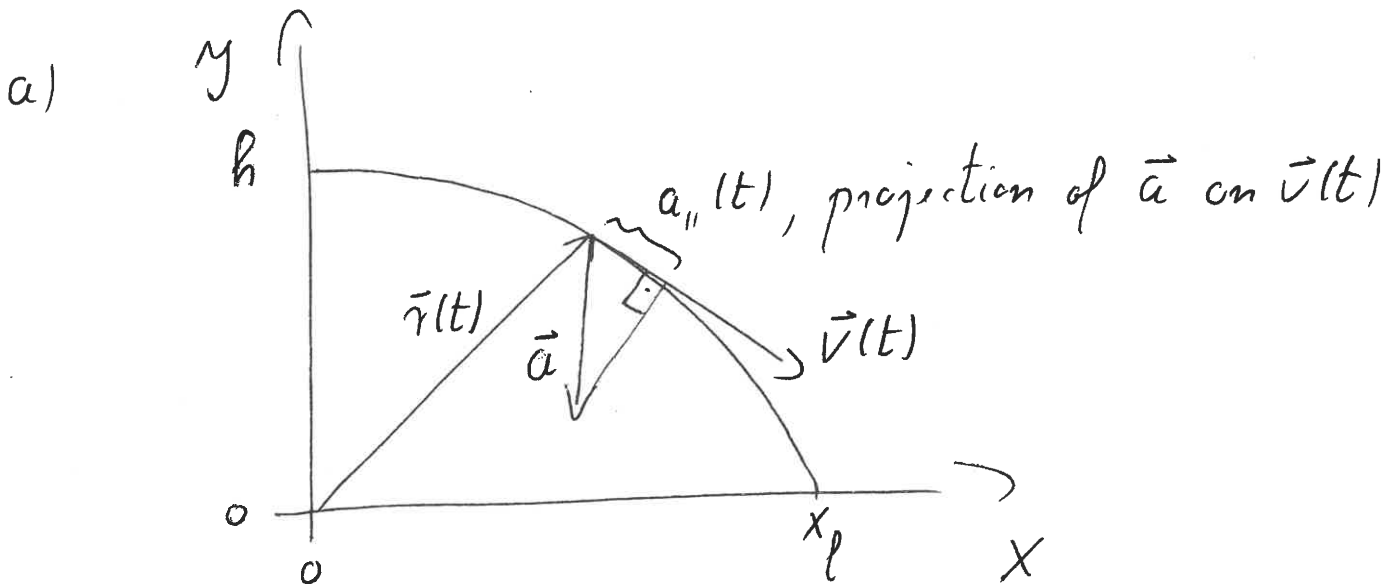


# Problem Set 2 - Solutions

## Problem 1



b)  $\vec{r}(t) = x(t)\hat{x} + y(t)\hat{y} = \omega t\hat{x} + (h - \frac{1}{2}gt^2)\hat{y}$

$\Rightarrow y(t) = h - \frac{1}{2}gt^2$

$y(t_f) = h - \frac{1}{2}gt_f^2 = 0 \Rightarrow t_f = \underline{\underline{\sqrt{\frac{2h}{g}}}}$

c)  $\vec{v} = \frac{d}{dt}\vec{r} = \omega\hat{x} - gt\hat{y}$

$v = |\vec{v}| = \sqrt{\omega^2 + g^2t^2} \quad (1), \quad \hat{v} = \frac{\vec{v}}{v}$

d)  $\vec{a} = \frac{d}{dt}\vec{v} = -g\hat{y} =: \vec{g} = \text{const}$

e)  $\frac{dv}{dt} \stackrel{1)}{=} \frac{1}{2} \frac{1}{\sqrt{\omega^2 + g^2t^2}} g^2 2t = \frac{g^2 t}{v}, \quad a_{||} = \vec{a} \cdot \hat{v} = \frac{g^2 t}{v}$

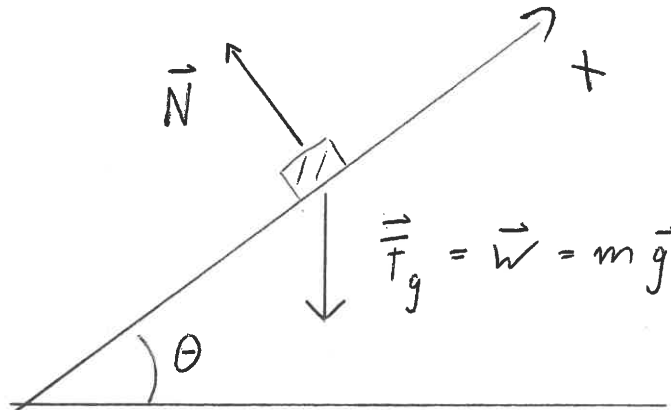
$\Rightarrow \frac{dv}{dt} = a_{||}, \quad a_{||} \text{ speeds up the projectile}$

