

INTRODUCTION AND APPROACH

- For a typical fusion device $Q \sim \mathcal{O}(\beta^2 B^4)$. Given a value of B it implies that high beta operation is favorable.
- The ballooning(or high wave-number) mode is useful for determining an upper limit of β that a tokamak plasma can achieve.
- Conventional methods assume that the least stable modes lie at the outboard mid-plane i.e. $\theta_0 = 0$.
- Our aim is to use $s - \alpha$ model for concentric flux surfaces and compare its results to gyrokinetic simulations of an equilibrium with finite θ_0

The two *equivalent* forms of the Ideal ballooning equation are

$$\underline{B} \cdot \nabla \left(\frac{|\nabla S|^2}{B^2} \underline{B} \cdot \nabla F \right) + 2 \frac{dp}{d\psi} \frac{(\underline{B} \times \nabla S) \cdot \underline{K}}{B^2} F, \quad \underline{K} = \frac{\nabla(p + B^2/2)}{B^2}, \quad \underline{B} \cdot \nabla S = 0 \quad (1)$$

$$\frac{1}{\mathcal{J}} \frac{\partial}{\partial y} \left[\frac{1}{\mathcal{J} R^2 B_\chi^2} \left\{ 1 + \left(\frac{R^2 B_\chi^2}{B} \int_{y_0}^y \nu' dy \right)^2 \right\} \frac{\partial F}{\partial y} \right] + F \left[\frac{2\mathcal{J} p'}{B^2} \frac{\partial}{\partial \psi} \left(p + \frac{B^2}{2} \right) - \frac{I p'}{B^4} \left(\int_{y_0}^y \nu' dy \right) \frac{\partial B^2}{\partial y} \right] = 0 \quad (2)$$

We apply Mercier-Luc formulation to these equations to obtain their lowest order form in ρ around a flux surface. For a $\mathcal{O}(\beta) \gg 1$, large aspect ratio $\mathcal{O}(\epsilon^{-1}) \gg 1$ equilibrium, the orderings of various the figures of merit are as follows

$$\beta_\varphi = \mathcal{O}\left(\frac{p}{B_\varphi^2}\right) \sim \epsilon, \quad \beta_p = \mathcal{O}\left(\frac{p}{B_p^2}\right) \sim \frac{1}{\epsilon}, \quad \frac{B_p}{B_\varphi} \sim \mathcal{O}(\epsilon), \quad q \sim \mathcal{O}(1), \quad \frac{\rho \psi_1}{\psi_0} \sim \mathcal{O}(\epsilon)$$

Applying these orderings to equation (1), we get the $s - \alpha$ equation

$$\frac{d}{d\theta} \left\{ 1 + \left(s(\theta - \theta_0) - \alpha (\sin(\theta) - \sin(\theta_0)) \right)^2 \right\} \frac{dF}{d\theta} + \alpha \left\{ \cos(\theta) + \sin(\theta) \left(s(\theta - \theta_0) - \alpha (\sin(\theta) - \sin(\theta_0)) \right) \right\} F = 0 \quad (3)$$

- We integrate the $s - \alpha$ equation for a concentric circle test profile and use Newcomb's criterion to determine the regions of marginal stability.
- Then we take a Grad-Shafranov stable high β equilibrium profile and plot $s - \alpha$ diagram and also use GS2 to perform Greene-Chance analyses to determine marginal stability

