# Chapter 9

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Im	nportant ideas:	
	• Compactness of $\mathcal{L}(x,y)$ .	

- $\partial^2 = 0$
- Using broken trajectories to compactify
- Gluing

# **1** | Background from Chapter 8

- $(M, \omega)$  with  $\omega \in \Omega^2(M)$  is a symplectic manifold with an almost complex structure J.
- $H \in C^{\infty}(M;\mathbb{R})$  a Hamiltonian with  $X_H$  the corresponding symplectic gradient.
  - Defined by how it acts on tangent vectors in  $T_xM$ :

$$\omega_x(\cdot, X_H(x)) = (dH)_x(\cdot).$$

– Zeros of vector field  $X_H$  correspond to critical points of H:

$$X_H(x) = 0 \iff (dH)_x = 0.$$

- Take the associated flow  $\psi^t: M \to M$ , assumed 1-periodic so  $\psi^1(x) = x$ : critical points of H are periodic trajectories.
- $u \in C^{\infty}(\mathbb{R} \times S^1; M)$  is a solution to the Floer equation.
- The Floer equation and its linearization:

$$\mathcal{F}(u) = \frac{\partial u}{\partial s} + J \frac{\partial u}{\partial t} + \operatorname{grad}_{u}(H) = 0$$
$$(d\mathcal{F})_{u}(Y) = \frac{\partial Y}{\partial s} + J_{0} \frac{\partial Y}{\partial t} + S \cdot Y$$

$$Y \in u^*TW, \ S \in C^{\infty}(\mathbb{R} \times S^1; \operatorname{End}(\mathbb{R}^{2n})).$$

- $\mathcal{L}M$  is the free loop space of M, i.e. space of contractible loops on M, i.e.  $C^{\infty}(S^1; M)$  with the  $C^{\infty}$  topology
  - Loops in  $\mathcal{L}M$  can be viewed as maps  $S^2 \to M$ , since they're maps  $I \times S^1 \to M$  with the boundaries pinched:



Figure 1: Loops in  $\mathcal{L}M$ 

- Elements  $x \in \mathcal{L}M$  can be viewed as maps  $S^1 \to M$ .
- Can extend to maps from a closed disc,  $u: \overline{\mathbb{D}}^2 \to M$ .
- The action functional is given by

$$\mathcal{A}_H : \mathcal{L}W \to \mathbb{R}$$
 
$$x \mapsto -\int_{\mathbb{D}} u^* \omega + \int_0^1 H_t(x(t)) dt$$

- Example:  $W = \mathbb{R}^{2n} \implies A_H(x) = \int_0^1 (H_t dt p dq).$
- Correspondence between trajectories of the gradient of  $\mathcal{A}_H$  and solutions to Floer equations.
- Assumption of symplectic asphericity, i.e. the symplectic form is zero on spheres. Statement: for every  $u \in C^{\infty}(S^2, M)$ ,

$$\int_{S^2} u^* \omega = 0 \quad \text{or equivalently} \quad \langle \omega, \ \pi_2 M \rangle = 0.$$

• Assumption of symplectic trivialization: for every  $u \in C^{\infty}(S^2; M)$  there exists a symplectic trivialization of the fiber bundle  $u^*TM$ , equivalently

$$\langle c_1 TM, \ \pi_2 M \rangle = 0.$$

Locally a product of base and fiber, transition functions are symplectomorphisms.

- x, y periodic orbits of H (nondegenerate, contractible), equivalently critical points of  $\mathcal{A}_H$ .
- Maslov index: used the fact that
  - $\operatorname{Sp}(2n,\mathbb{R})$  retracts onto U(n): use a polar decomposition S=PQ as a PSD times orthogonal, then homotope P to I.
  - $-\pi_1 U_n = \mathbb{Z}$ : use  $U(n,\mathbb{C}) \simeq SU(n,\mathbb{C}) \times S^1$  by the determinant, and  $\pi_1 SU(n,\mathbb{C}) = 0$ .
  - Thus every path in  $\gamma: I \to \operatorname{Sp}(2n,\mathbb{R})$  can be assigned an integer by getting a map  $\tilde{\gamma}: I \to S^1$  and taking (approximately) its winding number.
- $\mathcal{M}(x,y)$ , the moduli space of contractible finite-energy solutions to the Floer equation connecting x,y.
  - Showed that after perturbing H to get transversality, get a manifold of dimension  $\mu(x) \mu(y)$ .
  - How did we do it: describe as zeros of a section of a vector bundle over  $\mathcal{P}^{1,p}(x,y)$  (Banach manifold modeled on the Sobolev spaces  $W^{1,p}$ ), apply Sard-Smale to show  $\mathcal{M}(x,y)$  is the inverse image of a regular value of some map.
  - Needed tangent maps to be Fredholm operators, proved in Ch. 8 and used to show transversality. Followed from showing  $(d\mathcal{F})_u$  is a Fredholm operator of index  $\mu(x) \mu(y)$ .

