

# **Radio Interferometric Studies of Cool Evolved Stellar Outflows**

A dissertation submitted to the University of Dublin  
for the degree of Doctor of Philosophy

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Trinity College Dublin, September 2013

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SCHOOL OF PHYSICS  
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## **Summary**

You should write a nice summary here...

*A dedication if you wish...*

## **Acknowledgements**

Some sincere acknowledgements...

# List of Publications

## Refereed

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# 1

## Introduction

Here is the introduction of the thesis, complete with a few references ([Prothero & Buell, 2007](#); [Sagan, 1997](#)). Section 1.1 contains Equation 1.1, Section 1.2 has Figure 1.1 and Section 2.3 has Table ???. Chapter 2 has pretty much nothing in it.

### 1.1 First Section

This section has an equation. Here it is:

$$L_{\odot} = 4\pi R_{\odot}^2 \sigma T_e^4 \quad (1.1)$$

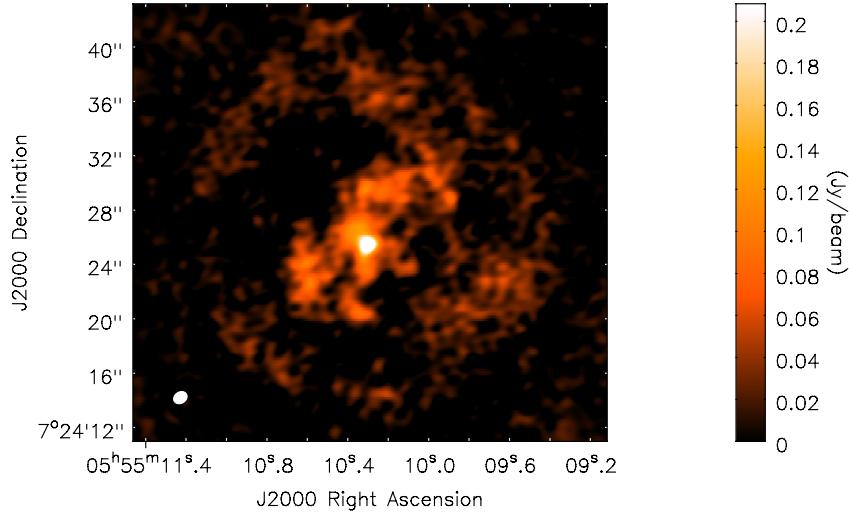
which is a nice way of describing the luminosity.

### 1.2 Second Section

So this section has a figure in it<sup>1</sup>. That figure depicts the basic structure of a red giant.

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<sup>1</sup>And also a footnote.



**Figure 1.1:** Red Giant and Asymptotic Giant Branch Stars. The left side of the figure shows the basic structure of a star on the giant branch of the HR diagram, while the right side shows a similar star after it has evolved to ascend the asymptotic giant branch. *Image Credit: Australian Telescope National Facility.*

### 1.3 Second Section

# 2

## Introduction to Radio Interferometry

The poor spatial resolution provided by a single dish radio antenna can cause difficulties in obtaining accurate flux density measurements of radio astronomical sources, especially at long wavelengths. A single dish radio antenna is unable to distinguish against background radio emitters located in the primary beam, and therefore the observed flux density can contain emission from unrelated sources. This limitation can be overcome through interferometry. An interferometer acts as a spatial filter, and can discriminate against smooth backgrounds, while its higher resolution allows separation of the target from nearby confusing sources. This chapter describes the basic elements of a radio interferometer, what happens to the radio signal before it is sent for correlation, introduces the general concept of interferometry, and explains the process of synthesis imaging.

### 2.1 Radio Antenna Fundamentals

The quality and properties of the final radio image produced from a synthesis array are partially dependent on the properties of the individual antennas in the array. The most important such properties are discussed in the following sections and include aperture size, aperture efficiency, pointing accuracy and sidelobe

level. We define the radio antenna as the piece of equipment which converts the electromagnetic waves emitted from the observed source into an electric current ready to be input into to the first low noise amplifier where the signal is at the radio/sky frequency,  $\nu_{\text{RF}}$ . What happens to the signal after this will be discussed in Section 2.2.

### 2.1.1 Properties of a Radio Antenna

The power gain of a transmitting antenna is a measure of the antenna's capability of converting power into radio waves in a specific direction. In radio astronomy, the receiving counterpart of transmitting power gain is the effective collecting area of an antenna,  $A(\nu, \theta, \phi)$ , where  $\nu$  is frequency and  $\theta$  and  $\phi$  are direction coordinates. An ideal radio antenna would collect all incident radiation from a distant point source and convert it to electrical power. The total spectral power  $P_\nu$  collected by it would then be a product of its geometric area and the incident spectral power per area, or flux density  $F_\nu$ . By analogy then, the effective area of a real radio antenna is defined

$$A(\nu, \theta, \phi) = \frac{P_\nu}{F_\nu} = \frac{P}{I(\nu, \theta, \phi)\Delta\nu\Delta\Omega} \quad (2.1)$$

where  $I(\nu, \theta, \phi)$  is the source brightness in units  $\text{W m}^{-2} \text{ Hz}^{-1} \text{ sr}^{-1}$  that the antenna is pointing at and  $P$  is the power (in Watts) received by the antenna in bandwidth  $\Delta\nu$  from element  $\Delta\Omega$  of solid angle. The normalized antenna reception pattern  $\mathcal{A}$ , often referred to as the power pattern due to the duality between receiving and transmitting, is defined as

$$\mathcal{A}(\nu, \theta, \phi) = \frac{A(\nu, \theta, \phi)}{A_0} \quad (2.2)$$

where  $A_0$  ( $\text{m}^2$ ) is often referred to as the effective area of the antenna and is the response at the center of the main lobe of  $A(\nu, \theta, \phi)$  [i.e.,  $A(\nu, 0, 0)$ ]. Then the beam solid angle,  $\Omega_A$ , of the primary beam is

$$\Omega_A = \iint_{\text{all sky}} \mathcal{A}(\theta, \phi) d\Omega \quad (2.3)$$

and is a measure of the field of view of the antenna.

## 2.1 Radio Antenna Fundamentals

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In the case of an isotropic antenna [i.e.,  $\mathcal{A}(\nu, \theta, \phi) = 1$ ], it can be shown that the product of the effective area and the primary beam solid angle is equal to the square of the wavelength ([Kraus et al., 1986](#))

$$A_0 \Omega_A = \lambda^2. \quad (2.4)$$

$\Omega_A$  has its maximum possible value of  $4\pi$  if  $\mathcal{A}$  is everywhere equal to 1. This means that the primary antenna can see the whole sky with equal sensitivity. Even though a large field of view is usually desirable in radio astronomy, [Equation 2.4](#) ensures that for any given wavelength, when  $\Omega_A$  is a maximum, the power received is a minimum and therefore the sensitivity is also at a minimum. To improve sensitivity, one could increase the collecting area of the antenna, but [Equation 2.4](#) then ensures that the field of view must decrease. Thus, when deciding on the primary antenna size in a synthesis array, there is always a trade-off between field of view and sensitivity.

In reality, an antenna cannot radiate isotropically and will radiate preferentially in one or more directions. A Fourier transform relationship exists between the complex voltage distribution of the field,  $f(u, v)$ , in the aperture of the antenna and the complex far-field voltage radiation pattern,  $F(l, m)$ , of the antenna ([Kraus et al., 1986](#))

$$F(l, m) = \iint_{\text{aperture}} f(u, v) e^{2\pi i(ul+vm)} dudv \quad (2.5)$$

and

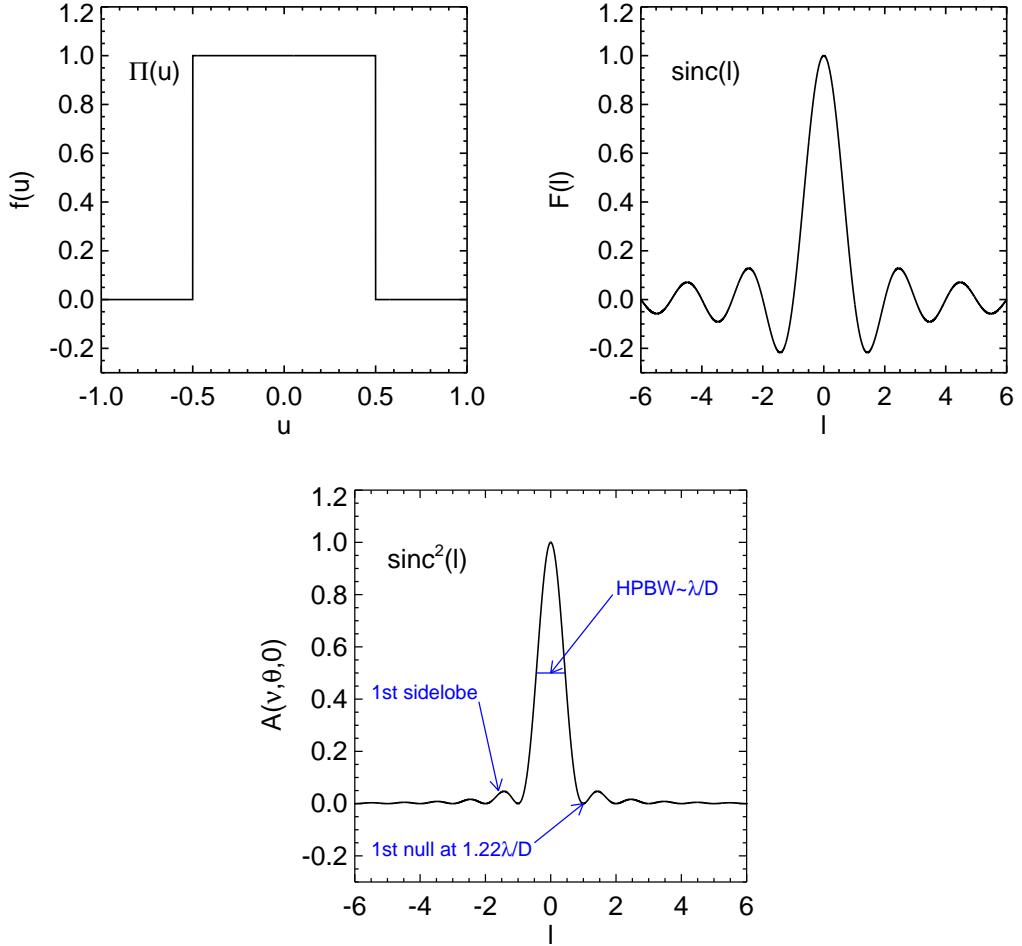
$$f(u, v) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} F(l, m) e^{-2\pi i(ul+vm)} dl dm \quad (2.6)$$

where

$$u = \sin\theta \cos\phi \quad \text{and} \quad v = \sin\theta \sin\phi \quad (2.7)$$

## 2.1 Radio Antenna Fundamentals

---



**Figure 2.1:** *Top Left:* A uniformly illuminated 1-D aperture  $f(u)$ . *Top Right:* The Fourier transform of  $f(u)$  gives the antenna radiation pattern in the far-field,  $F(l)$ . *Bottom:* The power pattern of the antenna is given by  $\mathcal{A} = |F(l)|^2$ .

are the antenna coordinates and  $l$  and  $m$  are their Fourier counterparts. The from of  $f(u, v)$  is determined by the manner in which the antenna feed illuminates the aperture. Therefore Equations 2.6 and 2.7 tell us that the radiation pattern in the far-field of a two-dimensional aperture is the two-dimensional Fourier trans-form of the aperture field illumination. For a uniformly illuminated 1-D aperture shown in Figure 2.1, the radiation pattern in the far-field is the *sinc* function. The radiation pattern in the far-field,  $F(l, m)$ , of such an antenna is related to the antenna power pattern by  $\mathcal{A} = |F(l, m)|^2$ . This power pattern is known as the Airy pattern if the antenna is uniformly illuminated and is also shown in Figure

[2.1](#). The central peak of this power pattern is called the main beam while the smaller secondary peaks are called sidelobes. The antenna is maximally sensitive to radiation from the direction of the peak of the beam, but is also slightly sensitive to radiation in the direction of the side lobes. The half-power beamwidth (HPBW) of the main beam  $\theta_{\text{HPBW}}$  is a term commonly used in the literature to describe the field of view of an antenna/interferometer and satisfies

$$\theta_{\text{HPBW}} \propto \frac{\lambda}{D} \quad (2.8)$$

where  $D$  is the diameter of the antenna. The constant of proportionality varies slightly with the illumination taper and can be shown to be equal to  $\sim 0.89$  for a uniformly illuminated 1-D aperture and  $\sim 1.2$  for most single dish radio antennas. When the sky is scanned with a single dish antenna, then this HPBW is the resolution of the resulting map.

### 2.1.2 Antenna Structural Design

The design of the primary antenna element of an interferometric array will depend on the wavelength range to be observed. In general, dipoles are used for wavelengths longer than  $\sim 1$  m, while reflector antennas are typically used at shorter wavelengths. The reason why the more simple and less expensive dipoles are not used at all wavelengths is given by Equation [2.4](#). For an isotropic antenna, this equation tells us that the effective area is just

$$A_0 = \frac{\lambda^2}{4\pi}. \quad (2.9)$$

Therefore, at short wavelengths a non-directional antenna such as a dipole, will have a small effective collecting area, giving it poor sensitivity for reception. Thus, dipoles can be used at long wavelengths as they have sufficient collecting area, but cannot be used at shorter wavelengths as an impractical amount would be needed to produce useful collecting areas. Since the interferometric arrays used in this thesis use reflector antennas, the rest of this section will focus on them.

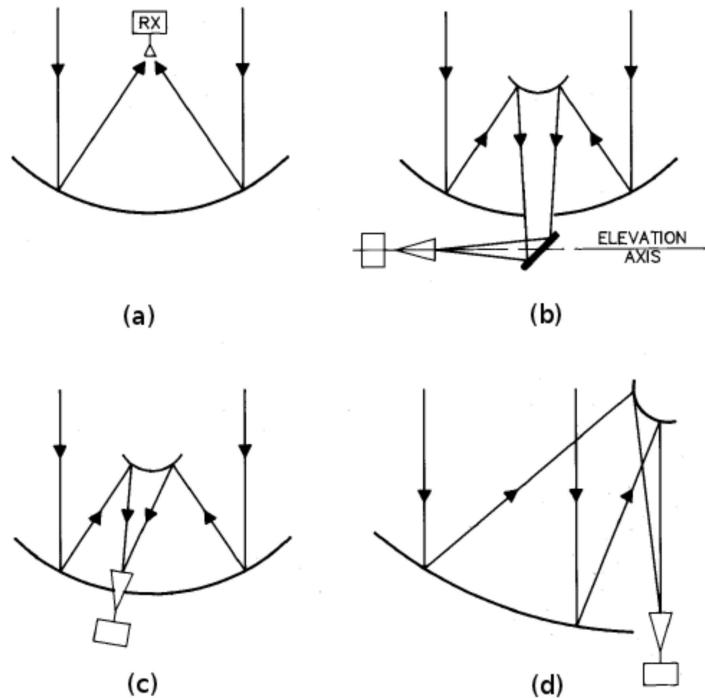
## 2.1 Radio Antenna Fundamentals

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*Choice of Antenna Mount.* Nearly all interferometric arrays consist of antennas which have altitude over azimuth (alt-azimuth) mounts. These antennas lie on a horizontal azimuth track on which the antenna can turn in azimuth, and on a horizontal elevation axle about which the antenna can change in zenith angle. The main advantage of such a design is simplicity and thus lower cost. Gravity always acts on the reflector in the same plan thus reducing the problem of keeping the reflector profile accurate during the duration an observation. However, sources close to the zenith can usually not be observed due to the high rate of azimuth rotation required. Also, the beam rotates with respect to the source for long duration observations which can affect the dynamic range of total intensity images of very large sources. The other type of mount occasionally used is the equatorial mount. Its polar axis is aligned parallel to the axis of rotation of the Earth and therefore only needs to rotate about the declination axis to observe a source. Its beam also doesn't have the beam rotation problem encountered by the alt-azimuth design and can track sources close to the zenith. Its major disadvantage and the reason for its scarce usage is the complexity of its design and resulting increased cost.

*Choice of Antenna Optics.* In Figure 2.2 we show the main optical systems which can be used to feed a large radio reflector. The prime focus system (e.g. used in the Giant Meter Radio Telescope) has the advantage that it can be used at long wavelengths where the use of secondary focus feeds (i.e., sub-reflectors) become impractical. However, access to and space for the feeds and receivers are limited, and sensitivity can be lost due to spillover noise from the ground. The other designs have the advantage of easier access to the feeds and receivers and less spillover noise from the ground. The off-axis Cassegrain (e.g. used in the Very Large Array) also has the advantage of increased frequency capability as many feeds can be located in a circle around the center of the reflector and a slight rotation of the sub-reflector is all that is required to change observing frequency. The receivers and feeds in the Naysmith geometry (e.g. used in the Combined Array for Research in Millimeter wave Astronomy) are located externally to the antenna structure. Finally, the offset Cassegrain (e.g. used in the Green Bank Telescope)

has no blockage and will have a circularly symmetric beam with low sidelobes although the increase complexity of its structure leads to increased costs.



**Figure 2.2:** Common optical systems used for radio antennas. (a) Prime focus, (b) Naysmith, (c) Off-axis Cassegrain, (d) Offset Cassegrain [Figure adapted from Taylor *et al.* (1999)].

### 2.1.3 Antenna Performance Parameters

*Aperture Efficiency.* The geometric collecting area of a parabolic antenna  $A_{\text{geo}}$  ( $= \pi D^2/4$ ) is related to the effective area (i.e., the collecting area when pointing directly at a source) via the dimensionless quantity  $\eta$  ( $\eta < 1$ ) known as the aperture efficiency where

$$\eta = \frac{A_0}{A_{\text{geo}}}. \quad (2.10)$$

The aperture efficiency directly impacts on the sensitivity of the interferometric array and can be defined as the product of a number of different efficiency loss

## 2.1 Radio Antenna Fundamentals

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factors,

$$\eta = \eta_{\text{sf}} \eta_{\text{bl}} \eta_s \eta_t \eta_{\text{misc}}. \quad (2.11)$$

The surface efficiency  $\eta_{\text{sf}}$  accounts for the aperture efficiency loss as a result of reflector profile inaccuracies. Such inaccuracies result in the electric field from various parts of the aperture not adding together in phase at the feed leading to a decrease in power. The aperture blockage efficiency  $\eta_{\text{bl}}$  accounts for the fact that the sub-reflector (or feed) and its support structure result in a reduction in the incident radiation on the antenna. The feed spillover efficiency  $\eta_{\text{bl}}$  is best understood if the antenna is considered in transmission rather than in reception mode, and is defined as the fraction of power radiated by the feed that is intercepted by the reflector for a prime focus system, or by the sub-reflector for a Cassegrain system. The illumination taper efficiency  $\eta_t$  accounts for the fact that the feed pattern does not illuminate the primary reflector uniformly but illuminates the outer part of the reflector at a lower level than the inner part. Finally, the miscellaneous efficiency losses such as reflector diffraction and feed position phase errors are accounted for in  $\eta_{\text{misc}}$ . As an example, the total aperture efficiency of the VLA antennas can vary between 0.65 and 0.2 at 6 and 0.7 cm, respectively.

*Pointing Accuracy.* The main lobe of an antenna's power pattern will usually not point exactly in the desired direction due to gravity deformations, wind pressure deformations, and mechanical inaccuracy. The angular offset,  $\Delta\theta$ , between the actual and desired pointing direction is called the pointing error. Usually, the desirable pointing error of an antenna at the highest operational frequency is  $\Delta\theta < \theta_{\text{HPBW}}/20$  ([Taylor et al., 1999](#)). With this specification reached, an antenna pointing at a compact source will suffer negligible intensity variations as  $\mathcal{A}(\theta_{\text{HPBW}}/20) > 0.99$ . However, this pointing error of only  $\theta_{\text{HPBW}}/20$  will still have a substantial effect on the accuracy of the outer image. For example, a source located at the half power point will suffer a substantial fractional intensity variation of  $2\mathcal{A}(\theta_{\text{HPBW}}/2 + \theta_{\text{HPBW}}/20) \simeq 0.86$ . The blind pointing of a VLA antenna is only about  $10''$  and can be much worst in daytime, occasionally exceeding  $1'$ . This means that at Q-band (45 GHz; 0.7 cm), which is the highest observing frequency on the VLA, the pointing error is only at best  $\theta_{\text{HPBW}}/6$ , and

---

## 2.2 Receiving System for an Interferometer

at worst  $> \theta_{\text{HPBW}}$ , meaning that the target may lie outside of the primary beam. To overcome this problem of large antenna pointing errors at high frequencies with the VLA, a technique known as referenced pointing is implemented. This technique will be discussed further in Chapter 3.

## 2.2 Receiving System for an Interferometer

A radiometer (a radio receiver) is a device used to measure the timed-averaged power of the noise coming from a radio telescope within a well-defined radio frequency (RF) range,  $\nu_{\text{RF}} - \Delta\nu_{\text{RF}}/2 \rightarrow \nu_{\text{RF}} + \Delta\nu_{\text{RF}}/2$ , where  $\Delta\nu_{\text{RF}}$  is the bandwidth of the receiver and  $\Delta\nu_{\text{RF}} < \nu_{\text{RF}}$ . The simplest radiometer carries out the following tasks:

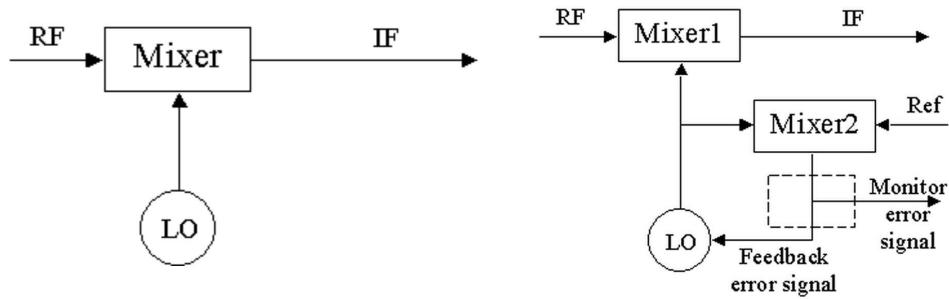
1. Filters the broadband noise coming from the antenna via a bandpass filter.
2. Multiplies the filtered voltage by itself (i.e., its output voltage is proportional to its input power).
3. Smooths out the rapidly fluctuating output of the detected voltage via a signal averager or integrator.
4. Measures the smoothed voltage.

In practice, radiometers are never as simple as those described above and nearly all practical radiometers are *superheterodyne* receivers which incorporate a number of additional steps to produce an output voltage.

The RF front end is the term used to describe all the circuitry between the feed horn and the lower intermediate frequency (IF) stage. The first task of the front end is to amplify the received signal. The radio signals we want to measure are generally very weak and therefore need to be initially amplified by many orders of magnitude so they are above the noise level in succeeding stages. However, the front end electronic components produce random electrical noise which will also get amplified by this large amount. Therefore the role of the pre-amplifier is to amplify the incoming signal from the antenna while adding as little noise as possible. For this reason, the pre-amplifier is often called a low noise amplifier

## 2.2 Receiving System for an Interferometer

(LNA) and are often cooled to very low temperatures to minimize the amount of noise contributed by the components.



**Figure 2.3:** *Left:* Block diagram of a simple superheterodyne receiver. The amplified RF signal is mixed with a signal from a local oscillator to convert the signal to the more manageable intermediate frequency. *Right:* For interferometry, a phase lock loop is used to ensure all antennas are locked to the same frequency.

The amplified RF signal is then sent through a mixer which multiplies the RF signal by a sine wave of frequency  $\nu_{\text{LO}}$  which is generated by a local oscillator (LO), as shown in Figure 2.3. The effect of this is that RF signal is changed to a lower frequency which can be easier handled by the IF amplifier and improves frequency selectivity. Mathematically, the mixer does the following

$$2\sin(2\pi\nu_{\text{LO}}t) \times \sin(2\pi\nu_{\text{RF}}t) = \cos[2\pi(\nu_{\text{LO}} - \nu_{\text{RF}})t] - \cos[2\pi(\nu_{\text{LO}} + \nu_{\text{RF}})t], \quad (2.12)$$

and produces two additional outputs, one at the input signal frequency minus the local oscillator frequency, and one at the sum of these frequencies. The lower of the two outputs called the intermediate frequency (IF) is taken by passing the mixer output through a filter in the IF amplifier. In interferometry, where the signals from two antennas are correlated, it is crucial that the receivers from both antennas are operating at the same frequency to control the phase difference between them. This is achieved by using a phase lock system, whose block diagram is laid out in Figure 2.3. In this system another mixer compares the LO to a reference frequency, which is the same for all antennas. Any existing phase error results in an error signal that is sent back to the oscillator so that its frequency can be adjusted to maintain exact frequency tuning. After this, the IF can finally be sent to the radiometer and recorded.

## 2.3 Fundamentals of Radio Interferometry

The angular resolution  $\Delta\theta$  of a radio antenna is the minimum angular separation which two point sources can have in order to be recognized as separate objects. The *Rayleigh criterion* is the generally accepted criterion for defining the angular resolution of a filled circular aperture of diameter  $D$ , at the observational wavelength  $\lambda$  and is given as

$$\Delta\theta = 1.22 \frac{\lambda}{D} \text{ rad.} \quad (2.13)$$

The Rayleigh criterion states that two objects are resolved when the first null of the diffraction pattern of one object coincides with the maximum of the diffraction pattern of the other. An immediate consequence of Equation 2.13 is that at large wavelengths, the angular resolution becomes large unless the diameter of the aperture can be increased substantially. In order to achieve modest angular resolution at radio wavelengths with a single radio antenna then, the diameter becomes impractically large. For example, in order to achieve an angular resolution of  $1''$  at 6 cm a 12 km aperture would be required. Radio interferometry is a technique used in radio astronomy to overcome this problem of poor angular resolution at long wavelengths.