RESULTS

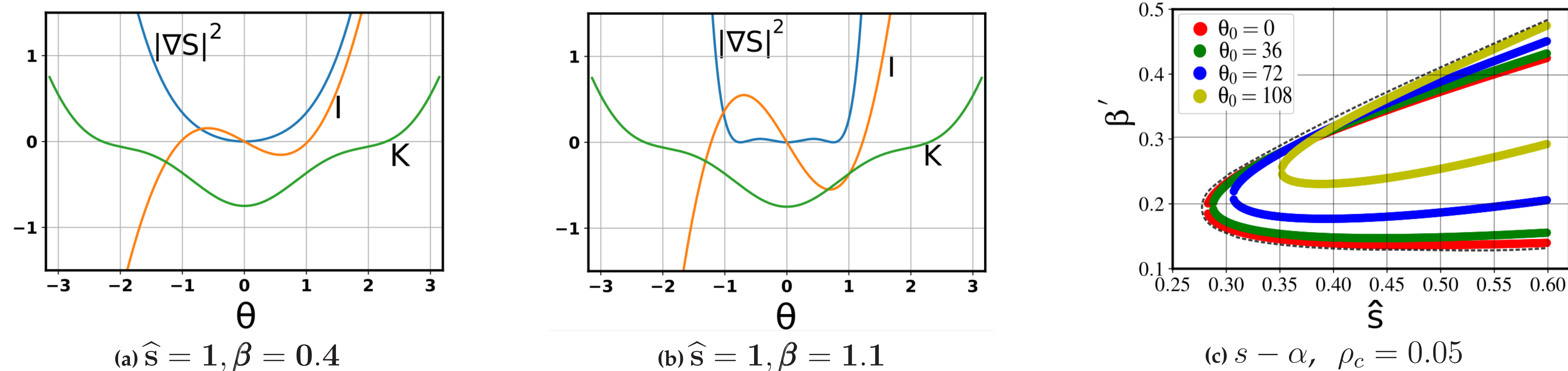


Figure 1. a) and b) Qualitative variation of the stretching($|\nabla S|^2$), total shear(I and curvature(K) terms) terms, c) $s - \alpha$ diagrams for various values of θ_0

