The strong superadditivity conjecture holds for the quantum depolarizing channel in any dimension

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Given a quantum channel Φ in a Hilbert space H put $\hat{H}_{\Phi}(\rho) = \min_{\rho_{av} = \rho} \sum_{j=1}^{k} \pi_{j} S(\Phi(\rho_{j}))$, where $\rho_{av} = \sum_{j=1}^{k} \pi_{j} \rho_{j}$, the minimum is taken over all probability distributions $\pi = \{\pi_{j}\}$ and states ρ_{j} in H, $S(\rho) = -Tr\rho\log\rho$ is the von Neumann entropy of a state ρ . The strong superadditivity conjecture states that $\hat{H}_{\Phi\otimes\Psi}(\rho) \geq \hat{H}_{\Phi}(Tr_K(\rho)) + \hat{H}_{\Psi}(Tr_H(\rho))$ for two channels Φ and Ψ in Hilbert spaces H and K, respectively. We have proved the strong superadditivity conjecture for the quantum depolarizing channel in any dimensions.

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I. INTRODUCTION

A linear trace-preserving map Φ on the set of states (positive unit-trace operators) $\mathfrak{S}(H)$ in a Hilbert space H is said to be a quantum channel if Φ^* is completely positive ([7]). The channel Φ is called bistochastic if $\Phi(\frac{1}{d}I_H) = \frac{1}{d}I_H$. Here and in the following we denote by d and I_H the dimension of H, $dim H = d < +\infty$, and the identity operator in H, respectively.

Given a quantum channel Φ in a Hilbert space H put ([10])

$$\hat{H}_{\Phi}(\rho) = \min_{\rho_{av} = \rho} \sum_{j=1}^{k} \pi_j S(\Phi(\rho_j)), \tag{1}$$

where $\rho_{av} = \sum_{j=1}^{k} \pi_{j} \rho_{j}$ and the minimum is taken over all probability distributions $\pi = \{\pi_i\}$ and states $\rho_i \in$ $\mathfrak{S}(H)$. Here and in the following $S(\rho) = -Tr(\rho \log \rho)$ is the von Neumann entropy of a state ρ . The strong superadditivity conjecture states that

$$\hat{H}_{\Phi \otimes \Psi}(\rho) \ge \hat{H}_{\Phi}(Tr_K(\rho)) + \hat{H}_{\Psi}(Tr_H(\rho)), \tag{2}$$

 $\rho \in \mathfrak{S}(H \otimes K)$ for two channels Φ and Ψ in Hilbert spaces H and K, respectively. Start-ex.

- (1)Siewissennichts. (GO FICTION 007)
 - youb. But don'tknowanything. (ETrans FICTION 007)
- (2)a. Die Frauen hat das nicht gerade zimperlich gemacht. (GO FICTION 007)

- b. The womenweren't exactly prudes.(ETrans FICTION 007)
- (3)a. Wenn wir also in diesem Sinne unseren Interessen und Werten dienen wollen, dann muss Europa erstens wachsam gegenüber den neuen Bedrohungen sein, denen die freien und offenen Gesellschaften ausgesetzt sind. (GO SPEECH 010)
 - b. So if we want to serve our interests and values in line with this definition, Europe must: firstly, be vigilant to the new threats to which the free and open societies are exposed. (ETrans SPEECH 010)
- a. Every country has its own political issues and this makes resolution of our disputes increasingly difficult. (EO SPEECH 009)
 - b. Jedes Landhatseineeigenenpoli-Anliegen, wodurchStreitschlichtung zunehmend erschwert wird. (GTrans _SPEECH_009)
- a. And if the EU does as it has in the past, and provides financing to Airbus at belowmarket rates of return, we could be facing a very large and highly contentious fight in the WTO. (EO SPEECH 009)
 - b. Und wenn die EU sich wie in der Vergangenheit verhält und dem Airbus Finanzierung zu Zinssätzen unter den auf dem Markt aültigen bietet, könnte uns ein großer und sehr kontroverser Kampf in der WTO bevorstehen. (GTrans SPEECH 009)

The infimum of the output entropy of a quantum chan-

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nel Φ is defined by the formula

$$S_{min}(\Phi) = \inf_{\rho \in \mathfrak{S}(H)} S(\Phi(\rho)). \tag{3}$$

The additivity conjecture for the quantity $S_{min}(\Phi)$ states ([9])

$$S_{min}(\Phi \otimes \Psi) = S_{min}(\Phi) + S_{min}(\Psi) \tag{4}$$

for an arbitrary quantum channel Ψ . It was shown in ([10]) that if the strong superadditivity conjecture holds, then the additivity conjecture for the quantity S_{min} holds too. Nevertheless the conjecture (2) is stronger than (3).

In the present paper we shall prove the strong superadditivity conjecture for the quantum depolarizing channel for all dimensions of H.

II. THE ESTIMATION OF THE OUTPUT ENTROPY

Our approach is based upon the estimate of the output entropy proved in [13]. Combining formulae (111) and (112) in [13] we get the lemma formulated below.