1.37 \* (a) The two forces on the puck are its weight  $mg$  and the normal force  $N$  of the incline. If we choose axes with  $x$  measured up the slope,  $y$  along the outward normal, and  $z$  horizontally across the slope, then  $N = (0, N, 0)$  and  $g = (-g \sin \theta, -g \cos \theta, 0)$ . Thus Newton's second law reads

$$m\ddot{\mathbf{r}} = \mathbf{N} + m\mathbf{g} \quad \text{or} \quad \begin{cases} m\ddot{x} = -mg \sin \theta \\ m\ddot{y} = N - mg \cos \theta \\ m\ddot{z} = 0 \end{cases}$$

Since  $\dot{z} = 0$  initially, it remains so and hence  $z = 0$  for all  $t$ . The normal force adjusts itself so that  $\ddot{y} = 0$ , and  $y = 0$  for all  $t$ . Finally,  $\ddot{x} = -g \sin \theta$ , which can be integrated twice to give  $x = v_0 t - \frac{1}{2}gt^2 \sin \theta$ .

(b) Solving for the times when  $x = 0$ , we find that  $t = 0$  (at launch) or  $t = 2v_0/(g \sin \theta)$  (the answer of interest).

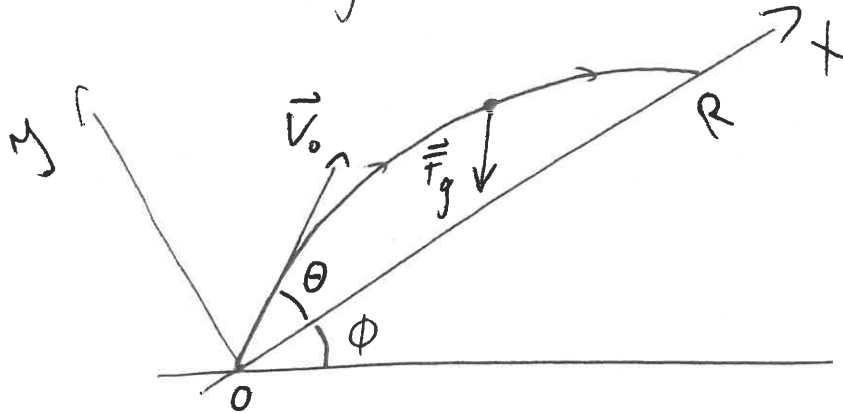


(a)

1.39 \*\*  $x = v_0 t \cos \theta - \frac{1}{2} g t^2 \sin \phi$ ,  $y = v_0 t \sin \theta - \frac{1}{2} g t^2 \cos \phi$ ,  $z = 0$ . When the ball returns to the plane,  $y$  is 0, which implies that  $t = 2v_0 \sin \theta / (g \cos \phi)$ . Substituting this time into  $x$  and using a couple of trig identities yields the claimed answer for the range  $R$ . To find the maximum range, differentiate  $R$  with respect to  $\theta$  and set the derivative equal to zero. This gives  $\theta = (\pi - 2\phi)/4$ , and substitution into  $R$  (plus another trig identity) yields the claimed value of  $R_{\max}$ .

(b)

more explicitly:



force on ball:  $\vec{F}_g = m\vec{g} = -mg \sin(\phi) \hat{x} - mg \cos(\phi) \hat{y}$

(N2), (1.35):

$$m\ddot{x} = F_x = -mg \sin(\phi), \quad m\ddot{y} = F_y = -mg \cos(\phi)$$

initial velocity:  $\vec{v}_0 = v_0 \cos(\theta) \hat{x} + v_0 \sin(\theta) \hat{y}$

$$\Rightarrow x(t) = v_0 \cos(\theta) \cdot t - \frac{1}{2} g \sin(\phi) t^2 \quad (1)$$

$$y(t) = v_0 \sin(\theta) \cdot t - \frac{1}{2} g \cos(\phi) t^2 \quad (2)$$

$\Rightarrow$

(a)

