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Problem 1. What is the type of (elliptic, parabolic, hyperbolic) the following second-order PDEs:

(a)
$$u_{xx} - 3u_{xy} + 2u_{yy} + u_y + 5u = 0$$

(b)
$$9u_{xx} + 6u_{xy} + u_{yy} + u_x = 0$$

Solution.

(a)

$$D = b^{2} - 4ac$$

$$= (-3)^{2} - 4(1)(2)$$

$$= 9 - 8$$

$$= 1$$

$$> 0$$

Since D > 0, the PDE is hyperbolic.

(b)

$$D = b^{2} - 4ac$$

$$= (6)^{2} - 4(9)(1)$$

$$= 36 - 36$$

$$= 0$$

$$= 0$$

Since D = 0, the PDE is parabolic.

Problem 2. Sketch the characteristic lines, and solve the transport equation

$$\begin{cases} u_t + 2u_x = 0 \\ u(x,0) = x^3 \end{cases}$$

Sketch u(x, 0), u(x, 1), u(x, 2) on the same graph and convince yourself that the solutions are moving to the right with speed 2.

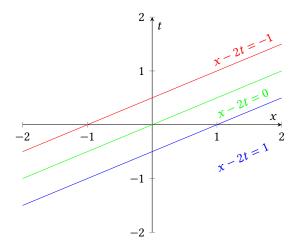
Solution.

$$(u_x, u_t) \cdot (2, 1) = 0$$

u is constant along the lines parallel to (2,1). These characteristic lines are given by

$$t = \frac{x}{2} + C_1 \iff x - 2t = C_2$$

for some constants C_1 , C_2 . We may graph a few characteristic lines:



We can express u(x, y) as a function of x - 2t:

$$u(x,t) = f(x-2t)$$

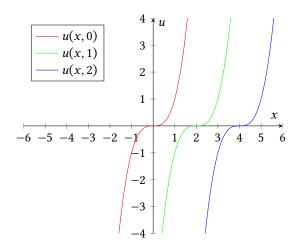
We can solve for f using the initial condition:

$$u(x,0) = f(x) = x^3$$

Thus, the solution to the transport equation is

$$u(x,t) = (x-2t)^3$$

We can graph u(x, 0), u(x, 1), u(x, 2) on the same graph:



We can see that the solutions are moving to the right with speed 2.

Problem 3. We define

$$u(x,t) := \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}}$$

Verify that the function u(x, t) is a solution of the heat equation

$$u_t = Du_{xx}$$

for all (x, t) with $-\infty < x < \infty$ and t > 0.

Solution.

$$\begin{split} u_t &= \frac{1}{\sqrt{4\pi D}} \left(-\frac{1}{2} t^{-\frac{3}{2}} e^{-\frac{x^2}{4Dt}} + t^{-\frac{1}{2}} e^{-\frac{x^2}{4Dt}} \frac{x^2}{4Dt} \right) \\ &= \frac{1}{2t \sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}} \left(\frac{x^2}{2Dt} - 1 \right) \\ &= \frac{1}{2t} \left(\frac{x^2}{2Dt} - 1 \right) u(x, t) \\ u_x &= \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}} \left(-\frac{x}{2Dt} \right) \\ u_{xx} &= \frac{1}{\sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}} \left(\frac{x^2}{4D^2t^2} - \frac{1}{2Dt} \right) \\ &= \frac{1}{2t \sqrt{4\pi Dt}} e^{-\frac{x^2}{4Dt}} \left(\frac{x^2}{2Dt} - 1 \right) \frac{1}{D} \\ &= \frac{1}{2Dt} \left(\frac{x^2}{2Dt} - 1 \right) u(x, t) \end{split}$$

We can see that

$$u_t = Du_{xx}$$

for all (x, t) with $-\infty < x < \infty$ and t > 0.

Problem 4. For Q > 0, compute

$$\int_{-\infty}^{\infty} e^{-Qx^2} dx$$

Hint: Multiply by the identical integral but replace x variable to y, and change vari- ables to polar coordinates.