### 2.3.1 Young's Slits

The basic principles of interferometry can be understood through Young's double-slit experiment. If coherent radiation emitted from a distant point source propagates through two slits, an illumination pattern composed of bright and dark fringes is observed. The phenomenon is a result of the constructive and destructive interference between the secondary waves produced by the slits. The fringe separation is  $\lambda/B$ , where  $B$  is the projected separation of the slits and is called the baseline. The fringe contrast which is historically known as the fringe visibility,  $V$ , can be written as

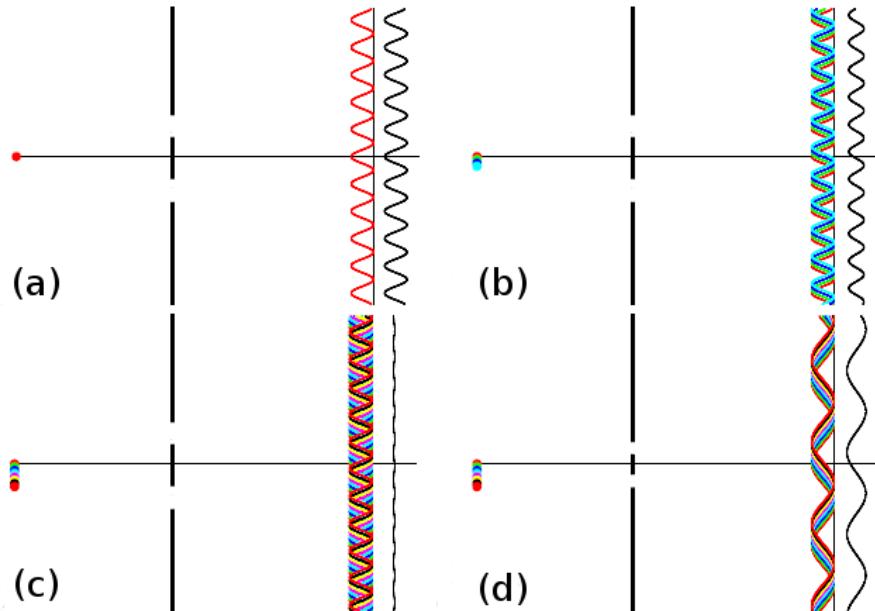
$$|V| = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}} \quad (2.14)$$

where  $I_{\max}$  and  $I_{\min}$  are the maximum and minimum intensity of the fringes, respectively. In other words, the fringe visibility is the fringe amplitude normalized

## 2.3 Fundamentals of Radio Interferometry

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by the sum of the maximum and minimum intensity.



**Figure 2.4:** The resulting fringe pattern produced by Young's slits under various conditions. The source is shown on the left of the slits in each panel, while the separate fringe patterns (colors) along with the added fringe pattern (black) is shown on the right of the slits. (a) Point source at infinity, (Visibility = 1). Fringes are separated by an angular distance of  $\lambda/B$ . (b) An increase in source size results in a drop in visibility. (c) When the source size is equal to  $\lambda/B$ , the visibility is zero. (d) If the source size remains the same as in (c) and the slit spacing is reduced, then the fringes re-appear. Figure adapted from [Jackson \(2008\)](#).

In the simple case shown in Figure 2.4a, the angular size of the source is  $\ll \lambda/B$  and the fringe visibility is 1. In interferometry, this equates to the situation in which the source size is smaller than the synthesized beam and only an upper limit of the source size can be obtained (i.e., the source is unresolved). In Figure 2.4b, the angular size of the source is now larger and can be thought of as a sequence of point sources each emitting radiation which is uncorrelated with emission from the others. An angular shift of  $\phi$ , called phase, in the sources position results in a shift in the corresponding fringe pattern by the same angle the other way. The total interference intensity pattern is then just the sum of these individual patterns and the visibility is reduced. When the extension of the

## 2.3 Fundamentals of Radio Interferometry

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source equals  $\lambda/B$ , the fringes disappear and give a constant illumination pattern. In this case the fringe visibility is zero and the source is completely resolved as shown in Figure 2.4c. Finally, if the source size is the same as that in Figure 2.4c but the slit separation is reduced, then the fringe separation  $\lambda/B$  will again increase as shown in Figure 2.4d. This is because the source now produces much less displacement of the fringe patterns as a fraction of the fringe separation. In interferometry, this result means that extended sources can only be probed with short baselines.

Visibility and phase are often expressed together as the complex visibility  $V = |V|e^{i\phi}$ , which completely defines a pattern of interference fringes. Young's double-slit experiment demonstrates a fundamental property of interferometry, namely that the contrast of fringes is a function of the geometry of the source. The results of the experiment are summarized in Figure 2.5. The top row shows that a large source (i.e., one whose intensity distribution extends out to a large angle on the sky) has a fringe visibility pattern which falls off quickly as projected baseline length increases. The bottom row shows that for compact sources the fringe visibility remains high out to large baselines. In the following section we will show that the relationship between the sky brightness distribution  $I(\theta)$  and the visibility  $V(B)$  is a Fourier transform.

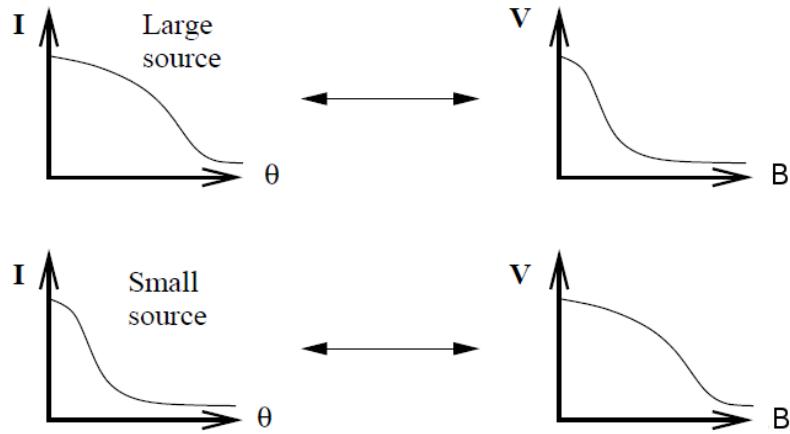
### 2.3.2 The Two-element Interferometer

Interferometers with  $N$  antennas can be treated as  $N(N - 1)/2$  independent interferometer pairs so it is worthwhile studying the simplest case of the two-element interferometer. A simplified block diagram of the components of such an interferometer is shown in Figure 2.6. The figure shows two identical antennas separated by a baseline vector  $\mathbf{b}$  pointing towards a distant radio source in a direction indicated by the unit vector  $\mathbf{s}$ . The plane waves from the distant radio source reach antenna 1 at a time  $\tau_g$  later than they reach antenna 2.  $\tau_g$  is called the geometric delay and is given by

$$\tau_g = \frac{\mathbf{b} \cdot \mathbf{s}}{c} = \frac{b \cos \theta}{c} \quad (2.15)$$

## 2.3 Fundamentals of Radio Interferometry

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**Figure 2.5:** *Left column:* Intensity distribution as a function of sky angle for an extended source (top) and for a more compact source (bottom). *Right column:* The corresponding fringe visibility as a function of slit separation or baseline. Figure adapted from [Jackson \(2008\)](#).

where  $c$  is the speed of light. If we assume that the interferometer only responds to a very narrow band centered on frequency  $\nu = \omega/2\pi$ , then the output voltages of antennas 1 and 2 at time  $t$  can be written as

$$V_1(t) = V \cos[\omega(t - \tau_g)] \quad \text{and} \quad V_2(t) = V \cos(\omega t). \quad (2.16)$$

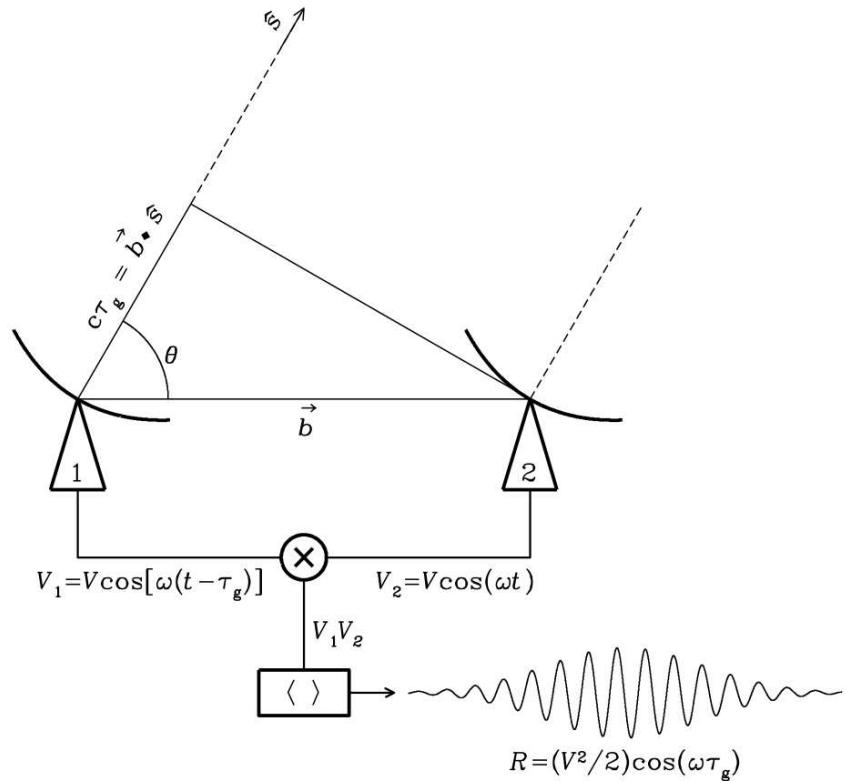
The signals are then passed through a correlator which first multiplies these voltages to give

$$V_1(t)V_2(t) = \frac{V^2}{2} [\cos(2\omega t - \omega\tau_g) + \cos(\omega\tau_g)] \quad (2.17)$$

and then averages them over a time interval  $\Delta t$  which is long enough such that  $\Delta t \gg (2\omega)^{-1}$  to give the final output  $R$ :

$$R = \langle V_1(t)V_2(t) \rangle = \frac{V^2}{2} [\cos(\omega\tau_g)]. \quad (2.18)$$

As the Earth rotates,  $\tau_g$  varies slowly with time and the resultant oscillations in the correlator output voltage represent the motion of the source. These



**Figure 2.6:** Simplified schematic diagram of a two-element interferometer. The correlator multiplies and averages the voltage outputs  $V_1$  and  $V_2$  of the two dishes and yields an output amplitude  $V^2/2$  which is proportional to the point-source flux density  $F_\nu$ . *Image Credit:* National Radio Astronomy Observatory.

sinusoidal oscillations are called fringes, and the fringe phase is

$$\phi = \omega\tau_g = \frac{\omega b \cos\theta}{c} \quad (2.19)$$

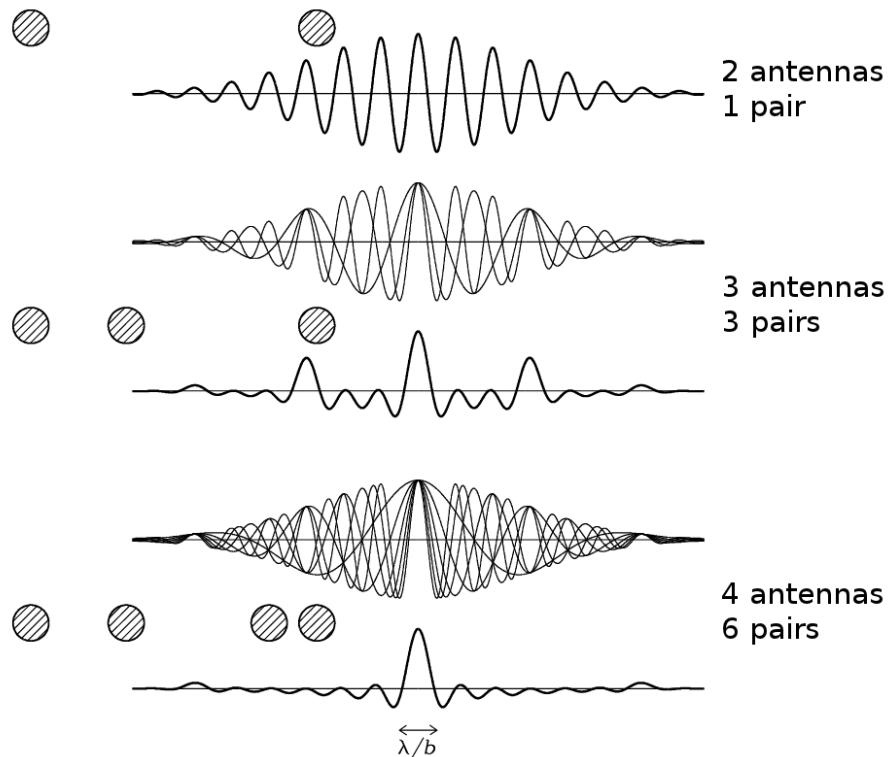
which changes with source direction as follows

$$\frac{d\phi}{d\theta} = \frac{\omega b \sin\theta}{c} = 2\pi \left( \frac{b \sin\theta}{\lambda} \right). \quad (2.20)$$

The fringe phase completes a full period (i.e.,  $\Delta\phi = 2\pi$ ) when an angular change  $\Delta\theta = (\lambda/b \sin\theta)$  occurs. This tells us that the fringe phase is an extremely sensitive measure of source position if the projected baseline  $b \sin\theta$  is many wavelengths long

## 2.3 Fundamentals of Radio Interferometry

and is the reason why interferometers can determine the positions of compact radio sources with exquisite accuracy.



**Figure 2.7:** The instantaneous point source responses of an interferometer with two, three and four elements is indicated by the thick curves. The individual responses of the three pairs of two-element interferometers of the three-element interferometer and the six pairs of two-element interferometers of the four-element interferometer are plotted as thin curves. The main beam of the four-element interferometer is nearly Gaussian and has a width of  $\sim \lambda/B$ . This is known as the instantaneous synthesized beam of the interferometer. *Image Credit:* National Radio Astronomy Observatory.

If the antennas in an interferometric array are isotropic, then the point-source response of the interferometer would be a sinusoid spanning the entire sky and the interferometer would be only sensitive to one Fourier component of the sky brightness distribution having angular period  $\lambda/b\sin\theta$ . The response of a two-element interferometer  $R$  with non-isotropic antennas is this sinusoid multiplied by the product of the voltage patterns (i.e., defined as  $f(u, v)$  in Section 2.1.1)

## 2.3 Fundamentals of Radio Interferometry

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of the individual antennas. If the antennas are identical then this product is the power pattern of the individual antennas called the primary beam. The primary beam is usually a Gaussian that is much wider than the fringe period as  $D \ll b\sin\theta$ . The result is that an interferometer with directive antennas responds to a finite range of angular frequencies centered on  $b\sin\theta/\lambda$ . The instantaneous point source response of an interferometer is known as the synthesized beam and is the point source response obtained by averaging the outputs of all antenna pairs. The synthesized beam of an interferometer is an important quantity as it defines the maximum angular resolution of the instrument. The synthesized beams produced by an interferometer with a various number of antennas arranged in 1-D is shown in Figure 2.7. The figure shows that the synthesized beam can be improved by acquiring more Fourier components (i.e., baselines) and rapidly approaches a Gaussian as  $N$  increases. However, sidelobes are still significant and a broad negative “bowl” exists between the main beam and the first sidelobes due to the absence of short spacings.

### 2.3.3 Complex Visibility

The interferometer output can be expressed in terms of the radio brightness over the sky, which is sometimes also called specific intensity and has units  $\text{W m}^{-2} \text{Hz}^{-1} \text{sr}^{-1}$ . If the radio brightness of a spatially incoherent extended source in the direction of unit vector  $\mathbf{s}$  is  $I(\mathbf{s})$ , then the response of the two-element interferometer with “cosine” correlator output near frequency  $\nu = \omega/2\pi$  is obtained by treating the extended source as the sum of independent point sources:

$$R_c = \int_{\Omega} \mathcal{A}(\mathbf{s}) I_{\nu}(\mathbf{s}) \cos\left(\frac{2\pi\mathbf{b}\cdot\mathbf{s}}{\lambda}\right) d\Omega \quad (2.21)$$

where  $\mathcal{A}$  is the normalized antenna reception pattern defined in Section 2.1.1 and we call  $\mathcal{A}(\mathbf{s}) I_{\nu}(\mathbf{s})$  the modified brightness distribution. However, the cosine function in the “cosine” correlator output is only sensitive to the even part of the sky brightness distribution, which can be written as the sum of even and odd parts:

$$I(\mathbf{s}) = I_e(\mathbf{s}) + I_o(\mathbf{s}). \quad (2.22)$$

## 2.3 Fundamentals of Radio Interferometry

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A “sine” correlator whose output is odd, is needed to detect the odd part of  $I(\mathbf{s})$  and this is implemented by inserting a  $90^\circ$  phase delay into the signal of one of the antennas to give

$$R_s = \int_{\Omega} \mathcal{A}(\mathbf{s}) I_{\nu}(\mathbf{s}) \sin\left(\frac{2\pi\mathbf{b}\cdot\mathbf{s}}{\lambda}\right) d\Omega \quad (2.23)$$

It is convenient to write the cosines and sines as complex exponentials using the identity

$$e^{i\phi} = \cos(\phi) + i\sin(\phi) \quad (2.24)$$

and so the combination of “cosine” and “sine” correlators is called a “complex” correlator. The term *visibility* was first introduced by [Michelson \(1890\)](#) to describe the relative amplitudes of the optical fringes that he observed. The visibility is a complex quantity in radio astronomy and has dimensions of spectral power flux density ( $\text{W m}^{-2} \text{ Hz}^{-1}$ ). The complex visibility is the response of a two-element interferometer with a complex correlator to an extended source with brightness distribution  $I(\mathbf{s})$  and is defined as

$$V \equiv R_c - iR_s, \quad (2.25)$$

i.e.,

$$V_{\nu} = \int_{\Omega} \mathcal{A}(\mathbf{s}) I_{\nu}(\mathbf{s}) e^{-2\pi i \nu \mathbf{b} \cdot \mathbf{s} / c} d\Omega. \quad (2.26)$$

### 2.3.4 Coordinate Systems for Imaging

The baseline vector  $\mathbf{b}$  has coordinates  $(u, v, w)$  in three dimensions shown in Figure 2.8 where  $w$  points in the directions of interest, i.e., towards a direction  $s_0$  that becomes the center of the synthesized image.  $u$ ,  $v$  and  $w$  are measured in wavelengths (i.e., the components of  $\mathbf{b}/\lambda$ ) and have directions towards the East, the North and the phase tracking center, respectively. An arbitrary unit vector  $\mathbf{s}$  has components  $(l, m, n)$  called direction cosines, where  $n = \cos\theta =$

## 2.4 Synthesis Imaging

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$(1 - l^2 - m^2)^{1/2}$ . Using these coordinates the parameters in Equation 2.26 become

$$\begin{aligned} \frac{\nu \mathbf{b} \cdot \mathbf{s}}{c} &= ul + vm + wn, \\ d\Omega &= \frac{dl dm}{n} = \frac{dl dm}{\sqrt{1 - l^2 - m^2}}. \end{aligned} \quad (2.27)$$

Therefore Equation 2.26 can be defined in terms of the coordinate system laid out in Figure 2.8 as

$$V_\nu(u, v, w) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathcal{A}_\nu(l, m) I_\nu(l, m) e^{-2\pi i [ul + vm + w(\sqrt{1 - l^2 - m^2})]} \frac{dl dm}{\sqrt{1 - l^2 - m^2}} \quad (2.28)$$

which is not a three-dimensional Fourier transform. This equation becomes a two-dimensional Fourier transform if  $w = 0$  which is a good approximation for small field imaging, i.e., when  $|l|$  and  $|m|$  are small. In this case Equation 2.28 can be inverted to find the modified sky brightness distribution:

$$\mathcal{A}_\nu(l, m) I_\nu(l, m) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} V_\nu(u, v) e^{2\pi i (ul + vm)} du dv. \quad (2.29)$$

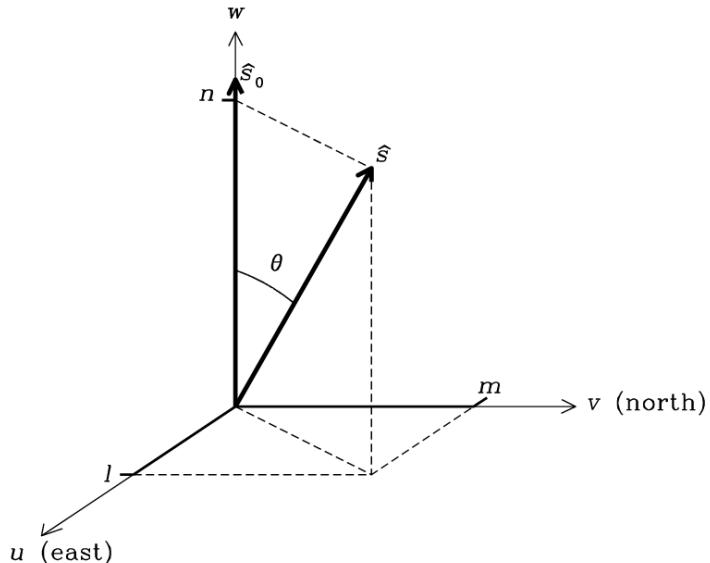
Therefore Equation 2.29 demonstrates the important Fourier Transform relationship between the sky brightness distribution and the complex visibility (i.e., the interferometer response).

## 2.4 Synthesis Imaging

A synthesis imaging telescope consists of a number radio antennas fixed on the ground and uses the Earth's rotation to vary the projected baseline coverage to increase the sampling of the  $u - v$  plane. In this section we describe how Earth-rotation aperture synthesis is used to convert the complex visibilities outputted from the correlator to a final radio image of the observed sky.

### 2.4.1 Visibility Sampling

An example of how a radio interferometer samples the  $u - v$  plane is shown in Figure 2.9. The left panel of this figure shows the overhead view of the VLA in its



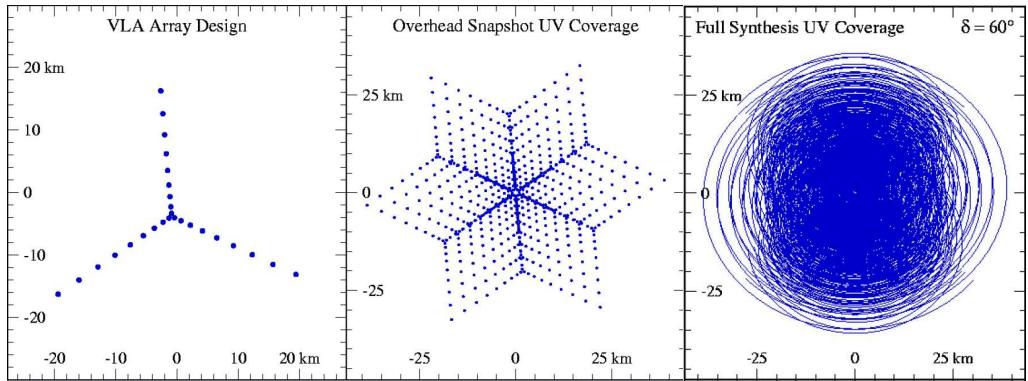
**Figure 2.8:** The  $(u, v, w)$  interferometric coordinate system.  $l$ ,  $m$ , and  $n$  are the projections of the unit vector  $\mathbf{s}$  onto the  $u$ ,  $v$ , and  $w$  axes, respectively.

most extended configuration while the other two panels shows the corresponding  $u - v$  coverage for different periods of time. We define  $u$  and  $v$  as the east-west and north-south components of the projected baseline in wavelengths, respectively. As the Earth rotates, the projected baseline of every two-element pair in the array changes, thus sampling a different part of the  $u - v$  plane. The middle panel shows that the total  $u - v$  coverage of the VLA for a very short duration track (i.e., a snapshot) results in a ‘snowflake’ like pattern, with more dense coverage in the direction of the arms of the array due to the larger number of baselines. Most radio interferometers have their own unique array configuration layout and thus produce a different ‘snapshot’  $u - v$  coverage to that shown in Figure 2.9. Over many hours, the  $u - v$  points trace out portions of ellipses and eventually after a full Earth rotation the points can trace out full ellipses as shown in the right panel of Figure 2.9.

### 2.4.2 Imaging (Making a Dirty Map)

For every sky brightness distribution  $I(l, m)$  there exists a continuous complex visibility function  $V(u, v)$  that is its Fourier Transform. An array of antennas

## 2.4 Synthesis Imaging



**Figure 2.9:** *Left:* The VLA in A-configuration is an example of an ‘Y’ shaped array design. *Middle:* The corresponding overhead snapshot  $u - v$  coverage results in ‘snowflake’ pattern. *Right:* The corresponding  $u - v$  coverage after a 12 hour track of a source at a declination of  $60^\circ$ . Note the more intense  $u-v$  coverage in the direction of the three straight arms of the VLA for a snapshot track compared to the more uniform coverage over a longer duration track.

will only ever measure a certain set of values of this visibility function where the measured set is called the sampling function  $S(u, v)$ . This function is zero where no data have been taken. The actual data provided by the array is known as the sampled visibility function,  $S(u, v)V(u, v)$ . If we take the inverse Fourier transform of this function we get what is known as the *dirty image*:

$$I_\nu^D(l, m) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S(u, v)V_\nu(u, v)e^{2\pi i(ul+vm)} du dv. \quad (2.30)$$

where we have used  $I_\nu(l, m)$  to denote the modified sky brightness  $\mathcal{A}(l, m)I_\nu(l, m)$  as the correction for primary beam can be made at the final stage of data processing. Using the convolution theorem, the relationship between the dirty image and the desired intensity distribution  $I_\nu(u, v)$  is

$$I_\nu^D(u, v) = I_\nu(u, v) * B(l, m) \quad (2.31)$$

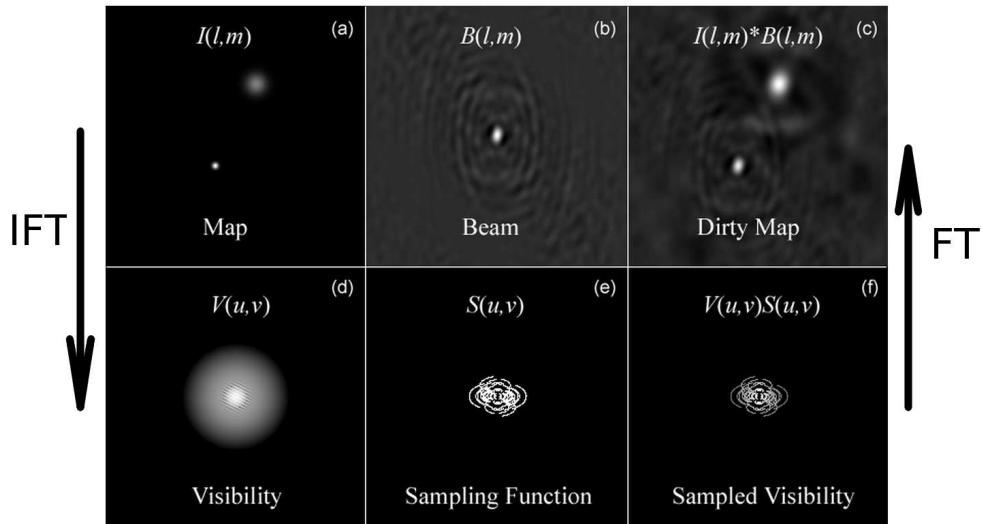
where the asterisk implies convolution and

$$B(l, m) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} S(u, v)e^{2\pi i(ul+vm)} du dv \quad (2.32)$$

## 2.4 Synthesis Imaging

is the *point spread function* (PSF), or *synthesized beam*, or *dirty beam* (i.e., the inverse Fourier Transform of the sampling function  $S$ ). Equation 2.31 says that the dirty image  $I^D$  is the true intensity distribution  $I$ , convolved with the synthesized beam  $B$ .