Goals

• Construct Floer homology and prove the Arnold Conjecture ("Symplectic Morse Inequalities?"):

# {1-Periodic trajectories of 
$$X_H$$
}  $\geq \sum_{k \in \mathbb{Z}} HM_k(w; \mathbb{Z}/2\mathbb{Z}).$ 

Steps

- 1. Define the action functional  $A_H$ .
- 2. Construct the chain complex (graded vector space)  $CF_*$ .
- 3. Define  $X_H$ , which will be used to define  $\partial$  later.

- 4. Count trajectories.
- 5. Show finite-energy trajectories connect critical points of  $\mathcal{A}_H$ .
- 6. Show compactness property for space of trajectories of finite energy.
- 7. Define  $\partial$  (uses a compactness property in 9.1c)
- 8. Show space of trajectories is a manifold (plus genericity, "Smale property")
- 9. Show that  $\partial^2 = 0$ .
- 10. Show that  $HF_*$  doesn't depend on  $A_H$  or  $X_H$
- 11. Show  $HF_* \cong HM_*$ , and compare dimensions of the vector spaces  $CM_*$  and  $CF_*$ .

# $\mathbf{2}$ | 9.1 and Review

• Defined moduli space of (parameterized) solutions:

 $\mathcal{M}(x,y) = \{\text{Contractible finite-energy solutions connecting } x,y\}$ 

 $\mathcal{M} = \{\text{All contractible finite-energy solutions to the Floer equation}\} = \bigcup_{x,y} \mathcal{M}(x,y).$ 

• Defined the moduli space of (unparameterized) **trajectories** connecting x to y:

$$\mathcal{L}(x,y) := \mathcal{M}(x,y)/\mathbb{R}.$$

- Use the quotient topology, define sequentially:

$$\tilde{u}_n \stackrel{n \to \infty}{\longrightarrow} \tilde{u} \iff \exists \{s_n\} \subset \mathbb{R} \text{ such that } u_n(s_n + s, \cdot) \stackrel{n \to \infty}{\longrightarrow} u(s, \cdot).$$

- When  $|\mu(x) - \mu(y)| = 1$ , get a compact 0-manifold, so the number of trajectories

$$n(x,y) \coloneqq \#\mathcal{L}(x,y)$$

is well-defined.

- $C_k(H) := \mathbb{Z}/2\mathbb{Z}[S]$  where S is the set of periodic orbits of  $X_H$  of Maslov index k.
  - Finitely many since they are nondegeneracy implies they are isolated.
- Defined a differential

$$\partial: C_k(H) \to C_{k-1}(H)$$
$$x \mapsto \sum_{\mu(y)=k-1} n(x,y)y$$

 $n(x,y) := \# \{ \text{Trajectories of grad } \mathcal{A}_H \text{ connecting } x,y \} \mod 2$ =  $\# \mathcal{L}(x,y) \mod 2$ . • Examined  $\partial^2$ :

$$\partial^{2}: C_{k}(H) \to C_{k-2}(H)$$

$$x \mapsto \partial(\partial(x))$$

$$= \partial \left(\sum_{\mu(y)=\mu(x)-1} n(x,y)y\right)$$

$$= \sum_{\mu(y)=\mu(x)-1} n(x,y) \partial(y)$$

$$= \sum_{\mu(y)=\mu(x)-1} \sum_{\mu(z)=\mu(y)-1} n(y,z)z$$

$$= \sum_{\mu(y)=\mu(x)-1} \sum_{\mu(z)=\mu(y)-1} n(x,y)n(y,z)z$$

$$= \sum_{\mu(z)=\mu(y)-1} \left(\sum_{\mu(y)=\mu(x)-1} n(x,y)n(y,z)\right)z \qquad \text{(finite sums, swap order),}$$

so it suffices to show

$$\sum_{\mu(y)=\mu(x)-1} n(x,y) n(y,z) = 0 \text{ when } \mu(z) = \mu(x) - 2.$$

Easier to examine parity, so we'll show it's zero mod 2.

