M1-IRELE

ELEC-H415 Communication Channels

V2V communication project

Author :

Colot Emmeran

Professor :

De Doncker Philippe

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Introduction

TODO

Might want to check all 3d plots names

Write an introduction

Resize figures to fill pages

Theoretical answers
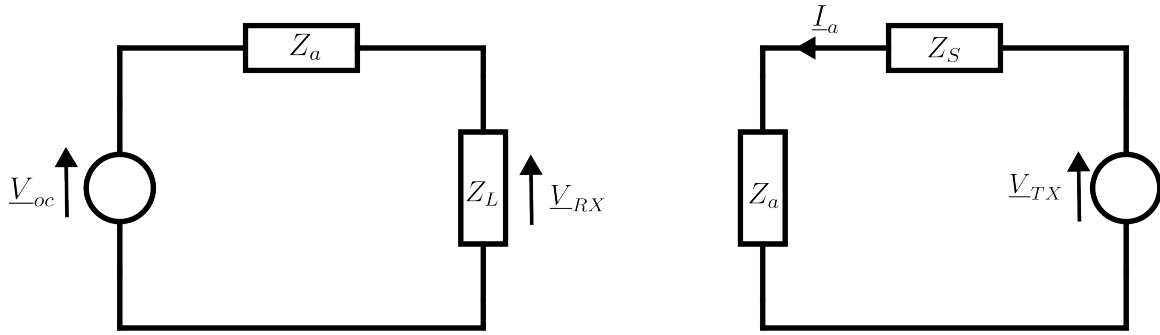
2.1 $\lambda/2$ **antennas**

Figure 2.1: Equivalent electric circuit of an antenna in RX (left) and TX (right)

Fig 2.1 shows the equivalent electrical circuit at RX and TX where $\underline{V}_{oc}$ is the induced voltage, $\underline{V}_{RX}$ the voltage at the output of the RX antenna, $\underline{V}_{TX}$ at the input of the TX antenna and $\underline{I}_a$ the current entering the TX antenna.

As both transmitting and receiving antenna are vertical $\lambda/2$ dipoles, their equivalent heights can be analytically computed:

$$\vec{h}_e(\theta,\phi) = \frac{\lambda}{\pi} \frac{\cos(\frac{1}{2}\cos\theta)}{\sin^2\theta} \vec{1}_z$$
$$\vec{h}_{e\perp}(\theta,\phi) = -\frac{\lambda}{\pi} \frac{\cos(\frac{\pi}{2}\cos\theta)}{\sin\theta} \vec{1}_\theta$$

As θ and ϕ are spherical coordinates, the horizontal plane (in which the simulation is done) corresponds to $\theta = \pi/2$. The transverse equivalent height is then reduced to hte following, where it does not depend on the azimuthal angle ϕ:

$$\vec{h}_{e\perp}(\phi) = -\frac{\lambda}{\pi}\vec{1}_\theta$$

The transverse part of the equivalent height gives rise to an expression for the emitted electric field.

$$\begin{aligned}
\underline{\vec{E}} &= -j\omega\underline{I}_a\frac{\mu_0}{4\pi}\frac{e^{-j\beta r}}{r}\vec{h}_{e\perp}(\theta,\phi) \\
&= j\omega\underline{I}_a\frac{\mu_0\lambda}{4\pi^2}\frac{e^{-j\beta r}}{r}\vec{1}_\theta \\
&= j\frac{\underline{I}_a\mu_0 c}{2\pi}\frac{e^{-j\beta r}}{r}\vec{1}_\theta \\
&= j\frac{\underline{I}_a Z_0}{2\pi}\frac{e^{-j\beta r}}{r}\vec{1}_\theta
\end{aligned}$$

To make the transmission parameters appear in the electric field expression, β and $\underline{I}_a$ must be replaced with the formulas given below. Z_0, Z_a, Z_S can also be replaced by respectively 120π, $\frac{720\pi}{32}$ and $\frac{720\pi}{32}$ (assuming the impedances are matching).

$$\beta = \frac{2\pi f_c}{c}$$
$$\underline{V}_{\text{TX}} = (Z_a + Z_S)\underline{I}_a = \frac{720\pi}{16}\underline{I}_a$$

Wich yields

$$\begin{aligned}
\underline{\vec{E}} &= j\frac{120\pi\underline{I}_a}{2\pi}\frac{e^{-\frac{j2\pi f_c r}{c}}}{r}\vec{1}_\theta \\
&= j\frac{4\underline{V}_{\text{TX}}}{3\pi}\frac{e^{-\frac{j2\pi f_c r}{c}}}{r}\vec{1}_\theta
\end{aligned}$$

The last step is to replace the travel distance r with $c\tau$ as the wave is propagating at the speed of light in free space. The electric field thus becomes:

$$\underline{\vec{E}} = j\frac{4\underline{V}_{\text{TX}}}{3\pi}\frac{e^{-j2\pi f_c\tau}}{c\tau}\vec{1}_\theta$$

The voltage at the output of the antenna $\underline{V}_{\text{RX}}$ can be deduced from the equivalent electric circuit 2.1 as it is a simple voltage divider. Assuming a matching between the antenna and the load, we have:

$$\underline{V}_{\text{RX}} = \frac{Z_L}{Z_a + Z_L}\underline{V}_{oc} = \frac{1}{2}\underline{V}_{oc}$$

$$\underline{V}_{oc} = -\left.\vec{h}_{e\perp}(\theta,\phi)\right|_{\theta=\pi/2} \cdot \underline{\vec{E}}_i$$

Where $\underline{\vec{E}}_i = -\underline{\vec{E}}$ due to the change of coordinates origin. It is here assumed that there was no reflexion, refraction or transmission through another material. This results in:

$$\underline{V}_{oc} = \frac{\lambda}{\pi}\cdot\underline{E}_i$$

$$= -j\frac{4\lambda\underline{V}_{\text{TX}}}{3\pi^2}\frac{e^{-j2\pi f_c\tau}}{c\tau}$$

$$\underline{V}_{\text{RX}} = -j\frac{2\lambda}{3\pi^2}\frac{e^{-j2\pi f_c\tau}}{c\tau}\underline{V}_{\text{TX}}$$

For a practical use in step 3, the time of flight τ is replaced by the traveled distance d to get:

$$\underline{V}_{\text{RX}} = -j\frac{2\lambda}{3\pi^2}\frac{e^{-j\frac{2\pi f_c d}{c}}}{d}\underline{V}_{\text{TX}} \tag{2.1}$$

2.2 LOS channel

Assuming the communication takes place via a LOS ray only, the channel impulse response $h(\tau)$ is defined as follows:

$$h(\tau) = \frac{\alpha_1 e^{-j\frac{2\pi f_c d_1}{c}}}{d_1}\delta(\tau - \tau_1)$$

Where d_1 is the distance of propagation and τ_1 he propagation delay of the direct ray. α_1 (which might be complex) takes into account a phase change or attenuation due for example to reflections whereas the imaginary exponential next to it corresponds to the phase change due to the propagation delay. In the case of LOS transmission α_1 is equal to 1.