In Figure 2.10 we graphically summarize what has been said above. The panels in the upper row show the sky plane representations of the true image, the point spread function and the dirty image, while the panels in the lower row show the corresponding  $u - v$  plane representations of the true visibility, the sampling function and the sampled visibility. In other words, Equation 2.29 is summarized graphically by the relationship between panels (a) and (d), Equation 2.32 by (b) and (e), and Equations 2.30 and 2.31 by (c) and (f).



**Figure 2.10:** The Fourier transform pairs in synthesis imaging. (a) and (d): The true sky brightness and the visibility. (b) and (e): The dirty beam and the sampling function. (c) and (f): The dirty image and the sampled visibility. *Image Credit:* Prof. Dale E. Gary (New Jersey Institute of Technology).

Before the dirty image is computed, a weighting system is often applied to the visibilities to control the PSF. The two most common types of weighting system used are:

$$D_k = 1, \quad \text{natural weighting} \quad (2.33)$$

$$D_k = \frac{1}{N_s(k)}, \quad \text{uniform weighting} \quad (2.34)$$

where  $D_k$  is the weight to be applied to cell  $k$ , and  $N_s(k)$  is the number of data samples falling into cell  $k$  of characteristic width  $s$ . Natural weighting treats all points alike and gives the best signal-to-noise ratio for detecting weak sources. However, it produces a beam with a broad low-level plateau which is undesirable when imaging sources with both large and small scale structure. Uniform gridding produces fewer artifacts in the final map, while keeping the full resolution of the array but gives poorer signal-to-noise than natural weighting.

The ‘direct Fourier transform’ method can then be used to solve for the dirty image Equation 2.30. However, if this method is evaluated at every point on a  $N \times N$  grid, then the number of multiplications required goes as  $N^4$ . The fast Fourier transform (FFT) algorithm can also be used to solve Equation 2.30 but requires interpolating the data onto a regular grid (i.e., a process known as *gridding*). This method is widely used for large data volumes as it requires only a few times  $N^2 \log_2 N$  operations - not  $\mathcal{O}(N^4)$ , and the total time taken for gridding and FFT is usually a lot less than would take using the direct Fourier transform method.

### 2.4.3 Deconvolution (Making a CLEAN map)

The solution to the inverse Fourier transform given in Equation 2.29 is not unique, because the unmeasured points in the  $u - v$  plane could have *any* value without violating the data constraints. The ‘principle solution’ is the one in which all missing  $u - v$  measurements are set to zero and gives the dirty image discussed in the previous section. The dirty image is usually not a satisfactory representation of the sky as one would expect a more continuous distribution of visibilities than that provided by the array. The goal of the deconvolution process is to find a method that determines more reasonable values for the unmeasured  $u - v$  data. A priori information is the key to choosing ‘reasonable’ values. For example, we know that the Stokes parameter  $I$  must be positive and that radio sources generally do not have sidelobe patterns.

## 2.4 Synthesis Imaging

The CLEAN algorithm ([Högbom, 1974](#)) is the most widely used technique in radio interferometry to deconvolve the true sky intensity from the dirty beam. It assumes that the radio source can be represented by a number of point sources in an otherwise empty field and a simple iterative process is used to find the strengths and positions of these point sources. The final CLEAN image (i.e., the deconvolved image) is the sum of these point sources convolved with a CLEAN beam, which is usually an elliptical Gaussian of the same size and shape as the inner part of the dirty beam. The CLEAN algorithm obeys the following steps:

1. Find the strength and position of the brightest point in the dirty image. It may also be desirable to search for peaks in specified areas of the image, called CLEAN windows or regions.
2. At this position in the dirty image, subtract the dirty beam multiplied by the peak strength and a damping factor  $g$  ( $g \leq 1$ , usually called the loop gain).
3. Record the position and the subtracted flux in a model.
4. Iterate between (1), (2), and (3) until the peak is below some user specified level. The remainder of the dirty image is now termed the residuals.
5. Convolve the accumulated point source model with an idealized CLEAN beam (usually an elliptical gaussian of the same size and shape as the central lobe of the dirty beam).
6. Add the residuals to the image in (5) to create the final CLEAN image.

A problem with CLEAN is that the final CLEANed image is somewhat dependent upon the various control parameters such as CLEAN boxes, the loop gain and the number of CLEAN subtractions. For example, using too high a gain tends to make extended, weak emission undetectable and noisy. This problem is unavoidable, and input values must be chosen on a case by case basis, depending on the source and data quality.

Another deconvolution algorithm used in radio synthesis imaging, albeit less often, is the Maximum Entropy Method (MEM) which operates by minimizing a

## 2.4 Synthesis Imaging

smoothness function ('entropy') in an image. To conclude this section, we briefly discuss the practical differences between CLEAN and MEM:

1. CLEAN is nearly always faster than MEM, unless the image contains more than 1 million pixels.
2. MEM images are nearly always smoother than CLEAN images. This is because for CLEAN, what happens at one pixel is not coupled to what happens to its neighbours, while MEM couples pixels together by minimizing the spread in pixel values.
3. CLEAN sometimes makes extended emission look blotchy and may introduce artificial stripes into the image while MEM copes very poorly with point sources in extended emission. (Multi-scale CLEAN which is discussed in Chapter 3 is now becoming a popular choice in the radio community as an alternative deconvolution algorithm for images containing extended emission.)
4. For MEM, it is necessary to know the noise level quite well and it also helps to know the total flux density of the image. Knowledge of these are not required for CLEAN.

# 3

## Targets, Instrumentation, and Observations

Two unique data sets were analyzed as part of this thesis. The first data set was a millimeter interferometric multi-configuration study of Betelgeuse’s circumstellar envelope at 1.3 mm. In the first half of this chapter we give a detailed introduction to Betelgeuse and discuss our current understanding of the star and its complex stellar atmosphere. We give a description of the millimeter interferometer CARMA, which we used to study its complex outflow region on a number of spatial scales and also describe our observations which spanned  $\sim 2.5$  yr. The second data set consisted of a multi-wavelength centimeter study of two non-dusty red giants, Arcturus and Aldebaran. In the second half of this chapter we give a detailed introduction to both of these stars, including a description of their stellar properties, and their existing atmospheric models which we test against our data in Chapter 6. We give a short description of the VLA, which was the centimeter interferometer we used to study these stars, and finally describe our VLA observations.

### 3.1 Betelgeuse

Betelgeuse ( $\alpha$  Ori: M2 Iab) is the second closet red supergiant [ $D = 197$  pc (Harper *et al.*, 2008), barely losing that honor to Antares ( $\alpha$  Sco: M1.5 Iab + B2.5 V)], and subtends the largest angular diameter of any star in the northern sky apart from the Sun. It is by far the best studied red supergiant and has been observed with various techniques from the radio to the UV. According to Harper *et al.* (2001), Betelgeuse probably was a runaway star from the star-formation region Ori OB1 (see, e.g., Hoogerwerf *et al.*, 2000), and was a spectral type O9 V star while on the main sequence where it had a mass of  $\sim 20 M_{\odot}$ . The evolutionary models of Meynet & Maeder (2003) suggest that its current mass is  $18 M_{\odot}$  which corresponds to a surface gravity of  $0.5 \text{ cm s}^{-2}$  (i.e.,  $2 \times 10^{-5} g_{\odot}$ ). Its mass loss is characterized by a low velocity wind ( $10 - 17 \text{ km s}^{-1}$ ) and a mass loss rate of  $\sim 2 \times 10^{-6} M_{\odot} \text{ yr}^{-1}$  (Harper *et al.*, 2001). Like most late-type evolved stars, Betelgeuse's terminal wind velocity  $v_{\infty}$  is smaller than the surface escape speed  $v_{\text{esc}}$  (values given in Table 3.1). This means that most of the energy and momentum are deposited into its flow within the first few stellar radii.

Betelgeuse has an extremely slow rotation period of about 17 yr (Uitenbroek *et al.*, 1998) which implies that it is probably experiencing very limited action of a solar-like dynamo. Its large radius ( $950 R_{\odot}$ ) means that any possible remnant of a strong main sequence magnetic field is likely negligible due to a large dilution factor. In comparison to Sun like stars, its large radius also means that it has a large photospheric scale height ( $H_{\star} \sim 0.01 R_{\star}$ ). The expected result of this is the presence of no more than a few giant and stable convection cells in the photosphere Schwarzschild (1975). Such features should enhance the detectability of the magnetic fields generated through a turbulent dynamo (Vögler & Schüssler, 2007). Bedecarrax *et al.* (2013) have recently monitored Betelgeuse over a three year period using high resolution spectropolarimetry, and find a longitudinal magnetic field strength which varies in values ranging from  $-3 \text{ G}$  to  $+2 \text{ G}$  with temporal changes in the field strength occurring over timescales as short as a few weeks.

Various studies have assumed that Betelgeuse has a warm ( $T_e \sim 8000 \text{ K}$ ) extended atmosphere where ionized hydrogen dominates the electron density (e.g.,

### 3.1 Betelgeuse

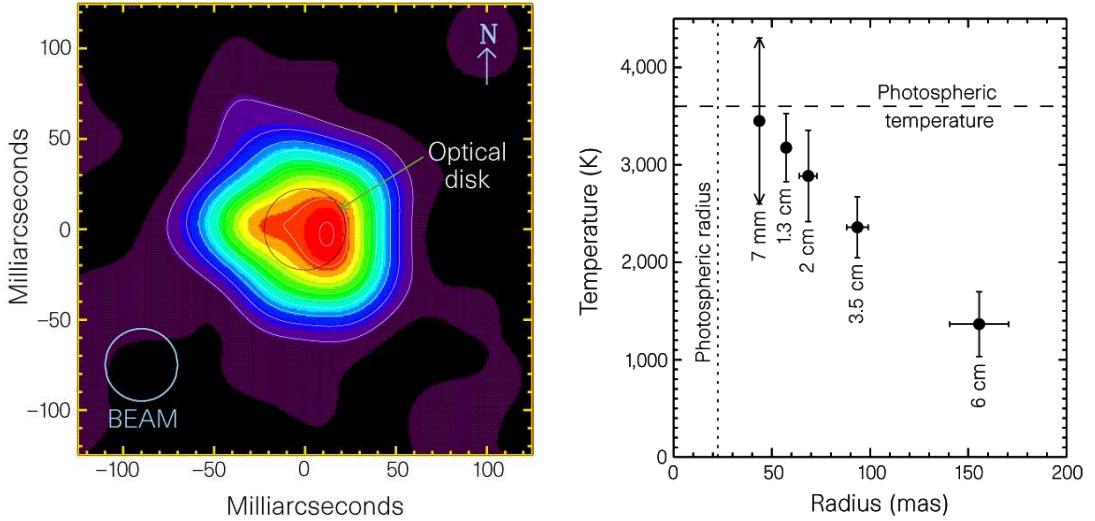
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**Table 3.1:** Physical Properties of  $\alpha$  Ori.

Property	Value	Reference
HD Number	39801	...
Spectral Type	M2 Iab	Perryman <i>et al.</i> (1997)
App. Mag. (V)	0.45	Perryman <i>et al.</i> (1997)
RA (ICRS: ep=J2000)	05 <sup>h</sup> 55 <sup>m</sup> 10.305 <sup>s</sup>	van Leeuwen (2007)
dec (ICRS: ep=J2000)	+07°24'25.430"	van Leeuwen (2007)
Proper motion-RA (mas yr <sup>-1</sup> )	27.54	van Leeuwen (2007)
Proper motion-dec (mas yr <sup>-1</sup> )	11.30	van Leeuwen (2007)
$\pi$ (mas)	5.07 ± 1.1	Harper <i>et al.</i> (2008)
d (pc)	197 ± 45	Harper <i>et al.</i> (2008)
Main Sequence $M_*$ ( $M_\odot$ )	~ 20	Meynet & Maeder (2003)
Main Sequence Spectral Type	O9 V	Harper <i>et al.</i> (2008)
Current $M_*$ ( $M_\odot$ )	~ 18	Meynet & Maeder (2003)
$\theta_{UD}$ (mas)	43.33 ± 0.04	Perrin <i>et al.</i> (2004)
$R_*$ ( $R_\odot$ )	950	Harper <i>et al.</i> (2008)
$T_{\text{eff}}$ (K)	3650	Levesque <i>et al.</i> (2005)
Log $L_*/L_\odot$	5.10 ± 0.22	Harper <i>et al.</i> (2008)
$g_*$ (cm s <sup>-2</sup> )	0.5	Meynet & Maeder (2003)
Heliocentric $v_{\text{rad}}$ (km s <sup>-1</sup> )	20.7 ± 0.4	Harper <i>et al.</i> (2008)
$v_{\text{esc}}$ (km s <sup>-1</sup> )	85	...
$v_\infty$ (km s <sup>-1</sup> )	10, 17	Bernat <i>et al.</i> (1979)
$T_{\text{wind}}$ (K)	< 4000	Lim <i>et al.</i> (1998)
$\dot{M}_*$ ( $M_\odot$ yr <sup>-1</sup> )	3 × 10 <sup>-6</sup>	Harper <i>et al.</i> (2001)
Photospheric $H_*$ ( $R_*$ )	0.006	...
Fe/H	0.05 ± 0.14	Ramírez <i>et al.</i> (2000)
O/C	2.5	Lambert <i>et al.</i> (1984)
Rotational period (yr)	17	Uitenbroek <i>et al.</i> (1998)
$B_*$ (G)	~ 1	Aurière <i>et al.</i> (2010)
Chromosphere Model	...	Harper <i>et al.</i> (2001)
Wind Model	...	Harper <i>et al.</i> (2001)
Outer CSE	...	Rodgers & Glassgold (1991)
$^{12}\text{C}/^{13}\text{C}$	6 ± 1	Harris & Lambert (1984)
$^{16}\text{O}/^{17}\text{O}$	525 ± 250	Harris & Lambert (1984)
$^{16}\text{O}/^{18}\text{O}$	700 ± 300	Harris & Lambert (1984)

### 3.1 Betelgeuse

Hartmann & Avrett, 1984; Newell & Hjellming, 1982). In fact, high resolution UV imaging with the HST partially resolved the *hot* chromosphere revealing  $6000 - 8000$  K plasma extending to more than  $2 R_\star$  (Gilliland & Dupree, 1996). However, Lim *et al.* (1998) spatially resolved Betelgeuse with the ‘old’ VLA at 5 different wavelengths and showed that the mean electron temperature is actually much cooler, i.e.,  $T_e = 2000 - 4000$  K and extends out to several stellar radii as shown in Figure 3.1. In fact, the mid-IR [Fe II] line studies of Harper *et al.* (2009) suggest that the temperature of the bulk of the plasma where the chromosphere has its largest filling factor is  $< 2500$  K. Therefore, it appears that the hot chromospheric plasma has a small filling factor and co-exists with the much more abundant cool plasma within several stellar radii. Harper *et al.* (2001) constructed a detailed physical model for the inner atmosphere of Betelgeuse which reproduces the radio fluxes of Lim *et al.* (1998) and consists of electron temperature and gas densities as a function of radius.



**Figure 3.1:** Left: The spatially resolved Q-band VLA image of Betelgeuse showing asymmetry in its atmosphere. Right: The multi-frequency spatially resolved radio maps revealed that the mean electron temperature profile was lower than expected (Lim *et al.*, 1998).

In comparison to other well studied red supergiants such as VY CMa, Betelgeuse has a modest mass loss rate, a low gas-to-dust mass ratio of  $200 - 1700$ , and low molecular abundances (Harper *et al.*, 2001). However, it is known to possess

### 3.1 Betelgeuse

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a molecular shell known as a MOLsphere above its photosphere (Tsuji, 2000), and recently Verhoelst *et al.* (2006) and Perrin *et al.* (2007) have used the VLTI to establish some of its properties. They find the MOLsphere to have a geometrical thin extent ( $\sim 0.1 R_\star$ ) with a temperature of  $\sim 1500$  K at  $\sim 1.4 R_\star$ , and contains molecules such as H<sub>2</sub>O, SiO, and Al<sub>2</sub>O<sub>3</sub>. Notably, they suggest that this molecular layer is where dust nucleation commences in support of a dust-driven wind scenario. However, the existence of Al<sub>2</sub>O<sub>3</sub> is disputed by Kamiński *et al.* (2013) who argue that the absence of AlO (which is required to form Al<sub>2</sub>O<sub>3</sub>) in the photospheric spectrum is puzzling, and verification by independent studies is needed. In addition, Woitke (2006) finds that the average temperature between  $1 - 3 R_\star$  is too high for silicates to form, and indeed recent studies suggest that silicates are not formed inside  $\sim 23 R_\star$  (Skinner *et al.*, 1997; Tatebe *et al.*, 2007). For Betelgeuse, it therefore appears that radiation pressure on dust grains is not the main mechanism driving its mass loss (Harper, 2010).

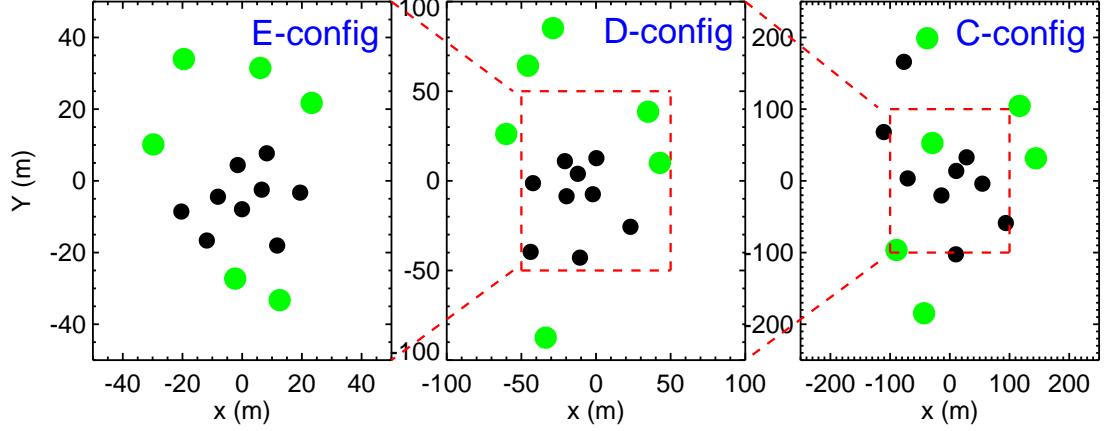
In the past decade or so, a number of sensitive multi-wavelength studies of Betelgeuse have revealed a complex non-spherically symmetric circumstellar environment. Starting at its surface, a *H* band interferometric image was reconstructed using data from the Infrared Optical Telescope Array (IOTA) and showed a non-uniform brightness distribution across its surface (Haubois *et al.*, 2009). This data was modelled with 3-D hydrodynamical simulations by Chiavassa *et al.* (2010), resulting in the detection of a granulation pattern on the surface. They concluded that the surface contains a large ( $\sim 30$  mas) convective cell and a few small to medium scale (5 – 15 mas) convection-related surface structures. Adaptive optics images with the Very Large Telescope (VLT) at  $1.04 - 2.17 \mu m$  revealed an irregular circumstellar envelope at a few  $R_\star$  with a few bright plumes extending out to  $6 R_\star$  (Kervella *et al.*, 2009). These plumes have been attributed to the action of giant convection cells. Thermal infrared VLT ( $\lambda = 8 - 20 \mu m$ ) imaging revealed oxygen-rich dust out to  $100 R_\star$  (2.5'') (Kervella *et al.*, 2011), while Herschel images show a chaotic dust distribution far out in the circumstellar envelope, i.e., beyond  $600 R_\star$  (i.e.  $> 15''$ ) (Decin *et al.*, 2012). A conclusion that can be drawn from these studies is the constant presence of inhomogeneities in the circumstellar environment which probably play a role in

the formation of dust and molecular species. We defer a discussion on previous observations of CO in Betelgeuse's circumstellar environment until Chapter 5.

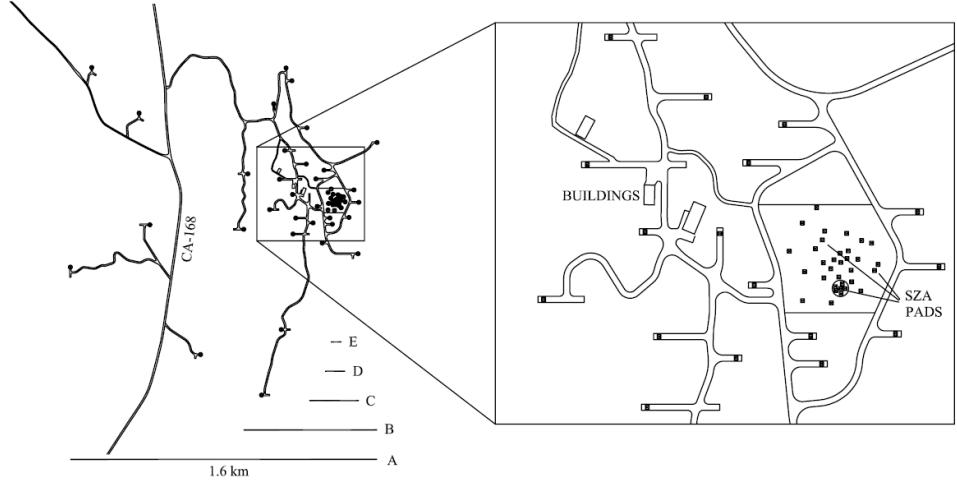
## 3.2 CARMA

The Combined Array for Research in Millimeter-wave Astronomy (CARMA) (Bock, 2006) is a millimeter interferometer located at Cedar Flat in eastern California at an elevation of 2200 m. The array consists of nine 6.1 m antennas and six 10.4 m antennas formerly from the Berkeley Illinois Maryland Association (BIMA) and the Owens Valley Radio Observatory (OVRO) arrays respectively, and operates at 85-115 GHz (3 mm) and 215-270 GHz (1.3 mm). Eight additional 3.5 m antennas known as the Sunyaev-Zel'dovich Array (SZA) can also be added to CARMA for continuum observations at 26-36 GHz (1 cm) and 85-115 GHz (3 mm). The different sizes of the CARMA antennas makes it a heterogeneous array with a total collecting area equivalent to a single 32 m dish antenna. Despite some technical difficulties associated with image restoration for a heterogeneous array (e.g., the 15-element CARMA array has 3 different primary beams), there are a number of advantages. Such an array samples shorter  $u-v$  spacings directly when the smallest antennas are in a compact configuration. This results in more of the total flux density being recovered and also recovers more of the large scale structure and is the reason why the smaller diameter antennas are usually left in the center of the array as shown in Figure 3.2. In fact, it has been shown that a heterogeneous CARMA array produces better image fidelity than a homogeneous array with the same number of antennas and collecting area (Wright, 1999).

CARMA is a reconfigurable array and has 5 different configurations providing baselines ranging between 8 m and 2 km; A configuration being the most extended and E configuration being the most compact. Figure 3.3 gives an overview of how these configurations are achieved at Cedar Flat. As the resolution is set by the array configuration and frequency of observation, we list the various achievable resolutions (i.e. the HPBW of the synthesized beam) for each of CARMA's 5 configurations at 230 GHz, in Table 3.2. We also list the largest angular scale that can be imaged at 230 GHz as this is an important property of each configuration, especially when imaging extended emission (i.e., Chapter 5). This limitation



**Figure 3.2:** The three CARMA array configurations used to study the CSE of Betelgeuse. The most compact CARMA configuration is E configuration (*left*) which has  $B_{\max} = 66$  m, D configuration (*middle*) has  $B_{\max} = 148$  m, while C configuration (*right*) has  $B_{\min} = 370$  m and was the most extended configuration used in our study. The 10.4 m antennas are marked green while the 6.1 m antennas are marked black.



