### Chapter 9

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Im	aportant Theorems:	
	<ul><li>9.1.7</li><li>9.2.1</li><li>9.2.3</li></ul>	
Im	portant ideas:	

- Compactness of  $\mathcal{L}(x,y)$ .
- Using broken trajectories to compactify
- Gluing

## **Background from Chapter 8**

- $(M,\omega)$  with  $\omega \in \Omega^2(M)$  is a symplectic manifold with an almost complex structure J.
- $H \in C^{\infty}(M; \mathbb{R})$  a Hamiltonian with  $X_H$  the corresponding vector field.
  - Recall, could be time-dependent. If so, assumed 1-periodic.
- x, y periodic orbits of H (nondegenerate, contractible)
- $u \in C^{\infty}(\mathbb{R} \times S^1; M)$  is a solution to the Floer equation.

• The Floer equation and its linearization:

$$\mathcal{F}(u) = \frac{\partial u}{\partial s} + J \frac{\partial u}{\partial t} + \text{grad } u(H) = 0$$
$$(d\mathcal{F})_u(Y) = \frac{\partial Y}{\partial s} + J_0 \frac{\partial Y}{\partial t} + S \cdot Y$$

$$Y \in u^*TW, S \in C^{\infty}(\mathbb{R} \times S^1; \operatorname{End}(\mathbb{R}^{2n})).$$

- $\mathcal{L}M$  is the free loop space of M, i.e. space of contractible loops on M, i.e.  $C^{\infty}(S^1; M)$  with the  $C^{\infty}$  topology
  - Loops in  $\mathcal{L}M$  can be viewed as maps  $S^2 \to M$ , since they're maps  $I \times S^1 \to M$  with the boundaries pinched:

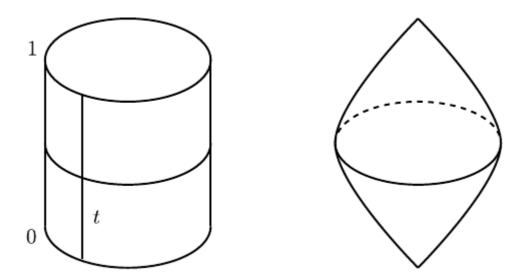


Figure 1: Loops in  $\mathcal{L}M$ 

- Elements  $x \in \mathcal{L}M$  can be viewed as maps  $S^1 \to M$ .
- Can extend to maps from a closed disc,  $u: \overline{\mathbb{D}}^2 \to M$ .
- The action functional is given by

$$\mathcal{A}_H: \mathcal{L}W \to \mathbb{R}$$
 
$$x \mapsto -\int_{\mathbb{D}} u^*\omega + \int_0^1 H_t(x(t)) \ dt$$
 
$$- \text{ Example: } W = \mathbb{R}^{2n} \implies A_H(x) = \int_0^1 \left(H_t \ dt - p \ dq\right).$$

• Assumption of symplectic asphericity, i.e. the symplectic form is zero on spheres. Statement: for every  $u \in C^{\infty}(S^2, M)$ ,

$$\int_{S^2} u^* \omega = 0 \quad \text{or equivalently} \quad \langle \omega, \ \pi_2 M \rangle = 0.$$

• Assumption of symplectic trivialization: for every  $u \in C^{\infty}(S^2; M)$  there exists a symplectic trivialization of the fiber bundle  $u^*TM$ , equivalently

$$\langle c_1 TM, \ \pi_2 M \rangle = 0.$$

• Maslov index 6.2.2

#### Results

Locally a product of base and fiber, transition functions are symplectomorphisms.

•  $(d\mathcal{F})_u$  is a Fredholm operator of index  $\mu(x) - \mu(y)$ .

### $\mathbf{2}$ | 9.1 and Review

- $\int_{S^2} u^* \omega = \sigma_1$  where  $u \in C^{\infty}(S^2, W)$ .
- $\langle c_1(TW), \pi_2(TW) \rangle = 0$ ?
- $C_k(H) := \mathbb{Z}/2\mathbb{Z}[S]$  where S is the set of periodic orbits of  $X_H$  of Maslov index k.
- x, y critical points of  $\mathcal{A}_H$  with  $\mathcal{M}(x, y)$  the moduli space of contractible solutions of finite energy connecting x, y.

