

# Chapter 9

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Sunday 20<sup>th</sup> September, 2020

## Contents

<b>1</b>	<b>Background from Chapter 8</b>	<b>1</b>
<b>2</b>	<b>9.1 and Review</b>	<b>3</b>
2.1	Review Last Time . . . . .	3
<b>3</b>	<b>9.2</b>	<b>4</b>
3.1	Three steps to gluing theorem . . . . .	4
<b>4</b>	<b>9.3: Pre-gluing</b>	<b>4</b>
<b>5</b>	<b>9.4: Construction of <math>\psi</math>.</b>	<b>5</b>

Important Theorems:

- 9.1.7
- 9.2.1
- 9.2.3

Important ideas:

- Compactness of  $\mathcal{L}(x, y)$ .
- $\partial^2 = 0$ .
- Using broken trajectories to compactify
- Gluing

## 1 | Background from Chapter 8

- $(M, \omega)$  with  $\omega \in \Omega^2(M)$  is a symplectic manifold
- $H \in C^\infty(M; \mathbb{R})$  a Hamiltonian with  $X_H$  the corresponding vector field.
  - Recall, could be time-dependent. If so, assumed 1-periodic.

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- The Floer equation and its linearization:

$$\begin{aligned}\mathcal{F}(u) &= \frac{\partial u}{\partial s} + J \frac{\partial u}{\partial t} + \text{grad}_u(H) = 0 \\ (d\mathcal{F})_u(Y) &= \frac{\partial Y}{\partial s} + J_0 \frac{\partial Y}{\partial t} + S \cdot Y\end{aligned}$$

$$Y \in u^*TW, \quad S \in C^\infty(\mathbb{R} \times S^1; \text{End}(\mathbb{R}^{2n})).$$

- $\mathcal{LM}$  is the *free loop space* of  $M$ , i.e. space of contractible loops on  $M$ , i.e.  $C^\infty(S^1; M)$  with the  $C^\infty$  topology
  - Loops in  $\mathcal{LM}$  can be viewed as maps  $S^2 \rightarrow M$ , since they're maps  $I \times S^1 \rightarrow M$  with the boundaries pinched:

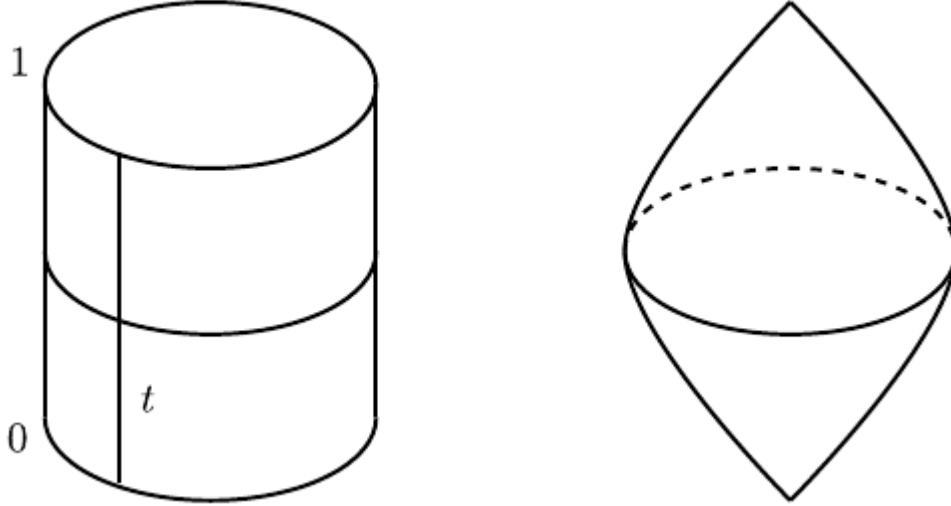


Figure 1: Loops in  $\mathcal{LM}$

- Elements  $x \in \mathcal{LM}$  can be viewed as maps  $S^1 \rightarrow M$ .
- Can extend to maps from a closed disc,  $u : \mathbb{D}^2 \rightarrow M$ .
- The action functional is given by

$$\begin{aligned}\mathcal{A}_H : \mathcal{LW} &\rightarrow \mathbb{R} \\ x &\mapsto - \int_{\mathbb{D}} u^* \omega + \int_0^1 H_t(x(t)) \, dt\end{aligned}$$

$u : \mathbb{D} \rightarrow W$  is an extension of  $x : S^1 \rightarrow W$  to the disc with  $u(\exp(it)) = x(t)$ .

- Example:  $W = \mathbb{R}^{2n} \implies \mathcal{A}_H(x) = \int_0^1 (H_t \, dt - p \, dq)$ .