**Figure 3.3:** The layout of antenna pads for CARMA and a visual of the extent of each configuration (Bock, 2006).

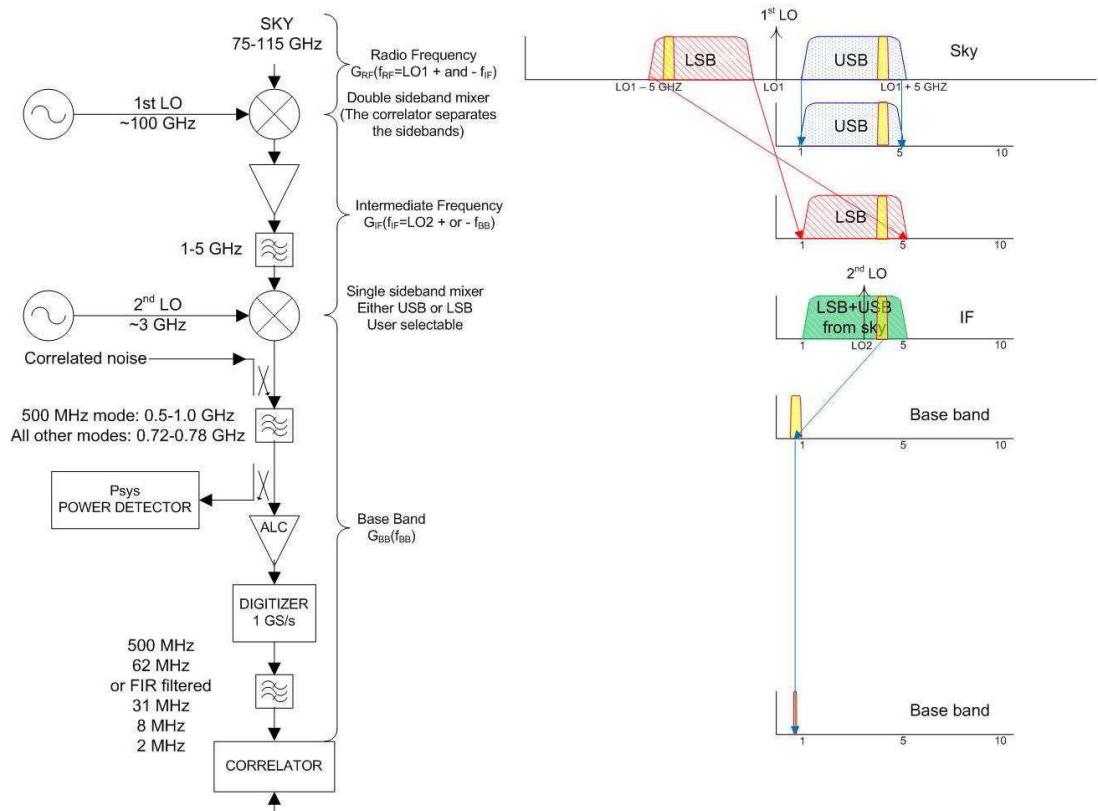
is unique to interferometers and means that structures on angular scales significantly larger than the fringe spacing formed by the shortest baseline are not measured. This information can only be obtained by observing in a smaller array configuration or by using the mosaicing method. The half-power beamwidth of

### 3.2 CARMA

the primary beam (i.e., the FOV) at a frequency  $\nu$  is  $50'' \times (230 \text{ GHz}/\nu)$  for the 6.1 m antennas, while is  $30'' \times (230 \text{ GHz}/\nu)$  for the 10.4 m antennas.

**Table 3.2:** Properties of the 5 CARMA configurations.

	A	B	C	D	E
$B_{\max}$ (m)	1883	946	370	148	66
$B_{\min}$ (m)	150	82	26	11	8.5
Synthesized Beam $\theta_{\text{HPBW}}$ ('')	0.15	0.38	0.86	2.1	4.4
Largest Angular Scale $\theta_{\text{LAS}}$ ('')	1.1	2.0	6.2	14.6	18.9
6.1 m Primary Beam $\theta_{\text{HPBW}}$ ('')			$50'' \times (230 \text{ GHz}/\nu)$		
10.4 m Primary Beam $\theta_{\text{HPBW}}$ ('')			$30'' \times (230 \text{ GHz}/\nu)$		



**Figure 3.4:** Major components in the signal path through the CARMA system (Wei *et al.*, 2008).

The 10.4 meter and 6.1 meter antennas are equipped with circular polarization SIS receivers for the 1mm band, and single-polarization SIS receivers for the 3 mm band. The tuning ranges of the 1mm and 3mm receivers are 215-265 GHz and 85-116 GHz, respectively. Figure 3.4 shows the main components in the signal path through the CARMA system. Upon passing through the atmosphere and antenna Naysmith optics, the RF signal is sent through the cooled low noise double sideband mixer where it is combined with a local oscillator  $\nu_{\text{LO}1}$  to produce the intermediate frequency  $\nu_{\text{IF}}$ . The  $\nu_{\text{LO}1}$  can be either below or above the  $\nu_{\text{RF}}$ ; known as the upper and lower sidebands, respectively, giving:

$$\nu_{\text{IF}} = \nu_{\text{RF}} \pm \nu_{\text{LO}1}. \quad (3.1)$$

The radio frequency gain  $G_{\text{RF}}(\nu_{\text{RF}})$ , results from components before this mixer (i.e., the sky and antenna optics) and depends only on the sky frequency. The IF signals then pass through amplifiers and attenuators in the antenna, are converted to light and sent along fibers to the control building, where they are then converted back to microwaves, amplified and attenuated some more before reaching the second mixer. Even though many of the components in the IF path are common across all bands, the IF gain  $G_{\text{IF}}(\nu_{\text{IF}})$ , will vary per band as there are components in the second mixer which are separate for each band. The second mixer is a single sideband mixer and its baseband output contains either the upper sideband signals above the second LO or the lower sideband:

$$\nu_{\text{BB}} = \nu_{\text{IF}} \pm \nu_{\text{LO}2}. \quad (3.2)$$

A noise signal is passed through all of the remaining correlator configuration dependent components and the selectable analog filters are measured by this noise source. The power measurement  $P_{\text{sys}}$ , used for the system noise temperature measurement is calculated after the analog filters and before the automatic level control (ALC) amplifier which ensures that the digitizers operate at a constant level.

## 3.3 CARMA Observations of Betelgeuse

We used CARMA to obtain on-source profiles of the rotational transition line  $^{12}\text{C}^{16}\text{O}(J = 2\text{-}1)$  which has a rest frequency of 230.538 GHz (1.3 mm). Table 3.3 summarizes our complete set of multi-configuration observations which span the period 2007 May - 2009 November. Our first set of observations took place in June 2007 when CARMA was in D configuration providing a spatial resolution of 2.1''. A total of 5 D configuration tracks were carried out and one which took place in 18 May 2007 was excluded from the final analysis due to large levels of noise in the final map as a result of poor weather conditions. We obtained five E configuration (the most compact configuration available) tracks during the summer of 2009 which provided a relatively low spatial resolution of 4.4'' but gave the best sensitivity to extended emission. However, only one track was found to be of sufficient quality to be included in the final analysis. Our C configuration observations which provide sub-arcsecond resolution at 230 GHz (i.e., 0.9'') were due to take place in December 2008 but were postponed for a full year due to very poor weather conditions at Cedar Flat. We obtained a total of 4 usable tracks in C configuration amounting to 8.4 hours on source. The FOV of the individual 10.4 m antennas is  $\sim 32''$  at the observed frequency.

CARMA can now take measurements in eight bands yielding a maximum bandwidth of 4 GHz per sideband. However, back when our data was obtained, the CARMA correlator took measurements in just three separate bands, each having an upper and lower sideband. One band was set to the low resolution 468 MHz bandwidth mode (15 channels of 31.25 MHz each) to observe continuum emission and was centered on the line. The other two bands were configured with 62 MHz and 31 MHz bandwidth across 63 channels (with a resolution of  $1.3 \text{ km s}^{-1}$  and  $0.65 \text{ km s}^{-1}$  respectively) and were also centered on the line. The line was measured in the upper sideband in the C and E array and in the lower sideband in the D array.

Bandpass and phase calibration were performed using 3C120 and 0530+135. 0532+075 was used as a secondary phase calibrator to determine the quality of the phase transfer from the primary phase calibrator. The observing sequence was to integrate on the primary phase calibrator for  $\sim 2.5$  minutes, the target

### **3.3 CARMA Observations of Betelgeuse**

for  $\sim$  18 minutes, and the secondary phase calibrator for  $\sim$  2.5 minutes. The cycle was repeated for each track which lasted between 1.5 hours and 5 hours. Absolute flux calibration was carried out with 0530+135 and 3C120 using the continuously updated CARMA flux catalog to obtain their flux values at each observation.

**Table 3.3:** CARMA Observations of  $\alpha$  Ori between June 2007 and November 2009.

Date	Configuration	Time on Source (hr)	Flux Calibrator	Phase Calibrator	Image Cube Dynamic Range
2007 Jun 18	D	0.9	0530+135	0530+135, 0532+075	22.8
2007 Jun 21	D	3.0	0530+135	0530+135, 0532+075	22.7
2007 Jun 24	D	2.1	0530+135	0530+135, 0532+075	26.1
2007 Jun 25	D	2.4	0530+135	0530+135, 0532+075	30.2
2009 Jul 07	E	3.2	3C120	3C120, 0532+075	30.1
2009 Nov 05	C	1.2	3C120	3C120, 0532+075	17.3
2009 Nov 09	C	3.0	3C120	3C120, 0532+075	27.2
2009 Nov 15	C	1.0	3C120	3C120, 0532+075	17.8
2009 Nov 16	C	3.2	3C120	3C120, 0532+075	32.0
All	C	8.4	...	...	43.8
All	D	8.4	...	...	31.9
All	Multi-configuration	20.0	...	...	52.3

## 3.4 Arcturus and Aldebaran

Currently the most detailed spatial information about the atmospheres of K and early M spectral type evolved stars is obtained from eclipsing binaries such as the  $\zeta$  Aurigae and symbiotic systems (e.g., Wright 1970; Baade *et al.* 1996; Eaton 2008; Crowley *et al.* 2008). Even though these systems offer us the best opportunity to obtain information on the dynamics and thermodynamics at various heights in the evolved star's atmosphere, the very nature of the binary system may introduce further complexities. For example, the orbital separation is often within the wind acceleration region and one could expect flow perturbations to be present (e.g., Chapman 1981). Using the old VLA, Harper *et al.* (2005) find a slow wind acceleration for  $\zeta$  Aurigae and confirm that its velocity structure is not typical of single stars with similar spectral types, such as  $\lambda$  Velorum (Carpenter *et al.*, 1999). In order to avoid the assumed additional complexities of a companion, we have selected two single luminosity class III red giants: Arcturus ( $\alpha$  Boo: K2 III) and Aldebaran ( $\alpha$  Tau: K5 III). These nearby red giants have been extensively studied at other wavelengths and their stellar parameters, which are summarized in Table 3.4, are accurately known. These stars are predicted to be point sources at all frequencies between 1 and 50 GHz in all VLA configurations so our radio observations measure their total flux density,  $F_\nu$ . For example, our observations of these stars which are discussed in Section 3.6, were taken in B configuration providing a maximum spatial resolution of  $\sim 0.14''$  at 45 GHz which is equivalent to  $\sim 7 R_*$  for both stars. The radio emission from these stars at 45 GHz is expected to be chromospheric in origin (Harper *et al.*, 2013), and as the spatial extent of red giant chromospheres is expected to be less than  $1.5 R_*$  (Berio *et al.*, 2011), then we can be assured that these targets will be unresolved even at the highest VLA frequencies. Moreover, both stars have existing semi-empirical 1-D chromospheric and wind models which we can directly compare our data against.

### *Arcturus ( $\alpha$ Boo: K2 III)*

Arcturus ( $\alpha$  Boo: K2 III) is the nearest ( $d = 11.3$  pc) and brightest ( $V = -0.04$  mag) noncoronal red giant and is probably the best example of a red giant whose

atmosphere can be studied in detail with the VLA. It is the leader of a group of stars that share a similar  $V$  space velocity (the component of stellar motion relative to the LSR in the direction of rotation), age ( $\geq 10$  Gyr), and metallicity ( $[\text{Fe}/\text{H}] \sim -0.5$ ), known as the Arcturus moving group (Eggen, 1971). The group has traditionally been regarded as the remains of a dissolved open cluster (e.g., Eggen, 1971, 1996) but it has also been suggested to be the debris of a metal-poor accreted satellite some billions of years ago (Navarro *et al.*, 2004). Recent analysis of chemical abundances are consistent with the former hypothesis but do not entirely rule out a merger one (Williams *et al.*, 2009). Arcturus is ascending the red giant branch and being a single star, its mass is relatively poorly constrained but is similar to that of the Sun ( $0.8 \pm 0.2 M_{\odot}$  by Kallinger *et al.*, 2010). During the last decade there has been a large dispersion in the reported values of Arcturus' effective temperature (a nice graph summarizing this is presented in Griffin, 1996) but nowadays it is generally accepted to be about 4300 K (di Benedetto, 1993). A number of interferometric measurements of the limb-darkened angular diameter of the star are available in the literature with most values agreeing within their uncertainties. The weighted mean value of these values is  $\theta_{\text{LD}} = 21.06 \pm 0.17$  mas (Ramírez & Allende Prieto, 2011) giving the star a radius of  $25.4 \pm 0.3 R_{\odot}$ .

Arcturus is an important target for high spatial and spectral resolution calibration. Thus when the Hipparcos catalog flagged Arcturus as a two component object (Perryman *et al.*, 1997) it caused quite a stir in the community. However, the uncertainties in the Hipparcos data (Soderhjelm & Mignard, 1998) along with a non-detection in adaptive optics observations (Turner *et al.*, 1999) and sensitive interferometric techniques (Lacour *et al.*, 2008) suggest that Arcturus is single and can still be used as a calibrator. Nevertheless, Lacour *et al.* (2008) still do not rule out the possibility of a planetary companion of a few Jovian masses as suggested by its long period radial velocity variations of  $\sim 233$  days (Brown, 2007; Hatzes & Cochran, 1993). Variations in the order of a few days period are also seen in its radial velocity (Merline, 1999) as well as photometry (Retter *et al.*, 2003). The photometric amplitude oscillations can vary by up to a percent and may be the manifestation of convection such as large-scale granulation, or solar-like oscillations (Dziembowski *et al.*, 2001).

### 3.4 Arcturus and Aldebaran

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The first detailed model of the atmosphere of Arcturus was the 1-D time-independent semi-empirical model of Ayres & Linsky (1975) which was based on diagnostics observable from the ground (i.e. Ca II  $h$  and  $k$  and the Ca II IR triplet) and early satellite observations of the Mg II  $h$  and  $k$  emission lines from *Copernicus*. They were able to calculate temperature and mass column densities

**Table 3.4:** Basic Properties of  $\alpha$  Boo and  $\alpha$  Tau.

Property	$\alpha$ Boo	$\alpha$ Tau	Reference
HD Number	124897	29139	...
Spectral Type	K2 III	K5 III	1
App. Mag. (V)	-0.5	0.86v	1
RA (ICRS: ep=J2000)	14 <sup>h</sup> 15 <sup>m</sup> 39.672 <sup>s</sup>	04 <sup>h</sup> 35 <sup>m</sup> 55.239 <sup>s</sup>	2
dec (ICRS: ep=J2000)	+19°10'56.673"	+16°30'33.489"	2
Proper motion-RA (mas yr <sup>-1</sup> )	-1093.39 ± 0.44	63.45 ± 0.84	2
Proper motion-dec (mas yr <sup>-1</sup> )	-2000.06 ± 0.39	-188.94 ± 0.65	2
$\pi$ (mas)	88.83 ± 0.54	48.94 ± 0.77	2
d (pc)	11.3 ± 0.1	20.4 ± 0.3	2
$M_\star$ ( $M_\odot$ )	0.8 ± 0.2	1.3 ± 0.3	3, 4
$\theta_{LD}$ (mas)	21.06 ± 0.17	20.58 ± 0.03	5, 6
$R_\star$ ( $R_\odot$ )	25.4 ± 0.3	44.2 ± 0.9	...
$T_{\text{eff}}$ (K)	4294 ± 30	3970 ± 49	7
$L_\star/L_\odot$	198 ± 3	442 ± 11	...
$g_\star$ (cm s <sup>-2</sup> )	34	18	...
Heliocentric $v_{\text{rad}}$ (km s <sup>-1</sup> )	+5.19 ± 0.04	+54.11 ± 0.04	8
$v_{\text{esc}}$ (km s <sup>-1</sup> )	110	106	...
$v_\infty$ (km s <sup>-1</sup> )	35 – 40	30	9, 10
$T_{\text{wind}}$ (K)	~10,000	~10,000	9, 10
$\dot{M}_\star$ ( $M_\odot$ yr <sup>-1</sup> )	$2 \times 10^{-10}$	$1.6 \times 10^{-11}$	9, 10
Photospheric $H_\star$ ( $R_\star$ )	0.005	0.005	...
Fe/H	-0.5 ± 0.2	-0.15 ± 0.2	11
Rotational period (yr)	2.0 ± 0.2	1.8	12, 13
$B_\star$ (G)	0.65 ± 0.26	unknown	14
Chromosphere Model	Drake (1985)	McMurtry (1999)	...
Wind Model	Drake (1985)	Robinson <i>et al.</i> (1998)	...

References.-(1) Perryman *et al.* (1997); (2) van Leeuwen (2007); (3) Kallinger *et al.* (2010); (4) Lebzelter *et al.* (2012) (5) Ramírez & Allende Prieto (2011) (6) Richichi & Roccatagliata (2005); (7) di Benedetto (1993); (8) Massarotti *et al.* (2008); (9) Drake (1985); (10) Robinson *et al.* (1998); (11) Decin *et al.* (2003); (12) Gray *et al.* (2006); (13) Hatzes & Cochran (1993); (14) Sennhauser & Berdyugina (2011)

for the upper photosphere and chromosphere and estimated the temperature at the top of the chromosphere to be  $\sim 8000$  K. They also found that  $T_{\min}/T_{\text{eff}} \sim 0.77$  which is similar to that of the Sun. The important IUE cool star survey by Linsky & Haisch (1979) placed Arcturus on the right of the dividing line. This meant that its atmosphere showed lines formed at temperatures no hotter than  $10,000 - 20,000$  K, suggestive of a chromosphere only. However, they also developed a hot coronal wind model for the star and found it to be consistent with the absence of transition region material in the IUE data; that is, the high temperature transition region emission line fluxes could be below the detection limit. Drake (1985) developed a semi-empirical chromosphere and wind model for Arcturus based on the Mg II  $k$  emission line from the IUE and showed that its wind is very extended and estimated a mass loss rate of  $2 \times 10^{-10} M_{\odot} \text{ yr}^{-1}$ .

Evidence began to emerge that Arcturus actually falls into the class of late type stars known as hybrids when deeply exposed IUE echelle spectrograms showed the weak presence of the Si III]  $\lambda 1892.0$  feature, indicating the existence of a small amount of plasma at temperatures as hot as  $6 \times 10^4$  K. Its hybrid status was confirmed when C IV and N V (indicative of temperatures  $\sim 1 \times 10^5$  K) were detected with the HST STIS, and also with a tentative  $3\sigma$  X-ray detection made with the Chandra X-Ray Observatory (Ayres *et al.*, 2003). It appears that Arcturus has been able to sustain a modest level of magnetic activity. In fact, three possible values for the mean longitudinal magnetic field (albeit weak:  $B = 0.65 \pm 0.26$ ,  $0.43 \pm 0.16$ , and  $-0.23 \pm 0.20$  G) have reportedly been detected on the star via the Zeeman effect (Sennhauser & Berdyugina, 2011), and a magnetic cycle with a period of  $\geq 14$  years has also been reported (Brown *et al.*, 2008).

Arcturus appears to have a thermally bifurcated chromosphere which consists of material within the chromosphere that is cooler than the surrounding chromospheric temperature minimum [i.e a CO-mosphere, Wiedemann *et al.* (1994)]. This suggests that the assumption in current semi-empirical chromospheric models of the *hot* UV emitting material having a filling factor of unity may not be correct. The more recent spectroscopic analysis of CO and H<sub>2</sub>O transitions have confirmed the existence of cooler molecular *clouds* in the outer atmosphere of Arcturus (Ryde *et al.*, 2002; Tsuji, 2009).

#### *Aldebaran ( $\alpha$ Tau: K5 III)*

At a distance of 20.4 pc, Aldebaran ( $\alpha$  Tau: K5 III) is a nearby red giant and is one of the most easily recognizable stars from the northern and most of the southern hemisphere. Even though it is almost twice as far away as Arcturus, its large stellar radius ( $R_\star = 44.2 \pm 0.9 R_\odot$ ) gives it a comparable angular diameter ( $\theta_{\text{LD}} = 20.58 \pm 0.03$ ), and is therefore another excellent candidate for multi-frequency studies with the VLA. Its effective temperature ( $T_{\text{eff}} = 3970 \pm 49$  K) is slightly lower than Arcturus', and it also has a slightly higher mass ( $M = 1.3 \pm 0.3 M_\odot$ ). Using these values for its radius and mass gives a surface gravity of  $g_\star = 18$  (in units of  $\text{cm s}^{-2}$ ), about 1500 times lower than the Sun's. The metallicity of Aldebaran is marginally subsolar with  $[\text{Fe}/\text{H}] = -0.15 \pm 0.2$  (Decin *et al.*, 2003). Using high spectral resolution in the H, K, and L bands, Tsuji (2008) derived the carbon, nitrogen, and oxygen abundances in Aldebaran which suggests the mixing of the CN-cycled material in the first dredge-up. A 643 day period in the radial velocity for Aldebaran was reported by Hatzes & Cochran (1993), and Hatzes & Cochran (1998) find evidence to support the hypothesis that this variability comes from the reflex motion of the central star due to a planetary companion having a mass of  $11 M_{\text{Jup}}$  although this has not been confirmed to date.

Aldebaran has been extensively studied at UV wavelengths. Early IUE observations placed the star well to the noncoronal side of the Linsky-Haisch dividing line (Linsky & Haisch, 1979). The first chromospheric model of the star was developed in the late 1970s (Kelch *et al.*, 1978) and was based on both optical (mainly Ca II H and K) and UV (Mg II *h* and *k*) emission line fluxes. Later, GHRS spectra revealed the presence of significant flux in the C IV resonance lines around 1550 Å (Carpenter & Robinson, 1996), indicating the presence of some *hot* transition region plasma. In light of the new GHRS findings, a new model of the chromosphere and transition region of Aldebaran was developed, with temperatures reaching up to  $T_e \sim 1 \times 10^5$  K. Modelling the GHRS optically thick Mg II and O I resonance lines (which show typical stellar wind absorption features), Robinson *et al.* (1998) found evidence for the acceleration of a slow wind and derived a mass-loss rate of  $1.6 \times 10^{-11} M_\odot \text{ yr}^{-1}$  and a terminal wind velocity of 30 km  $\text{s}^{-1}$ . Unexpectedly, FUSE spectra revealed the presence (albeit weak)

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of the coronal proxy O VI 1032 and 1038 Å emission lines (Dupree *et al.*, 2005) although Ayres *et al.* (2003) failed to detect any X-ray emission from the star. Like Arcturus, Aldebaran appears to harbour small levels of magnetic activity.

Wiedemann *et al.* (1994) observed Aldebaran's infrared ro-vibration absorption lines ( $v = 2 - 1$ ) and matched them with synthetic spectra based on a model containing the semi-empirical photosphere of Kelch *et al.* (1978) and the best fitting temperature and column density profiles. Like Arcturus, they found a steady decrease in temperature with height, with the chromospheric temperature being constantly below the temperature minimum ( $T_e^{\min} \sim 2800$  K) of the Kelch model. Again, they explained their results by suggesting the existence of an extra *cool* component with a large (> 99%) filling factor in the outer atmosphere (i.e. a thermally bifurcated CO-mosphere). Recently, Ohnaka observed the CO first overtone lines ( $v = 2 - 0$ ) of Arcturus near 2.3 μm with the Very Large Telescope Interferometer (VLTI) and discovered a CO layer extending out to  $2.5 R_\star \pm 0.3 R_\star$ . They were unable to constrain the geometrical thickness of this CO layer from their data but we will show how our VLA radio data can constrain this value in Chapter 6. H<sub>2</sub>O was also confirmed in its MOLsphere when Tsuji (2001) detected a H<sub>2</sub>O absorption feature at ∼ 6.6 μm. A narrow absorbtion feature in the midst of the wind absorption of the GHRS Mg II h and k lines has been interpreted and modelled as a feature of α Tau's wind-ISM interaction region also known as its astrosphere (Wood *et al.*, 2007).

## 3.5 The Karl G. Jansky Very Large Array

The NRAO<sup>1</sup> Karl G. Jansky Very Large Array (VLA) is an aperture synthesis radio telescope located on the Plains of San Agustin, New Mexico, USA and is capable of producing radio images with a spatial resolution greater than that of the HST. It is the product of a program to modernize the electronics of the old VLA which had been in operation at the same site since the late 1970's. One of the main upgrades to the VLA is the addition of the Wideband Interferometric Digital Architecture (WIDAR) correlator which allows the digital correlation of

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<sup>1</sup>The National Radio Astronomy Observatory is a facility of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc.

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**Table 3.5:** Improved Performance Parameters of the VLA.

Parameter	old VLA	VLA	Improvement Factor
Continuum sensitivity ( $1\sigma$ , 9 hr)	$10 \mu\text{Jy}$	$1 \mu\text{Jy}$	10
Bandwidth per polarization	0.1 GHz	8 GHz	80
Coarsest frequency resolution	50 MHz	2 MHz	25
Finest frequency resolution	381 Hz	0.12 Hz	3180
Channels at max. bandwidth	16	16,384	1024
Maximum number of channels	512	4,194,304	8192

very wideband signals. WIDAR digitally filters and splits the data into sub-bands which are then separately cross correlated and integrated before being stitched together again to yield the final wideband spectrum. The new WIDAR correlator and its superior bandwidth capability provides the VLA with greater sensitivity, allowing the detection of lower flux density sources than was previously possible with the old VLA. A comparison of the performance parameters of the VLA with those of the old VLA is shown in Table 3.5. The three major new observational abilities of the VLA are:

1. Complete frequency coverage between 1 and 50 GHz opening up new regions of the electromagnetic spectrum to astronomy.
2. An increase in continuum sensitivity by an order of magnitude at some frequencies, by increasing the bandwidth to 8 GHz per polarization.
3. Process the large bandwidth with a minimum of 16,384 spectral channels per baseline.

