

Last Time:

- Simulation with Constraints
- DAEs
- Baumgarte Stabilization

Today:

- Variational Integrators
 - Momentum
 - Legendre Transform
 - Hamiltonian Mechanics
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Variational Integrators:

- Rather than dealing with generic $\dot{x} = f(x)$, we specialize to Lagrangian systems
- Key idea: discretize/approximate Lagrangian/least-action principle rather than ODE.
- By construction, automatically obey constraints and conserve momentum, energy, etc.

* Least-Action:

$$\min_{q(t)} S(q(t)) = \int_{t_0}^{t_f} L(q, \dot{q}) dt$$

- Let's break this into a bunch of tiny pieces (exact so far):

$$\min_{q(t)} S(q(t)) = \sum_{n=1}^{N-1} \int_{t_n}^{t_{n+1}} L(q, \dot{q}) dt$$

$t_n = t_0 + kh$

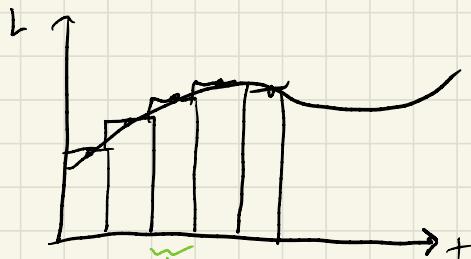
- Now let's approximate the tiny integrals:

$$\int_{t_n}^{t_{n+1}} L(q, \dot{q}) dt \approx h L\left(\frac{q_n + q_{n+1}}{2}, \frac{\dot{q}_{n+1} - \dot{q}_n}{h}\right)$$

midpoint approximation

"Discrete Lagrangians"

$$L_d(q_n, \dot{q}_{n+1})$$



- Approximating an integral like this is called "quadrature".
- Well-developed theory very similar to RK methods (i.e. order conditions etc.)
- Now we get a discrete "Action Sum":

$$\min_{q_{1:N}} S_d(q_{1:N}) = \sum_{n=1}^{N-1} L_d(q_n, \dot{q}_{n+1})$$

- This is exactly the same scenario we saw with the hanging cable. We can get local FON conditions by looking at a very short segment with only 3 knot points:

$$\frac{\partial}{\partial q_2} [L_d(q_1, \dot{q}_2) + L_d(q_2, \dot{q}_3)] = 0$$

$$\Rightarrow D_2 L_d(q_1, \dot{q}_2) + D_1 L_d(q_2, \dot{q}_3) = 0$$

"Discrete Euler-Lagrange Equation"

- We can use the D_{EL} equation as an implicit integrator
- Given $q_1, q_2 \Rightarrow$ can solve for q_3 using Newton's method

* Pendulum Example:

- Energy oscillates about true value \Rightarrow Lyapunov stable
- Energy behavior is independent of step size
- Very good for long-term simulations
- Turns out accuracy is 2nd order. In general accuracy of the quadrature rule used in Ld = accuracy of variational integrator.
- Easy to construct higher-order methods.

Momentum + Hamiltonian:

- We needed \dot{q}_1 and \dot{q}_2 to initialize our variational integrator.
- How do we do this with q_1 and U , instead?
- How do we get U_n in general?

- Remember the finite-diff velocities in Ld are defined at the midpoints, not knot points.
- We can trivially rewrite the least-action principle as:

$$\min_{\begin{smallmatrix} t(q) \\ \mathcal{J}(q) \end{smallmatrix}} \int_{t_0}^{t_f} L(q, v) dt$$

s.t. $\dot{q} = v \leftarrow \text{kinematic constraints}$

- Enforcing the kinematic constraint with a Lagrange multiplier gives KKT conditions:

$$\frac{\partial}{\partial q(t)} \left[\int_{t_0}^{t_f} L(q, v) + \rho^T (q - v) dt \right] = 0$$

$$\frac{\partial}{\partial v(t)} \left[\int_{t_0}^{t_f} L(q, v) + \rho^T (q - v) dt \right] = 0$$

$$\dot{q} - v = 0$$

- Doing the standard integration by parts trick:

$$\frac{\partial L}{\partial q} - \dot{\rho} = 0$$

$$\frac{\partial L}{\partial v} - \rho = 0$$

$$\dot{q} = v$$

- Let's plug in $L = \frac{1}{2}v^T M v - U(q)$:

$$\Rightarrow \begin{cases} \frac{\partial L}{\partial q} = -\nabla U(q) = \dot{p} \\ \frac{\partial L}{\partial v} = Mv = p \end{cases} \Rightarrow \begin{matrix} \text{Lagrange multiplier} \\ = \text{Momentum} \end{matrix}$$

$\underbrace{-\nabla U(q)}_{F}$ $\underbrace{\dot{p}}_{M\dot{v}}$
 $\underbrace{Mv}_{\text{momentum}}$

- This version of the least-action principle is called the Hamilton-Pontryagin principle because of its connection to optimal control.
- We can replace v with p in these equations if we want:

$$\frac{\partial L}{\partial v} = M(q)v = p \Rightarrow v = M^{-1}p$$

$$\Rightarrow \int_{t_0}^{t_f} \frac{1}{2} v^T M(q)v - U(q) + p^T (\dot{q} - v) dt$$

$$= \int_{t_0}^{t_f} \frac{1}{2} p^T M(q)^{-1} p - U(q) + p^T (\dot{q} - M(q)^{-1}p) dt$$

$$= \int_{t_0}^{t_f} -\frac{1}{2} p^T M(q)^{-1} p - U(q) + p^T \dot{q} dt$$

- We can define the Hamiltonian:

$$H(q, p) = -\frac{1}{2} p^T M(q)^{-1} p + U(q)$$

Such that:

$$-S = \int_{t_0}^{t_f} H(q, p) - p^T \dot{q} dt$$

- In these variables, standard FON conditions are:

$$\frac{\partial}{\partial q(t)} \left[\int_{t_0}^{t_f} H(q, p) - p^T \dot{q} dt \right] = 0$$

$$\frac{\partial}{\partial p(t)} \left[\quad \quad \quad \right] = 0$$

$$\boxed{\begin{aligned} \frac{\partial H}{\partial q} + \dot{p} &= 0 \\ \frac{\partial H}{\partial p} - \dot{q} &= 0 \end{aligned}} \leftarrow \text{"Hamilton's Equations"}$$

- In optimization terms, Hamiltonian is "dual form" of least-action problem (Lagrangian is "primal form").
- The change of variables $v \rightarrow p$ is called "Legendre Transform" in classical mechanics.