- When  $\mu(z) = \mu(x) 2$ ,  $\mathcal{L}(x, z)$  is a non-compact 1-manifold, so we compactify by adding in broken trajectories to get  $\overline{\mathcal{L}}(x, y)$ .
- We'll then have

$$\overline{\mathcal{L}}(x,z) = \mathcal{L}(x,z) \cup \partial \overline{\mathcal{L}}(x,z), \qquad \partial \overline{\mathcal{L}}(x,z) = \bigcup_{\mu(y) = \mu(x) - 1} \mathcal{L}(x,y) \times \mathcal{L}(y,z),$$

which "space-ifies" the equation we want.

• We'll show  $\partial \overline{\mathcal{L}}(x,z)$  is a 1-manifold, which must have an even number of points, and thus

$$\sum_{\mu(y)=\mu(x)-1} n(x,y) n(y,z) = \#\Big(\partial \overline{\mathcal{L}}(x,z)\Big) \equiv 0 \mod 2.$$

#### 2.1 Three Important Theorems

• Shown last time: a sequence of trajectories can converge to a broken trajectory.

### Theorem 2.1(9.1.7).

Let  $\{u_n\}$  be a sequence in  $\mathcal{M}(x,y)$ , then there exist

- A subsequence  $\{u_{n_j}\}$
- Critical points  $\{x_0, x_1, \dots, x_{\ell+1}\}$  with  $x_0 = x$  and  $x_{\ell+1} = y$
- Sequences  $\left\{s_n^1\right\}, \left\{s_n^2\right\}, \cdots, \left\{s_n^{\ell}\right\}.$
- Elements  $u^k \in \mathcal{M}(x_k, x_{k+1})$  such that for every  $0 \le k \le \ell$ ,

$$u_n \cdot s_n^k \stackrel{n \to \infty}{\longrightarrow} u^k$$
.

### **Definition 2.1.1** (Regular Pair).

For an almost complex structure J and a Hamiltonian H, the pair (H, J) is **regular** if the Floer map  $\mathcal{F}$  is transverse to the zero section in the following vector bundle:

#### Theorem 2.2(9.2.1).

Let (H, J) be a regular pair with H nondegenerate and x, z be two periodic trajectories of H such that

$$\mu(x) = \mu(z) + 2.$$

Then  $\overline{\mathcal{L}}(x,z)$  is a compact 1-manifold with boundary with

$$\partial \overline{\mathcal{L}}(x,z) = \bigcup_{y \in \mathcal{I}(x,z)} \mathcal{L}(x,y) \times \mathcal{L}(y,z) \quad \text{where} \quad \mathcal{I}(x,z) = \left\{ y \mid \mu(x) < \mu(y) < \mu(z) \right\}.$$

Note: possibly a typo in the book?

#### Remark 1.

- As a corollary,  $\partial^2 = 0$ .
- Most of chapter 9 is spent proving this theorem.

#### Remark 2.

Some notation:

$$\mathbb{R} \longrightarrow \mathcal{M}(x,z)$$

$$\downarrow^{\pi}$$

$$\mathcal{L}(x,z)$$

Hats will generally denote maps induced on quotient.

#### Theorem 2.3(9.2.3: Gluing).

Let x, y, z be three critical points of  $\mathcal{A}_H$  with three consecutive indices

$$\mu(x) = \mu(y) + 1 = \mu(z) + 2.$$

and let

$$(u,v) \in \mathcal{M}(x,y) \times \mathcal{M}(y,z) \quad \leadsto \quad (\widehat{u},\widehat{v}) \in \mathcal{L}(x,y) \times \mathcal{L}(y,z).$$

