

题 1

- 1 分别用直接摄动法和 Lindstedt-Poincaré 摄动法求下述系统自由振动的二次近似解，并对结果进行一致有效展开检验

$$\begin{cases} \ddot{u} + u + \varepsilon u^2 = 0 \\ u(0) = a, \quad \dot{u}(0) = 0 \end{cases}$$

直接摄动法

二次近似

$$u(\varepsilon, t) = u_0(t) + \varepsilon u_1(t) + \varepsilon^2 u_2(t) \quad (1)$$

代入方程

$$(\ddot{u}_0 + \varepsilon \ddot{u}_1 + \varepsilon^2 \ddot{u}_2) + (u_0 + \varepsilon u_1 + \varepsilon^2 u_2) + \varepsilon(u_0 + \varepsilon u_1 + \varepsilon^2 u_2)^2 = 0 \quad (2)$$

可得

$$\begin{cases} \ddot{u}_0 + u_0 = 0 \\ \ddot{u}_1 + u_1 = -u_0^2 \\ \ddot{u}_2 + u_2 = -2u_0 u_1 \end{cases} \quad (3)$$

迭代

解

$$\begin{cases} \ddot{u}_0 + u_0 = 0 \\ u_0(0) = a \\ \dot{u}_0(0) = 0 \end{cases} \quad (4)$$

得

$$u_0 = a \cos t \quad (5)$$

代入

$$\begin{cases} \ddot{u}_1 + u_1 = -u_0^2 \\ u_1(0) = 0 \\ \dot{u}_1(0) = 0 \end{cases} \quad (6)$$

解得

$$u_1(t) = \frac{a^2}{6}(\cos 2t + 2 \cos t - 3) \quad (7)$$

代入到

$$\begin{cases} \ddot{u}_2 + u_2 = -2u_0 u_1 \\ u_2(0) = 0 \\ \dot{u}_2(0) = 0 \end{cases} \quad (8)$$

解得

$$u_2(t) = \frac{a^3}{144}(3 \cos 3t + 16 \cos 2t + 29 \cos t + 60t \sin t - 48) \quad (9)$$

二次近似

$$\begin{aligned} u(t) &= u_0(t) + \varepsilon u_1(t) + \varepsilon^2 u_2(t) \\ &= a \cos t + \frac{\varepsilon a^2}{6}(\cos 2t + 2 \cos t - 3) \\ &\quad + \frac{\varepsilon^2 a^3}{144}(3 \cos 3t + 16 \cos 2t + 29 \cos t + 60t \sin t - 48) \end{aligned} \quad (10)$$

LP 摄动法

做二次近似

$$\begin{aligned} u(\varepsilon, t) &= u_0(t) + \varepsilon u_1(t) + \varepsilon^2 u_2(t) \\ \omega(\varepsilon)^2 &= \omega_0^2 + \varepsilon b_1 + \varepsilon^2 b_2 \end{aligned} \quad (11)$$

则有

$$\omega_0^2 = \omega^2 - \varepsilon b_1 - \varepsilon^2 b_2 \quad (12)$$

代入方程合并同类项有:

$$\begin{cases} \ddot{u}_0 + \omega^2 u_0 = 0 \\ \ddot{u}_1 + \omega^2 u_1 = b_1 u_0 - u_0^2 \\ \ddot{u}_2 + \omega^2 u_2 = b_2 u_0 + b_1 u_1 - 2u_0 u_1 \end{cases} \quad (13)$$

解

$$\begin{cases} \ddot{u}_0 + \omega u_0 = 0 \\ u_0(0) = a \\ \dot{u}_0(0) = 0 \end{cases} \quad (14)$$

得

$$u_0 = a \cos \omega t \quad (15)$$

代入有

$$\begin{cases} \ddot{u}_1 + \omega u_1 = b_1 u_0 - u_0^2 \\ u_1(0) = 0 \\ \dot{u}_1(0) = 0 \end{cases} \quad (16)$$

其中 $\cos \omega t$ 的系数为 b_1 , 为避免永年项, 取 $b_1 = 0$

解得

$$u_1(t) = \frac{1}{6\omega^2}(2a^2 \cos \omega t + a^2 \cos 2\omega t - 3a^2) \quad (17)$$

代入

$$\begin{cases} \ddot{u}_2 + \omega^2 u_2 = b_2 u_0 + b_1 u_1 - 2u_0 u_1 \\ u_2(0) = 0 \\ \dot{u}_2(0) = 0 \end{cases} \quad (18)$$