#### 2.1 Review Last Time

- $\mathbb{R} \curvearrowright \mathcal{M}_{x,y}$ , so we quotient to define  $\mathcal{L}(x,y) := \mathcal{M}_{x,y}/\mathbb{R}$  with the quotient topology.
- Topology defined by when sequences converge:

$$\tilde{u}_n \stackrel{n \to \infty}{\to} \tilde{u} \iff \exists \{s_n\} \subseteq \mathbb{R} \text{ such that } u_n(s_n + s, \cdot) \stackrel{n \to \infty}{\to} u(s, \cdot).$$

#### Proposition 2.1(?).

 $\mathcal{L}(x,y)$  is Hausdorff.

- Want to show  $\mathcal{L}(x,y)$  is a compact 0-dimensional manifold.
- Have a differential

$$\partial: C_k(H) \longrightarrow C_{k-1}(H)$$

$$\partial(x) = \sum_{\text{Ind}(y)=k-1} n(x,y)y.$$

with n(x,y) the number (mod 2) of trajectories of grad  $\mathcal{A}_H$  connecting x,y, i.e solutions to the Floer equation.

• Want to prove that the following is a 1-dimensional manifold:

$$M := \overline{\mathcal{L}}(x, z) = \mathcal{L}(x, z) \cup_{\mu(y) = \mu(x) + 1} \mathcal{L}(x, y) \times \mathcal{L}(y, z).$$

and show that M is compact with  $\partial M$  equal to the last union.

- Last time: closure of space of trajectories connecting x, y contains "broken" trajectories.
- Last time: toward proving that M is compact

# $\mathbf{3}$ | 9.2

- Wanted to compactify  $\mathcal{L}(x,y)$ , needed to go to space of broken trajectories.
- Main theorem of chapter 9: 9.2.1.

#### Theorem 3.1(9.2.1).

Let (H, J) be a regular pair with H nondegenerate.

Let x, z be two periodic trajectories of H such that  $\mu(x) = \mu(z) + 2$ .

Then  $\overline{\mathcal{L}}(x,y)$  is a compact 1-manifold with boundary satisfying

$$\partial \overline{\mathcal{L}}(x,y) = \bigcup_{\mu(x) < \mu(y) < \mu(z)} \mathcal{L}(x,y) \times \mathcal{L}(y,z).$$

As a corollary,  $\partial^2 = 0$ .

- Know  $\overline{\mathcal{L}}(x,y)$  is compact and  $\mathcal{L}(x,y)$  is a 1-manifold
- Now suffices to study in a neighborhood of boundary points ("gluing theorem")

#### 3.1 Three steps to gluing theorem

- 1. Pre-gluing: Get a function  $w_p$  which interpolates between u and v (not exactly a solution itself, but will be approximated by one later).
- 2. Constructing  $\psi$  a "true solution" from  $w_p$  using the Newton-Picard method. We'll have

$$\psi(p) = \exp_{w_p}(\gamma(p)) \qquad \gamma(p) \in W^{1,p}(w_p^* TW) = T_{w_p} \mathcal{P}(x, z).$$

where  $\mathcal{P} = ?$ .

- 3. Get a lift  $\hat{\psi} = \pi \circ \psi$  where  $\pi = ?$  satisfying
- $\widehat{\psi}(p) \stackrel{n \to \infty}{\to} (\widehat{u}, \widehat{v})$
- $\widehat{\varphi}$  is an embedding
- $\widehat{\psi}$  is unique in the following sense (the last point)

#### Theorem 3.2(9.2.3 (Gluing Theorem)).

Let x, y, z be critical points of the action functional  $\mathcal{A}_H$  such that  $\mu(x) = \mu(y) + 1 = \mu(z) + 2$ .