The transfer function $H(f)$ of the channel is found by taking the Fourier transform of $h(\tau)$

$$H(f) = \int_{-\infty}^{\infty} h(\tau) e^{-j2\pi f\tau} d\tau$$

$$= \frac{e^{-j\frac{2\pi f_c d_1}{c}}}{d_1} e^{-j2\pi f\frac{d_1}{c}}$$

$$= \frac{e^{-j\frac{2\pi(f+f_c)d_1}{c}}}{d_1}$$

The narrowband impulse response of the channel $h_{\text{NB}}(\tau)$ is, in the case of a single ray, simply equal to $h(\tau)$ and the signal received at the output of the channel $y(t)$ is given by:

$$y(t) = h_{\text{NB}}(t)x(t)$$

$$\text{where} \qquad h_{\text{NB}}(t) = h(t) = \frac{e^{-j\frac{2\pi f_c d_1}{c}}}{d_1}\delta(t-\tau_1)$$

Where $x(t)$ is the transmitted signal.

The ratio between the received power P_{RX} and the transmitted power P_{TX} is found with:

$$\frac{P_{\text{received}}}{P_{\text{transmitted}}} = \frac{|h(\tau)|^2}{2}$$

$$= \frac{1}{d_1^2} \tag{2.2}$$

This result can be compared with the Friis formula, given by:

$$P_{\text{RX}}(d) = P_{\text{TX}}G_{\text{TX}}(\theta_{\text{TX}},\phi_{\text{TX}})G_{\text{RX}}(\theta_{\text{RX}},\phi_{\text{RX}})\left(\frac{\lambda}{4\pi d}\right)^2 \tag{2.3}$$

The reason for the big difference between the two formulas can be easily explained: in eq 2.2, $P_{\text{transmitted}}$ and P_{received} are the powers of the waves, not the one of the signal before/after passing through the antennas. To correct it, they are replaced by the Poynting vectors $\vec{S}_{\text{TX}}$ and $\vec{S}_{\text{RX}}$:

$$\frac{\left|\vec{S}_{\text{RX}}\right|}{\left|\vec{S}_{\text{TX}}\right|} = \frac{1}{d_1^2} \tag{2.4}$$

To compare the received and the injected power, the Poynting vectors should be replaced by

$$\left|\vec{\mathbf{S}}_{\text{TX}}\right| = G_{\text{TX}} P_{\text{TX}}$$

$$\left|\vec{\mathbf{S}}_{\text{RX}}\right| A_{eRX} = P_{\text{RX}}$$

$$\text{where} \quad A_{eRX} = G_{\text{RX}} \left(\frac{\lambda}{4\pi}\right)^2$$

When placed back in eq 2.4, the expression matches the Friis formula (eq 2.3).

To further simplify and as the antennas are considered to be lossless dipoles, equation 2.6 replaces their gain by $\frac{16}{3\pi}$, the theoretical gain of such antennas.

$$\frac{P_{\text{RX}}}{P_{\text{TX}}} = G_{\text{TX}} G_{\text{RX}} \left(\frac{\lambda}{4\pi d_1}\right)^2 \tag{2.5}$$

$$= \frac{16}{9\pi^2} \left(\frac{\lambda}{\pi d_1}\right)^2 \tag{2.6}$$

Two deductions can be made of this result:

- The power loss on the channel depends on the square of the distance between the emitter and the receiver. It implies that doubling the range of an antenna would need a multiplication of the transmitting power by a factor 4.

- A more surprising result is the wavelength squared appearing at the numerator. It shows that there is always a compromise between data rate and power consumption as a larger wavelength will indeed be less attenuated but the maximal data rate will then be lowered in order for the channel to still be considered as narrowband.

3.1 Raytracing

After having implemented a raytracing algorithm, the simulation is able to compute every path from the transmitter to the receiver. Figure 3.1 shows the result of the simulation for cars spaced by 100m on a 20m wide road surrounded by buildings. For a better understanding, the angles are computed between the buildings and the rays. This is only done on this graph as the angles used to compute the reflection coefficients $\Gamma_\perp$ are the ones between the normal to the obstacle and the ray.
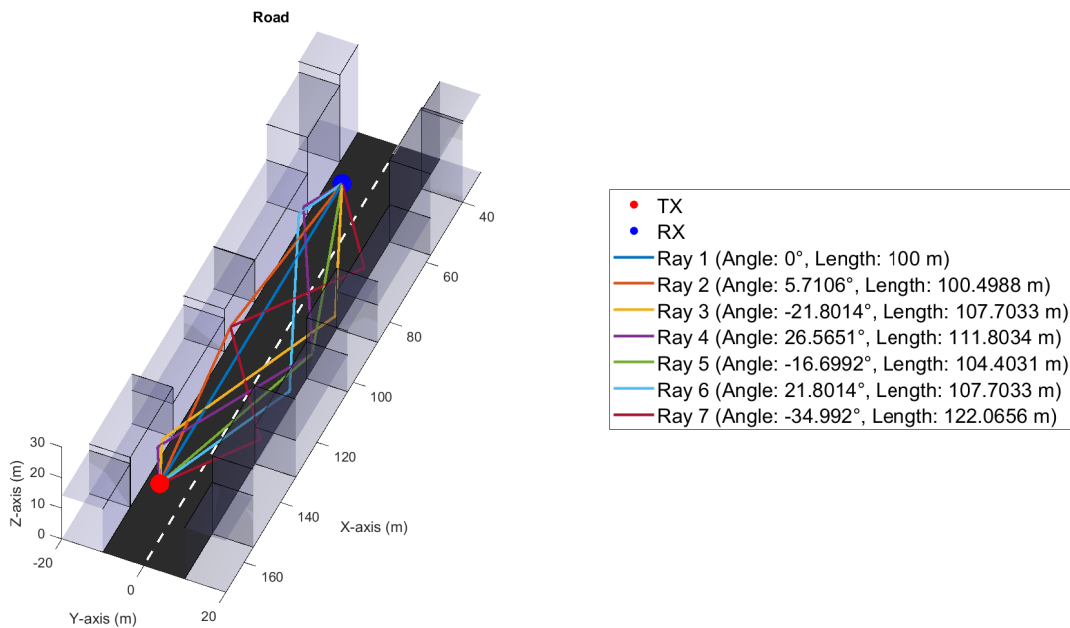
Figure 3.1: Raytracing simulation of a road surrounded by buildings

The received voltage has been computed on the same setup using equation 2.1 to which the reflection coefficients $\Gamma_\perp$ have been added.