令 $\cos \omega t$ 的系数为 0

$$\frac{5a^3 + 6a\omega^2 b_2}{6\omega^2} = 0 \quad (19)$$

解得

$$b_2 = -\frac{5}{6} \frac{a^2}{\omega^2} \quad (20)$$

解得方程为

$$u_2(t) = \frac{1}{144\omega^2} (3a^3 \cos 3\omega t + 16a^3 \cos 2\omega t + 29a^3 \cos \omega t - 48a^3) \quad (21)$$

可得

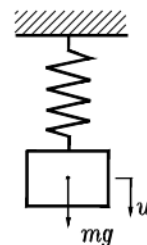
$$\begin{aligned} \omega &= \sqrt{\omega_0^2 + \varepsilon b_1 + \varepsilon^2 b_2} \\ &= \sqrt{1 - \varepsilon^2 \frac{5a^2}{6\omega^2}} \end{aligned} \quad (22)$$

题 2

- 2 考察图示重力场中的单自由度无阻尼系统，其非线性弹簧恢复力 q 与变形 δ 间关系为

$$q(\delta) = k\delta + \varepsilon k\delta^3, \quad k > 0, \quad 0 < \varepsilon \ll 1$$

- (1) 以系统静平衡位置为位移原点，建立系统的振动微分方程；
(2) 用 Lindstedt-Poincaré 摄动法分析重力对系统自由振动频率的影响。



题 2 图

1 建立方程

由牛顿定律

$$\begin{cases} mg - q(u_g) = 0 \\ mg - q(u_g + u) = m\ddot{u} \end{cases} \quad (23)$$

消掉 mg 得

$$m\ddot{u} + ku + \varepsilon k(u^3 + 3u_g u^2 + 3u u_g^2) = 0 \quad (24)$$

标准化

$$\ddot{u} + \omega_0^2 u + \varepsilon \omega_0^2 (u^3 + 3u_g u^2 + 3u u_g^2) = 0 \quad \omega_0 = \sqrt{\frac{k}{m}} \quad (25)$$

LP 摄动法

做二次近似

$$\begin{aligned} u(\varepsilon, t) &= u_0(t) + \varepsilon u_1(t) + \varepsilon^2 u_2(t) \\ \omega(\varepsilon)^2 &= \omega_0^2 + \varepsilon b_1 + \varepsilon^2 b_2 \end{aligned} \quad (26)$$

则有

$$\omega_0^2 = \omega^2 - \varepsilon b_1 - \varepsilon^2 b_2 \quad (27)$$

代入方程合并同类项有:

$$\begin{cases} \ddot{u}_0 + \omega^2 u_0 = 0 \\ \ddot{u}_1 + \omega^2 u_1 = -3u_g^2 \omega^2 u_0 + b_1 u_0 - 3u_g \omega^2 u_0^2 - \omega^2 u_0^3 \\ \ddot{u}_2 + \omega^2 u_2 = 3u_g^2 b_1 u_0 + b_2 u_0 + 3u_g b_1 u_0^2 + b_1 u_0^3 - 3u_g^2 \omega^2 u_1 + b_1 u_1 - 6u_g \omega^2 u_0 u_1 - 3\omega^2 u_0^2 u_1 \end{cases} \quad (28)$$

解

$$\begin{cases} \ddot{u}_0 + \omega u_0 = 0 \\ u_0(0) = a \\ \dot{u}_0(0) = 0 \end{cases} \quad (29)$$

得

$$u_0 = a \cos \omega t \quad (30)$$

代入有

$$\begin{cases} \ddot{u}_1 + \omega^2 u_1 = -3u_g^2 \omega^2 u_0 + b_1 u_0 - 3u_g \omega^2 u_0^2 - \omega^2 u_0^3 \\ u_1(0) = 0 \\ \dot{u}_1(0) = 0 \end{cases} \quad (31)$$

令 $\cos \omega t$ 的系数为 0

$$\frac{1}{4}(3a^3 \omega^2 + 12au_g^2 \omega^2 - 4ab_1) = 0 \quad (32)$$

得

$$b_1 = \frac{3a^3 + 12au_g^2}{4a} \omega^2 \quad (33)$$

解得

$$u_1(t) = \frac{1}{32}(-48a^2 u_g - a^3 \cos \omega t + 32a^2 u_g \cos \omega t + 16a^2 u_g \cos 2\omega t + a^3 \cos 3\omega t) \quad (34)$$