Ball returns to plane:

$$y(t) = 0 \Rightarrow v_0 \sin \theta \cancel{t} - \frac{1}{2} g \cos \phi \cdot \cancel{t^2} = 0$$

$$\Rightarrow \boxed{t} = \frac{2v_0}{g} \frac{\sin \theta}{\cos \phi}$$

In (1):

$$R = x(t) = \frac{2v_0^2}{g} \underbrace{\cos \theta \frac{\sin \theta}{\cos \phi}}_{\frac{\sin \theta}{\cos^2 \phi} \cos \theta \cos \phi}$$

$$- \underbrace{\frac{1}{2} g \frac{4v_0^2}{g^2}}_{\frac{2v_0^2}{g}} \underbrace{\sin \phi \frac{\sin^2 \theta}{\cos^2 \phi}}_{\frac{\sin \theta}{\cos^2 \phi} \sin \theta \sin \phi}$$

$$R = \frac{2v_0^2}{g} \frac{\sin \theta}{\cos^2 \phi} \underbrace{(\cos \theta \cos \phi - \sin \theta \sin \phi)}_{\cos(\theta + \phi)}$$

$$= \frac{2v_0^2}{g} \frac{\sin \theta}{\cos^2 \phi} \cos(\theta + \phi) \quad (3)$$

result quoted in question

$$R_{\max} \text{ from } \frac{dR}{d\theta} = 0 :$$

$$\frac{dR}{d\theta} = \frac{2v_0^2}{g \cos^2 \phi} \underbrace{\left\{ \cos \theta \cos(\theta + \phi) - \sin \theta \sin(\theta + \phi) \right\}}_{\cos(2\theta + \phi)}$$

$$= 0$$

$$\Rightarrow 2\theta + \phi = \frac{\pi}{2} \hat{=} 90^\circ$$

$$2\theta + \phi = \frac{\pi}{2} \quad (4)$$

$$\Rightarrow \theta + \phi = \frac{\pi}{2} - \theta$$

$$\text{in (3): } \cos(\theta + \phi) = \cos\left(\frac{\pi}{2} - \theta\right) = \sin \theta$$

$$\Rightarrow R_{\max} = \frac{2v_o^2}{g} \frac{\sin^2 \theta}{\cos^2 \phi} \quad (5)$$

$$(4) \Rightarrow \theta = \frac{1}{2} \left( \frac{\pi}{2} - \phi \right) \quad (b)$$

$$\Rightarrow \sin^2 \theta = \sin^2 \left\{ \frac{1}{2} \left( \frac{\pi}{2} - \phi \right) \right\}$$

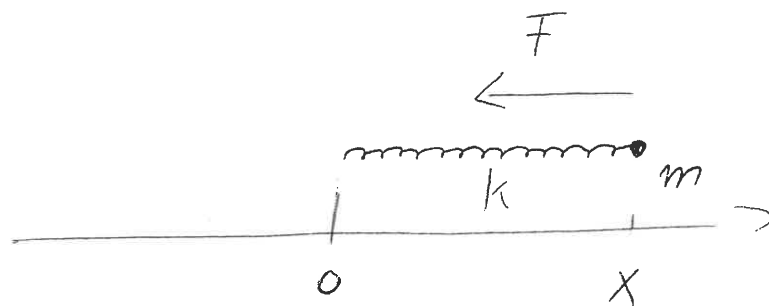
$$= \frac{1}{2} \left\{ 1 - \underbrace{\cos\left(\frac{\pi}{2} - \phi\right)}_{\sin \phi} \right\} \text{ trig. relation}$$

$$\text{in (5): } \cos^2 \phi = 1 - \sin^2 \phi = (1 + \sin \phi)(1 - \sin \phi)$$

$$\Rightarrow R_{\max} = \frac{v_o^2}{g} \frac{1 - \cancel{\sin \phi}}{(1 + \sin \phi)(1 - \cancel{\sin \phi})}$$

$$= \frac{v_o^2}{g(1 + \sin \phi)} \quad \text{quoted in question}$$

# Problem 4



$$F = -kx \quad (1)$$

$$a) \quad F = ma = m \frac{d^2x}{dt^2} \quad (N2)$$

$$\stackrel{(1)}{\Rightarrow} m \frac{d^2x}{dt^2} = -kx \quad | \cdot \frac{1}{m}$$

$$\Rightarrow \frac{d^2x}{dt^2} = -\frac{k}{m}x \quad (2)$$

(2) contains 2nd derivative  $\frac{d^2x}{dt^2}$

$\hat{=}$  2nd order differential equation (DE)  
for position variable  $x(t)$

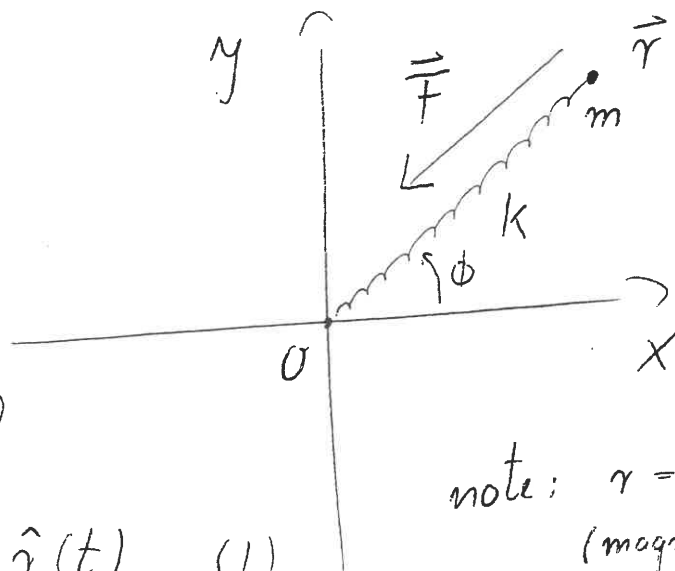
b) Show:  $x(t) = a \cos(\omega t) + b \sin(\omega t)$   
solves the DE (2):

