# Dave-Haller's Weyl law and the tangent groupoid

Edward McDonald (Penn State University)

University of Wollongong

September 4, 2024

#### Introduction

This talk is based on none of my own work, instead I want to advertise the following two papers:

E. Van Erp and R. Yuncken, A groupoid approach to pseudodifferential calculi. *J. Reine Angew. Math.* **756** (2019), 151–182.

and

S. Dave and S. Haller, The heat asymptotics on filtered manifolds *J. Geom. Anal.* **30** (2020), no. 1, 337–389.

## Plan for this talk

- Carnot-Caratheodory geometry
- The tangent groupoid (of Connes)
- The H-tangent groupoid (of van Erp and Yuncken)
- The Volterra calculus
- Dave-Haller's Weyl law for H-elliptic operators.

In the classic game of asteroids, a player controls a spaceship moving on a two dimensional toroidal space,  $\mathbb{R}/\mathbb{Z} \times \mathbb{R}/\mathbb{Z}$ . There are two controls available:

- The spaceship can be rotated
- The spaceship can be moved forward along the direction it is facing.

The configuration space of the game is the three-dimensional torus  $\mathbb{T}^3 = (\mathbb{R}/\mathbb{Z})^3$ , with coordinates  $(x,y,\theta)$ , where (x,y) is the position of the spaceship and  $\theta$  is its angle.

The controls of the game allow us to move along the vector fields

$$X = \cos(\theta)\partial_x + \sin(\theta)\partial_y, \quad Y = \partial_\theta.$$

The configuration space of the game is the three-dimensional torus  $\mathbb{T}^3 = (\mathbb{R}/\mathbb{Z})^3$ , with coordinates  $(x,y,\theta)$ , where (x,y) is the position of the spaceship and  $\theta$  is its angle.

The controls of the game allow us to move along the vector fields

$$X = \cos(\theta)\partial_x + \sin(\theta)\partial_y, \quad Y = \partial_\theta.$$

The player moves the spaceship along a path which is parallel to the span of X and Y. That is, the path of the spaceship in configuration space is  $\{\gamma(t)\}_{t\geq 0}$ , where

$$\dot{\gamma}(t) \in \operatorname{span}\{X_{\gamma(t)}, Y_{\gamma(t)}\}, \quad t \ge 0.$$

The configuration space of the game is the three-dimensional torus  $\mathbb{T}^3 = (\mathbb{R}/\mathbb{Z})^3$ , with coordinates  $(x,y,\theta)$ , where (x,y) is the position of the spaceship and  $\theta$  is its angle.

The controls of the game allow us to move along the vector fields

$$X = \cos(\theta)\partial_x + \sin(\theta)\partial_y, \quad Y = \partial_\theta.$$

The player moves the spaceship along a path which is parallel to the span of X and Y. That is, the path of the spaceship in configuration space is  $\{\gamma(t)\}_{t\geq 0}$ , where

$$\dot{\gamma}(t) \in \operatorname{span}\{X_{\gamma(t)}, Y_{\gamma(t)}\}, \quad t \ge 0.$$

Despite there being only two available directions, we can reach any point  $(x, y, \theta)$  from any other point by travelling parallel to X and Y.

In general, if we can travel parallel to X and Y then we can approximate paths along [X,Y], by the Lie-Kato-Trotter product formula

$$\exp(t[X,Y]) = \lim_{n \to \infty} (\exp(\frac{t}{n}X) \exp(\frac{t}{n}Y) \exp(-\frac{t}{n}(X+Y)) \exp(-\frac{t}{n}Y))^n.$$

But moving along [X, Y] is harder than moving along X and Y. In the asteroids example,

$$[X, Y] = \sin(\theta)\partial_x - \cos(\theta)\partial_y$$

so  $\{X,Y,[X,Y]\}$  form a basis for the tangent space to  $\mathbb{T}^3$  at every point.

Thinking about X and Y as derivations (not just as directions), we should think of X, Y as being order 1 and [X, Y] as being order 2. The operator

$$\Xi = X^2 + Y^2 = \partial_{\theta}^2 + (\cos(\theta)\partial_x + \sin(\theta)\partial_y)^2$$

is homogeneous of order 2.

## Carnot manifolds

The plane bundle  $\mathrm{span}(X,Y)\subset T\mathbb{T}^3$  is an example of a contact structure, and this leads us to what is in general called a Carnot manifold.

#### Definition

A Carnot manifold is a manifold X equipped with a filtration of sub-bundles of TX. That is, there are subbundles  $(H^j)_{j=0}^N$  such that

$$0 = H^0 < H^1 < \dots < H^N = TX$$

and if  $E \in \Gamma H^j$ ,  $F \in \Gamma H^j$ , then  $[E, F] \in \Gamma H^{j+k}$ .