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$$\Gamma_\perp = \frac{\cos\theta_i - \sqrt{\epsilon_r - \sin^2\theta_i}}{\cos\theta_i + \sqrt{\epsilon_r - \sin^2\theta_i}}$$

The voltage carried by each ray is shown in figure 3.2 using a colormap and the exact values are given in table 3.1. The total received voltage is computing by summing the voltages of each ray and taking the phase into account. This gave the following result:

$$\text{Total voltage} = 145.6\mu V \angle -48.8°$$

Ray Index	Ray physical angle (°)	Ray length (m)	Carried voltage (μV)	Phase (°)
1	180.0	100.0	129.1	30.0
2	174.3	100.5	114.6	-81.2
3	-163.3	104.4	88.9	-3.7
4	-158.2	107.7	51.2	-149.3
5	158.2	107.7	51.2	-149.3
6	153.4	111.8	24.9	161.93
7	-145.0	122.1	15.0	-134.1

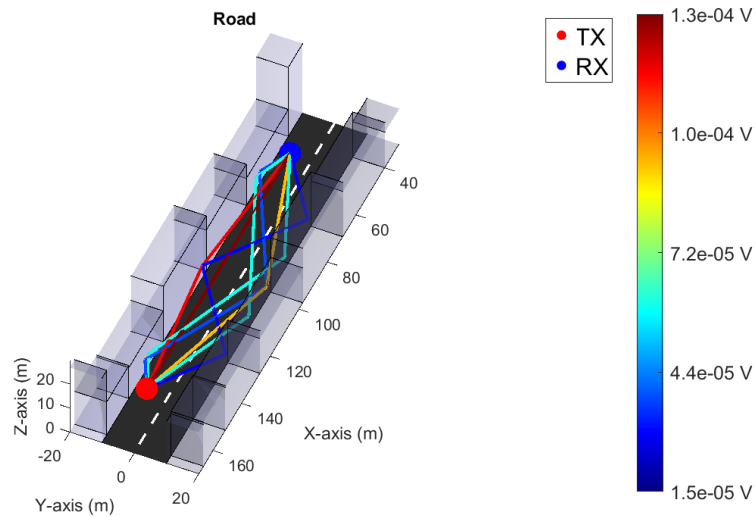Table 3.1: Ray properties: index, angle, length, and voltage amplitude

Figure 3.2: Received voltage for each ray

3.2 Power analysis

By varying the distance between the two communicating cars between 0 and 1000m, a comparison is made between the raytracing simulation and the theoretical model on figure 3.3. The simulated power is computed by taking the square of V_{OC} divided by $Z_L \parallel Z_a$ (see figure 2.1) and the theoretical power is computed using equation 2.6.
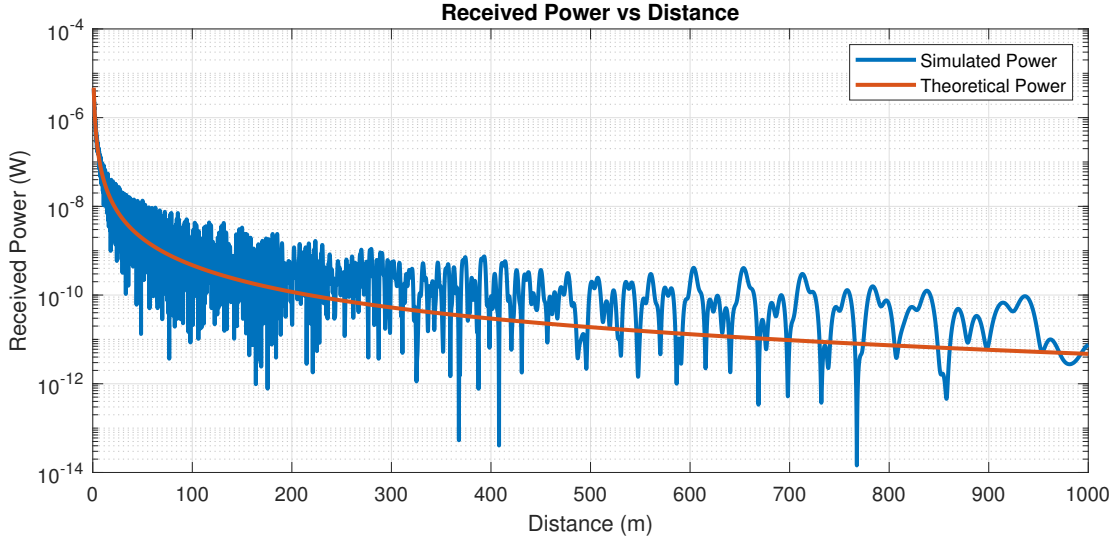

Figure 3.3: Received power as a function of the distance between the two cars

For short distances, the LOS ray is the only one contributing to the received power which explains the good match of the simulation with theory.

After a few meters, the multi path components (MPC) start to have an impact on the received power and the resulting interferences make the simulated power oscillate around the true power.

At greater distances and because of the limited number of bounces, one can clearly see the power is a multisine signal. This is because the reflection angles don't vary much with the distance at that point.

The oscillation around the expected power can be modeled with a Rice distribution as the LOS ray is the strongest one. The Rice factor K is defined as the ratio between the power of the LOS ray and the power of the other rays:

$$K = \frac{a_0^2}{\sum_{n=1}^{N} a_n^2}$$

where a_n is the amplitude of the nth ray with a_0 the amplitude of the LOS ray. As those amplitudes vary with distance, figure 3.4 shows the rice factor as a function of the distance between the two cars.
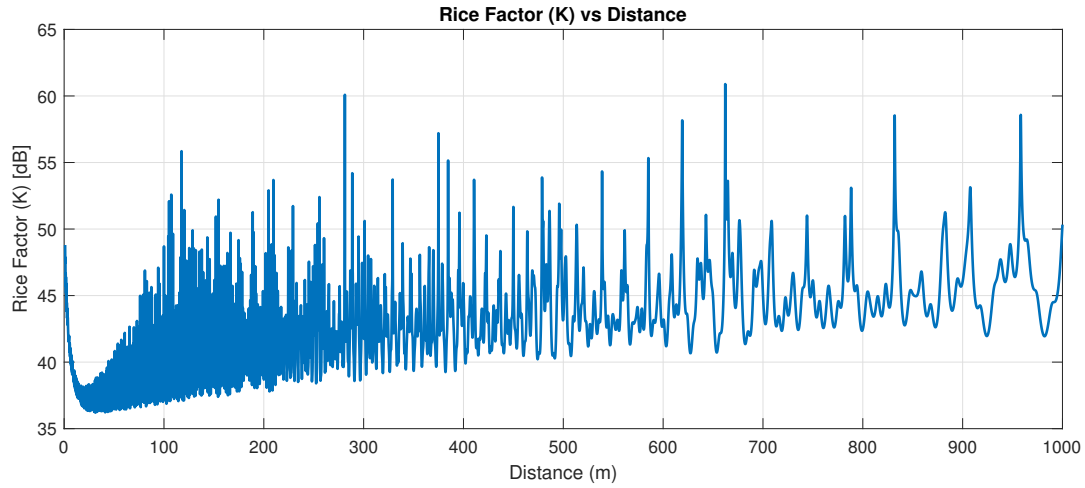
Figure 3.4: Rice factor as a function of the distance between the two cars