Apart from the addition of more feeds at the center of the reflector, the structural design of the VLA has not changed during its recent upgrade. As before it consists of 27 fully steerable alt-azimuth antennas arranged along the arms of an upside-down ‘Y’ as shown in Figure 2.9. The array is re-configurable and can vary its resolution by over a factor of  $\sim 50$  through movement of its component antennas along twin railroad tracks. Four standard configurations of antennas along the arms of the array are possible whose scales vary by the ratios 1 : 3.28

### **3.5 The Karl G. Jansky Very Large Array**

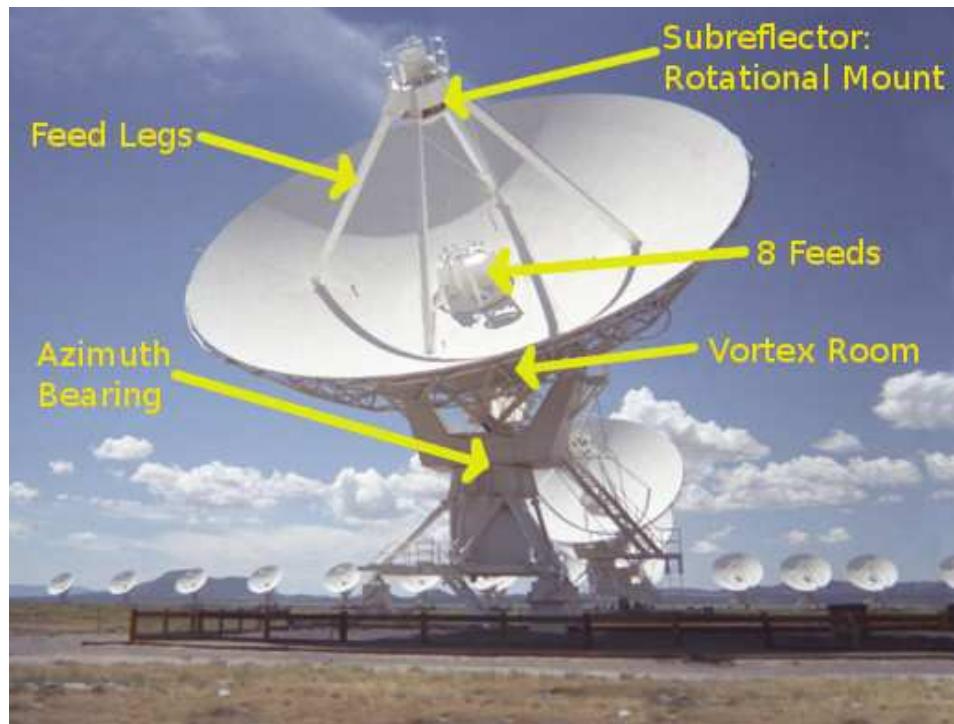
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: 10.8 : 35.5 from smallest to largest. These are called D, C, B, and A configurations, with A having the longest baselines ( $\sim 36$  km) giving the best spatial resolution, but lacking short baselines needed for imaging extended structure. In each configuration, the distance of each antenna from the center of the ‘Y’ is equal to  $m^{\ln 2}$  where  $m$  is the antenna location number, counting outwards from the center of each arm. With this design, the  $m$ ’th station in any configuration coincides with the  $2m$ ’th station in the next smaller configuration. This means that only 72 stations are needed to handle all four configurations. Additionally, there are 3 hybrid configurations called DnC, CnB, and BnA, which are well suited for sources with low declination. In these configurations, the North arm antennas are deployed in the next larger configuration than the SE and SW arm antennas resulting in a more circular synthesized beam for these sources.

Each antenna is 25 m in diameter giving the array a total collecting area equivalent to a single dish of 130 m in diameter. Each antenna has an off-axis Cassegrain design with a rotatable sub-reflector at the prime focus of the main reflector and is supported by four feed legs as shown in Figure 3.5. All feeds are located on a feed ring at the Cassegrain focus and the observing feed is changed by rotating the asymmetric sub-reflector about the main reflector axis so that the secondary focal point moves to the desired feed. The standard observing mode for all feeds is circular polarization. RF signals from each feed are sent via a waveguide to the antenna vortex room located directly underneath the main reflector where they are feed into low noise front ends. The vortex room is temperature controlled and also contains cryogenic cooling systems for the front end, portions of the LO, and IF equipment. IF signals from each antenna are sent by cable to a shielded room where the signals are cross correlated.

All VLA antennas are outfitted with eight receivers providing continuous frequency coverage between 1 and 50 GHz. As shown in Table 3.6, the frequency ranges of 1-2 GHz, 2-4 GHz, 4-8 GHz, 8-12 GHz, 12-18 GHz, 18-26.5 GHz, 26.5-40 GHz, and 40-50 GHz, are commonly referred to as L, S, C, X, Ku, K, Ka, and Q bands, respectively. Additionally, the VLA is currently being outfitted with even lower frequency receivers, P-band (230-470 MHz) and 4-band (54-86 MHz). The VLA’s spatial resolution is set by the maximum baseline  $B_{\max}$  and frequency of observation. This means that structures smaller than the diffraction

### 3.5 The Karl G. Jansky Very Large Array



**Figure 3.5:** Main features of a VLA antenna. The sub-reflector is located at prime focus on a rotational mount is supported by four feed legs. The 8 feeds are located in a ring at the secondary focus. The feeds send the RF signal to the front end located in the vortex room directly beneath the main reflector.

limit ( $\theta_{\text{HPBW}} \sim \lambda/B_{\max}$ ) will be smoothed to the resolution of the array. Table 3.6 also summarizes the maximum resolution for each of the four main configurations at each wavelength. The resolution is defined here as the HPBW of the synthesized beam, using uniform weighting, over a full 12 hour synthesis observation of a source which passes near the zenith. For completeness, we also give the field of view (FOV) at each observing frequency in Table 3.6, defined as the HPBW of the primary beam, which for the VLA antennas can be approximated using the formula:  $\text{FOV}(') = 45/\nu_{\text{GHz}}$ .

### **3.6 VLA Observations of Arcturus and Aldebaran**

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**Table 3.6:** Frequency coverage, primary beam, and spatial resolution of the VLA.

	L	S	C	X	Ku	K	Ka	Q
$\nu$ (GHz)	1.5	3.0	6.0	10	15	22	33	45
$\lambda$ (cm)	20	13	6.0	3.0	2.0	1.3	1.0	0.7
$\nu$ Range (GHz)	1-2	2-4	4-8	8-12	12-18	18-26.5	26.5-40	40-50
A config: $\theta_{\text{HPBW}}$ ('")	1.3	0.65	0.33	0.20	0.13	0.089	0.059	0.043
B config: $\theta_{\text{HPBW}}$ ('")	4.3	2.1	1.0	0.6	0.42	0.28	0.19	0.14
C config: $\theta_{\text{HPBW}}$ ('")	14	7.0	3.5	2.1	1.4	0.95	0.63	0.47
D config: $\theta_{\text{HPBW}}$ ('")	46	23	12	7.2	4.6	3.1	2.1	1.5
FOV(')	30	15	7.5	4.5	3.0	2.0	1.4	1.0

## **3.6 VLA Observations of Arcturus and Aldebaran**

The Open Shared Risk Observing (OSRO) program at the VLA existed during its commissioning phase to provide observers with early access to a number of VLA correlator capabilities and observing modes. This represented a considerable improvement over the capabilities of the old VLA correlator as observes were provided with increased bandwidth capability at existing VLA bands, increased spectral resolution capabilities, and access to new spectral bands. In September 2010 our proposal (PI: G. M. Harper, Program ID: 10C-105) to observe two archetypical red giants at multiple frequencies was allocated the requested 15.5 hours of observing time with the VLA as part of NRAO’s OSRO Science Program 2010C. A number of observing scripts called scheduling blocks (SBs) were prepared during December 2010 and their duration were kept to  $\leq 2.5$  hours to increase their likelihood of being scheduled. The VLA now uses dynamic scheduling for deciding which SBs are executed at any time. This takes into account many factors like the scheduling priority assigned by the time allocation committee, weather constraints, and SB duration. Dynamic scheduling means that the observer does not know when their observations will occur but in general, the chances of observations being scheduled are increased if the duration of the SB is kept short.

**Table 3.7:** VLA Observations of  $\alpha$  Boo and  $\alpha$  Tau obtained in February 2011 and July 2012.

Star	Date	Band	$\nu$ (GHz)	$\lambda$ (cm)	Time on Star (hr)	Restoring Beam ("×")	Bandwidth (GHz)	Number of Antennas	Phase Calibrator
$\alpha$ Boo	2011 Feb 22	Q	43.3	0.7	0.3	$0.19 \times 0.15$	0.256	22	J1357+1919
	2011 Feb 22	Ka	33.6	0.9	0.2	$0.25 \times 0.20$	0.256	23	J1357+1919
	2011 Feb 22	K	22.5	1.3	0.4	$0.35 \times 0.28$	0.256	24	J1357+1919
	2011 Feb 11	X	8.5	3.5	0.3	$1.14 \times 0.70$	0.256	18	J1415+1320
	2011 Feb 11	C	5.0	6.0	0.5	$2.02 \times 1.30$	0.256	21	J1415+1320
	2011 Feb 13	S	3.1	9.5	1.8	$2.57 \times 2.08$	0.256	12	J1415+1320
	2012 Jul 19	S	3.0	10.0	0.7	$2.82 \times 2.30$	2.0	23	J1415+1320
	2012 Jul 20	L	1.5	20.0	1.6	$4.46 \times 3.94$	1.0	23	J1415+1320
$\alpha$ Tau	2011 Feb 11	Q	43.3	0.7	0.3	$0.18 \times 0.16$	0.256	22	J0431+1731
	2011 Feb 11	Ka	33.6	0.9	0.2	$0.22 \times 0.20$	0.256	19	J0449+1121
	2011 Feb 11	K	22.5	1.3	0.4	$0.35 \times 0.31$	0.256	21	J0449+1121
	2011 Feb 13	X	8.5	3.5	0.5	$0.85 \times 0.78$	0.256	25	J0449+1121
	2011 Feb 13	C	5.0	6.0	1.2	$1.48 \times 1.32$	0.256	21	J0449+1121
	2011 Feb 12	S	3.1	9.5	1.8	$2.74 \times 2.02$	0.256	11	J0431+2037

### **3.6 VLA Observations of Arcturus and Aldebaran**

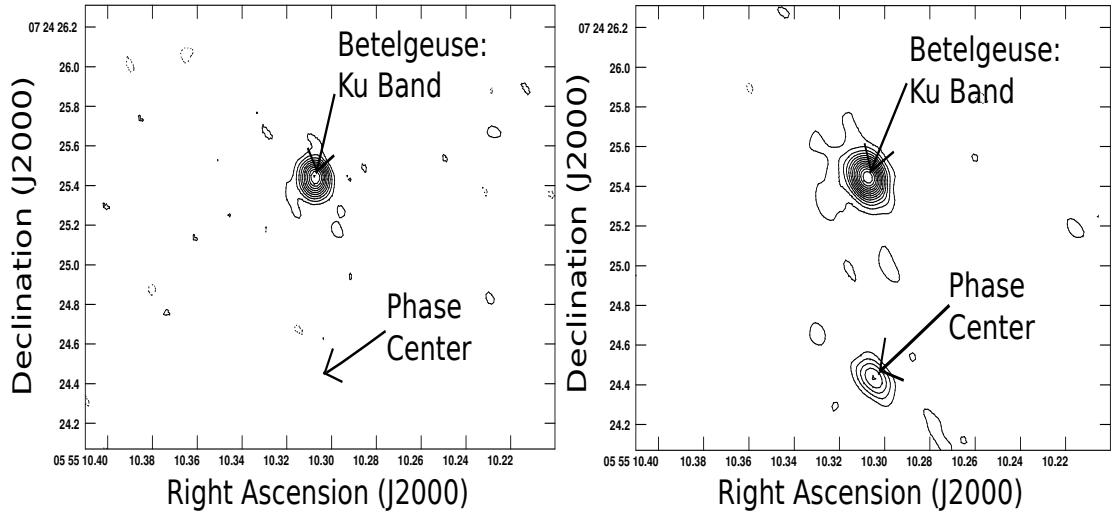
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Our main set of observations took place in February 2011 while the VLA was in B configuration. All observations were taken in continuum mode and the correlator was set up with two 128 MHz sub-bands centered on the frequencies listed in Table 3.7. Each sub-band had sixty-four channels of width 2 MHz and four polarization products (RR, LL, RL, LR). We obtained all our requested observations of  $\alpha$  Tau in just two days between the 11<sup>th</sup> and 13<sup>th</sup> of February 2011 which consisted of Q, Ka, K, X, C, and S band observations of the star. We did not request L band (i.e. 1.5 GHz) observations of  $\alpha$  Tau as it was believed that the star would be too faint to be observable at this frequency. There was also insufficient Ku band receivers available at the time to carry out observations at 15 GHz. We obtained Q, Ka, K, X, C, and S band observations of  $\alpha$  Boo in eleven days between the 11<sup>th</sup> and 22<sup>nd</sup> of February 2011. We also had prepared a 2.5 hour SB for  $\alpha$  Boo at L-band but this SB was never executed.

For this reason we applied for (and were awarded) 3 additional hours of directors discretionary time (DDT) in early 2012 (PI: E. O’Gorman, Program ID: 12A-472) to observe  $\alpha$  Boo at S and L band. We decided to include a short observation at S band even though we already had an observation at this band to make sure that the stars flux density had not significantly changed over that period and so any possible L band detection could be included in the analysis of the main set of data from the previous year. Our DDT observations took place in July 2012 when the VLA was again in B configuration with details given in Table 3.7. The capabilities of the VLA had greatly increased in the  $\sim 1.5$  years since the main set of observations and we now could utilize the full 1 and 2 GHz of bandwidth at L and S band, respectively. The 1-2 GHz and 2-4 GHz frequency ranges were both divided into 16 sub-bands, each with sixty-four channels. The channel width was 2 and 1 MHz for S and L-band, respectively, and each sub-band had four polarization products (RR, LL, RL, LR).

In radio interferometry, baseline-dependent additive errors in the visibilities can occasionally lead to artifacts occurring at phase center of the final image. Such errors may be caused by unflagged low level interference picked up by some antennas baselines. An example of this is demonstrated in Figure 3.6 in which two radio images of Betelgeuse at 15 GHz are shown (A. Brown, priv. comm.). The left panel was taken on 2<sup>nd</sup> February 2002 and shows a  $30\sigma$  detection of Betelgeuse

### 3.6 VLA Observations of Arcturus and Aldebaran

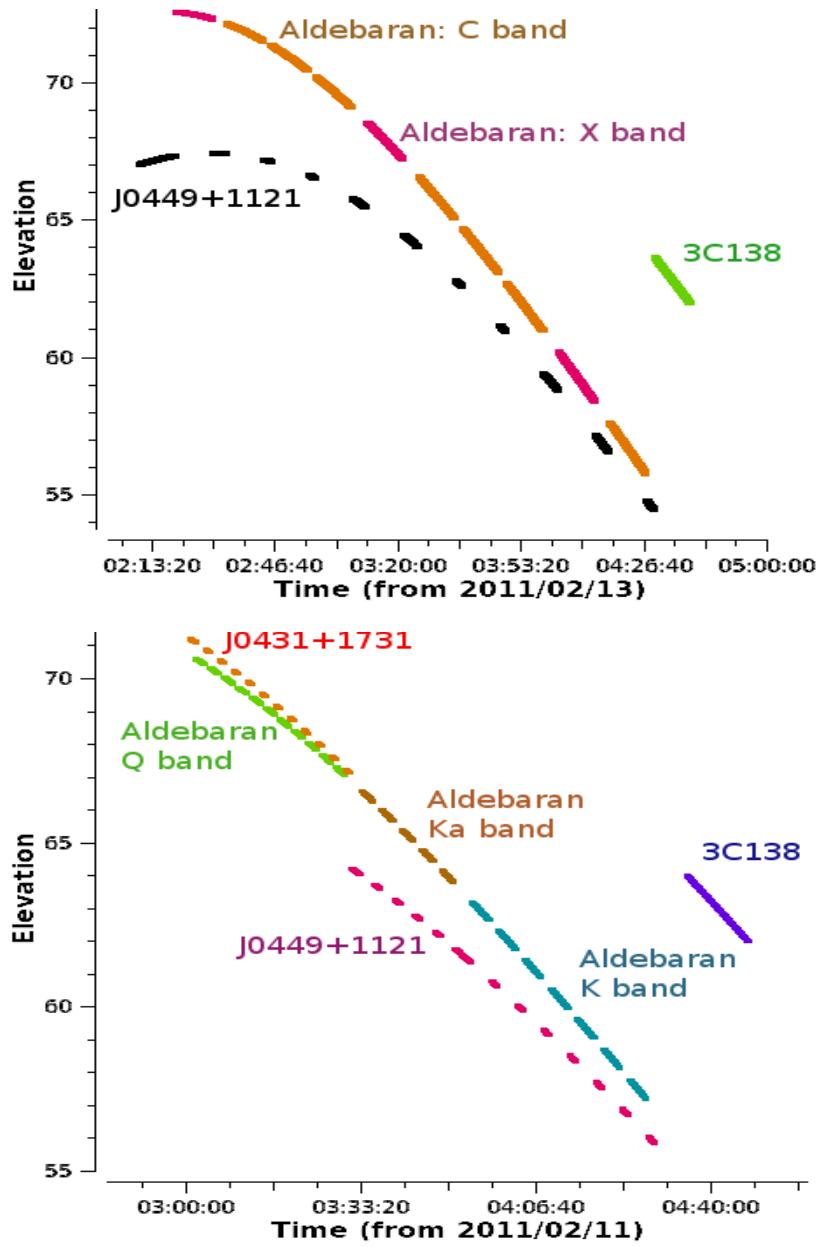


**Figure 3.6:** An example to highlight the importance of offsetting the source from phase center. *Left:* Old VLA image of Betelgeuse at 15 GHz taken on showing no source at the phase center of the image. Contour levels at  $\sigma(-6,-3,3,\dots,30)$  where  $\sigma = 84 \mu\text{Jy}$ . *Right:* Two months later Betelgeuse was again imaged at 15 GHz but now shows a strong artifact at phase center. Contour levels at  $\sigma(-6,-3,3,\dots,45)$  where  $\sigma = 90 \mu\text{Jy}$  (A. Brown, priv. comm.).

with low level background noise. The right panel which was taken two months later on 2<sup>nd</sup> April 2002 shows a  $45\sigma$  detection of the star but also now shows a  $15\sigma$  artifact at phase center. If the target had been observed at phase center in this case, then this artifact would lead to an incorrect flux density measurement for the star. For our VLA observations, both  $\alpha$  Boo and  $\alpha$  Tau were slightly offset from the phase-center by  $\sim 5$  synthesized beam widths in order to avoid source contamination by these rare but possible phase center artifacts and therefore avoiding spurious detections.

The main problem at low VLA frequencies (L and S bands) is disturbances in the ionosphere caused by solar activity. At L-band, solar flares can be as strong as  $1 \times 10^6$  Jy and are a major source of interference, with their effects sometimes being impossible to remove from the data. To avoid such problems the S and L band SBs were scheduled for night time observing only. The low to intermediate frequency observations (L - X bands) were composed of repeatedly interleaved observations of the target and a nearby phase calibrator with cycle times of 12 minutes; 10 minutes on the target and 2 minutes on the phase calibrator. For  $\alpha$

### 3.6 VLA Observations of Arcturus and Aldebaran



**Figure 3.7:** Overview of a low and high frequency VLA SB for  $\alpha$  Tau. *Top Panel:* The low frequency observations consisted of interleaved observations of the target and a nearby phase calibrator with cycle times of 12 minutes. The X band observations were interspersed between C band observations to obtain a good spread in  $u-v$  coverage. *Bottom Panel:* The high frequency observations had short cycle times to compensate for tropospheric effects. Q band observations were taken first to ensure the best pointing solutions were applied to them.

### **3.6 VLA Observations of Arcturus and Aldebaran**

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Boo, the point source J1415+1320 which is located  $6^\circ$  away was used as the phase calibrator at these frequencies. For  $\alpha$  Tau, the point source J0449+1121 located  $6^\circ$  away was used at C and X band, while the brighter J0431+2037 located  $4^\circ$  away was used at S band due to it being unresolved at this frequency (but is resolved at C and X band). The primary calibration sources 3C286 and 3C138 were observed at the end of all low and intermediate frequency SBs and were used to measure the bandpass and set the absolute flux for  $\alpha$  Boo and  $\alpha$  Tau, respectively. The top panel in Figure 3.7 shows a plot of elevation against time for all the sources observed in the SB of the C and X band observations of  $\alpha$  Tau. In this SB, observations of  $\alpha$  Tau at X band were interspersed between C band observations throughout the track to obtain a good spread in  $u - v$  coverage.

At high VLA frequencies (i.e. K, Ka, and Q band) the troposphere can cause major phase variations to incoming radio waves. To reduce this problem, the old VLA used a technique called *fast switching* which reduced the setup overhead and slewing compared to traditional iterating between source and calibrator scans. This overhead is sufficiently reduced with the new VLA and no special fast switching mode is necessary. Instead regular but short duration source-calibrator loops are implemented. As a result, the calibration overheads for high frequency observing are typically considerably larger than for lower frequency observations. For both stars, the total cycle times for the Q, Ka, and K-band observations were 160, 230, and 290 s, respectively. These high frequency observations were combined into a single 2 hour SB for each star and commenced with X-band reference pointing with solutions being applied on-line. As mentioned in Section 2.1.3, the blind pointing errors of the VLA antennas can occasionally be as large as the HPBW of the primary beam at high frequencies. Thankfully, the pointing can be calibrated for using a technique known as *reference pointing* whereby a nearby known calibrator is observed in interferometer pointing mode every hour or so. The measured local pointing corrections are then be applied to subsequent target observations. Reference pointing can reduce the rms pointing errors to as little as  $2''$  if the reference source is with  $10^\circ$  of the target source. After X-band pointing the target source was observed at Q-band to ensure the best pointing solutions were used as shown in the lower panel of Figure 3.7 for  $\alpha$  Tau.

## 3.7 VLA-Pie Town Observations of Betelgeuse

As part of our study into Betelgeuse’s extended atmosphere we decided to re-visit some archival high spatial resolution VLA data from the early 2000’s. This was motivated by the recent findings of [Richards \*et al.\* \(2013\)](#) who detected two hot chromospheric features with the very long baseline interferometer e-MERLIN. The two features, also referred to as ‘hot spots’, were found to have brightness temperatures well above the mean brightness temperature measured at the same radio wavelength, and were separated from each other by 90 mas. We wanted to see if these features were present at other wavelengths and if so, how they varied over time.

Asymmetries in Betelgeuse’s atmosphere were first detected in the radio by [Skinner \*et al.\* \(1997\)](#) and later confirmed by [Lim \*et al.\* \(1998\)](#). To see how these asymmetries evolved with time, a multi-wavelength study into the temporal evolution of Betelgeuse’s atmosphere was carried out with the old VLA over a number of years in the early 2000’s in its most extended A configuration plus the Pie Town link<sup>1</sup>. This study focused on fitting various models to the visibility data as was done by [Lim \*et al.\* \(1998\)](#), i.e., the analysis was mainly carried out in the visibility plane. The detection of the two unique hot spots in the e-MERLIN radio maps by [Richards \*et al.\* \(2013\)](#) inspired us to revisit this data which, because of the addition of the Pie Town link to the VLA, gave spatial resolution comparable to or exceeding that of e-MERLIN at some wavelengths. Our goal was to carry out the analysis in the image plane to make a direct comparison with the findings of [Richards \*et al.\* \(2013\)](#). In Table 3.8 we summarize this unique data set which consists of almost four years of high resolution multi-wavelength observations of Betelgeuse. To increase our observational time period we also include the short wavelength high resolution (i.e., A configuration) observations of the star from 1998. Betelgeuse was resolved at all wavelengths listed in Table 3.8 except at L band.

#### ***The VLA-Pie Town Link***

The 25 m antenna located at Pie Town, New Mexico is one of 10 antennas spread

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<sup>1</sup>This data was carefully calibrated by Dr. Alexander Brown.

### 3.7 VLA-Pie Town Observations of Betelgeuse

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**Table 3.8:** Multi-wavelength VLA + Pie Town Observations of Betelgeuse

Date	Band	$\lambda$ (cm)	Time on Star (hr)	Restoring Beam (mas)	rms noise (mJy/beam)
2004 Oct 21,30	Q	0.7	1.2	$39 \times 26$	0.37
	K	1.3	2.6	$80 \times 41$	0.09
	U	2.0	3.3	$121 \times 91$	0.08
	X	3.5	3.8	$208 \times 126$	0.02
	C	6.2	2.9	$377 \times 265$	0.02
	L	16.7	2.4	$1262 \times 889$	0.03
2003 Aug 10,12	Q	0.7	0.9	$40 \times 27$	0.46
	K	1.3	2.6	$80 \times 42$	0.18
	U	2.0	3.3	$119 \times 96$	0.10
	X	3.5	3.0	$204 \times 139$	0.03
	C	6.2	2.8	$378 \times 297$	0.03
	L	16.7	2.6	$1247 \times 931$	0.04
2002 Feb 17,18	K	1.3	2.0	$82 \times 47$	0.14
	U	2.0	2.0	$128 \times 90$	0.11
	X	3.5	2.1	$200 \times 135$	0.03
	C	6.2	1.6	$372 \times 273$	0.3
	L	16.7	1.1	$1312 \times 951$	0.05
2002 Apr 12,13	K	1.3	2.0	$91 \times 59$	0.18
	U <sup>a</sup>	2.0	2.0	$130 \times 98$	0.39
	X	3.5	2.1	$224 \times 154$	0.03
	C <sup>b</sup>	6.2	1.6	$406 \times 295$	0.3
	L	16.7	0.9	$1397 \times 1146$	0.07
2001 Jan 02	K	1.3	6.0	$78 \times 42$	0.08
2000 Dec 23	Q	0.7	6.0	$44 \times 20$	0.18
1998 Mar 29,30	Q <sup>c</sup>	0.7	5.3	$39 \times 36$	0.38
	K <sup>c</sup>	1.3	6.3	$114 \times 89$	0.25

<sup>a</sup> Spurious source seen at phase center of image.

<sup>b</sup> The star was not detected in this data set.

<sup>c</sup> No Pie Town link was used in this data.

across the United States of America which together make up the Very Long Baseline Array (VLBA). It was successfully connected in real time via a 104 km long fiber optic line to the VLA in 1999 and was operational at all old VLA frequencies

### **3.7 VLA-Pie Town Observations of Betelgeuse**

prior to its upgrade to the Karl G. Jansky VLA. The Pie Town antenna is located approximately 52 km west of the center of the VLA, effectively doubling its resolution while in A configuration when connected. It is worth noting that the Pie Town link was only functional when the VLA was in A configuration. The extra collecting area together with 27 additional baselines also increases the sensitivity of the array. The Pie Town antenna generates IF signals in VLBA format and these are converted, at Pie Town, to a 50 MHz wide VLA signal where they are multiplexed and then modulated onto a single laser before being sent over 104 km of fiber to the VLA. When the signal arrives at the VLA site it is received and converted back into an RF signal. The received RF signal is then amplified and filtered before its is sent on to the backend receiver.