代入

$$\begin{cases} \ddot{u}_2 + \omega^2 u_2 = 3u_g^2 b_1 u_0 + b_2 u_0 + 3u_g b_1 u_0^2 + b_1 u_0^3 - 3u_g^2 \omega^2 u_1 + b_1 u_1 - 6u_g \omega^2 u_0 u_1 - 3\omega^2 u_0^2 u_1 \\ u_2(0) = 0 \\ \dot{u}_2(0) = 0 \end{cases} \quad (35)$$

令 $\cos \omega t$ 的系数为 0

$$\begin{aligned}
& -\frac{3}{64}a^5\omega^2 + \frac{9}{4}a^4u_g\omega^2 - \frac{243}{32}a^3u_g^2\omega^2 + 3a^2u_g^3\omega^2 + \frac{23}{128}a^2(-3a^3\omega^2 - 12au_g^2\omega^2) \\
& + \frac{1}{4}au_g(-3a^3\omega^2 - 12au_g^2\omega^2) + \frac{3}{4}u_g^2(-3a^3\omega^2 - 12au_g^2\omega^2) - ab_2 = 0
\end{aligned} \quad (36)$$

解得

$$b_2 = -\frac{75a^5\omega^2 - 192a^4u_g\omega^2 + 1536a^3u_g^2\omega^2 + 1152au_g^4\omega^2}{128a} \quad (37)$$

解得方程为

$$\begin{aligned}
u_2(t) = \frac{1}{1024} & (2016a^4u_g - 3072a^3u_g^2 + 4608a^2u_g^3 + 23a^5\cos\omega t - 1120a^4u_g\cos\omega t \\
& + 1952a^3u_g^2\cos\omega t - 3072a^2u_g^3\cos\omega t - 1024a^4u_g\cos 2\omega t \\
& + 1024a^3u_g^2\cos 2\omega t - 1536a^2u_g^3\cos 2\omega t - 24a^5\cos 3\omega t + 96a^4u_g\cos 3\omega t \\
& + 96a^3u_g^2\cos 3\omega t + 32a^4u_g\cos 4\omega t + a^5\cos 5\omega t)
\end{aligned} \quad (38)$$

可得

$$\begin{aligned}
\omega &= \sqrt{\omega_0^2 + \varepsilon b_1 + \varepsilon^2 b_2} \\
&= \sqrt{\omega_0^2 + \varepsilon \frac{3a^3 + 12au_g^2}{4a}\omega^2 + \varepsilon^2 \frac{75a^5\omega^2 - 192a^4u_g\omega^2 + 1536a^3u_g^2\omega^2 + 1152au_g^4\omega^2}{128a}}
\end{aligned} \quad (39)$$

可知, 重力大, u_g 大, ω 也会增大

题 3

3 用平均法求下述保守系统周期振动的一阶近似解

$$(1) \quad \ddot{u} + \sin u = 0$$

$$(2) \quad \ddot{u} + u + \varepsilon(u^2 + u^3) = 0, \quad 0 < \varepsilon \ll 1$$

1

近似

$$\sin u \approx u - \frac{1}{3!}u^3 \quad (40)$$

代入则有

$$\ddot{u} + u = \frac{1}{6}u^3 \quad (41)$$

令

$$\begin{cases} u = a(t) \cos(\omega_0 t + \varphi(t)) \\ \dot{u} = -\omega_0 a(t) \sin(\omega_0 t + \varphi(t)) \end{cases} \quad (42)$$

求导得

$$\begin{cases} \dot{u} = \dot{a} \cos(\omega_0 t + \varphi) - a \sin(\omega_0 t + \varphi)(\omega_0 + \dot{\varphi}) \\ \ddot{u} = -\omega_0 [\dot{a} \sin(\omega_0 t + \varphi) + a \cos(\omega_0 t + \varphi)(\omega_0 + \dot{\varphi})] \end{cases} \quad (43)$$

则有

$$\begin{cases} \dot{a} \cos(\omega_0 t + \varphi) - a \sin(\omega_0 t + \varphi) \dot{\varphi} = 0 \\ \dot{a} \sin(\omega_0 t + \varphi) + a \cos(\omega_0 t + \varphi) \dot{\varphi} = -\frac{1}{\omega_0} \frac{1}{6} a^3 \cos^3(\omega_0 t + \varphi) \end{cases} \quad (44)$$

解得

$$\begin{cases} \dot{a} = -\frac{1}{\omega_0} \frac{1}{6} a^3 \cos^3(\omega_0 t + \varphi) \sin(\omega_0 t + \varphi) \\ \varphi(a) = -\frac{1}{\omega_0} \frac{1}{6} a^3 \cos^3(\omega_0 t + \varphi) \cos(\omega_0 t + \varphi) \end{cases} \quad (45)$$