The average received power in local areas is computed using:

$$< P_{\text{RX}} > = \frac{1}{8Z_a} \sum_{n=1}^{N} |V_{OC}|^2$$
$$= \frac{1}{45\pi} \sum_{n=1}^{N} |V_{\text{RX}}|^2$$

Where the sum is done over all the received rays. Figure 3.5 shows the average received power in local areas with a width of 5m for an emitter placed at the intersection of the two roads. The power is in dBm as it quicly becomes very small.

Another simulation has been done with 2m wide local areas and an emitter placed a bit further from the intersection. The result is shown in figure 3.6 and it can be clearly seen that it is really difficult to get any power when no LOS ray is present.

It must be mentioned that the simulation does not take any transmission ray into account. This is under the hypothesis that the walls of the buildings are thick enough to not let any ray pass through.
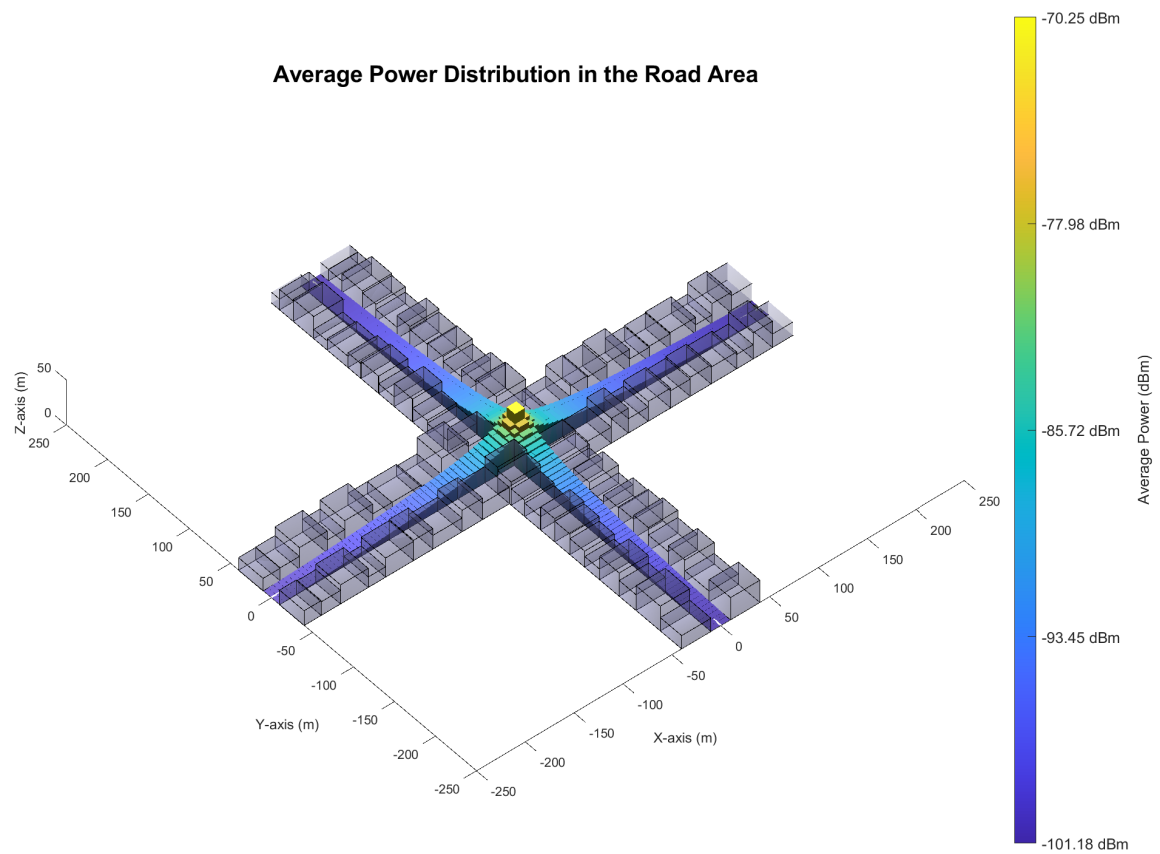
Average Power Distribution in the Road Area

-70.25 dBm

-77.98 dBm

-85.72 dBm

Average Power (dBm)

-93.45 dBm

-101.18 dBm

Z-axis (m)

50

0

250

200

150

100

50

0

-50

-100

-150

-200

-250

Y-axis (m)

