} \cdot \vec{b}|^2 + |\hat{u} \times \vec{b}|^2$$

$$= |\hat{u}|^2 |\hat{b}|^2 \cos^2 \theta + |\hat{u}|^2 |\hat{b}|^2 \sin^2 \theta$$

$$= |\hat{b}|^2 (\cos^2 \theta + \sin^2 \theta) = |\hat{b}|^2$$

1.17 ** (a) Let us start with the x component of $r \times (u + v)$. From the definition (1.9), we see that

$$[\mathbf{r} \times (\mathbf{u} + \mathbf{v})]_{\mathbf{x}} = r_y(u_z + v_z) - r_z(u_y + v_y) = (r_y u_z - r_z u_y) + (r_y v_z - r_z v_y) = (\mathbf{r} \times \mathbf{u})_x + (\mathbf{r} \times \mathbf{v})_x.$$

Since the y and z components follow in the same way, we conclude that $\mathbf{r} \times (\mathbf{u} + \mathbf{v}) = \mathbf{r} \times \mathbf{u} + \mathbf{r} \times \mathbf{v}$.

(b) Starting again from (1.9), we find for the x component

$$\frac{d}{dt}(\mathbf{r}\times\mathbf{s})_x = \frac{d}{dt}(r_y s_z - r_z s_y) = \left(r_y \frac{ds_z}{dt} - r_z \frac{ds_y}{dt}\right) + \left(\frac{dr_y}{dt} s_z - \frac{dr_z}{dt} s_y\right) = \left(\mathbf{r}\times\frac{d\mathbf{s}}{dt} + \frac{d\mathbf{r}}{dt}\times\mathbf{s}\right)_x.$$

This is the x component of the desired identity. Since the y and z components follow in exactly the same way, our proof is complete.

or, following hint:
(b)
$$\vec{r} \times \vec{s} = \sum_{ijk} \sum_{ijk} r_j s_k \hat{e}_i$$

= $\frac{1}{3t} (\vec{r} \times \vec{s}) = \sum_{ijk} \sum_{ijk} (\vec{r}_j \cdot \vec{r}_k + r_j \cdot \vec{s}_k) \hat{e}_i = \vec{r} \times \vec{s} + \vec{r} \times \vec{s}$
 $(\cdot = \frac{1}{3t})$ -2-

 $\begin{array}{ll} 1.19 \star \star & \frac{d}{dt}[\mathbf{a}\cdot(\mathbf{v}\times\mathbf{r})] = \frac{d\mathbf{a}}{dt}\cdot(\mathbf{v}\times\mathbf{r}) + \mathbf{a}\cdot\frac{d}{dt}(\mathbf{v}\times\mathbf{r}) = \dot{\mathbf{a}}\cdot(\mathbf{v}\times\mathbf{r}) + \mathbf{a}\cdot(\dot{\mathbf{v}}\times\mathbf{r} + \mathbf{v}\times\dot{\mathbf{r}}). \\ \text{The final term } \mathbf{a}\cdot(\mathbf{v}\times\dot{\mathbf{r}}) \text{ is zero because } \dot{\mathbf{r}} = \mathbf{v} \text{ and } \mathbf{v}\times\mathbf{v} = 0. \text{ The second to last term is } \mathbf{a}\cdot(\mathbf{a}\times\mathbf{r}) = \mathbf{0}, \text{ because } \mathbf{a}\times\mathbf{r} \text{ is perpendicular to } \mathbf{a}, \text{ so their scalar product is zero. This leaves us with the requested identity.} \end{array}$

$$\bar{a} \cdot (\bar{v} \times \bar{\gamma}) = Z \quad a_i \, v_j \, \gamma_k$$

=
$$\frac{1}{Jt} \left[\bar{a} \cdot (\bar{v} \times \bar{\gamma}) \right] = \frac{1}{ijk} \left(\dot{a}_i \dot{v}_j \gamma_k + a_i \dot{v}_j \gamma_k + a_i \dot{v}_j \gamma_k \right)$$

$$= \dot{\vec{a}} \cdot (\vec{v} \times \vec{\gamma}) + \dot{\vec{a}} \cdot (\vec{v} \times \vec{\gamma}) + \dot{\vec{a}} \cdot (\vec{v} \times \vec{\gamma})$$

$$= \vec{a} \cdot (\vec{a} \times \vec{\gamma}) = \vec{a} \cdot (\vec{v} \times \vec{v}) = 0$$

$$= 0 \frac{1}{c} (\vec{a} \times \vec{\gamma}) = \vec{a} \cdot (\vec{v} \times \vec{v}) = 0$$

$$= \dot{\bar{a}} \cdot (\bar{\nu} \times \bar{\gamma})$$

note: if
$$\bar{a} = const$$
, i.e., $\bar{a} = 0$:
$$= \frac{1}{2} \int_{\bar{a}} \left[\bar{a} \cdot (\bar{v} \times \bar{\tau}) \right] = 0$$
i.e., $\bar{a} \cdot (\bar{v} \times \bar{\tau}) = const$.
(constant of motion)