Solution. Let $I = \int_{-\infty}^{\infty} e^{-Qx^2} dx = \int_{-\infty}^{\infty} e^{-Qy^2} dy$. Then

$$I^{2} = \left(\int_{-\infty}^{\infty} e^{-Qx^{2}} dx \right) \left(\int_{-\infty}^{\infty} e^{-Qy^{2}} dy \right)$$
$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-Q(x^{2} + y^{2})} dx dy$$

We can change variables to polar coordinates:

$$I^2=\int_0^{2\pi}\int_0^\infty e^{-Qr^2}rdrd heta
onumber \ =2\pi\int_0^\infty e^{-Qr^2}rdr$$

Let $u = -Qr^2$, then du = -2Qrdr. Thus,

$$I^{2} = 2\pi \left[\frac{1}{-2Q} e^{-Qr^{2}} \right]_{0}^{\infty}$$
$$= \frac{\pi}{Q}$$

Thus,

$$I = \int_{-\infty}^{\infty} e^{-Qx^2} dx = \sqrt{\frac{\pi}{Q}}$$

Problem 5. We will derive the transport equation $u_t + cu_x = 0$ for $x, t \in \mathbb{R}$ using an approach via compartment models similar to what we used for the derivation of the heat equation in class: assume that u(x,t) is the concentration of particles at position x and time t. Let h and τ be positive and small. First, divide up the real line into intervals of length h centered at positions ..., x - 2h, x - h, x, x + h, x + 2h, Next, we assume that during the time interval from t to $t + \tau$ all particles in an interval move to the nearest interval to their right.

- (a) Write down a balance equation that reflects this situation. Please provide a clear description of the various terms in the balance equation.
- (b) Rearrange the terms in your balance equation and substitute them from Taylor expansions appropriately.
- (c) Choose an appropriate way to let h, t go to zero, and show that the limit is the transport equation we want.

Solution.

(a) Number of particles at time $t + \tau$ = Number of particles at time t + Change:

$$hu(x, t + \tau) = hu(x, t) + \text{Change}$$

Change = (Number of particles that moved from x - h to x) - (Number of particles that moved from x to x + h). = h(u(x - h, t) - u(x, t))

We have

$$hu(x, t + \tau) = hu(x, t) + h(u(x - h, t) - u(x, t))$$

(b) Rearranging the terms and dividing by $h\tau$, we have

$$\frac{u(x,t+\tau)-u(x,t)}{\tau} = \frac{h}{\tau} \frac{u(x-h,t)-u(x,t)}{h} \tag{1}$$

We can use a Taylor expansion to approximate u(x - h, t)

$$u(x - h, t) = u(x, t) - hu_x(x, t) + O(h^2)$$

so that we have

$$u(x - h, t) - u(x, t) = -hu_x(x, t) + O(h^2)$$

Substituting into (1), we have

$$\frac{u(x,t+\tau)-u(x,t)}{\tau}=\frac{h}{\tau}\frac{-hu_x(x,t)+O(h^2)}{h}$$

$$\frac{u(x,t+\tau)-u(x,t)}{\tau}=\frac{h}{\tau}\left(-u_x(x,t)+O(h)\right)$$

(c) Let $h \to 0$ and $\tau \to 0$. Let us also choose $\tau = \frac{h}{c}$ so that $\tau \to 0$ as $h \to 0$. Then

$$\lim_{h \to 0} \frac{u(x, t + \tau) - u(x, t)}{\tau} = \lim_{h \to 0} \frac{h}{\tau} \left(-u_x(x, t) + O(h) \right)$$

$$\lim_{\tau \to 0} \frac{u(x, t + \tau) - u(x, t)}{\tau} = \lim_{h \to 0} \frac{h}{\frac{h}{c}} \left(-u_x(x, t) + O(h) \right)$$

$$u_t(x,t) = -cu_x(x,t)$$

$$u_t + cu_x = 0$$

Thus, the limit is the transport equation.