250

200

150

100

50

0

-50

-100

-150

-200
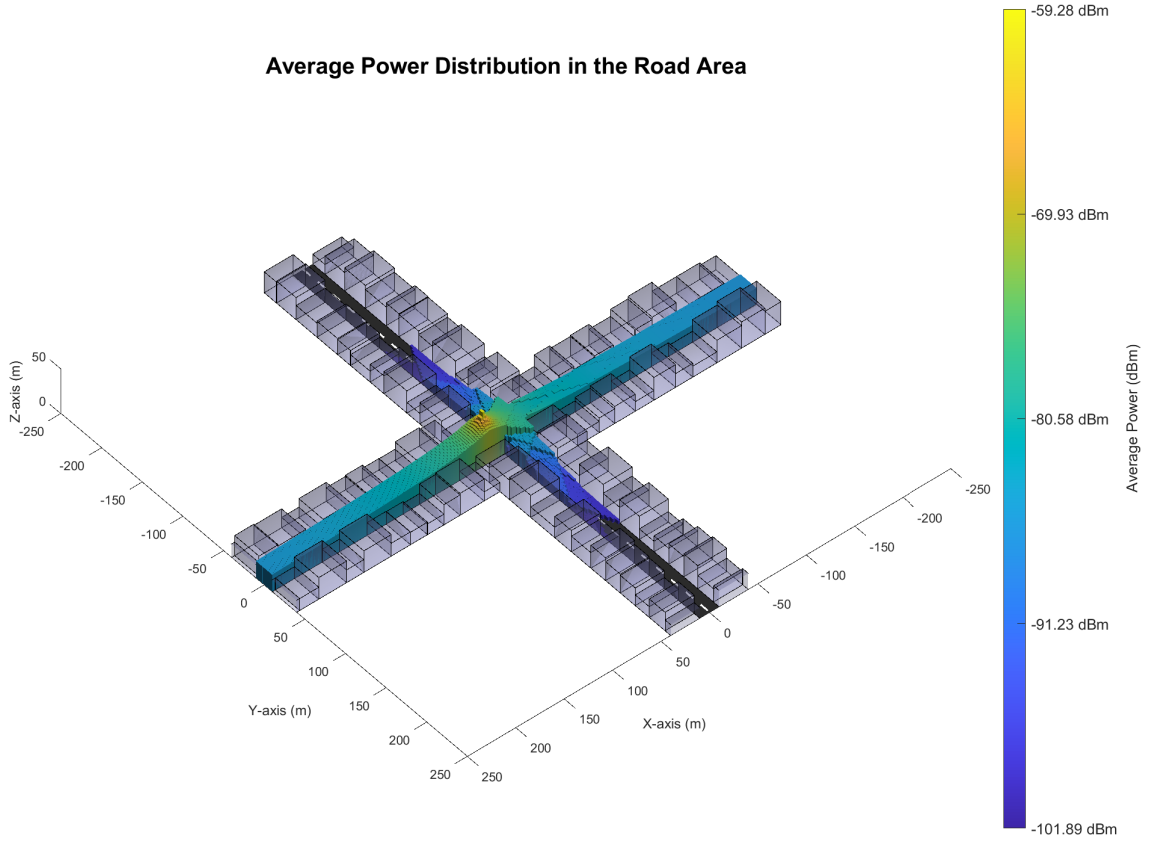
-250

X-axis (m)

Figure 3.5: Average received power in local areas of 5m

Figure 3.6: Average received power in local areas of 2m

3.3 Full channel, narrowband model

The pass loss of the channel $L(d)$ and the pass loss that does not depend on the antennas $L_0(d)$ are defined (with every term in dB) as:

$$L(d) = P_{TX} - P_{RX} \tag{3.1a}$$

$$L_0(d) = L(d) + G_{TX} + G_{RX} \tag{3.1b}$$

The canonical model for $L_0(d)$ is:

$$L_0(d) = L_0(d_0) + 10n log_{10}\left(\frac{d}{d_0}\right)$$

In which the parameters L_0 and n must be found to fit the values of the simulation. d_0 has been arbitrarily set to 1m. The result is shown in figure 3.7 and the resulting parameters give as path loss model:

$$L_0(d) = 58.501 + 10 \cdot 1.56 \cdot log_{10}\left(\frac{d}{1\mathrm{m}}\right) \tag{3.2}$$
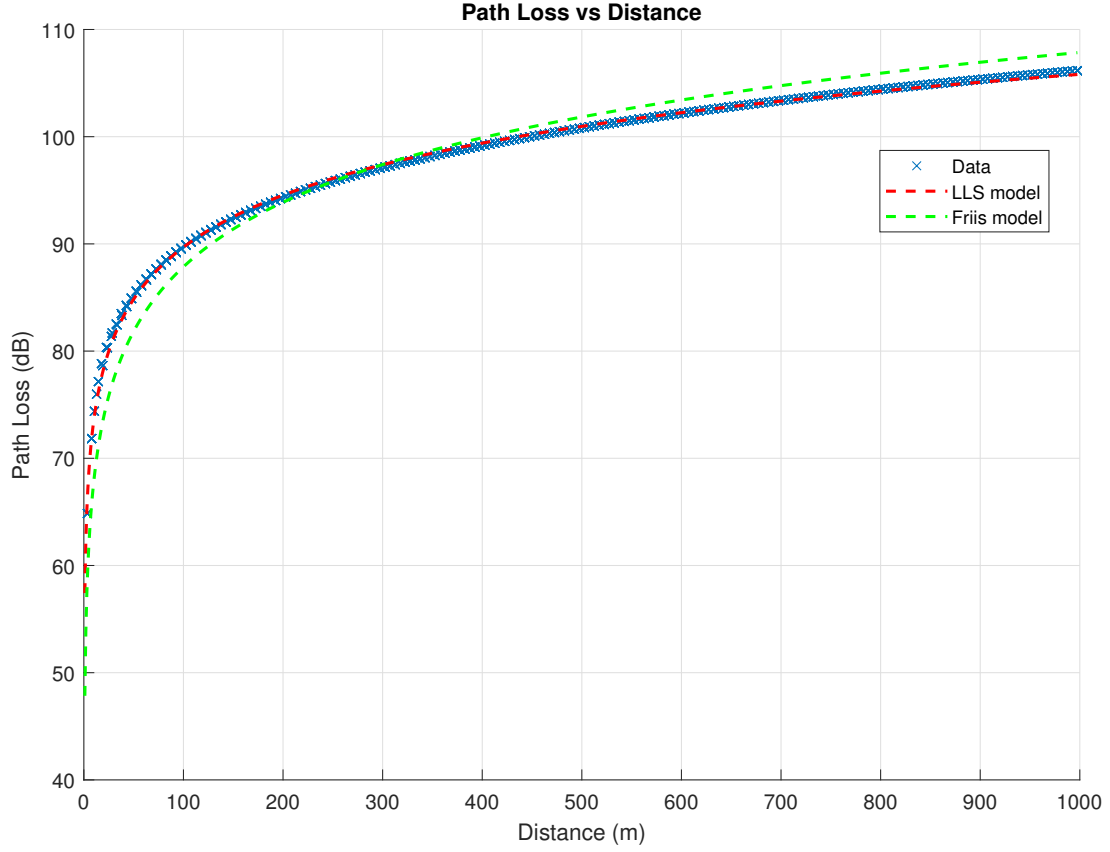

Figure 3.7: Pass loss of the channel and fitted model

As the LLS estimator only gives a "*best fit*" to the data, there is still a parameter to define: the variability σ_L. It is the standard deviation of the path loss around the fitted model and its value here is:

$$\sigma_L = 0.30236 \quad \mathrm{dB}$$

Regarding the communication reliability, it can be found under the assumption that the power variations due to shadowing is log-normal distributed. In this case, the probability of a successful communication is given by:

$$P_{\mathrm{success}} = \frac{1}{2}\mathrm{erfc}\left(\frac{M}{\sigma_L}\sqrt{2}\right)$$

Where M is the fade margin, defined as the difference between the received power and the minimum power needed to decode the signal. The fade margins required for a certain communication reliability are given in table 3.2. They have been found by plotting the error probability as a function of the fade margin (see figure 3.8). The cell range can also be computed for each value of P_{success} using equation 3.1 and 3.2.

$$L(d) = P_{\text{TX}} - P_{\text{RX}}$$

$$L_0(d_{\max}) = P_{\text{TX}} - P_{RX,\min} + G_{\text{TX}} + G_{\text{RX}}$$

$$58.501 + 10 \cdot 1.56 \cdot log_{10}\left(\frac{d_{\max}}{1\text{m}}\right) = P_{\text{TX}} - P_{RX,\min} + G_{\text{TX}} + G_{\text{RX}}$$

Communication reliability	Fade margin (dB)	Cell range (m)
50%	0	??
95%	0.4973	??
99%	0.7033	??