$$\frac{dx}{dt} = -a \sin(\omega t) \cdot \omega + b \cos(\omega t) \cdot \omega$$

$$\frac{d^2x}{dt^2} = -a \cos(\omega t) \cdot \omega^2 - b \sin(\omega t) \cdot \omega^2$$

$$= -\omega^2 x(t) \Rightarrow (2) \text{ fulfilled, with } \underline{\omega = \sqrt{\frac{k}{m}}}$$

# Problem 5



$$\vec{r}(t) = r(t) \hat{r}(t)$$

$$\vec{F}(t) = -k r(t) \hat{r}(t) \quad (1)$$

note:  $r = |\vec{r}|$   
(magnitude)

a)  $\vec{a}(t) = \frac{d^2}{dt^2} \vec{r}(t)$  in polar coordinates:  
(1.47),  $\cdot \triangleq \frac{d}{dt}$

$$\vec{a} = (\ddot{r} - r\dot{\phi}^2) \hat{r} + (r\ddot{\phi} + 2\dot{r}\dot{\phi}) \hat{\phi}$$

$\vec{F} = m\vec{a}$  with  $\vec{F}$  from (1):

$$-kr\hat{r} = m(\ddot{r} - r\dot{\phi}^2)\hat{r} + m(r\ddot{\phi} + 2\dot{r}\dot{\phi})\hat{\phi}$$

comparing components in  $\hat{r}, \hat{\phi}$ :

$$\Rightarrow -kr = m(\ddot{r} - r\dot{\phi}^2) \quad (2)$$

$$r\ddot{\phi} + 2\dot{r}\dot{\phi} = 0 \quad (3)$$

(2), (3) are two coupled 2nd order DEs  
for  $r(t), \phi(t)$ ; finding solutions is difficult  
in general!



b) Ansatz for  $r(t)$ ,  $\phi(t)$ :

$$r = R = \text{const.}$$

$$\phi(t) = \omega t, \quad \omega = \text{const} \quad (4)$$

= circular orbit with radius  $R$

$$\Rightarrow \dot{r} = 0, \quad \ddot{r} = 0$$

$$\dot{\phi} = \omega, \quad \ddot{\phi} = 0$$

$$\text{in (2):} \quad -kR \stackrel{!}{=} m(0 - R\omega^2)$$

$$-kR = -mR\omega^2$$

$$\text{fulfilled if } \underline{\omega = \sqrt{\frac{k}{m}}}$$

$$\text{in (3):} \quad R \cdot 0 + 2 \cdot 0 \cdot \omega \stackrel{!}{=} 0$$

$$0 = 0, \quad \text{fulfilled}$$

$\Rightarrow$  (4) is a solution of (2), (3)

with  $\omega = \sqrt{\frac{k}{m}}$ , same  $\omega$  as in P4!

explanation:

write  $\vec{F} = -k r \hat{r} = -k \vec{r}$

in Cartesian (x-y) coordinates:

$$\begin{aligned}\vec{F} &= -k \vec{r} = -k (x \hat{x} + y \hat{y}) \\ &= -k x \hat{x} - k y \hat{y}\end{aligned}$$

$$\vec{a} = \frac{d^2 \vec{r}}{dt^2} = \frac{d^2 x}{dt^2} \hat{x} + \frac{d^2 y}{dt^2} \hat{y}$$

comparing components in  $\hat{x}, \hat{y}$ :

$$\vec{F} = m \vec{a} \Rightarrow \vec{a} = \frac{1}{m} \vec{F}$$

$$\Rightarrow \frac{d^2 x}{dt^2} = -\frac{k}{m} x, \quad \frac{d^2 y}{dt^2} = -\frac{k}{m} y \quad (5)$$

(5) are two independent, one-dimensional DEs for  $x(t), y(t)$  as in P4!

$\Rightarrow$  harmonic motions for  $x(t), y(t)$  with  $\omega = \sqrt{\frac{k}{m}}$  as in P4; for a circular orbit,  $x(t), y(t)$  are  $90^\circ$  out of phase:

$$x(t) = R \cos(\omega t), \quad y(t) = R \cos(\omega t - \frac{\pi}{2}) = R \sin(\omega t)$$