We should think of the directions in  $H^{j}$  as having "order j".

# Differential operators on Carnot manifolds

We would like to understand operators coming from Carnot manifolds, such as

$$\Xi = X^2 + Y^2 = \partial_{\theta}^2 + (\cos(\theta)\partial_x + \sin(\theta)\partial_y)^2$$

from the asteroids example. What do we want to understand?

- $\Xi$  is not elliptic, but  $(1 \Xi)^{-1}$  does improve regularity somewhat. Why, and by how much?
- The spectrum of  $\Xi$  is discrete, with a sequence of eigenvalues  $0 \le \lambda(1,\Xi) \le \lambda(2,\Xi) \le \cdots$ . What is their asymptotic behaviour?

The usual way to analyse differential operators is to build a pseudodifferential calculus.

# The Connes tangent groupoid

The usual recipe for defining pseudodifferential operators on a manifold X is the following procedure:

- **1** Identify a class of symbols  $\sigma$  on  $\mathbb{R}^d$ .
- **②** Define pseudodifferential operators  $\mathrm{Op}(\sigma)$  by a quantisation formula such as

$$\operatorname{Op}(\sigma)u(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{i(x,\xi)} \sigma(x,\xi) \widehat{f}(\xi) d\xi$$

- Show that the class of pseudodifferential operators just defined is invariant under change of variables
- **3** A pseudodifferential operator on a manifold X is a linear operator  $T: C_c^\infty(X) \to C^\infty(X)$  which has smooth kernel away from the diagonal and which is pseudodifferential in every chart.

This is a little inelegant, is there a better way?

## Semiclassical quantisation

Often it is better to quantise a symbol  $\sigma$  into a whole family of operators  $\operatorname{Op}_{\hbar}(\sigma)$  depending on a parameter  $\hbar$ , by a formula such as

$$\operatorname{Op}_{\hbar}(\sigma)u(x) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{i(x,\xi)} \sigma(x,\hbar\xi) \widehat{f}(\xi) d\xi.$$

As  $\hbar\to 0,$  the noncommutative algebra of  $\hbar\text{-pseudodifferential}$  operators under operator composition is supposed to reduce to the commutative algebra of symbols under pointwise multiplication.

# Kernels of pseudodifferential operators

The Schwartz kernel of a pseudodifferential operators with symbol  $\sigma$  is given by the oscillatory integral

$$K(x,y) = (2\pi)^{-d} \int_{\mathbb{R}^d} e^{i(x-y)\cdot\xi} \sigma(x,\xi) \,d\xi$$

(in the distribution sense). If we consider the kernel  $K(\cdot,\cdot,\hbar)$  of the  $\hbar$ -quantisation, we should have

$$K(x,y,\hbar) = (2\pi\hbar)^{-d} \int_{\mathbb{R}^d} e^{i\frac{x-y}{\hbar}\cdot\xi} \sigma(x,\xi) d\xi.$$

In the limit as  $\hbar \to 0$ , what should happen is that this looks more and more like the kernel of a convolution operator. Really, the kernel K of a pseudodifferential operator on a manifold X should be thought of as a function (distribution) on the space

$$\mathbb{T}X = (TX \times \{0\}) \sqcup (X \times X \times (0, \infty)).$$

# The tangent groupoid

Let X be a manifold, and define the set

$$\mathbb{T}X = (TX \times \{0\}) \sqcup (X \times X \times \mathbb{R}^{\times}).$$

Connes invented a good topology for  $\mathbb{T}X$ , making it a manifold of dimension  $2\dim(X)+1$ . Better yet, this is a *Lie groupoid*, with range and source maps

$$r(x, y, \hbar) = (x, \hbar), \quad s(x, y, \hbar) = (y, \hbar), \quad r((x, z), 0) = s((x, z), 0) = (x, 0)$$

and composition law

$$(x, y, \hbar) \circ (y, w, \hbar) = (x, w, \hbar), \quad ((x, z), 0) \circ ((x, z'), 0) = (x, z + z', 0).$$

# The tangent groupoid

Elements of the convolution algebra of the groupoid  $\mathbb{T}X$  are distributions  $f,g\in\mathcal{D}'(\mathbb{T}X)$ , with convolution product

$$(f*g)(x,y,h) = \int_X f(x,w,h)g(w,y,h) dw, \quad (f*g)((x,z),0) = \int_{T_xX} f((x,z),0) dx$$

That is: the convolution of distributions on  $\mathbb{T}X$  looks like composition of kernels of pseudodifferential operators.

# Thank you for listening!