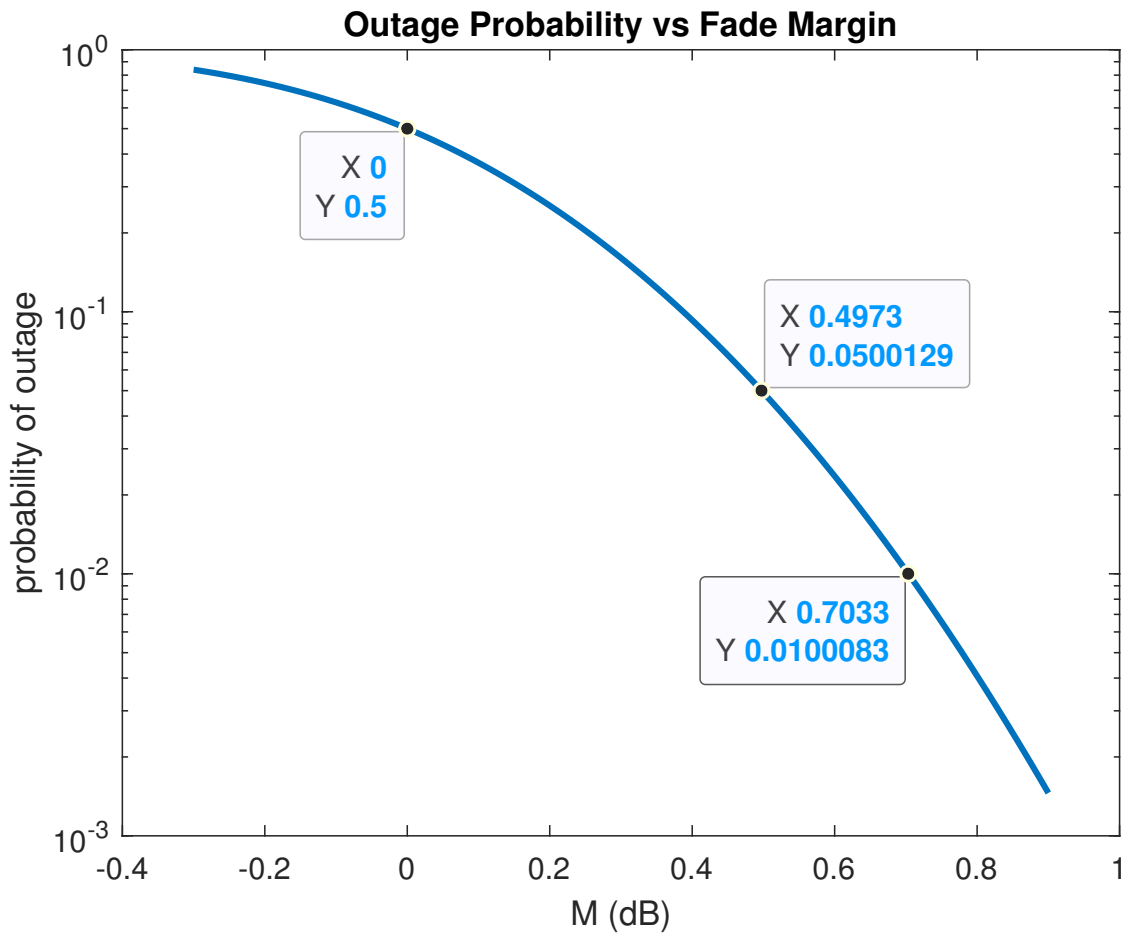Table 3.2: Fade margins for different communication reliability

Figure 3.8: Fade margin as a function of the communication reliability

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3.4   LOS channel, wideband model

For the wideband model and if there is only a single ray, the impulse response of the channel is given by:

$$h(t) = \frac{e^{-j\frac{2\pi f_c d_1}{c}}}{d_1} \delta(t - \tau_1)$$

The impulse response and the corresponding transfer function (computed by taking the DFT of $h(t)$) are shown in figure 3.9 and 3.10.
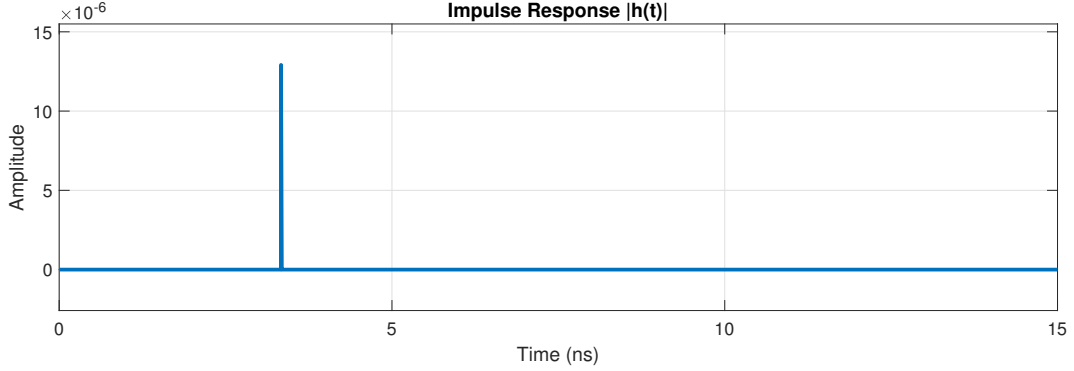

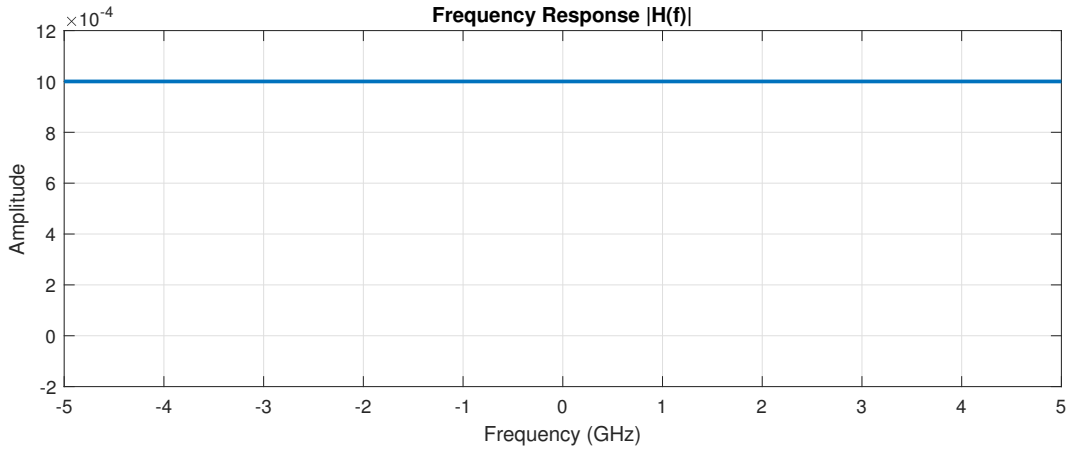Figure 3.9: Impulse response of the wideband channel with the LOS ray only