用平均值近似, 且由题意 $\omega_0 = 1$

$$\begin{cases} \dot{a} = -\frac{1}{2\pi\omega_0} \int_0^{2\pi} \frac{1}{6} a^3 \cos^3(\omega_0 t + \varphi) \sin(\omega_0 t + \varphi) dt = 0 \\ \varphi(a) = -\frac{1}{2\pi\omega_0 a} \int_0^{2\pi} \frac{1}{6} a^3 \cos^3(\omega_0 t + \varphi) \cos(\omega_0 t + \varphi) dt = \frac{1}{16} a^2 \end{cases} \quad (46)$$

所以

$$\begin{cases} a = a_0 \\ \varphi = \frac{1}{16} a^2 t \end{cases} \quad (47)$$

所以

$$u = a_0 \cos\left(t + \frac{1}{16} a^2 t\right) \quad (48)$$

2

令

$$\begin{cases} u = a(t) \cos(\omega_0 t + \varphi(t)) \\ \dot{u} = -\omega_0 a(t) \sin(\omega_0 t + \varphi(t)) \end{cases} \quad (49)$$

则有

$$\begin{cases} \dot{a} = -\frac{\varepsilon}{2\pi\omega} \int_0^\pi p(u, \dot{u}) \sin \psi d\psi \\ \dot{\varphi} = -\frac{\varepsilon}{2\pi\omega a} \int_0^\pi p(u, \dot{u}) \cos \psi d\psi \end{cases} \quad (50)$$

代入

$$\begin{cases} \dot{a} = \frac{\varepsilon}{2\pi} \int_0^\pi (a^2 \cos^3 \psi + a^3 \cos^3 \psi) \sin \psi d\psi = 0 \\ \dot{\varphi} = \frac{\varepsilon}{2\pi a} \int_0^\pi (a^2 \cos^3 \psi + a^3 \cos^3 \psi) \cos \psi d\psi = \frac{3}{8} \varepsilon a^2 \end{cases} \quad (51)$$

所以

$$\begin{cases} a = a_0 \\ \varphi = \frac{3}{8}\varepsilon a^2 t \end{cases} \quad (52)$$

所以

$$u = a_0 \cos\left(t + \frac{3}{8}\varepsilon a^2 t\right) \quad (53)$$

题 4

4 试论证可否用平均法求解下述立方非线性系统的周期振动

$$\ddot{u} + u^3 = 0$$

变形为

$$\ddot{u} + \omega_0^2 u = \varepsilon(\omega_0^2 u - u^3) \quad \varepsilon = 1 \quad (54)$$

则有

$$\begin{cases} \dot{a} = -\frac{\varepsilon}{2\pi\omega} \int_0^\pi p(u, \dot{u}) \sin \psi \, d\psi \\ \dot{\varphi} = -\frac{\varepsilon}{2\pi\omega a} \int_0^\pi p(u, \dot{u}) \cos \psi \, d\psi \end{cases} \quad (55)$$

代入则有

$$\begin{cases} \dot{a} = -\frac{1}{2\pi\omega} \int_0^\pi (\omega_0^2 a \cos \psi - a^3 \cos \psi^3) \sin \psi \, d\psi = 0 \\ \dot{\varphi} = -\frac{1}{2\pi\omega a} \int_0^\pi (\omega_0^2 a \cos \psi - a^3 \cos \psi^3) \cos \psi \, d\psi = -\frac{1}{2}\omega_0 + \frac{3}{8}\frac{a^2}{\omega_0} \end{cases} \quad (56)$$

得

$$\begin{cases} a = a_0 \\ \varphi = -\frac{1}{2}\omega_0 t + \frac{3}{8}\frac{a_0^2}{\omega_0} t \end{cases} \quad (57)$$

$$u = a_0 \cos\left(t + \frac{3}{8}\frac{a_0^2}{\omega_0} t - \frac{1}{2}\omega_0 t\right) \quad (58)$$

由于 $\varepsilon = 1$, 解与自由振动频率相差较大, 所以不用平均法

题 5

5 考察下述 Coulomb 摩擦阻尼系统

$$\ddot{u}(t) + \mu N \operatorname{sgn} \dot{u}(t) + \omega_0^2 u(t) = 0$$

其中摩擦系数 μ 为小参数, 用 KBM 法求解该系统自由振动的一次近似。

变形为

$$\ddot{u} + \omega_0^2 u = -\varepsilon N \operatorname{sgn} \dot{u} \quad \varepsilon = \mu \quad (59)$$