# 4

## Data Analysis

### 4.1 Measurement Equation

# 5

## Multi-wavelength Study of Betelgeuse's Outer Atmosphere

In this chapter we present the results from our multi-wavelength high resolution study of Betelgeuse's extended atmosphere. The majority of this chapter is based on the findings of O'Gorman *et al.* (2012) and detail the results of our sub-arcsecond resolution study into the distribution and kinematics of CO<sup>1</sup> around Betelgeuse using the CO( $J = 2 - 1$ ) rotational line as the main probe. The variation of the emission line profile with different CARMA array configurations provides us with a unique insight into the distribution and kinematics of the two previously known flows, S1 and S2, around Betelgeuse. We successfully image the star on spatial scales as small as  $0''.9$  ( $\sim 20 R_\star$ ) over a  $32''$  ( $\sim 750 R_\star$ ) field of view and we discuss these radio maps in detail. We compare these results to higher CO rotational lines obtained with other instruments. We also present the recent findings of Richards *et al.* (2013) who image two distinct chromospheric features at centimeter wavelengths. These surprising findings motivated us to re-analyze old VLA-Pie Town data to search for these features at shorter wavelengths

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<sup>1</sup>Throughout this chapter the usage of the term CO refers to the molecular isotope  $^{12}\text{C}^{16}\text{O}$ .

## 5.1 CO molecules in the CSE of Betelgeuse

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with comparable resolution, the results of which are presented at the end of this chapter.

### 5.1 CO molecules in the CSE of Betelgeuse

The CSE of Betelgeuse ( $\alpha$  Ori) is a proving ground for ideas and theories of mass loss from oxygen-rich M-type supergiants. Currently it is losing mass at a respectable rate  $\sim 2 \times 10^{-6} M_{\odot}$  yr $^{-1}$  (Harper *et al.*, 2001), as it has been over the past  $\sim 1000$  years. Most of the optically thin silicate dust lies beyond  $\sim 46$  stellar radii (Danchi *et al.*, 1994) and dust is, therefore, unlikely to be responsible for the bulk mass loss. This raises the important point that if the mass loss from Betelgeuse is not a result of dust then perhaps the same mechanisms that are responsible might also be active in the more dusty later M-type supergiants. Radiation pressure on atoms and molecules is another potential contributing candidate as a mass loss mechanism and so spatial and dynamical studies of molecules are a fruitful line of investigation, especially in relation to eventual formation of dust. Such studies also allow us to calculate the time scales on which certain mass loss episodes have occurred, and these can then be compared to the time scales of potential mass-loss initiators such as convection or magnetic dynamo cycles.

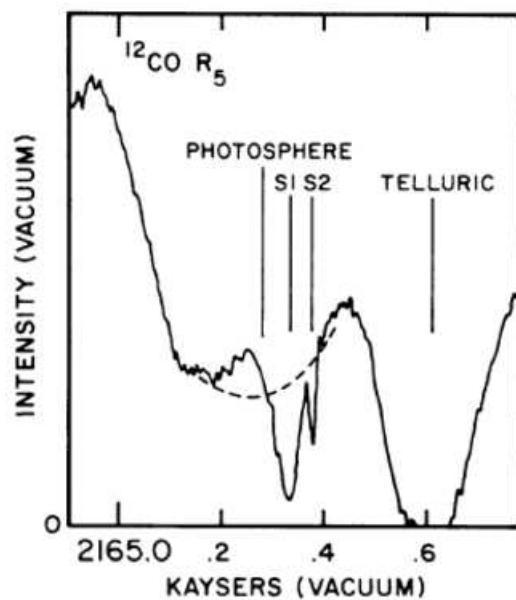
Despite the initial reports of an enhancement of C abundance in its atmosphere (Spinrad & Vardya, 1966), Betelgeuse is now known to be an oxygen-rich star. The comprehensive study of CO and OH ro-vibrational bands by Lambert *et al.* (1984) found  $\log \epsilon(\text{C}) = 8.4$  and  $\log \epsilon(\text{O}) = 8.8$ . For any species X

$$\log \epsilon(\text{X}) = \log (\text{X}/\text{H}) + 12, \quad (5.1)$$

which can be used to show O/C = 2.5. Lambert *et al.* (1984) also found  $\log \epsilon(\text{N}) = 8.6$  which indicates a contamination of the atmosphere with C-N processed material. This is in agreement with the theoretical work of Lamb *et al.* (1976) who predict the dredging up of processed material by the deep convective envelope which a supergiant like Betelgeuse is expected to possess. Betelgeuse is unlikely to have undergone a second dredge-up of material that would cause small changes in the CNO abundances due to its mass being greater than the upper mass limit

## 5.1 CO molecules in the CSE of Betelgeuse

for a second dredge-up to occur. This second dredge-up is predicted to occur after He core exhaustion for stars below  $8 - 10 M_{\odot}$  (Kaler *et al.*, 1978). The observed  $^{12}\text{C}/^{13}\text{C}$  ratio, which is an indication of the history of nuclear processing and mixing in the interior of stars (Eggleton *et al.*, 2007; Pavlenko *et al.*, 2003) has consistently found to be between 6 and 7 (Bernat *et al.*, 1979; Harris & Lambert, 1984; Hinkle *et al.*, 1976; Lambert *et al.*, 1974), well below the predicted value of 20 (Charbonnel, 1995; Pavlenko *et al.*, 2003), which is the case for most red giant stars (e.g., Boothroyd & Sackmann, 1999).

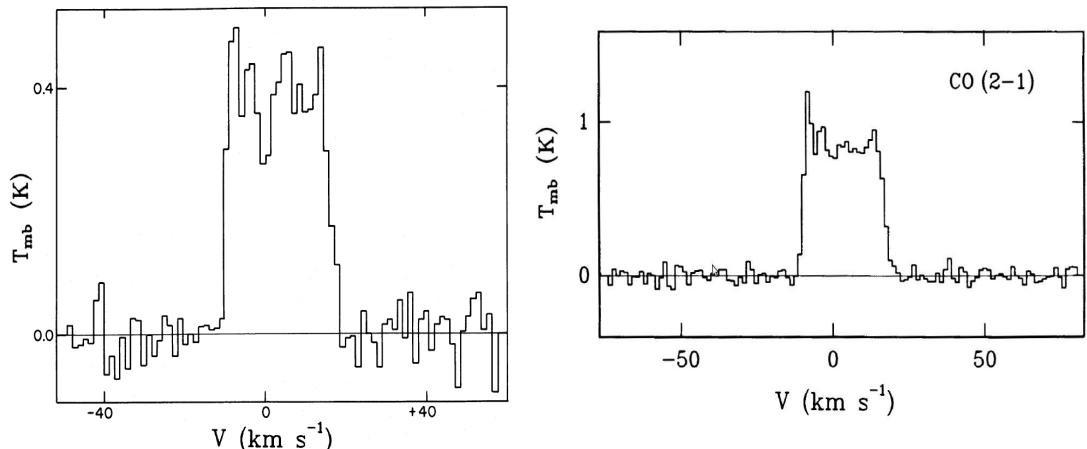


**Figure 5.1:** CO ro-vibrational fundamental line showing two sharp blueshifted absorption features originating from the expanding envelope (Bernat *et al.*, 1979). These two features became known as the S1 and S2 components of Betelgeuse's outflow having outflow velocities of  $\sim 9$  and  $\sim 16 \text{ km s}^{-1}$ , respectively. The units of the x-axis are  $\text{cm}^{-1}$ .

The study of CO molecules in the CSE of Betelgeuse began with the detection of  $4.6 \mu\text{m}$  ro-vibrational absorption lines of  $^{12}\text{C}^{16}\text{O}$  and  $^{13}\text{C}^{16}\text{O}$  by Bernat *et al.* (1979) who identified two absorption features shown in Figure 5.1, implying two distinct components within the overall outflow. One component, known as S1, has a Doppler shift of  $9 \text{ km s}^{-1}$  towards us with  $T_{\text{exc}} \simeq 200 \text{ K}$ ,  $v_{\text{turb}} \simeq 4 \text{ km s}^{-1}$ ,

## 5.1 CO molecules in the CSE of Betelgeuse

and  $N_{\text{CO}} = 4.7 \times 10^{17} \text{ cm}^{-2}$ . The second faster component, known as S2, has a Doppler shift of  $16 \text{ km s}^{-1}$  towards us with  $T_{\text{exc}} \simeq 70 \text{ K}$ ,  $v_{\text{turb}} \simeq 1 \text{ km s}^{-1}$ , and  $N_{\text{CO}} = 1.2 \times 10^{16} \text{ cm}^{-2}$ . [Bernat et al. \(1979\)](#) noted that a simple assumption of the dust and gas being in thermal equilibrium would infer that the S1 and S2 components have a spatial extent of 4 and  $55''$  respectively, using the dust model of [Tsuji \(1979\)](#). The S1 feature with its higher column density was well known from atomic absorption line studies (e.g. [Weymann, 1962](#)) and both features had been detected in high spectral resolution atomic Na and K absorption profiles ([Goldberg et al., 1975](#)).



**Figure 5.2:** Previous single dish CO( $J = 2 - 1$ ) rotational emission line profiles of Betelgeuse. *Left:* [Huggins \(1987\)](#) line profile using a HPBW of  $32''$ . They found some evidence for an S2 radius of about  $16''$ . *Right:* The line of [Huggins et al. \(1994\)](#) looks very similar even using a smaller HPBW in conflict with the findings of [Huggins \(1987\)](#). Velocities are plotted in the local standard of rest (LSR) frame.

CO was first detected at radio wavelengths by [Knapp et al. \(1980\)](#) who presented a low S/N ratio spectrum of the  $J = 2 - 1$  rotational emission line at 230 GHz. They used a single dish antenna with a HPBW of  $\sim 30''$  and detected one component expanding at  $15 \text{ km s}^{-1}$ . They assumed this component to be the S2 component detected by [Bernat et al. \(1979\)](#) and by assuming that they had resolved a flat top profile, predicted that its spatial extent didn't exceed  $10''$ . The weaker CO( $J = 1 - 0$ ) line was tentatively detected by [Knapp & Morris \(1985\)](#) with a 7 m dish which had a HPBW of  $100''$ . [Huggins \(1987\)](#) presented a higher

signal-to-noise CO( $J = 2 - 1$ ) profile that was obtained with a HPBW of  $32''$  and is shown in Figure 5.2. By comparing the  $(2 - 1)/(1 - 0)$  line intensities they found some evidence for an S2 radius of  $\sim 16''$ . However, a 30 m IRAM  $J = 2 - 1$  line profile was later presented by Huggins *et al.* (1994) and as can be seen in Figure 5.2 looked remarkably similar, even though it was observed with a smaller  $12''$  HPBW. The profile did not show the horned wing signature expected if it had been resolved as discussed in Chapter 1, apparently in conflict with the previous S2 radius estimate. These single dish line profiles also showed no signature of the slower S1 shell and so questions remained about the spatial extent of these two distinct outflow components in the CSE of Betelgeuse. A sensitive high spatial resolution study of its atmosphere was needed to untangle this puzzling evidence.

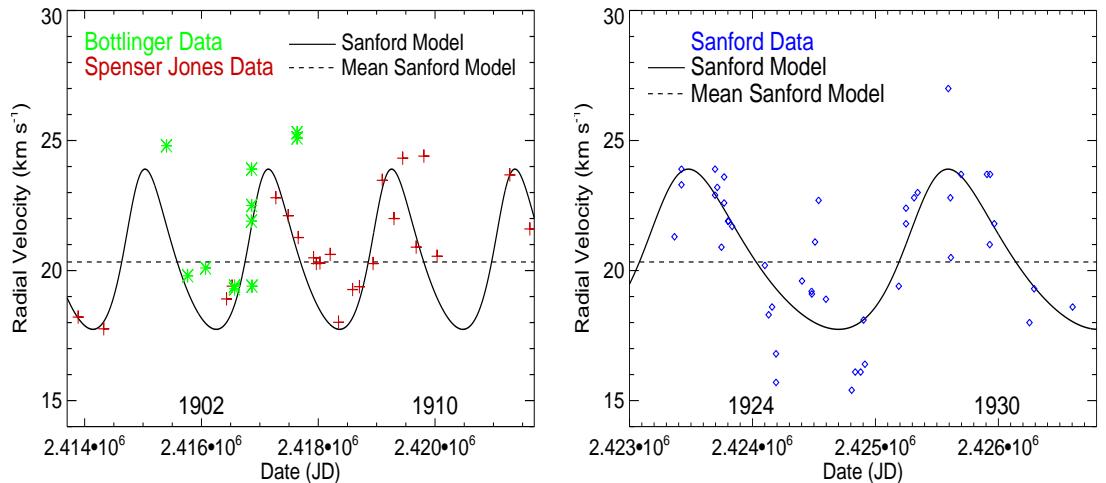
## 5.2 Adopted Radial Velocity

Before we discuss the CO( $J = 2 - 1$ ) spectra obtained from our multi-configuration CARMA observations which are summarized in Chapter 3, we begin by explaining which radial velocity value  $v_{\text{rad}}$ , for Betelgeuse was adopted. An accurate value of  $v_{\text{rad}}$  is necessary when plotting the spectra of any star with respect to its center of mass rest frame, as we do for all spectra in this chapter. Plotting spectra in this frame of reference is intuitive as it allows the reader to immediately see the gas expansion velocity relative to the photosphere. A positive  $v_{\text{rad}}$  denotes recession (i.e., redshifted lines) while a negative  $v_{\text{rad}}$  denotes advancement (i.e., blueshifted lines).

Betelgeuse is a semi-regular variable and its radial velocity exhibits variability on time scales ranging from short 1.5 year periods as suggested by Stebbins & Huffer (1931) to longer 5.8 year periods (Smith *et al.*, 1989; Spencer Jones, 1928). Stothers & Leung (1971) interpreted the long period as being the convective turnover time of giant convection cells on the stellar surface while Dupree *et al.* (1990) attribute the shorter period with pulsation. Betelgeuse's radial velocity amplitudes are also known to vary by at least  $\pm 3 \text{ km s}^{-1}$  (Smith *et al.*, 1989) making it difficult to determine a precise value for the stellar center of mass radial velocity. In Figure 5.3 we plot the radial velocity data and the corresponding model derived in the classical paper by Sanford (1933) which is based on

### 5.3 CARMA CO( $J = 2 - 1$ ) Spectra

observations spanning 1923 to 1931. We have also extrapolated the model back to show that it matches the earlier data of Bottlinger (1911) and Spencer Jones (1928) quite well, as shown in Figure 5.3. Goldberg (1984) has also shown that the model of Sanford (1933) can be extrapolated forward to give a reasonable fit to his measurements. In this study we adopt a heliocentric radial velocity of  $+20.7 \text{ km s}^{-1}$ ; a value used by Weymann (1962) and Harper *et al.* (2008) and is based on the mean values of Spencer Jones (1928) and Sanford (1933) radial velocity models.



**Figure 5.3:** The radial velocity model of Sanford (1933) along with measurements from Bottlinger (1911), Spencer Jones (1928), and Sanford (1933) spanning from 1897 – 1931. In this study we use  $+20.7 \text{ km s}^{-1}$  which is the mean value from the models of Spencer Jones (1928) and Sanford (1933). The radial velocity amplitude variations of  $\sim \pm 3 \text{ km s}^{-1}$  is clearly evident in these data.

### 5.3 CARMA CO( $J = 2 - 1$ ) Spectra

The spectrum for each individual configuration image cube (which are composed of all the appropriate configuration tracks listed in Table 3.3) along with the multi-configuration image cube (which is composed of the entire data set spanning  $\sim 2.5 \text{ yr}$ ) can be used to obtain information on the kinematics of the S1 and S2 flows. In particular, it is of interest to see how the CO( $J = 2 - 1$ ) line profile

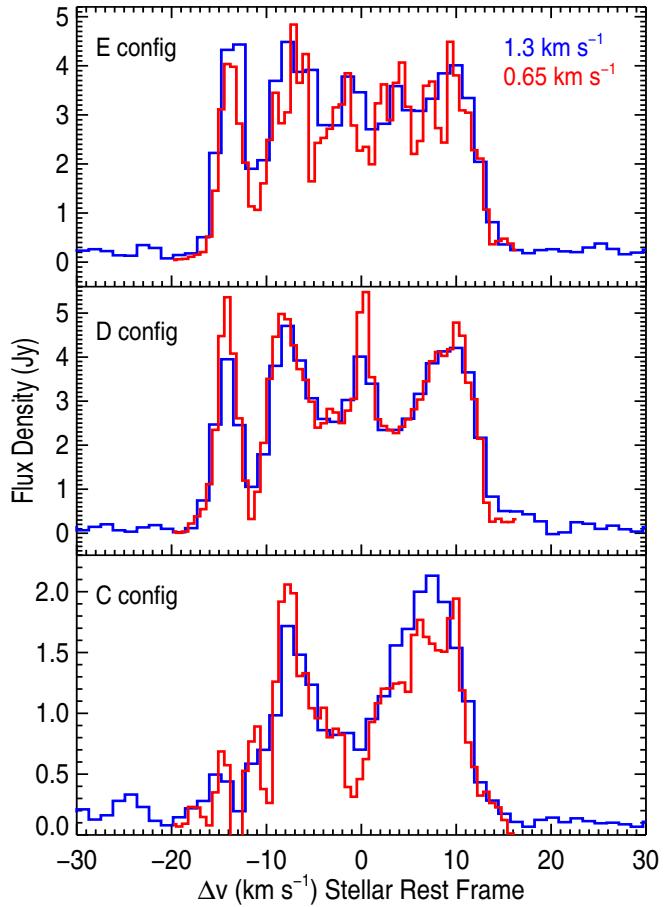
### **5.3 CARMA CO( $J = 2 - 1$ ) Spectra**

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changes when observed with different array configurations as each of these will be sensitive to emission on different spatial scales. . The spectra corresponding to the C, D, and E configuration image cubes are plotted in Figure 5.4 for both the high ( $0.65 \text{ km s}^{-1} \text{ bin}^{-1}$ ) and low ( $1.3 \text{ km s}^{-1} \text{ bin}^{-1}$ ) spectral resolution data and were obtained by integrating all emission within a circular area of radius  $5''$  centered on the source. The high and low spectral resolution modes allow two independent sets of spectra to be measured for each observation and thus provide a good check on the data quality. The high resolution spectra (channel width =  $0.65 \text{ km s}^{-1}$ ) give the best measure of the S1 and S2 flow kinematics and therefore all outflow velocities are derived from these spectra.

The E configuration image cube spectrum has a total line width of  $29.2 \text{ km s}^{-1}$  and the low spectral resolution profile contains a steep blue wing emission feature between  $-16.0 \text{ km s}^{-1}$  and  $-11 \text{ km s}^{-1}$  and a more flat-topped feature between  $-10.3 \text{ km s}^{-1}$  and  $-13.2 \text{ km s}^{-1}$ . This steep emission wing shows that the turbulence in the flow is less than or equal to the velocity bin size. The blue wing in the high resolution profile matches the lower resolution profile well but the remainder of the profile looks more complex than the flat-topped feature seen in the lower resolution profile. The line profile shape has been well documented by previous single dish observations (e.g., Figure 5.2) and, out of our three individual configuration spectra, we expect the most compact E configuration spectra to resemble these single dish measurements the closest due to its better sampling of the inner  $u - v$  plane and consequent sensitivity to extended structures. This indeed turns out to be the case when we compare our three individual configuration spectra to those previous single dish profiles. The blue wing emission feature appears again in the D configuration spectrum at the same velocities as those in the E configuration spectrum but the remainder of the profile appears quite different. Between  $-10.3 \text{ km s}^{-1}$  and  $+13.2 \text{ km s}^{-1}$  the D configuration spectrum is dominated by a blue wing at  $\sim -10.0 \text{ km s}^{-1}$ , a red wing at  $\sim +13.0 \text{ km s}^{-1}$  and an emission feature at  $\sim 0 \text{ km s}^{-1}$ . The large drop in emission at certain velocities in the D configuration spectrum is probably a result of the lower sensitivity to large scale structure in comparison to the E configuration as described in Table 3.2.

The line profile has a lower flux ( $\sim 50\%$ ) in the high spatial resolution C configuration spectrum due to its lack of sensitivity to extended structure. Notably,



**Figure 5.4:** Spectra integrated over a radius of  $5''$  for each array configuration image cube. The blueshifted emission component between  $-16.0 \text{ km s}^{-1}$  and  $-10.0 \text{ km s}^{-1}$  is almost resolved out in the C configuration image cube spectrum. The red and blue lines correspond to the high and low spectral resolution data respectively.

the blueshifted emission feature located between  $-16.0 \text{ km s}^{-1}$  and  $-11.0 \text{ km s}^{-1}$  in the E and D configuration spectra is almost completely spatially filtered by the extended C configuration. This component of the line has previously been associated with the outer S2 flow (Huggins, 1987) and as the majority of it has been spatially filtered by our C configuration we expect even less contribution from the S2 flow at lower absolute velocities still. For the redshifted line emission we again expect the majority of the S2 contribution to be spatially filtered, so we conclude that most, if not all of the emission in the C configuration spectrum

### 5.3 CARMA CO( $J = 2 - 1$ ) Spectra

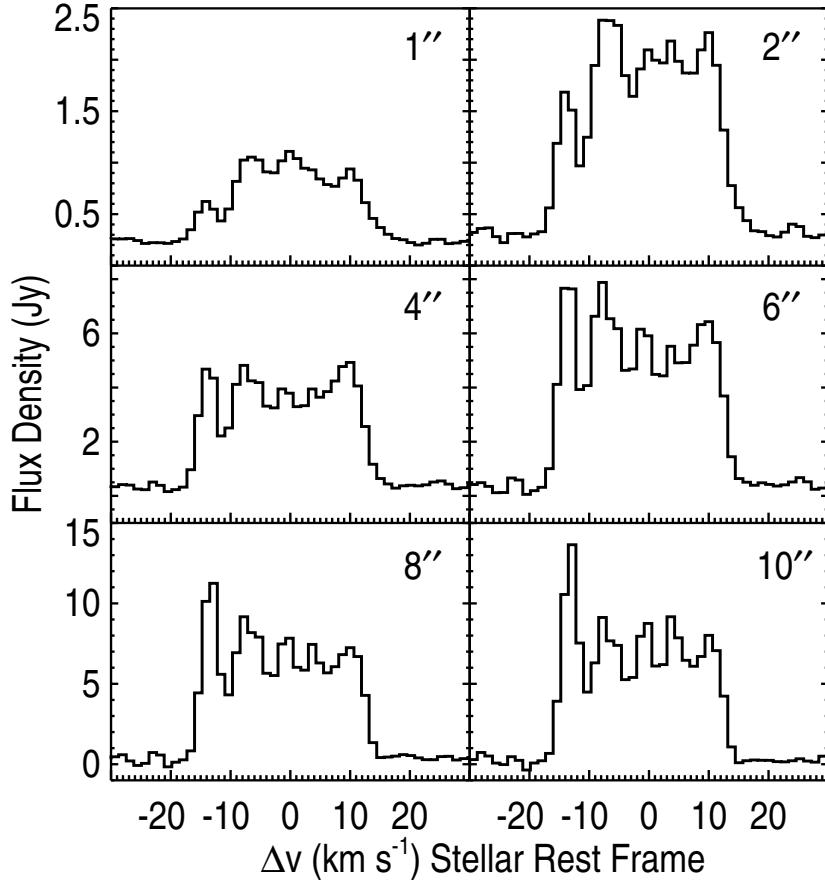
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emanates from the inner S1 flow. The spectrum is double peaked with the blue and redshifted wings extending to  $-9.0 \text{ km s}^{-1}$  and  $+10.6 \text{ km s}^{-1}$  respectively, and we define these as the outflow velocities of S1. As discussed in Chapter 3 the C configuration has a resolving out scale of  $\sim 6''$  at 1.3 mm and so is not sensitive to angular scales larger than this. If the emission between  $-9.0 \text{ km s}^{-1}$  and  $+10.6 \text{ km s}^{-1}$  in the C configuration spectrum appeared as a flat topped profile then we could conclude that the S1 flow lies within a radius of  $3''$  from the star. Clearly however, the profile has a more *horned* shaped appearance and so the lower absolute velocity components of this profile have been spatially filtered. We can therefore conclude that the radial extent of the S1 from the star is greater than  $3''$ . If we assume that the S1 flow would produce a  $19.6 \pm 0.7 \text{ km s}^{-1}$  wide top-hat line profile of 2 mJy were it not for the resolving out effects of the interferometer, its integrated line flux is then

$$F_{\text{tot S1}} = F_\nu \times \left( \frac{\Delta v}{\lambda} \right) = 3.1 \times 10^{-19} \text{ W m}^{-2}. \quad (5.2)$$

To obtain the most robust value for the S2 outflow velocities we examine the high spectral resolution multi-configuration image cube spectrum which is composed of all tracks from all three configurations. It is worth stressing that by analyzing the multi-configuration image cube we make the assumption that the physical properties of all three components (i.e.  $\alpha$  Ori, S1 and S2) have not changed over the total observation period (i.e.  $\sim 2.5$  yr). The profile is found to have a total linewidth of  $28.6 \pm 0.7 \text{ km s}^{-1}$ , which is in good agreement with previous single dish observations of the line where values of  $30.6 \pm 2.5 \text{ km s}^{-1}$  and  $28.6 \text{ km s}^{-1}$  were reported by Knapp *et al.* (1980) and Huggins (1987) respectively. The centroid velocity of the spectrum is  $-1.1 \pm 0.7 \text{ km s}^{-1}$  ( $v_{\text{lsr}} = 3.7 \pm 0.7 \text{ km s}^{-1}$ ) which again is in good agreement with Knapp *et al.* (1980) and Huggins (1987) values of  $v_{\text{lsr}} = 3.0 \pm 2.5 \text{ km s}^{-1}$  and  $v_{\text{lsr}} = 3.7 \pm 0.4 \text{ km s}^{-1}$ , respectively. We use Equation 5.2 to find the integrated total line flux to be  $1.5 \times 10^{-18} \text{ W m}^{-2}$ , of which at least 20% emanates from the S1 flow.