Then

1. There exists a  $\rho_0 > 0$  and a differentiable map

$$\psi: [\rho_0, \infty) \to \mathcal{M}(x, z)$$

such that  $\widehat{\psi}$ , the induced map on the quotient

$$[\rho_0, \infty) \xrightarrow{\psi} \mathcal{M}(x, z)$$

$$\widehat{\psi} \qquad \downarrow^{\pi}$$

$$\mathcal{L}(x, z)$$

is an embedding that satisfies

$$\widehat{\psi}(\rho) \stackrel{\rho \to \infty}{\longrightarrow} (\widehat{u}, \widehat{v}) \in \overline{\mathcal{L}}(x, z).$$

2. For any sequence  $\{\ell_n\} \subseteq \mathcal{L}(x,z)$ ,

$$\ell_n \stackrel{n \to \infty}{\longrightarrow} (\widehat{u}, \widehat{v}) \implies \ell_n \in \operatorname{im}(\widehat{\psi}) \text{ for } n \gg 0.$$

#### 2.2 Review Last Time

- $\mathbb{R} \curvearrowright \mathcal{M}_{x,y}$ , so we quotient to define  $\mathcal{L}(x,y) := \mathcal{M}_{x,y}/\mathbb{R}$  with the quotient topology.
- Topology defined by when sequences converge:

$$\tilde{u}_n \stackrel{n \to \infty}{\to} \tilde{u} \iff \exists \{s_n\} \subseteq \mathbb{R} \text{ such that } u_n(s_n + s, \cdot) \stackrel{n \to \infty}{\to} u(s, \cdot).$$

### Proposition 2.4(?).

 $\mathcal{L}(x,y)$  is Hausdorff.

- Want to show  $\mathcal{L}(x,y)$  is a compact 0-dimensional manifold.
- Have a differential

$$\partial: C_k(H) \longrightarrow C_{k-1}(H)$$

$$\partial(x) = \sum_{\text{Ind}(y)=k-1} n(x,y)y.$$

with n(x,y) the number (mod 2) of trajectories of grad  $\mathcal{A}_H$  connecting x,y, i.e solutions to the Floer equation.

• Want to prove that the following is a 1-dimensional manifold:

$$M := \overline{\mathcal{L}}(x, z) = \mathcal{L}(x, z) \cup_{\mu(y) = \mu(x) + 1} \mathcal{L}(x, y) \times \mathcal{L}(y, z).$$

and show that M is compact with  $\partial M$  equal to the last union.

- Last time: closure of space of trajectories connecting x, y contains "broken" trajectories.
- Last time: toward proving that M is compact

# $3 \mid 9.2$

- Wanted to compactify  $\mathcal{L}(x,y)$ , needed to go to space of broken trajectories.
- Main theorem of chapter 9: 9.2.1.

### Theorem 3.1(9.2.1).

Let (H, J) be a regular pair with H nondegenerate.

Let x, z be two periodic trajectories of H such that  $\mu(x) = \mu(z) + 2$ .

Then  $\overline{\mathcal{L}}(x,y)$  is a compact 1-manifold with boundary satisfying

$$\partial \overline{\mathcal{L}}(x,y) = \bigcup_{\mu(x) < \mu(y) < \mu(z)} \mathcal{L}(x,y) \times \mathcal{L}(y,z).$$

As a corollary,  $\partial^2 = 0$ .

- Know  $\overline{\mathcal{L}}(x,y)$  is compact and  $\mathcal{L}(x,y)$  is a 1-manifold
- Now suffices to study in a neighborhood of boundary points ("gluing theorem")

### 3.1 Three steps to gluing theorem

- 1. Pre-gluing: Get a function  $w_p$  which interpolates between u and v (not exactly a solution itself, but will be approximated by one later).
- 2. Constructing  $\psi$  a "true solution" from  $w_p$  using the Newton-Picard method. We'll have

$$\psi(p) = \exp_{w_p}(\gamma(p)) \qquad \gamma(p) \in W^{1,p}(w_p^* TW) = T_{w_p} \mathcal{P}(x, z).$$

where  $\mathcal{P} = ?$ .

- 3. Get a lift  $\hat{\psi} = \pi \circ \psi$  where  $\pi = ?$  satisfying
- $\widehat{\psi}(p) \stackrel{n \to \infty}{\to} (\widehat{u}, \widehat{v})$
- $\widehat{\varphi}$  is an embedding
- $\hat{\psi}$  is unique in the following sense (the last point)

#### Theorem 3.2(9.2.3 (Gluing Theorem)).

Let x, y, z be critical points of the action functional  $\mathcal{A}_H$  such that  $\mu(x) = \mu(y) + 1 = \mu(z) + 2$ .