Lemma. Let $\Phi_{dep}(\rho) = (1-p)\rho + \frac{p}{d}I_H$, $\rho \in \mathfrak{S}(H)$, $0 \le p \le \frac{d^2}{d^2-1}$, be the quantum depolarizing channel in the Hilbert space H of the dimension d. Then, for any quantum channel Ψ there exist the orthonormal basis $\{e_s, 1 \le s \le d\}$ in H and d states $\rho_s \in \mathfrak{S}(K)$, $1 \le s \le d$, such that

$$S((\Phi_{dep} \otimes \Psi)(\rho)) \ge -(1 - \frac{d-1}{d}p)\log(1 - \frac{d-1}{d}p) - (5)$$

$$\frac{d-1}{d}p\log\frac{p}{d} + \frac{1}{d}\sum_{s=1}^{d}S(\Psi(\rho_s))$$

and

$$\frac{1}{d}\sum_{s=1}^{d}\rho_{s}=Tr_{H}(\rho),$$

where $\rho \in \mathfrak{S}(H \otimes K)$, $\rho_s = dTr_H((|e_s> < e_s| \otimes I_K)\rho) \in \mathfrak{S}(K)$, $1 \leq s \leq d$.

In the present paper our goal is to prove the following theorem.

Theorem.Let Φ_{dep} be the quantum depolarizing channel in the Hilbert space of the dimension d. Then, for an arbitrary quantum channel Ψ in a Hilbert space K the strong superadditivity conjecture holds, i.e.

$$\hat{H}_{\Phi_{dep}\otimes\Psi}(\rho) \ge \hat{H}_{\Phi_{dep}}(Tr_K(\rho)) + \hat{H}_{\Psi}(Tr_H(\rho)).$$
 (6)

Proof. Suppose that

$$\rho = \sum_{j=1}^{k} \pi_j \rho_j \tag{7}$$

and the states ρ_j , $1 \le j \le k$, form the optimal ensemble for (1) in the sense that

$$\hat{H}_{\Phi_{dep} \otimes \Psi}(\rho) = \sum_{j} \pi_{j} S((\Phi_{dep} \otimes \Psi)(\rho_{j}))$$
 (8)

Applying (5) to each element of the sum in (8) we get

$$\hat{H}_{\Phi_{dep}\otimes\Psi}(\rho) \ge -\left(1 - \frac{d-1}{d}p\right)\log(1 - \frac{d-1}{d}p) - \tag{9}$$

$$\frac{d-1}{d}p\log\frac{p}{d} + \frac{1}{d}\sum_{i=1}^{k}\pi_{i}\sum_{s=1}^{d}S(\Psi(\rho_{js})),$$

where $\rho_{js} = dTr_H((|e_{js}\rangle < e_{js}| \otimes I_K)\rho_j) \in \mathfrak{S}(K), \ 1 \leq j \leq d$, and each the set $\{e_{js}, \ 1 \leq s \leq d\}$ forms the orthonormal basis of H for $1 \leq j \leq k$.

It follows from Lemma that

$$\frac{1}{d} \sum_{j=1}^{k} \pi_j \sum_{s=1}^{d} \Psi(\rho_{js}) = \sum_{j=1}^{k} \pi_j \Psi(Tr_H(\rho_j)) = \Psi(Tr_H(\rho)).$$
(10)

The equality (10) results in

$$\frac{1}{d} \sum_{j=1}^{k} \pi_{j} \sum_{s=1}^{d} S(\Psi(\rho_{js})) \ge \hat{H}_{\Psi}(Tr_{H}(\rho)). \tag{11}$$

Notice that the quantity (1) is always bounded from below by the quantity (3). For the quantum depolarizing channel Φ_{dep} (1) coincides with (3) for any state because (3) is achieved on any pure input state due to the covariance property of Φ_{dep} . Thus, we get

$$\hat{H}_{\Phi_{dep}}(\rho) = -(1 - \frac{d-1}{d}p)\log(1 - \frac{d-1}{d}p) - \tag{12}$$

$$\frac{d-1}{d}p\log\frac{p}{d} = S_{min}(\Phi_{dep})$$

for any state $\rho \in \mathfrak{S}(H)$. Taking into account (9),(11) and (12) we get

$$\hat{H}_{\Phi_{den}\otimes\Psi}(\rho) \geq \hat{H}_{\Phi_{den}}(Tr_K(\rho)) + \hat{H}_{\Psi}(Tr_H(\rho)),$$

 $\rho \in \mathfrak{S}(H \otimes K)$. Thus, the strong superadditivity conjecture for the quantum depolarizing channel is proved.

III. CONCLUSION

At the first time the additivity conjecture (4) for the quantum depolarizing channel was proved in [13]. The method was based upon the estimation of l_p -norms of the channel. On the other hand in the papers [1–3] it was shown that the decreasing property of the relative entropy also can be used to prove the additivity conjecture

for some partial cases at least. In the present paper we have proved that the estimation of the output entropy obtained in [13] allows to prove the strong superadditivity conjecture (2) for the quantum depolarizing channel. One of a possible basis for considering the strong superadditivity conjecture can be drawn from the paper [10]. There was presented the proof of the global equivalence of the additivity conjecture for the constrained channels and the strong superadditivity conjecture.

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