The outflow velocities of S2 are  $-15.4 \text{ km s}^{-1}$  and  $+13.2 \text{ km s}^{-1}$  which, like the S1 flow, are slightly asymmetric but in the opposite sense. Note that the S1 and S2 outflow velocities are dependent on the adopted radial velocity of Betelgeuse.



**Figure 5.5:** Spectral profiles of the low spectral resolution multi-configuration image cube for circular extraction areas of radius 1'', 2'', 4'', 6'', 8'', and 10''. The signal-to-noise of the line profile reduces at larger extraction areas as more noise is included from the outer regions of the channel maps.

If for instance, we instead adopt a radial velocity of 21.9 km s<sup>-1</sup> (Famaey *et al.*, 2005) then the S2 outflow velocities become even more asymmetric (-16.6 and +12.0 km s<sup>-1</sup>) while the S1 outflow becomes less so (-10.2 and +9.4 km s<sup>-1</sup>). Both S1 and S2 therefore cannot have spherically symmetric outflow velocities regardless of the adopted stellar radial velocity. Adopting a mass of  $18 M_{\odot}$  (Meynet & Maeder, 2003) and a radius of  $950 R_{\odot}$  (Harper *et al.*, 2008) then the escape velocity for Betelgeuse is 85 km s<sup>-1</sup> which is much greater than the S1 and S2 outflow velocities. This indicates that the majority of the stellar mass loss mechanism's energy goes into lifting the CO molecules out of the gravitational potential and

## 5.4 Individual Configuration Image Cubes

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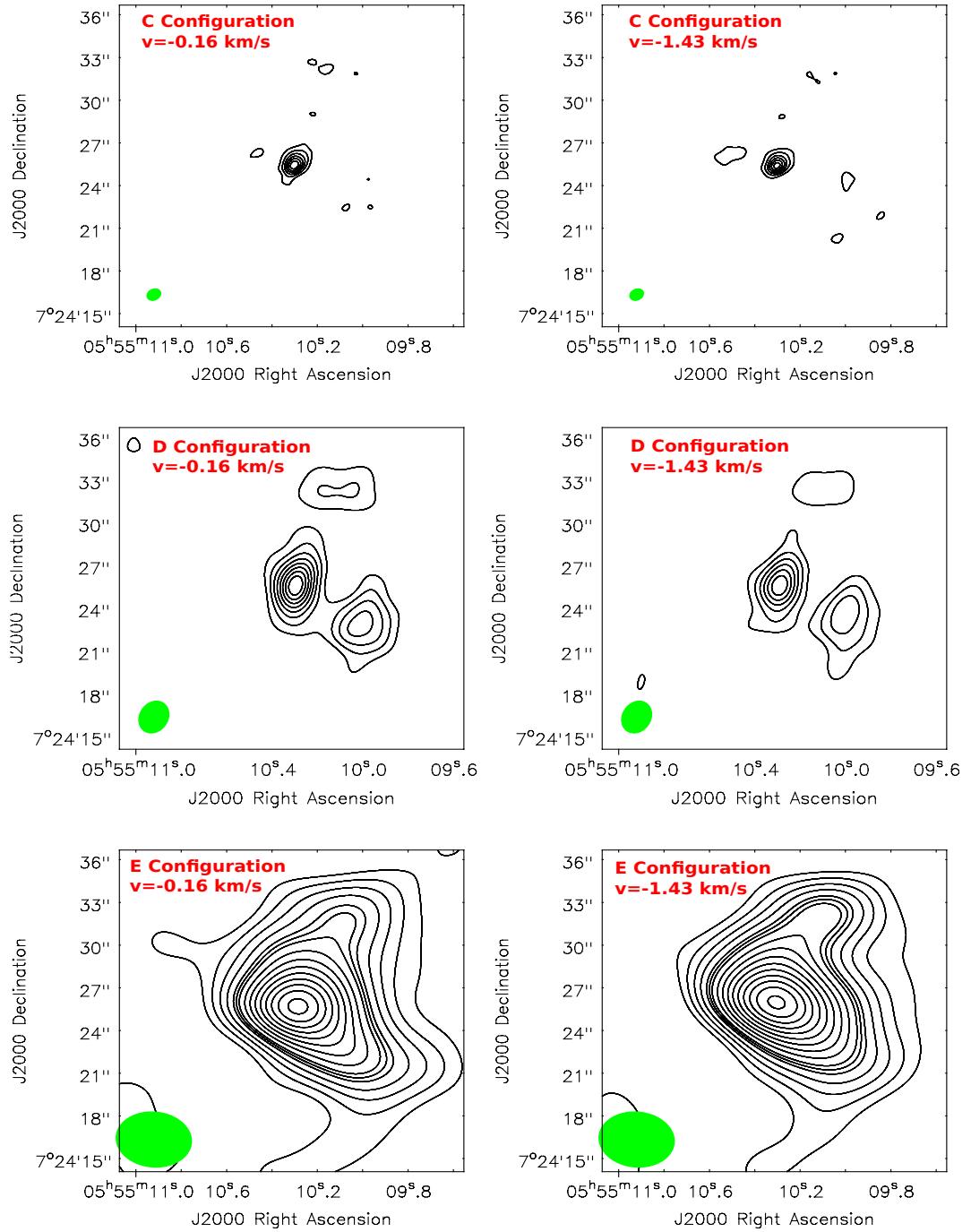
not into their outflow velocities. These outflow velocities are greater than the adiabatic hydrogen sound speed (i.e.,  $11.7\sqrt{10^{-4}T_e}$ ), which, if we assume that the gas temperature is the same as the excitation temperature, are  $1.7 \text{ km s}^{-1}$  and  $1 \text{ km s}^{-1}$  for S1 and S2 respectively.

The spectra in Figure 5.5 are taken from the low resolution multi-configuration image cube using circular extraction areas ranging in radius from  $1''$  to  $10''$  and demonstrates how the line profile changes over these different extraction areas. The most striking change in the line profile is the change in appearance of the extreme blue wing. At small extraction radii where we sample the most compact emission, the feature is weak in comparison to the rest of the line but becomes more dominant as we begin to sample more of the extended emission. This indicates that even the high absolute velocity components of the S2 flow have extended emission and this is why they are almost completely spatially filtered by CARMA's C configuration. The large reduction of flux at  $-11 \text{ km s}^{-1}$  suggests that there is more material moving towards the observer than at other lower absolute velocities indicating a non-isotropic (or non-spherical) S2 flow. This suggests a more sheet like (flatter) structure rather than a spherical cap.

## 5.4 Individual Configuration Image Cubes

We created three image cubes for each configuration by concatenating all good tracks together per configuration. The gradual formation of a ring type structure as one goes from high absolute velocities to lower absolute velocities indicative of a shell was not seen in any of these three image cubes. However, an additional spatially unresolved source was clearly detected in a number of the D configuration image cube maps (both high and low spectral resolution). The component is present in five continuous channel maps between  $\sim -4.0 \text{ km s}^{-1}$  and  $+2.4 \text{ km s}^{-1}$  and is located  $\sim 5''$  S-W of  $\alpha$  Ori. Its peak emission lies at  $\sim 0 \text{ km s}^{-1}$  and has a flux density of  $1.8 \text{ mJy}$  which is approximately 60% of the total source flux density at this velocity. The middle row in Figure 5.6 are two of these D configuration channel maps which clearly show this discrete source. The contour levels are set at the  $4\sigma$  level so the detection of the second source in the D configuration is significant. The corresponding channel maps in the E configuration image cube

## 5.4 Individual Configuration Image Cubes



**Figure 5.6:** Two channel maps from each of the three individual configuration image cubes. A spatially unresolved source is present in five continuous channel maps between  $\sim -4.0 \text{ km s}^{-1}$  and  $+2.4 \text{ km s}^{-1}$  and is located  $\sim 5''$  S-W of  $\alpha$  Ori. Contour levels begin at  $4\sigma$  where  $1\sigma$  is the rms channel noise. The restoring beam is shown in green in the lower left corner of each map.

## 5.5 Multi-configuration Image Cube Inspection

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show extended emission out to  $8''$  in the same S-W direction as can be seen in the lower panel of Figure 5.6. There is only weak emission detected in the same position in the corresponding C configuration channel maps, probably due to the lower sensitivity resulting from the smaller HPBW (i.e., the flux is diluted). This discrete second source thus has the effect of adding extra emission to the corresponding multi-configuration image cube spectrum at the low absolute velocities where it is present. The presence of this addition source at these low absolute velocities is probably one of the reasons why Huggins *et al.* (1994) did not detect a horned shaped spectrum with his  $12''$  HPBW. We also note that there is a weaker emission feature  $7''$  north north-west of  $\alpha$  Ori in these maps also which too can be seen in Figure 5.6.

## 5.5 Multi-configuration Image Cube Inspection

A subset of the blueshifted velocity channel maps of the low spectral resolution multi-configuration image cube is presented in Figure ???. The first channel map at  $-17.9 \text{ km s}^{-1}$  shows just the compact unresolved continuum emission with no extended emission present. Between  $-16.7 \text{ km s}^{-1}$  and  $-9.0 \text{ km s}^{-1}$  we see evidence for the development of a classical shell signature for the S2 flow. We first sample the highest velocity components where the emission is relatively compact (i.e. between  $-16.7 \text{ km s}^{-1}$  and  $-12.9 \text{ km s}^{-1}$ ) and then sample lower radial velocity components where S2 becomes a faint ring (i.e. between  $-11.6 \text{ km s}^{-1}$  and  $-9.0 \text{ km s}^{-1}$ ). At lower velocities again, these rings disappear into the noise of the maps and possibly extends out beyond the primary beam at zero velocity where S2 should have maximum spatial extent. The emission from the channel maps between  $-15.3 \text{ km s}^{-1}$  and  $-11.6 \text{ km s}^{-1}$  correspond to all the emission in the extreme blue wing component of the multi-configuration image cube line profile discussed in Section 5.3. We can see in Figure ?? that all of this emission is greater than the C configuration resolving out scale of  $\sim 6''$ , therefore confirming that our C configuration line profile is mainly composed of S1 emission. The shell formation signature of S2 is also apparent in the redshifted velocity channel maps between  $+7.5 \text{ km s}^{-1}$  and  $+13.8 \text{ km s}^{-1}$  but the emission appears weaker and the rings fainter therefore indicating that S2 is somewhat fragmented. The

multi-configuration maps also show the central compact emission from the S1 flow at velocities between  $-10.3 \text{ km s}^{-1}$  and  $+11.3 \text{ km s}^{-1}$ . This S1 emission can be seen in the final two maps of Figure ?? as a central slightly elongated emission feature surrounded by the fainter rings of the S2 flow. The spatial extent of the S1 flow varies from channel map to channel map but appears to be larger than the  $2''$  value given by [Smith et al. \(2009\)](#), who observed off-star wind scattered ro-vibration CO lines.

## 5.6 Spatial Extent of S2

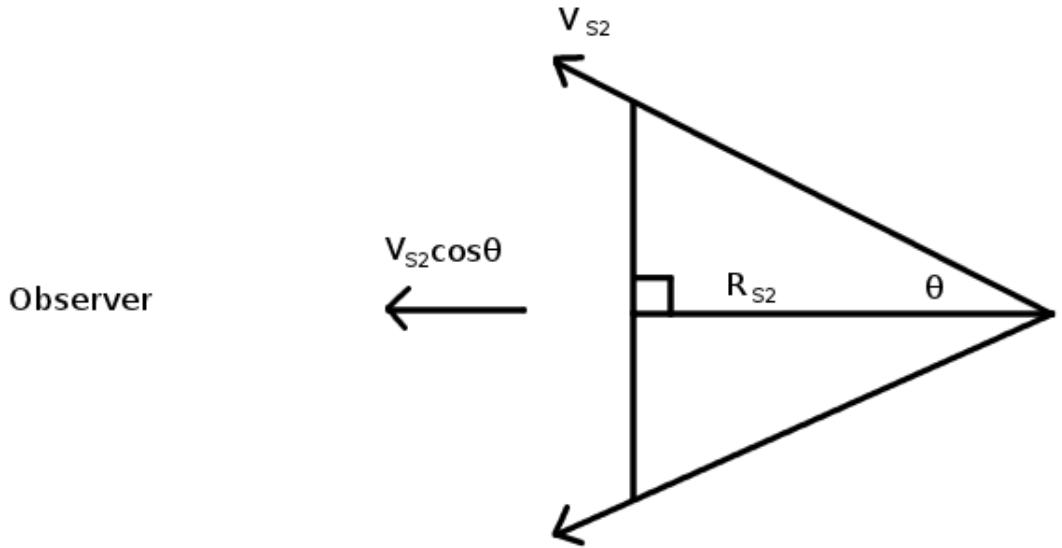
The spatial extent of the S2 flow around Betelgeuse was not directly determined from either the CO infrared absorption spectra of [Bernat et al. \(1979\)](#) or previous CO single dish radio observations ([Huggins, 1987](#); [Huggins et al., 1994](#); [Knapp et al., 1980](#)). Our low spectral resolution multi-configuration image cube has sufficient spatial resolution and S/N to make direct estimates of its maximum radius. We find that little or no signature of the S2 flow is present in the low absolute velocity channel maps where its spatial extent is maximum and either lies outside of the primary beam or is lost into the noise near the edge of the maps. Therefore we cannot make a simple direct measurement of its maximum size. However, by measuring how its size varies in the higher absolute velocity maps it is clearly present allows us to derive its maximum spatial extent.

If we assume that S2 is spherically symmetric with an outer radius  $R_{S2}$ , and is undergoing steady expansion with velocity  $V_{S2}$ , then its geometry is described by Figure 5.7. As we sample the line at different velocities, the spatial extent of S2 changes and gets progressively larger at lower absolute velocities as can be seen in Figure ???. We define the spatial extent of S2 in each map as

$$r_{\text{chan}} = R_{S2} \times \sin\theta \quad (5.3)$$

which can be directly measured in many of the channel maps. Each of these channel maps have an associated velocity  $V_{\text{chan}}$ , which we know, and is related to the actual velocity of the S2 flow by

$$V_{\text{chan}} = V_{S2} \times \cos\theta. \quad (5.4)$$

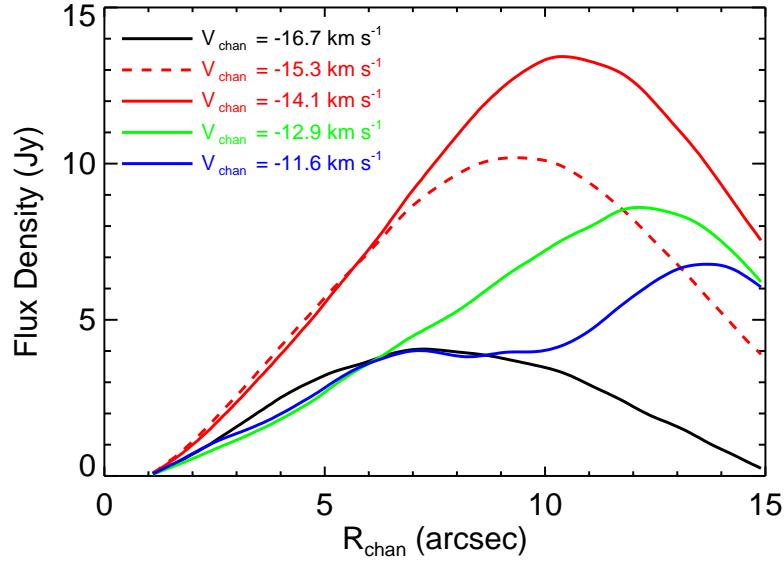


**Figure 5.7:** Geometry of a spherical symmetric flow used to derive the spatial extent of the S2 flow. By measuring the spatial extent of S2 (i.e.,  $R_{S2}\sin\theta$ ) as a function of channel/velocity (i.e.,  $V_{S2}\cos\theta$ ) the maximum spatial extent of the S2 flow can be derived.

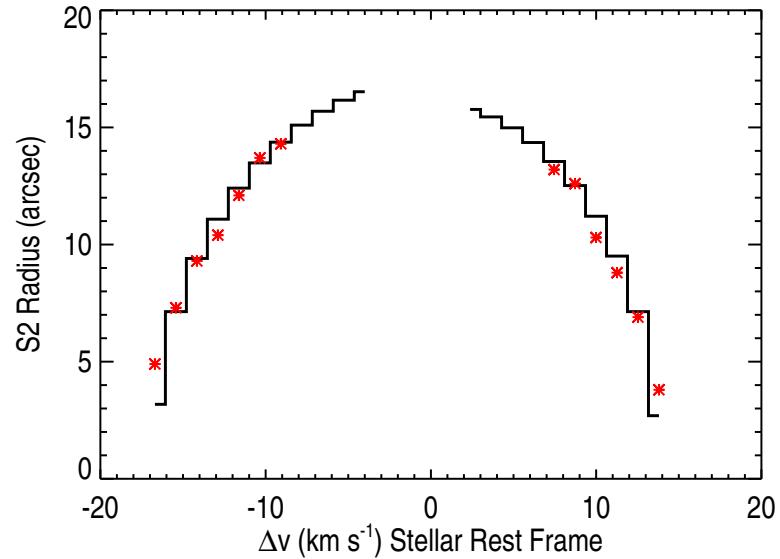
Combining Equations 5.3 and 5.4 via the angle dependence, gives the following equation which contains  $R_{S2}$  as the only unknown parameter,

$$r_{\text{chan}} = R_{S2} \times \sin \left[ \cos^{-1} \left( \frac{V_{\text{chan}}}{V_{S2}} \right) \right]. \quad (5.5)$$

An estimate of the S2 radius per channel ( $r_{\text{chan}}$ ) was found by creating annuli of increasing radius around the central emission in each relevant line channel map of the multi-configuration image cube, extracting all flux within each annulus and then plotting these fluxes against distance from the star for each channel map as demonstrated in Figure 5.8. The maximum of these resultant curves was then deemed to be the maximum radius of S2 per channel map. Figure 5.9 shows these data over-plotted with two model outflows which were created using Equation 5.5. The blueshifted data points were best fitted by a model outflow of maximum radius 17" and outflow velocity  $16.7 \text{ km s}^{-1}$ , while the redshifted data points were best fitted by a model outflow of maximum radius 16" and outflow velocity  $13.8 \text{ km s}^{-1}$ . It is worth mentioning that this estimate for the spatial



**Figure 5.8:** Deriving  $R_{\text{chan}}$  per channel map. The maximum of the curves representing the integrated flux as a function of distance from the star was deemed as the size of S2 in each channel map.



**Figure 5.9:** The derived S2 radius as a function of velocity (red points) overplotted with two model outflows. The blueshifted model (left) corresponds to an outflow with a maximum radius of 17'' and a velocity of  $16.7 \text{ km s}^{-1}$  while the redshifted model (right) corresponds to an outflow with a maximum radius of 16'' and a velocity of  $13.8 \text{ km s}^{-1}$ . Note: The line profile is  $1.9 \text{ km s}^{-1}$  wider in the low resolution image cube ( $\Delta v = 1.3 \text{ km s}^{-1}$ ) than in the high resolution image cube.

## 5.7 Intensity distribution of CO

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extent of S2 is only weakly dependent on our adopted radial velocity value for Betelgeuse and adopting a slightly different value would simply alter the outflow velocities.

## 5.7 Intensity distribution of CO

We first investigate how one would expect the intensity of CO to vary as a function of distance from the star for a spherically symmetric optically thin outflow. The CO( $J = 2 - 1$ ) emission line is formed over an extended region around Betelgeuse so a constant outflow velocity can be assumed when investigating the intensity distribution of the S1 and S2 outflow. This allows us to use the equation of mass continuity to find the number density of the molecules in the lower level of the transition,  $n_l$ , i.e.,  $n_l \propto r^{-2}$ . The Boltzmann formula can then be used to find the number density of CO molecules in the upper level,  $n_u$ ,

$$n_u(r) = n_l(r) \frac{g_u}{g_l} e^{-(E_u - E_l)/kT_{\text{exc}}(r)} \quad (5.6)$$

where  $g_u$  and  $g_l$  are the statistical weights of the two rotational levels, and  $E_u$  and  $E_l$  are their respective energies. Here,  $T_{\text{exc}}$  is the excitation temperature and equals to the gas temperature if collisions dominate the redistribution of the populations over the rotational ladder. To find the intensity we first need to know the emissivity,  $j_l$ , of the line photons which is just  $n_u(r)$  times the probability for spontaneous de-excitation per second,  $A_{ul}$ , times the energy of the emitted photons,  $h\nu$ , i.e.,

$$\begin{aligned} j_l(r) &= n_u(r) A_{ul} h\nu \\ &\propto r^{-2} e^{-(E_u - E_l)/kT_{\text{exc}}(r)}. \end{aligned} \quad (5.7)$$

If  $T_{\text{exc}}$  is constant then the line intensity,  $I_{\text{CO}}$ , can be written as

$$\begin{aligned} I_{\text{CO}} &= \int j_l(r) dr \\ &\propto r^{-1}. \end{aligned} \quad (5.8)$$

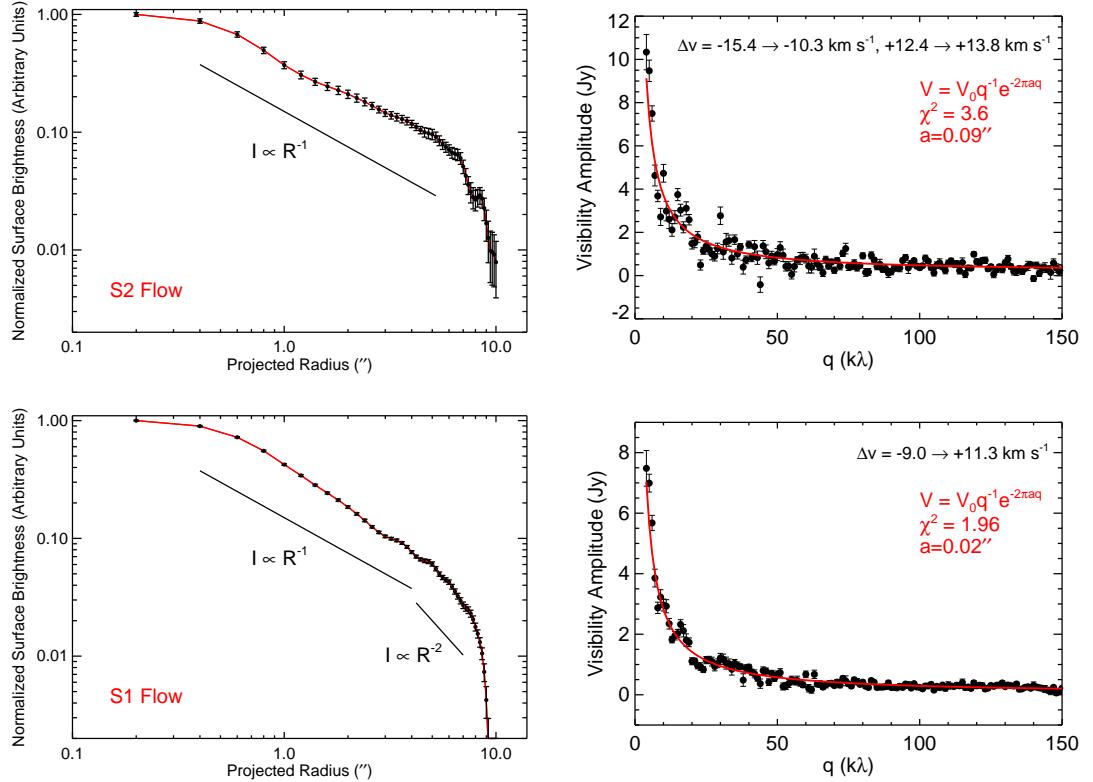
In the left column of Figure 5.10 we investigate the intensity distribution of

## 5.7 Intensity distribution of CO

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CO emission as a function of projected radius,  $R$ , for both the S1 and S2 flows. From our discussions in Section 5.3 we can assume that all line emission between  $-15.4 \rightarrow -10.3 \text{ km s}^{-1}$  and  $+12.4 \rightarrow +13.8 \text{ km s}^{-1}$  emanates solely from the S2 flow. Using the low spectral resolution multi-configuration image cube we integrate the surface brightness over these channels and find that the intensity fall-off is close to being proportional to  $R^{-1}$  (Figure 5.10: *top left*). To investigate the S1 flow intensity distribution around  $\alpha$  Ori we integrate the surface brightness over the channels between  $-9 \rightarrow +10.6 \text{ km s}^{-1}$ . Although these channels contain emission from both S1 and S2, most of the S2 emission here will have larger projected radii and thus the majority of the inner emission should emanate from the S1 flow. Between  $0.5''$  and  $4''$  from the star the intensity is again found to be close to proportional to  $R^{-1}$  (Figure 5.10: *bottom left*). We have just showed that such an intensity distribution is expected for an optically thin homogeneous constant velocity outflow with  $\rho \propto 1/R^2$ . Beyond  $\sim 6''$  the intensity fall off is more rapid and is steeper than a  $R^{-2}$  distribution and may be caused by factors such as a variation in the thermal or velocity profile, or even the absence of emitting material altogether. This region may mark the initiation of the current epoch of mass loss.