- $u \in C^\infty(\mathbb{R} \times S^1; W)$  is a solution to the Floer equation.
- $(d\mathcal{F})_u$  is a Fredholm operator of index  $\mu(x) - \mu(y)$ .

## 2 | 9.1 and Review

- $\int_{S^2} u^* \omega = \sigma_1$  where  $u \in C^\infty(S^2, W)$ .
- $\langle c_1(TW), \pi_2(TW) \rangle = 0$ ?
- $C_k(H) := \mathbb{Z}/2\mathbb{Z}[S]$  where  $S$  is the set of periodic orbits of  $X_H$  of Maslov index  $k$ .
- $x, y$  critical points of  $\mathcal{A}_H$  with  $\mathcal{M}(x, y)$  the moduli space of contractible solutions of finite energy connecting  $x, y$ .

### 2.1 Review Last Time

- $\mathbb{R} \curvearrowright \mathcal{M}_{x,y}$ , so we quotient to define  $\mathcal{L}(x, y) := \mathcal{M}_{x,y}/\mathbb{R}$  with the quotient topology.
- Topology defined by when sequences converge:

$$\tilde{u}_n \xrightarrow{n \rightarrow \infty} \tilde{u} \iff \exists \{s_n\} \subseteq \mathbb{R} \text{ such that } u_n(s_n + s, \cdot) \xrightarrow{n \rightarrow \infty} u(s, \cdot).$$

**Proposition 2.1 (?)**.  
 $\mathcal{L}(x, y)$  is Hausdorff.

- Want to show  $\mathcal{L}(x, y)$  is a compact 0-dimensional manifold.
- Have a differential

$$\begin{aligned} \partial : C_k(H) &\longrightarrow C_{k-1}(H) \\ \partial(x) &= \sum_{\text{Ind}(y)=k-1} n(x, y)y. \end{aligned}$$

with  $n(x, y)$  the number (mod 2) of trajectories of  $\text{grad } \mathcal{A}_H$  connecting  $x, y$ , i.e solutions to the Floer equation.

- Want to prove that the following is a 1-dimensional manifold:

$$M := \overline{\mathcal{L}}(x, z) = \mathcal{L}(x, z) \cup_{\mu(y)=\mu(x)+1} \mathcal{L}(x, y) \times \mathcal{L}(y, z).$$

and show that  $M$  is compact with  $\partial M$  equal to the last union.

- Last time: closure of space of trajectories connecting  $x, y$  contains “broken” trajectories.
- Last time: toward proving that  $M$  is compact

## 3 | 9.2

- Wanted to compactify  $\mathcal{L}(x, y)$ , needed to go to space of broken trajectories.
- Main theorem of chapter 9: 9.2.1.

### Theorem 3.1(9.2.1).

Let  $(H, J)$  be a regular pair with  $H$  nondegenerate.

Let  $x, z$  be two periodic trajectories of  $H$  such that  $\mu(x) = \mu(z) + 2$ .

Then  $\bar{\mathcal{L}}(x, y)$  is a compact 1-manifold with boundary satisfying

$$\partial \bar{\mathcal{L}}(x, y) = \bigcup_{\mu(x) < \mu(y) < \mu(z)} \mathcal{L}(x, y) \times \mathcal{L}(y, z).$$

As a corollary,  $\partial^2 = 0$ .

- Know  $\bar{\mathcal{L}}(x, y)$  is compact and  $\mathcal{L}(x, y)$  is a 1-manifold
- Now suffices to study in a neighborhood of boundary points (“gluing theorem”)

### 3.1 Three steps to gluing theorem

1. Pre-gluing: Get a function  $w_p$  which interpolates between  $u$  and  $v$  (not exactly a solution itself, but will be approximated by one later).
2. Constructing  $\psi$  a “true solution” from  $w_p$  using the Newton-Picard method. We’ll have

$$\psi(p) = \exp_{w_p}(\gamma(p)) \quad \gamma(p) \in W^{1,p}(w_p^* TW) = T_{w_p} \mathcal{P}(x, z).$$

where  $\mathcal{P} = ?$ .

3. Get a lift  $\hat{\psi} = \pi \circ \psi$  where  $\pi = ?$  satisfying
  - $\hat{\psi}(p) \xrightarrow{n \rightarrow \infty} (\hat{u}, \hat{v})$
  - $\hat{\varphi}$  is an embedding
  - $\hat{\psi}$  is unique in the following sense (the last point)

### Theorem 3.2(9.2.3 (Gluing Theorem)).

Let  $x, y, z$  be critical points of the action functional  $\mathcal{A}_H$  such that  $\mu(x) = \mu(y) + 1 = \mu(z) + 2$ .

Let  $(u, v) \in \mathcal{M}(x, y) \times \mathcal{M}(y, z)$  be trajectories, inducing  $(\bar{u}, \bar{v}) \in \mathcal{L}(x, y) \times \mathcal{L}(y, z)$ .