Let  $(u, v) \in \mathcal{M}(x, y) \times \mathcal{M}(y, z)$  be trajectories, inducing  $(\bar{u}, \bar{v}) \in \mathcal{L}(x, y) \times \mathcal{L}(y, z)$ .

- There exist a differentiable map  $\psi:(\rho_0,\infty)\to\mathcal{M}(x,z)$  for some  $\rho>0$  such that

- π ∘ ψ : (ρ<sub>0</sub>, ∞) → L(x, z) is an embedding
  ψ̂ <sup>ρ→∞</sup> (ū, v̄) ∈ L(x, z).
  If ℓ<sub>n</sub> ∈ L(x, z) with ℓ<sub>n</sub> <sup>n→∞</sup> (ū, v̄), then for n ≫ 1 we have ℓ ∈ ℑ(ψ̂).

## 9.3: Pre-gluing

- Choose a bump function  $\beta$  on  $\{0\}^c \subset \mathbb{R} \to [0,1]$  which is 1 on  $|x| \geq 1$  and 0 on  $|x| < \varepsilon$
- Split into positive and negative parts  $\beta^{\pm}$ :

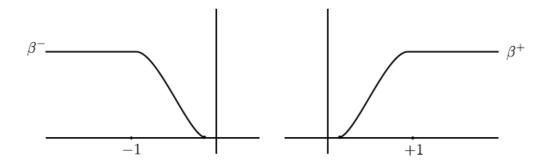


Figure 2: Figure 9.3

• Define the interpolation  $w_{\rho}$  from u to v in the following way:

$$w_{\rho}(s,t) = \begin{cases} u(s+\rho,t) & \text{if } s \leq -1\\ \exp_{y(t)} \left(\beta^{-}(s) \exp_{y(t)}^{-1}(u(s+\rho,t)) + \beta^{+}(s) \exp_{y(t)}^{-1}(v(s-\rho,t))\right) & \text{if } s \leq -1\\ v(s-\rho,t) & \text{if } s \geq 1 \end{cases}$$

• Why does this make sense?

$$|s| \le 1 \implies u(s \pm \rho, t) \in \left\{ \exp_{y(t)} Y(t) \mid \sup_{t \in S^1} ||Y(t)|| \le r_0 \right\}.$$

## **9.4**: Construction of $\psi$ .

- Have constructed  $w_{\rho} \in C_{\sim}^{\infty}(x,z)C^{\infty}(x,z)$  for every  $\rho \geq \rho_0$ , since there is exponential decay.
- Yields  $\psi_{\rho} \in \mathcal{M}(x,z)$  a true solution (to be defined).

• Need to check that  $\mathcal{F}(\psi_{\rho}) = 0$  where

$$\mathcal{F} = \frac{\partial}{\partial s} + J \frac{\partial}{\partial t} + \text{grad } Hx$$

in the weak sense.

- $\psi_{\rho}$  already continuous, and by elliptic regularity, makes it a strong solution.
- Trivialization
- Defining  $\mathcal{F}_{\rho}$ .

$$W^{1,p}\left(\mathbf{R}\times S^{1};\mathbf{R}^{2n}\right) \xrightarrow{\mathcal{F}_{\rho}} L^{p}\left(\mathbf{R}\times S^{1};\mathbf{R}^{2n}\right)$$
$$(y_{1},\ldots,y_{2n})\longmapsto \left[\left(\frac{\partial}{\partial s}+J\frac{\partial}{\partial t}+\operatorname{grad}H_{t}\right)\left(\exp_{w_{\rho}}\sum y_{i}Z_{i}^{\rho}\right)\right]_{Z_{i}}$$

where  $\mathcal{F}_{\rho} := \mathcal{F} \circ \exp_{w_{\rho}}$  written in the bases  $Z_i$ . sd - Newton-Picard method, general idea

• Original method and variant: find the limit of a sequence

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}, \qquad x_{n+1} = x_n - \frac{f(x_n)}{f'(x_0)}.$$

- Allows finding zeros of f given an approximate zero  $x_0$ .
- Linearize  $\mathcal{F}_{\rho}$ .