Figure 3.10: Transfer function of the wideband channel with the LOS ray only

The impulse response is a single peak at a time corresponding to the time of flight of the LOS ray. The transfer function has a constant value over the frequency spectrum. This means that there is no frequency selectivity for a single ray, a result easily understandable as no frequency is less attenuated than another in free space. This is of course valid only because some assumptions have been made: there is no multipath propagation and no parasitic effects (gaseous absorption, rain attenuation, tropospheric scintillation, etc.). Because of the limited bandwidth of the channel, the tapped delay line model is used to compute $h_{\text{TDL}}(t)$. The previous impulse response did not take into the fact that the receiver is only able to sample at the time resolution of the system $\Delta\tau = \frac{1}{B}$, where B is its bandwidth. The TDL model is found by:

$$h_{\text{TDL}}(t) = \sum_{l=0}^{L} h_l(t)\delta(t - l\Delta\tau) \tag{3.3}$$

$$\text{where} \qquad h_l(t) = \sum_{n=1}^{N} \alpha_n(t)\text{sinc}\,(B(\tau_n - l\Delta\tau))$$

In those equations, L is the maximal index of the temporal taps (linked with the maximal allowed delay) and N is the number of rays, which is simply 1 in this case. α_1 is different from the one defined in section 2.2, as it here takes the spacial dispersion into account:

$$\alpha_n = \prod_{b=1}^{B} \Gamma_{\perp,b} \cdot \frac{e^{\frac{-j2\pi f_c d_n}{c}}}{d_n} \qquad \text{where } B \text{ is the number of bounces of the } n^{th} \text{ ray}$$

When all of it is put together in equation 3.3, the impulse response of the channel becomes:

$$h_{\text{TDL}}(t) = \sum_{l=0}^{L} \frac{e^{\frac{-j2\pi f_c d_1}{c}}}{d_1}\text{sinc}\,(B(\tau_1 - l\Delta\tau))\,\delta(t - l\Delta\tau)$$

On figure 3.11, the effect of the sinc function is clearly visible: the impulse response is no longer a single peak but the adjacent time bins (of size $\Delta\tau$) are now different from zero. This is because the receiver is not able to sample the signal at the exact time of flight of the ray.
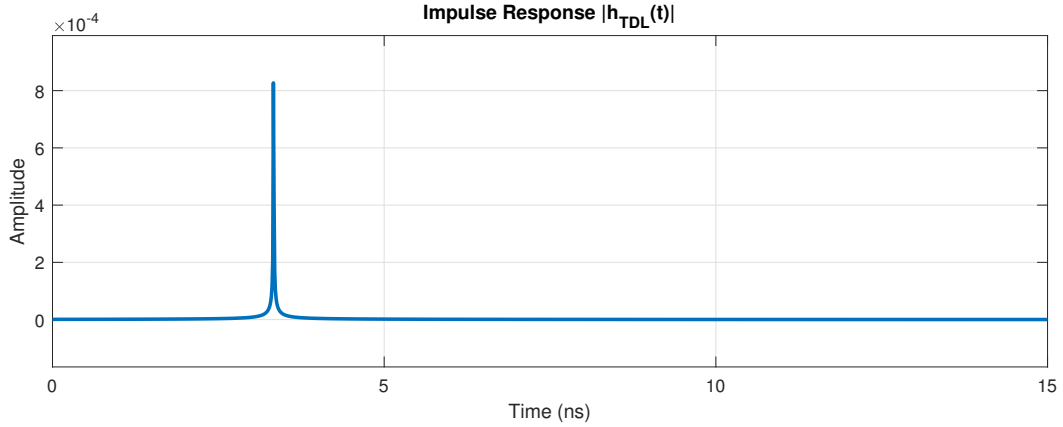
More explanations needed

Figure 3.11: TDL impulse response of the wideband channel with the LOS ray only

3.5   Full channel, wideband model

This section will do the same analysis as the previous one but for multiple rays. The impulse response and the corresponding transfer function are shown in figure 3.12 and 3.13.

The impulse response plot is zoomed in so the different peaks are visible. Because of those peaks, some

frequencies are more attenuated than others during their propagation through the channel. This is the reason why the transfer function is not constant anymore.

The transfer function plot is also only partially shown in order to focus around the carrier frequency f_c = 5.9GHz. This analysis shows that this choice for f_c is not optimal as the amplitude of $H(f)$ is low around this frequency. A wiser choice would have been to choose a frequency around f_c = 6.5GHz. Even though the transfer function is not maximal there, it does not vary much around it whereas a change of only 10 MHz above 5.9GHz could diminish the amplitude of the transfer function by half.
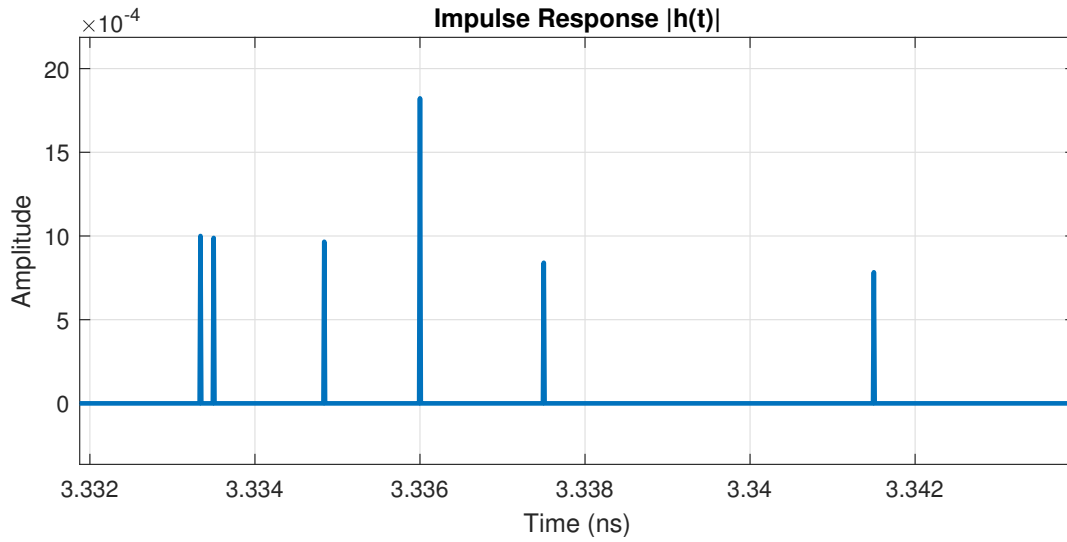