- There exist a differentiable map  $\psi : (\rho_0, \infty) \rightarrow \mathcal{M}(x, z)$  for some  $\rho > 0$  such that
- $\pi \circ \psi : (\rho_0, \infty) \rightarrow \mathcal{L}(x, z)$  is an embedding
- $\hat{\psi} \xrightarrow{\rho \rightarrow \infty} (\bar{u}, \bar{v}) \in \bar{\mathcal{L}}(x, z)$ .
- If  $\ell_n \in \mathcal{L}(x, z)$  with  $\ell_n \xrightarrow{n \rightarrow \infty} (\bar{u}, \bar{v})$ , then for  $n \gg 1$  we have  $\ell \in \mathfrak{I}(\hat{\psi})$ .

## 4 | 9.3: Pre-gluing

- Choose a bump function  $\beta$  on  $\{0\}^c \subset \mathbb{R} \rightarrow [0, 1]$  which is 1 on  $|x| \geq 1$  and 0 on  $|x| < \varepsilon$
- Split into positive and negative parts  $\beta^\pm$ :

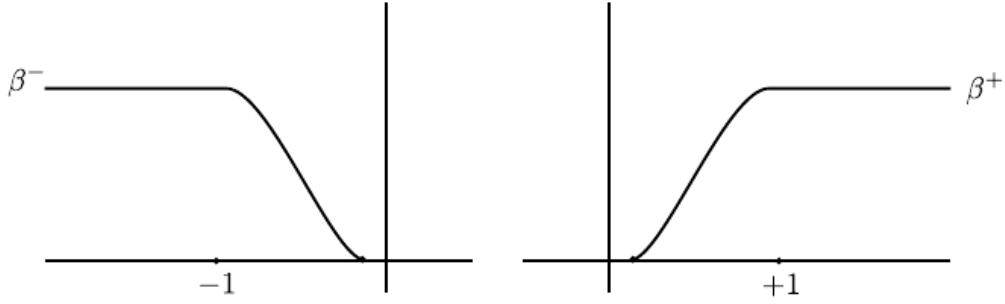


Figure 2: Figure 9.3

- Define the interpolation  $w_\rho$  from  $u$  to  $v$  in the following way:

$$w_\rho(s, t) = \begin{cases} u(s + \rho, t) & \text{if } s \leq -1 \\ \exp_{y(t)} \left( \beta^-(s) \exp_{y(t)}^{-1}(u(s + \rho, t)) + \beta^+(s) \exp_{y(t)}^{-1}(v(s - \rho, t)) \right) & \text{if } s \in [-1, 1] \\ v(s - \rho, t) & \text{if } s \geq 1 \end{cases}$$

- Why does this make sense?

$$|s| \leq 1 \implies u(s \pm \rho, t) \in \left\{ \exp_{y(t)} Y(t) \mid \sup_{t \in S^1} \|Y(t)\| \leq r_0 \right\}.$$

## 5 | 9.4: Construction of $\psi$ .

- Have constructed  $w_\rho \in C^\infty_\times(x, z)C^\infty(x, z)$  for every  $\rho \geq \rho_0$ , since there is exponential decay.
- Yields  $\psi_\rho \in \mathcal{M}(x, z)$  a true solution (to be defined).
- Need to check that  $\mathcal{F}(\psi_\rho) = 0$  where

$$\mathcal{F} = \frac{\partial}{\partial s} + J \frac{\partial}{\partial t} + \text{grad } Hx$$

in the weak sense.

- $\psi_\rho$  already continuous, and by elliptic regularity, makes it a strong solution.
- Trivialization
- Defining  $\mathcal{F}_\rho$ .

$$W^{1,p}(\mathbf{R} \times S^1; \mathbf{R}^{2n}) \xrightarrow{\mathcal{F}_\rho} L^p(\mathbf{R} \times S^1; \mathbf{R}^{2n})$$

$$(y_1, \dots, y_{2n}) \mapsto \left[ \left( \frac{\partial}{\partial s} + J \frac{\partial}{\partial t} + \text{grad } H_t \right) \left( \exp_{w_\rho} \sum y_i Z_i^\rho \right) \right]_{Z_i}$$

where  $\mathcal{F}_\rho := \mathcal{F} \circ \exp_{w_\rho}$  written in the bases  $Z_i$ . sd - Newton-Picard method, general idea

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- Original method and variant: find the limit of a sequence

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}, \quad x_{n+1} = x_n - \frac{f(x_n)}{f'(\textcolor{red}{x}_0)}.$$

- Allows finding zeros of  $f$  given an approximate zero  $x_0$ .
- Linearize  $\mathcal{F}_\rho$ .