

Last Time:

- Discrete Legendre Transform
- Variational Integrators w/ Constraints
- " " w/ External Forces

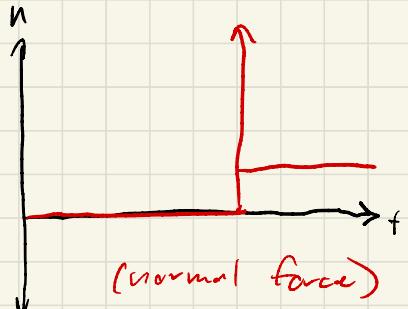
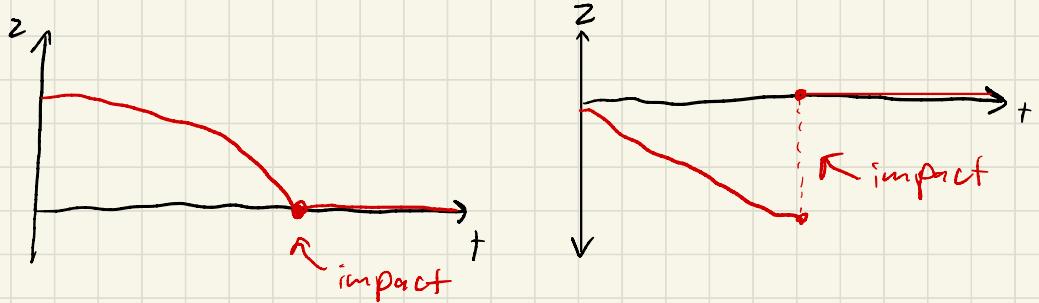
Today:

- Contact Dynamics
- Discrete Mechanics with Impacts

Contact Dynamics:

- Why is this hard?

- Imagine a falling particle that hits the ground (no friction):



- Impacts (and friction) cause velocity discontinuities.
For ideal rigid bodies this leads to infinite forces + accelerations.
- Strictly speaking, F_{norm} is not well defined in these scenarios
- In reality, things are not perfectly rigid, but practically this is still difficult.
- Makes modeling, simulation, control, learning hard.
- Three Broad Categories of simulation techniques:
 - 1) Smooth Contact Models
 - Surround all contact faces with e.g. non-linear spring-damper model
- Pros: Easy to implement, differentiable
- Cons: Not very accurate: interpenetration, creep, stiff ODEs
- MuJoCo does this

2) Hybrid / Event-Driven Methods

- Explicitly detect impact events while integrating ODEs. Perform discontinuous "jump" update, then proceed with smooth integration.
- Pros : Can use standard ODE tools, not stiff, captures qualitatively correct physics.
- Cons : Scales poorly with number of contacts, issues with simultaneous impacts, not differentiable.
- Common in control

3) Time-Stepping Methods

- Solve an optimization problem at each time step to compute contact forces that satisfy interpenetration constraints.
 - Pros : Scales well in number of contact points, No issues with simultaneous impacts, qualitatively correct physics.
 - Cons : Computationally expensive. Not differentiable.
 - Used by Bullet, Dart, Gazebo
- We're going to focus on 3)

Discrete Mechanics with Impacts:

- Review KKT conditions for equality constrained optimization problem:

$$\begin{array}{ll} \text{min}_x & f(x) \\ \text{s.t.} & C(x) \geq 0 \end{array} \quad \left\{ \Rightarrow L(x, \lambda) = f(x) - \lambda^T C(x) \right.$$

"Lagrangian" from optimization

- KKT Conditions:

$$\frac{\partial L}{\partial x} = \frac{\partial f}{\partial x} - \lambda^T \frac{\partial C}{\partial x} = 0 \quad (\text{stationarity})$$

"Complementarity condition"

$$\lambda \perp C(x) \quad \left\{ \begin{array}{l} C(x) \geq 0 \quad (\text{primal feasibility}) \\ \lambda \geq 0 \quad (\text{dual feasibility}) \\ \lambda^T C(x) = 0 \quad (\text{compl. slackness}) \end{array} \right.$$

↖ Hadamard product
(element-wise)

- Complementarity conditions define switching behavior for active vs. inactive constraints.
- Now let's apply this to the falling particle

- better define a signed-distance function that return distances between closest points:

$$\phi(q) = [0 \ 1] q$$

- Note: $\phi(q) = 0$ indicates contact
 $\phi(q) < 0$ indicates penetration (bad)

- Recall Lagrangian for a particle:

$$L(q, \dot{q}) = \frac{1}{2} m \dot{q}^T \dot{q} - m g [0 \ 1]^T q$$

$$L_d(q_u, q_{u+1}) = h L\left(\frac{q_{u+1} + q_u}{2}, \frac{q_{u+1} - q_u}{h}\right)$$

- Now we write down Least-Action s.t. interpenetration constraint:

$$\min_{q_{1:N}} \sum_{n=1}^{N-1} L_d(q_n, q_{n+1})$$

$$\text{s.t. } \phi(q_n) \geq 0 \quad \forall n$$

- Looks almost like joint constraints. KKT conditions:

$$D_2 L_d(q_{n-r}, q_n) + D_r L_d(q_n, q_{n+r}) + h n_n^T D_1 \phi(q_n) = 0$$

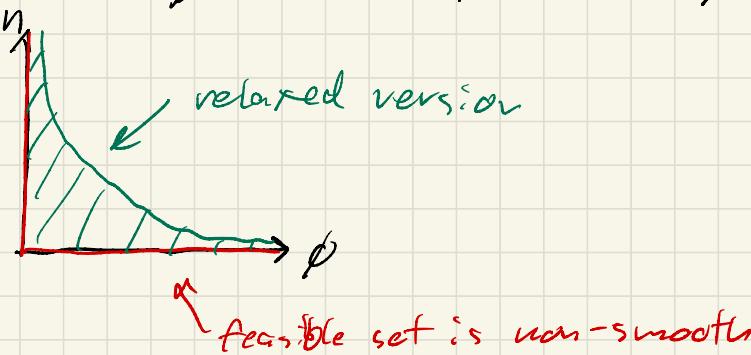
Stay above the ground $\phi(q_{n+r}) \geq 0$

Normal force can only push $n_n \geq 0$

No force off the ground $n_n \circ \phi(q_n) = 0$

* Quick notes on solving this:

- Feasible region for complementarity:



- Relax with a slack variable:

$$n \phi(q) = 0 \Rightarrow n \phi(q) \leq s, \quad s \geq 0$$