则有

$$\begin{cases} \dot{a} = -\frac{\varepsilon}{2\pi\omega} \int_0^\pi p(u, \dot{u}) \sin \psi \, d\psi \\ \dot{\varphi} = -\frac{\varepsilon}{2\pi\omega a} \int_0^\pi p(u, \dot{u}) \cos \psi \, d\psi \end{cases} \quad (60)$$

代入则有

$$\begin{cases} \dot{a} = \frac{\varepsilon}{2\pi\omega} \int_0^\pi (N \operatorname{sgn}(-\omega_0 a \sin \psi)) \sin \psi \, d\psi = -\frac{2\varepsilon N}{\pi\omega_0} \operatorname{sgn} a \\ \dot{\varphi} = \frac{\varepsilon}{2\pi\omega a} \int_0^\pi (N \operatorname{sgn}(-\omega_0 a \sin \psi)) \cos \psi \, d\psi = 0 \end{cases} \quad (61)$$

得

$$\begin{cases} a = a_0 - \frac{2\varepsilon N}{\pi\omega_0} \operatorname{sgn}(a)t \\ \varphi = 0 \end{cases} \quad (62)$$

所以

$$u = \left(a_0 - \frac{2\varepsilon N}{\pi\omega_0} \operatorname{sgn}(a)t \right) \cos \omega_0 t \quad (63)$$

题 6

设

$$\ddot{x} + \omega_0^2 x = \varepsilon f(x, \dot{x})$$

用矩阵变换推导平均方程

解

矩阵形式

$$\begin{pmatrix} \dot{x} \\ \ddot{x} \end{pmatrix} + \begin{pmatrix} \omega_0^2 & 0 \end{pmatrix} \begin{pmatrix} x \\ \dot{x} \end{pmatrix} = \begin{pmatrix} 0 \\ \varepsilon f(x, \dot{x}) \end{pmatrix} \quad (64)$$

令

$$\begin{cases} x = a(t) \cos \psi \\ \dot{x} = -\omega_0 a(t) \sin \psi \\ \psi = \omega_0 t + \varphi(t) \end{cases} \quad (65)$$

可得

$$\begin{pmatrix} \dot{x} \\ \ddot{x} \end{pmatrix} = \begin{pmatrix} \dot{a} \cos \psi - a \cos \psi (\omega_0 + \dot{\varphi}) \\ -\omega_0 (\dot{a} \sin \psi + a \cos \psi (\omega_0 + \dot{\varphi})) \end{pmatrix} \quad (66)$$

代入则有

$$\begin{pmatrix} \cos \psi & -\sin \psi \\ \sin \psi & \cos \psi \end{pmatrix} \begin{pmatrix} \dot{a} \\ a\dot{\varphi} \end{pmatrix} = \begin{pmatrix} 0 \\ -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \end{pmatrix} \quad (67)$$

知 $\begin{pmatrix} \cos \psi & -\sin \psi \\ \sin \psi & \cos \psi \end{pmatrix}$ 为旋转矩阵, 即正交矩阵, 有

$$A^{-1} = A^T \quad (68)$$

所以

$$\begin{aligned} \begin{pmatrix} \dot{a} \\ a\dot{\varphi} \end{pmatrix} &= \begin{pmatrix} \cos \psi & -\sin \psi \\ \sin \psi & \cos \psi \end{pmatrix}^{-1} \begin{pmatrix} 0 \\ -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \end{pmatrix} \\ &= \begin{pmatrix} \cos \psi & -\sin \psi \\ \sin \psi & \cos \psi \end{pmatrix}^T \begin{pmatrix} 0 \\ -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \end{pmatrix} \\ &= \begin{pmatrix} \cos \psi & \sin \psi \\ -\sin \psi & \cos \psi \end{pmatrix} \begin{pmatrix} 0 \\ -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \end{pmatrix} \\ &= \begin{pmatrix} -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \sin \psi \\ -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \cos \psi \end{pmatrix} \end{aligned} \quad (69)$$

所以

$$\begin{pmatrix} \dot{a} \\ \dot{\varphi} \end{pmatrix} = \begin{pmatrix} -\frac{\varepsilon}{\omega_0} f(x, \dot{x}) \sin \psi \\ -\frac{\varepsilon}{\omega_0 a} f(x, \dot{x}) \cos \psi \end{pmatrix} \quad (70)$$