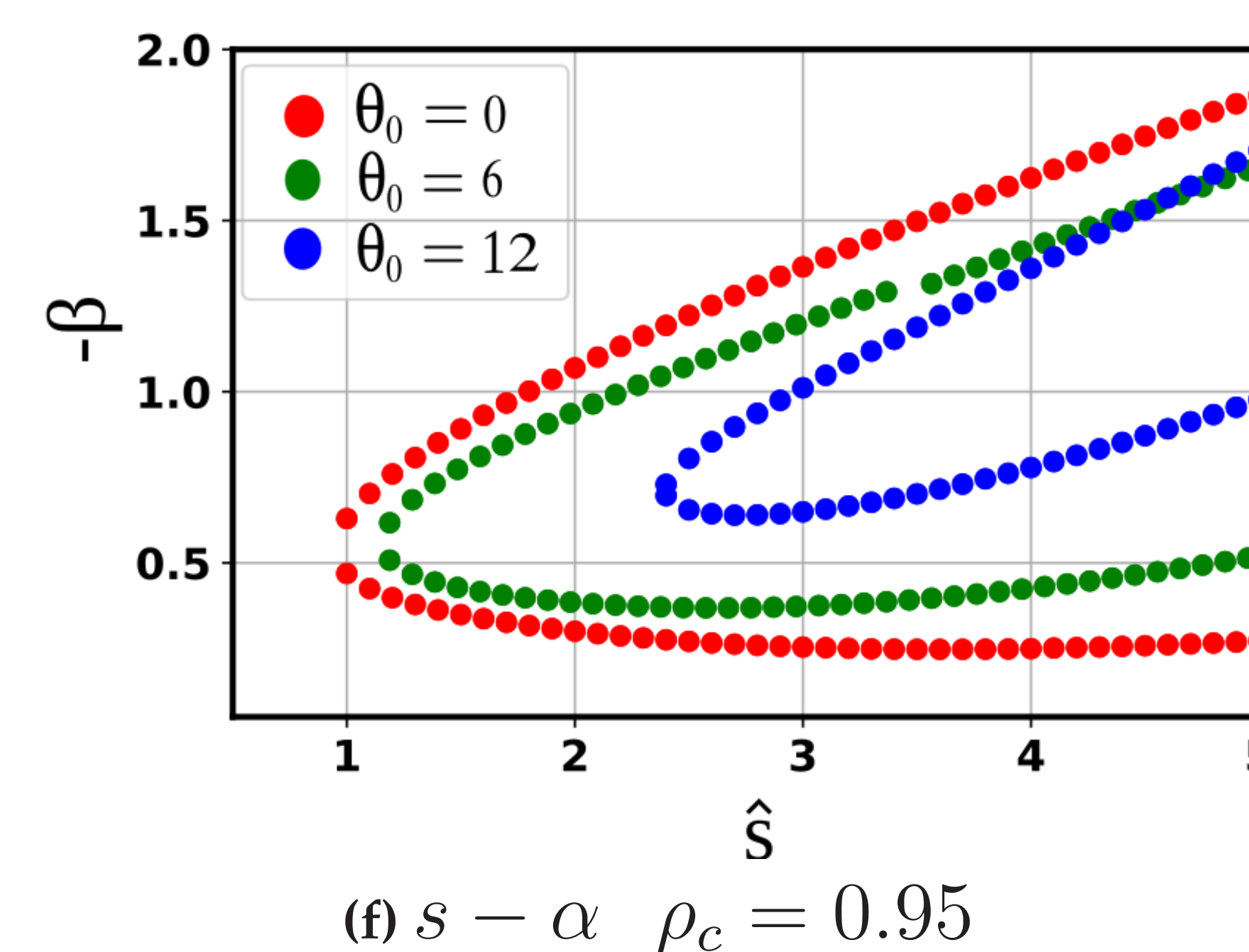
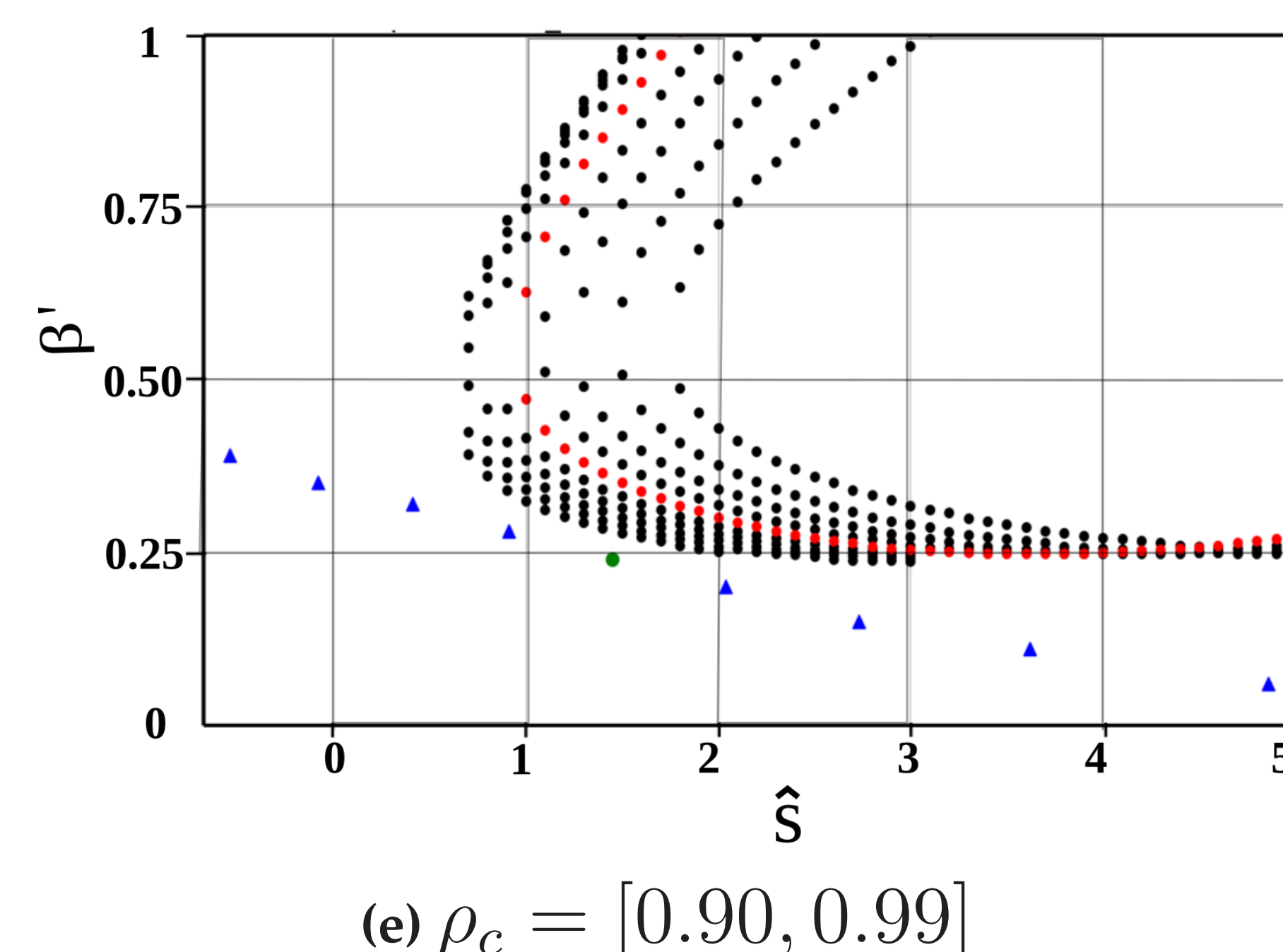
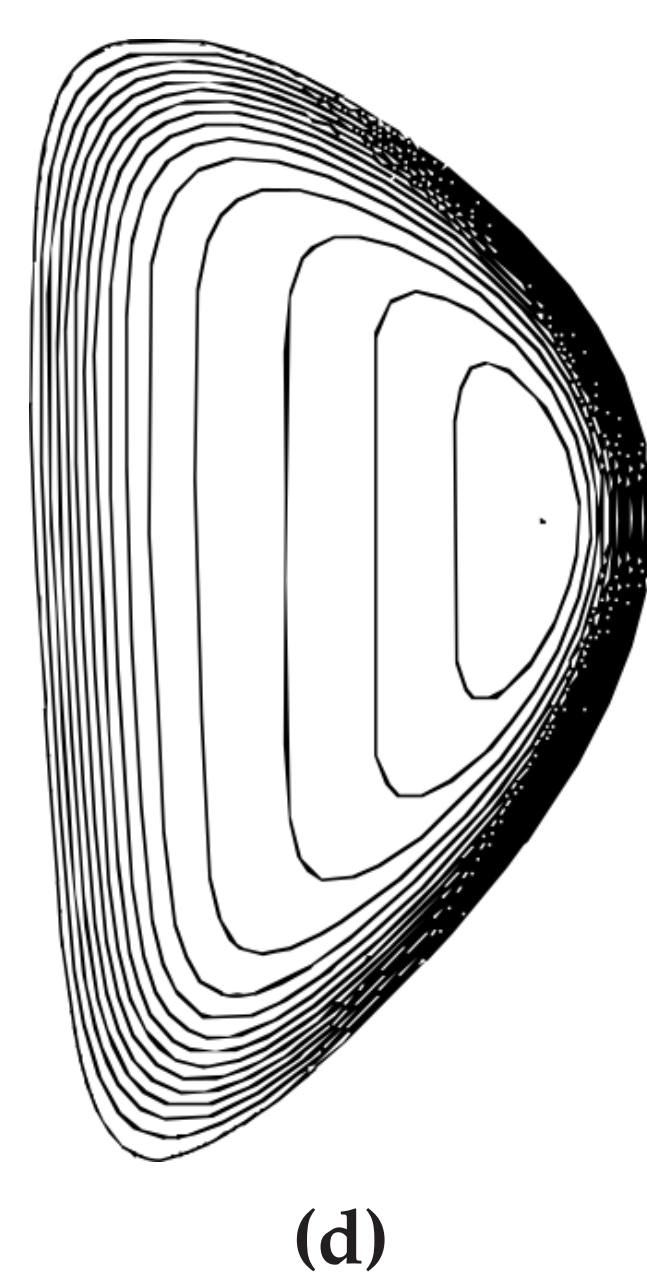


Figure 1. d) High beta equilibrium profile and e) $s - \alpha$ diagrams for different surfaces ρ_c near the LCFS and the stability boundary from GS2 f) $s - \alpha$ diagrams for various values of θ_0 for the surface $\rho_c = 0.95$

SUMMARY

- We have investigated the effect of finite θ_0 on the boundary of marginal stability to ballooning modes for a test profile and an actual high β equilibrium profile.
- Although the test profile shows significant changes to the $s - \alpha$ envelope, the actual profile doesn't change as much for the tested range of \hat{s} .
- Future work would involve checking the effect of large θ_0 and \hat{s} values and a broader range of flux surfaces(ρ_c).

REFERENCES

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