Figure 3.12: Impulse response of the wideband channel with multiple rays
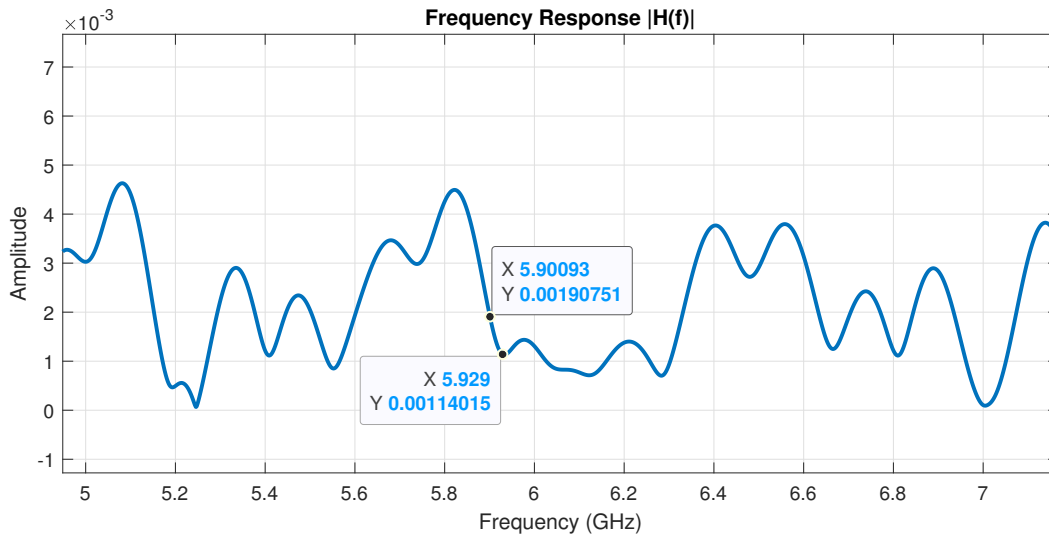

Figure 3.13: Transfer function of the wideband channel with multiple rays

The impulse response of the tapped delay line model is computed using the same method as in section 3.4 and it is plotted together with the one in with the LOS ray only on figure 3.14.
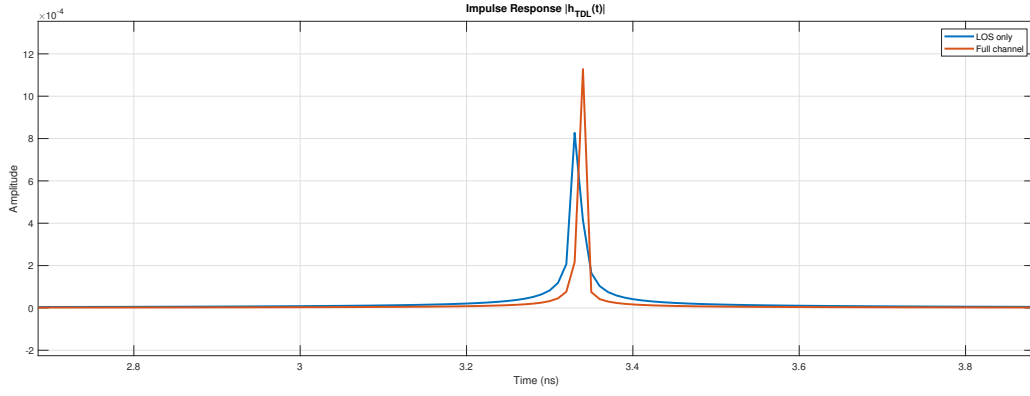
Figure 3.14: Comparison of h_{TDL} for the LOS ray only and the full channel

Taking the MPC into account makes h_{TDL} reach a higher value than the one with the LOS ray only. Because the bandwidth of the channel is lower than the coherence bandwidth (defined as the inverse of the delay between the first and last incoming ray), all the rays are arriving approximately at the same time.

To verify this, the same script has been used with 15 bounces instead of 3. The coherence bandwidth decreased from 123MHz to as low as 6MHz. The impulse response is shown on figure 3.15 and separate peaks are now visible.
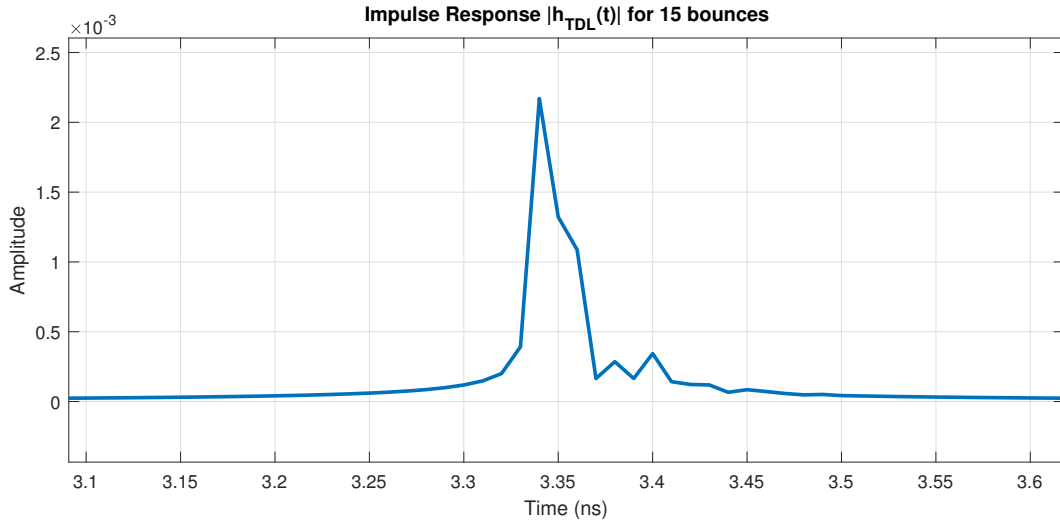

Figure 3.15: Impulse response of the wideband channel with multiple rays and 15 bounces

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Further analysis

- Take back step 4.1 and show coherence bandwidth vs distance to see whether the narrowband analysis is valid or not.

- try to do the whole analysis for different scenarios

- 3d plot with (distance, frequency, $|H(f)|$)