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Let  $(u, v) \in \mathcal{M}(x, y) \times \mathcal{M}(y, z)$  be trajectories, inducing  $(\bar{u}, \bar{v}) \in \mathcal{L}(x, y) \times \mathcal{L}(y, z)$ .

- There exist a differentiable map  $\psi:(\rho_0,\infty)\to\mathcal{M}(x,z)$  for some  $\rho>0$  such that

- π ∘ ψ : (ρ<sub>0</sub>, ∞) → L(x, z) is an embedding
  ψ̂ <sup>ρ→∞</sup> (ū, v̄) ∈ L(x, z).
  If ℓ<sub>n</sub> ∈ L(x, z) with ℓ<sub>n</sub> <sup>n→∞</sup> (ū, v̄), then for n ≫ 1 we have ℓ ∈ ℑ(ψ̂).

# 9.3: Pre-gluing

- Choose a bump function  $\beta$  on  $\{0\}^c \subset \mathbb{R} \to [0,1]$  which is 1 on  $|x| \geq 1$  and 0 on  $|x| < \varepsilon$
- Split into positive and negative parts  $\beta^{\pm}$ :

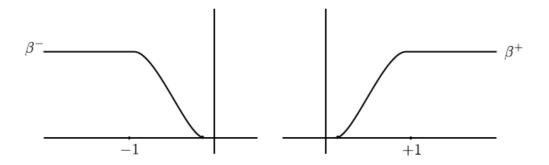


Figure 2: Figure 9.3

• Define the interpolation  $w_{\rho}$  from u to v in the following way:

$$w_{\rho}(s,t) = \begin{cases} u(s+\rho,t) & \text{if } s \leq -1\\ \exp_{y(t)} \left(\beta^{-}(s) \exp_{y(t)}^{-1} (u(s+\rho,t)) + \beta^{+}(s) \exp_{y(t)}^{-1} (v(s-\rho,t))\right) & \text{if } s \in [-1,1]\\ v(s-\rho,t) & \text{if } s \geq 1 \end{cases}$$

• Why does this make sense?

$$|s| \le 1 \implies u(s \pm \rho, t) \in \left\{ \exp_{y(t)} Y(t) \mid \sup_{t \in S^1} ||Y(t)|| \le r_0 \right\}.$$

# **9.4**: Construction of $\psi$ .

- Have constructed  $w_{\rho} \in C_{\sim}^{\infty}(x,z)C^{\infty}(x,z)$  for every  $\rho \geq \rho_0$ , since there is exponential decay.
- Yields  $\psi_{\rho} \in \mathcal{M}(x,z)$  a true solution (to be defined).

• Need to check that  $\mathcal{F}(\psi_{\rho}) = 0$  where

$$\mathcal{F} = \frac{\partial}{\partial s} + J \frac{\partial}{\partial t} + \text{grad } Hx$$

in the weak sense.

- $\psi_{\rho}$  already continuous, and by elliptic regularity, makes it a strong solution.
- Trivialization
- Defining  $\mathcal{F}_{\rho}$ .

$$W^{1,p}\left(\mathbf{R}\times S^{1};\mathbf{R}^{2n}\right) \xrightarrow{\mathcal{F}_{\rho}} L^{p}\left(\mathbf{R}\times S^{1};\mathbf{R}^{2n}\right)$$
$$(y_{1},\ldots,y_{2n})\longmapsto \left[\left(\frac{\partial}{\partial s}+J\frac{\partial}{\partial t}+\operatorname{grad}H_{t}\right)\left(\exp_{w_{\rho}}\sum y_{i}Z_{i}^{\rho}\right)\right]_{Z_{i}}$$

where  $\mathcal{F}_{\rho} := \mathcal{F} \circ \exp_{w_{\rho}}$  written in the bases  $Z_i$ . sd - Newton-Picard method, general idea

• Original method and variant: find the limit of a sequence

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}, \qquad x_{n+1} = x_n - \frac{f(x_n)}{f'(x_0)}.$$

- Allows finding zeros of f given an approximate zero  $x_0$ .
- Linearize  $\mathcal{F}_{\rho}$ .