Insight can also be gained into how the intensity varies on different size scales by conducting analysis in the  $u - v$  plane and plotting the visibility amplitude of our CO data against  $u - v$  distance. The result of this is shown in the right column of Figure 5.10 where the same channels corresponding to the S1 and S2 flows defined in the last paragraph have been used. The data are azimuthally averaged, and have been binned to produce one data point per  $k\lambda$ . The result for both the S1 and S2 data is a steep drop-off in visibility amplitude over a relatively short  $u - v$  distance, signaling that the sources are well resolved, as expected. Both sets of visibility data agree with an intensity proportional to  $(a^2 + R^2)^{-1/2}$ , where  $a$  is an inner spatial limit. This is because the Hankel transform of this function is  $q^{-1}e^{-2\pi aq}$  (Bracewell, 2000), where  $q$  is the  $u - v$  distance, and a vertically scaled version of this function is shown to match the visibility data very well in Figure 5.10. As analysis in both the sky and  $u - v$  plane indicate the intensity of both flows is proportional to  $R^{-1}$  we conclude that when azimuthally averaged, both outflows are consistent with an optically thin and quasi-steady



**Figure 5.10:** *Left Column:* Surface Brightness as a function of projected radius on sky,  $R$  (red line). The emission has been extracted from the low spectral resolution multi-configuration image cube and is integrated over the channels where S1 is present (bottom) and over the channels where only S2 is present (top). Intensity proportional to  $R^{-1}$  and  $R^{-2}$  is also shown for comparison. *Right Column:* The corresponding visibility amplitude as a function of  $u-v$  distance ( $q$ ) of both outflows can be modeled well by a  $R^{-1}$  fall off in intensity. The error bars in all plots represent the standard error of the mean.

flow which is in agreement with [Smith \*et al.\* \(2009\)](#) (i.e., S1) and [Plez & Lambert \(2002\)](#) (i.e., S2).

## 5.8 Spatial Extent of S1

An exact determination of the maximum spatial extent of the S1 flow is more difficult as we do not see the classical shell formation signature for it as we sample across velocities, like we do for S2. Instead its spatial extent varies over

the channel maps with evidence of discrete clumps being present in many of these maps. At 20% of maximum emission in the integrated intensity S1 map (i.e. composed of all channels between  $-9 \rightarrow +10.6 \text{ km s}^{-1}$ ) the S1 flow extends out to a mean distance of  $\sim 4''$  and is even more extended in the S-W direction due to the presence of the second emission feature in the compact configuration data sets. The HPBW of  $0.9''$  is not sufficient to determine whether the S1 flow is discrete or an extension of the current wind phase seen in ultraviolet spectra, e.g., ?, and cm-radio continuum interferometry (Harper *et al.*, 2001; ?).

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## 5.9 Continuum Flux Densities

Betelgeuse is known to show brightness variations at many continuum wavelengths. Goldberg (1984) reports a decrease of half a magnitude in visual brightness over a period of six years. Bookbinder *et al.* (1987) found stochastic 30%-40% variations in flux density at 6 cm over timescales as short as 10 days to as long as 8 months (i.e. the observational period). A more comprehensive study was carried out by Drake *et al.* (1992) who observed Betelgeuse with the VLA at centimeter wavelengths from 1986 to 1990 and found stochastic variability of 22%, 15%, and 21% at 6 cm, 3.6 cm, and 2 cm respectively, at a variety of different timescales down to less than one month (a possible source for this variability will be described in Section 5.11). The mm-continuum emission that we measure arises mainly from electron-ion and electron-atom bremsstrahlung (also the source of the longer wavelength continuum emission) and possibly dust emission, so it is not unreasonable to also expect variability at mm-wavelengths too.

In Table 5.1 we list the derived continuum flux densities for each of the three configuration image cubes and also the multi-configuration image cube. The high spectral resolution ( $\Delta v = 0.65 \text{ km s}^{-1}$ ) image cubes were just wide enough to image the CO line but were too narrow to make accurate estimates of the continuum flux density. Therefore, all continuum flux density estimates are derived from the lower spectral resolution ( $\Delta v = 1.3 \text{ km s}^{-1}$ ) image cubes from which we were able to take accurate measurements at both sides of the line. We fitted elliptical Gaussians to  $\sim 20$  continuum channels using CASA's *imfit* routine allowing the flux

## 5.10 Higher CO rotational lines

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**Table 5.1:** CARMA Continuum Fluxes at 230 GHz between 2007 → 2009.

Configuration	Date ("×")	Restoring Beam	Flux (mJy)	Uncertainty (mJy)
C	2007 Jun	$0.96 \times 0.76$	234	18
D	2009 Jul	$2.33 \times 1.87$	389	72
E	2009 Nov	$4.93 \times 3.84$	278	40
Multi-configuration	2007 → 2009	$1.05 \times 0.84$	289	21

and corresponding uncertainties to be calculated. The source was unresolved in most of these continuum channels. The D configuration data were acquired under adverse weather conditions and these data have the highest noise levels out of the three configurations. Its continuum emission measurement is approximately 50% greater than the C and E configuration continuum measurements which were also acquired approximately two years after the D configuration data. We believe the continuum emission value of  $289 \pm 21$  mJy derived from the multi-configuration image cube is a reasonable estimation of the mean mm-continuum flux density over the two year period, 2007 → 2009. It is in reasonably good agreement with the 250 GHz flux density of [Altenhoff \*et al.\* \(1994\)](#) who report a mean value of  $351 \pm 25$  mJy for 1986 → 1989.

## 5.10 Higher CO rotational lines

although a search at 86 GHz for SiO( $J = 2 - 1$ ) by [Lambert & Vanden Bout \(1978\)](#) had been unsuccessful

## 5.11 CARMA CO observations in Context

## 5.12 e-Merlin 6 cm Results

position, GMH model

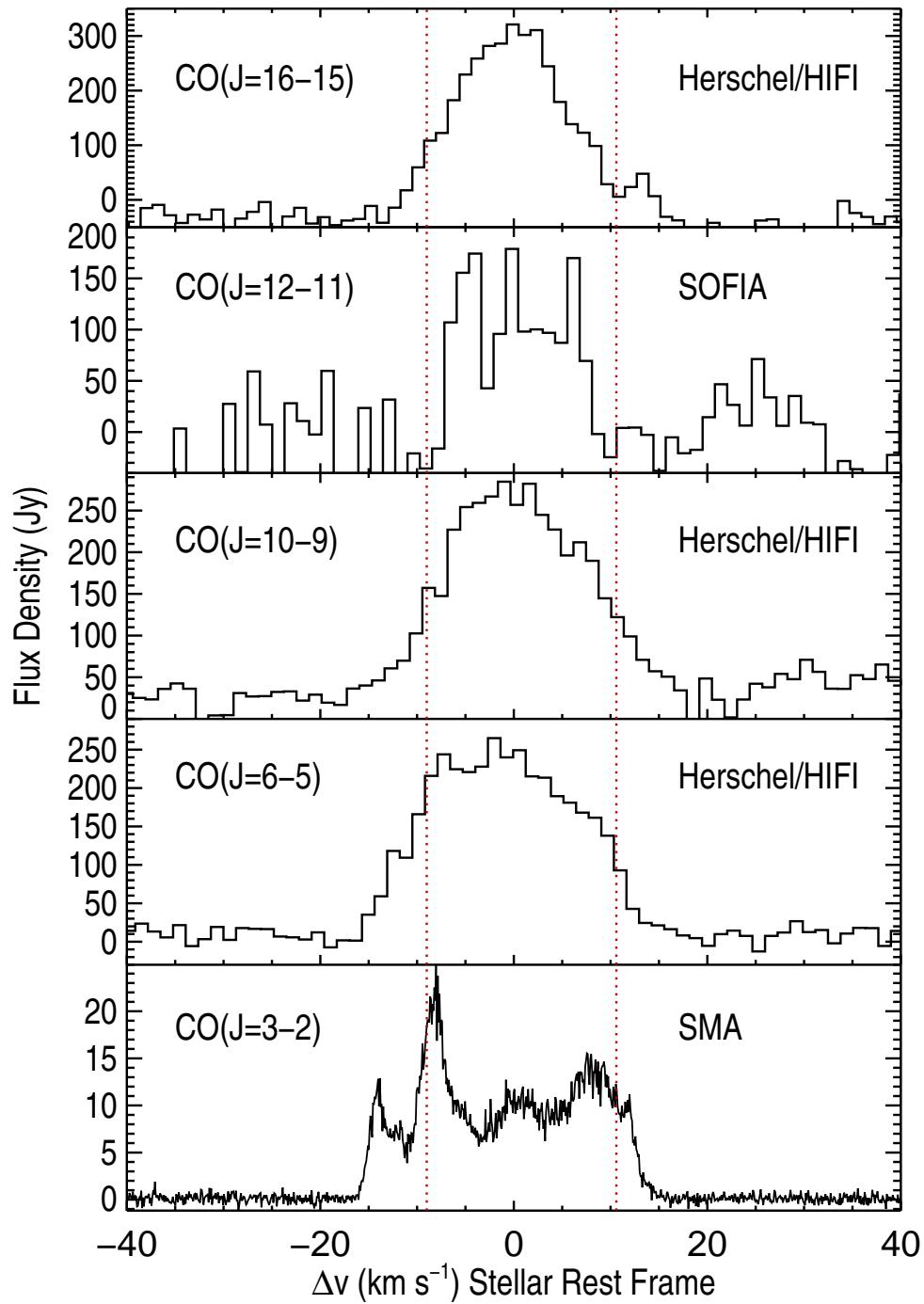
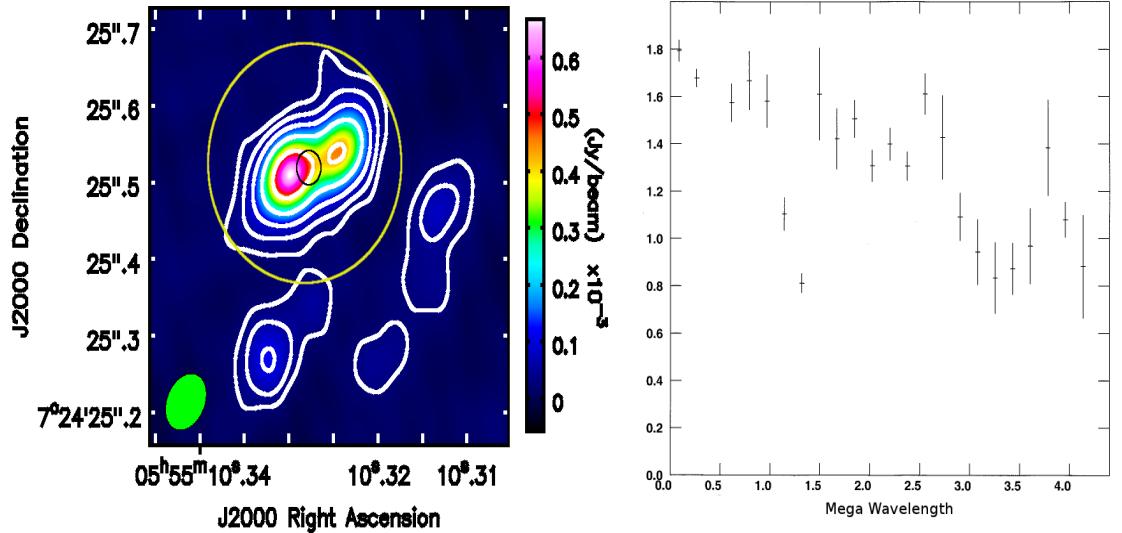


Figure 5.11

## 5.13 Multi-wavelength VLA-Pie Town Results



**Figure 5.12:** First e-MERLIN results for Betelgeuse. *Left:* 5.2 cm e-Merlin image of Betelgeuse taken in July 2012 showing two unresolved emission features separated by  $90 \pm 10$  mas. The image was restored with a  $80 \times 60$  mas $^2$  beam marked in green in the lower left corner and using uniform weighting. Contour levels are at  $(-3, 3, 6, 12, 24, 48) \times \sigma_{\text{rms}}$  where  $\sigma_{\text{rms}} = 9 \mu\text{Jy}$ . The photospheric radius is the small black circle which the large yellow circle defines the 6 cm radio disk of Lim *et al.* (1998). *Right:* The calibrated visibility amplitudes (in mJy) versus  $u - v$  distance measured in mega-wavelengths with the possible signature of two unresolved components with a brightness ratio differing from unity.

## 5.13 Multi-wavelength VLA-Pie Town Results

Confirmation of a low temperature wind

Time Variability

how does Anita's fit into this?

# 6

## Multi-wavelength Radio Continuum Emission Studies of Dust-free Red Giants

The vast increase in bandwidth coverage of the new VLA now allows possible detections of historically weak or undetectable red giants at multiple radio wavelengths. In this chapter we present and analyze the most comprehensive set of multi-wavelength thermal radio continuum observations of two standard luminosity class III red giants, Arcturus ( $\alpha$  Boo: K2 III) and Aldebaran ( $\alpha$  Tau: K5 III), to date. We report the first detections at several wavelengths for each star including a detection at 10 cm (3.0 GHz: S-band) for both stars and a 20 cm (1.5 GHz: L-band) detection for  $\alpha$  Boo. This is the first time single luminosity class III red giants have been detected at these continuum wavelengths. We compare our data with published semi-empirical atmospheric models based on UV data, and spectral indices are used to discuss the possible properties of the stellar atmospheres. Finally, we develop a simple analytical wind model for  $\alpha$  Boo based on our new long wavelength flux measurements. This chapter is based on the main findings of O’Gorman *et al.* (2013).

### 6.1 Why Study Red Giants with the VLA?

Although studies of wind-scattered UV and optical line profiles have provided clues to the mass loss rates and radial distribution of the mean and turbulent velocity fields, the thermal structure remains poorly constrained. In the UV, the source function  $S_\nu$ , of electron collisionally excited emission lines is sensitive to electron temperature  $T_e$ , (i.e.,  $S_\nu \propto e^{-h\nu/kT_e}/\sqrt{T_e}$ ). Therefore, a localized hot plasma component in a dynamic atmosphere can completely dominate the temporally and spatially averaged emission and hence not reflect the mean radial electron temperature distribution. At radio wavelengths however, the source function is thermal and is just the Rayleigh-Jeans tail of the Planck function, which is linear in electron temperature (i.e.,  $S_\nu = 2kT_e\nu^2/c^2$ ). This should give a more appropriate estimate of the mean radial electron temperature. It is this value that controls the atomic level populations and ionization of the mean plasma which is needed to quantify the implied thermal heating supplied to the wind by the unknown driving source/sources, allowing constraints on potential mass loss mechanisms to be derived.

In the cm-radio regime the radio opacity  $\kappa_\lambda$ , strongly increases with wavelength (i.e.,  $\kappa_\lambda \propto \lambda^{2.1}$ ) and so the longer wavelengths sample the extended layers of a stars atmosphere, thus providing us with spatial information about the star's mass outflow region. The VLA is sensitive to over three orders of magnitude in continuum optical depth  $\tau_\lambda$ , ( $\tau_{20 \text{ cm}}/\tau_{0.7 \text{ cm}} \approx 10^3$ ) and provides an area-averaged sweep through the wind acceleration zone of evolved late-type stars. The thermodynamic properties in this spatial region control the ionization in the far wind because the ionization balance, which also controls the cooling rates, becomes *frozen-in* at large radii due to advection. Furthermore, it is these outer extended regions of the star's atmosphere that contribute to the commonly seen P Cygni line profiles in the UV. In these profiles the line-of-sight absorption caused by the star's wind is superimposed on the blueshifted scattered emission. Thus, centimeter radio continuum observations can provide a test of models based on these UV profiles.

Despite the importance of radio continuum emission as a tool for studying cool evolved stellar outflows, red giants are feeble emitters at these wavelengths

however, and previous observations have provided only a small number of modest signal-to-noise measurements slowly accumulated over three decades. As most of our VLA observations were carried out over just a few days, we have obtained a unique snapshot of the different stellar atmospheric layers sampled at different wavelengths, independent of any long-term variability.

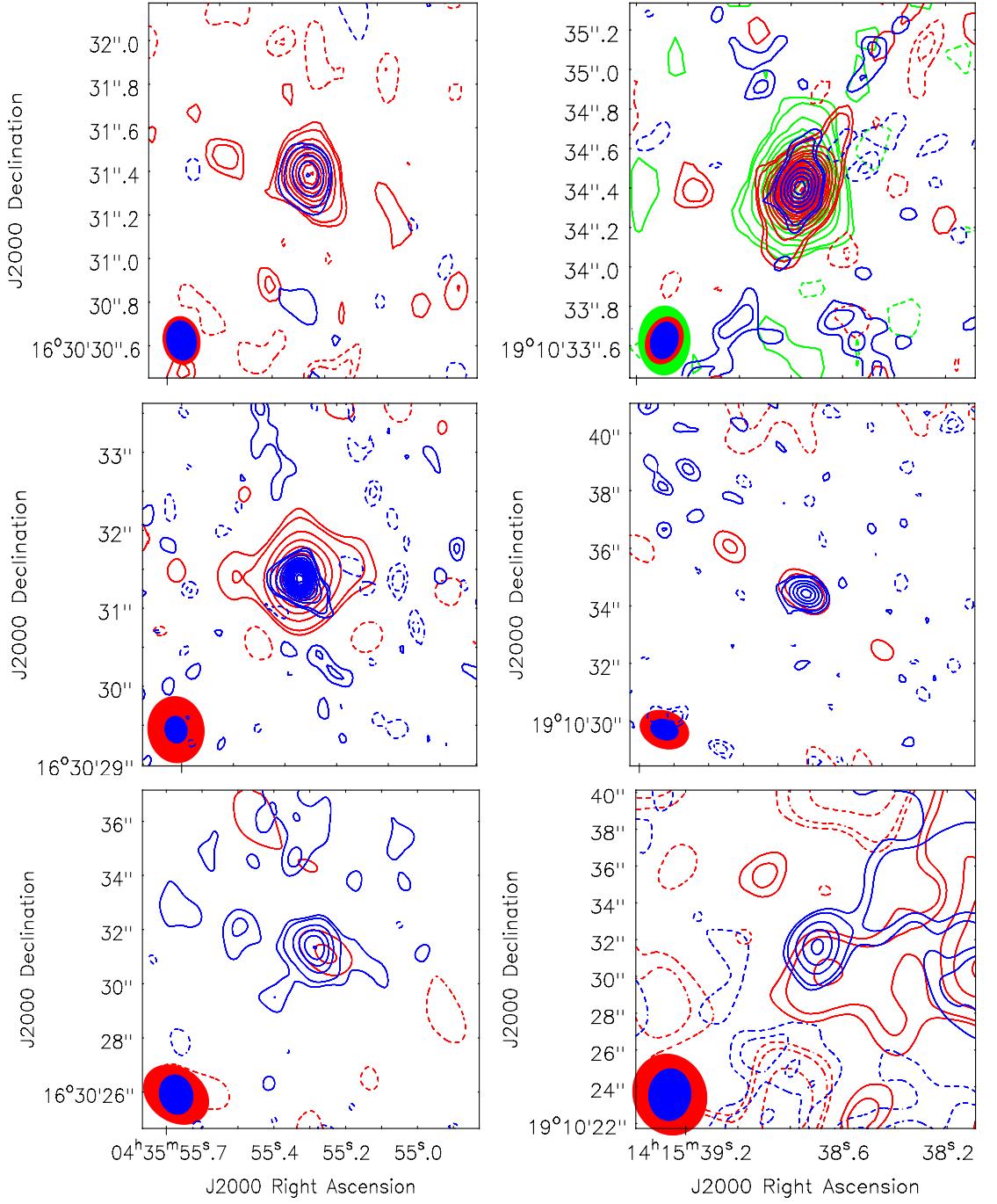
## 6.2 Radio Maps

Apart from  $\alpha$  Boo at C-band and  $\alpha$  Tau at S-band, detections were made in every sub-band for the 2011 data. For all other bands, the flux densities of the targets in both sub-bands were found to be the same within their uncertainties so we do not present separate values here. Instead we give the values from the radio maps produced by concatenating the two sub-bands. We present in Table 6.1 the target flux densities extracted from these concatenated radio maps. In each of these maps the flux density from the unresolved target source was calculated by, 1) taking the peak pixel value from the source, 2) manually integrating the flux density around the source, and 3) fitting an elliptical Gaussian model to the source and deriving the integrated flux density using the CASA *imfit* task. Each of these values along with the image rms noise measured from adjacent background regions and fitting error produced by *imfit* are given in Table 6.1 to indicate the quality of each radio map. Both sources are point sources at all frequencies so the peak flux value given in Table 6.1 will also be its total flux density value. For weak detections (i.e.,  $F_\nu \lesssim 5\sigma$ ) we avoid using the *imfit* task to obtain a flux density estimate, as this may produce biased parameter estimates (Taylor *et al.*, 1999). The flux density values used in the following Sections are the peak values listed in Table 6.1. We assume absolute flux density scale systematic uncertainties of 3% at all frequencies (Perley & Butler, 2013). In the following two sections we briefly discuss the properties of these maps for both targets which are shown in Figure 6.1.

### 6.2.1 $\alpha$ Boo Maps

High S/N detections ( $>19\sigma$ ) of  $\alpha$  Boo were made at 22.5, 33.6, and 43.3 GHz. Some residuals of the dirty beam remained in the CLEANed maps due to the

## 6.2 Radio Maps



**Figure 6.1:** Final VLA multi-wavelength radio maps of α Tau and α Boo. *Left column:* Q and Ka-band (top), K and X-band (middle), C and S-band (bottom) maps of α Tau. *Right column:* Q, Ka, and K-band (top), X and C-band (middle), S and L-band (bottom) maps of α Boo. Contour levels are set at  $(-6, -3, -2, +2, +3, +6, +9, \dots) \times \sigma$  where  $\sigma$  is the rms noise of each radio map. For α Boo, artifacts from the dirty beam can be seen in the Q-band map, while artifacts caused by the serendipitous strong source in the S and L-band maps are also present.

**Table 6.1:** VLA Flux Densities of  $\alpha$  Boo and  $\alpha$  Tau.

Star	Band	$\nu^a$ (GHz)	$\lambda$ (cm)	Peak $F_\nu$ (mJy beam $^{-1}$ )	Integrated $F_\nu$ (mJy)	<i>Imfit</i> Integrated $F_\nu$ (mJy)	Image rms (mJy beam $^{-1}$ )	<i>Imfit</i> Fitting Error (mJy)
$\alpha$ Boo	Q	43.28	0.7	5.94	6.09	6.42	0.30	0.26
	Ka	33.56	0.9	4.16	4.32	4.49	0.08	0.09
	K	22.46	1.3	1.83	1.78	1.81	0.04	0.05
	X	8.46	3.5	0.51	0.51	0.53	0.03	0.02
	C	4.90	6.1	0.21	0.14	0.16	0.04	0.01
	S	3.15	9.5	0.15	0.14	-	0.03	-
	S	2.87	10.4	0.13	0.12	0.12	0.01	0.02
	L	1.63	18.4	0.07	0.07	-	0.01	-
$\alpha$ Tau	Q	43.28	0.7	3.67	3.73	4.08	0.26	0.18
	Ka	33.56	0.9	2.19	1.96	2.13	0.09	0.07
	K	22.46	1.3	1.86	1.88	2.07	0.04	0.08
	X	8.46	3.5	0.30	0.29	0.28	0.01	0.02
	C	4.96	6.0	0.15	0.17	0.18	0.01	0.01
	S	3.15	9.5	0.06	0.04	-	0.02	-

<sup>a</sup>Frequency of the final image produced using the multi-frequency synthesis imaging mode within CASA's *clean* task.

paucity of uv-coverage in these short high frequency observations (see Figure 6.1, top right panel). At the lower frequencies, it was necessary to image confusing sources, notably a strong radio source located 186'' north-west of  $\alpha$  Boo. This non-thermal source was reported by [Drake & Linsky \(1983a\)](#) and their flux density of 25 mJy at 4.9 GHz is in close agreement with our measurement of 23.2 mJy at the same frequency. We find the source to have a spectral index  $\alpha$  ( $F_\nu \propto \nu^\alpha$ ) of -1.4 between 8.5 and 1.6 GHz; its flux density reaches 80.3 mJy at 1.6 GHz.

We detected  $\alpha$  Boo at  $6\sigma$  in the lower frequency sub-band of C-band, at 4.9 GHz. The noise was slightly higher and the images were poorer quality in the C-band higher frequency sub-band, with artifacts exceeding  $\pm 200 \mu\text{Jy}$ , and we cannot report a detection in this sub-band, so values given in Table 6.1 are taken from the lower frequency sub-band only. We obtain good detections ( $>5\sigma$ ) of the star for both epochs at  $\sim 3$  GHz (S-band) and the peak flux densities agree within their uncertainties. We can therefore safely assume that the 1.5 GHz (L-band) flux density has not changed significantly over that period either, and so can safely be included in any analysis. The map at L-band was highly contaminated by the sidelobes of the strong source north-west of  $\alpha$  Boo but the star is still detected at the  $5\sigma$  level. There is a slight positional offset of 1'' between the position of the peak flux density at 1.5 and at 3.0 GHz for the 2012 data, which were taken within 1 day of each other. However, the position uncertainties due to noise and phase uncertainties between the directions of the phase reference source and the target are at least 1'', and so we feel that it is highly likely that both detections are of  $\alpha$  Boo.

### 6.2.2 $\alpha$ Tau Maps

The final deconvolved radio maps of  $\alpha$  Tau were of excellent quality with the rms noise reaching the predicted noise levels in many cases. The target field at all frequencies was free from strong serendipitous radio sources and thus the final images were free of the sidelobe contamination that were present in the low frequency  $\alpha$  Boo images.  $\alpha$  Tau was the only source in the high frequency maps while the brightest source in the low frequency maps was located 106'' north north-east of  $\alpha$  Tau and had flux densities of 0.85, 1.35, and 1.7 mJy at 8.5, 5.0,

### 6.3 Results vs Previous Observations

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and 3.5 GHz, respectively. Strong detections ( $>14\sigma$ ) of  $\alpha$  Tau were made at all frequencies between 5.0 and 43.3 GHz. Due to the limited number of S-band receivers available at the time, a full 2.5 hr track was dedicated to  $\alpha$  Tau at 3.1 GHz in order to achieve the required sensitivity to give a possible detection. We report a tentative  $3\sigma$  detection of  $\alpha$  Tau at 3.1 GHz when we take its peak pixel value as its total flux density.

## 6.3 Results vs Previous Observations

Prior to and during the early operation of the ‘old’ VLA, a small number of single dish radio observations reported the detection of flares from single red giants (e.g., [Slee \*et al.\* 1989](#)). These transient radio events have never been re-observed however, even with more sensitive interferometers, suggesting that such detections were spurious (e.g., [Beasley \*et al.\* 1992](#)). The first definitive detection of thermal free-free emission from a luminosity class III single red giant at centimeter wavelengths was of  $\alpha$  Boo at 6 cm ([Drake & Linsky, 1986, 1983a](#)). Since then there has been a modest number of centimeter and millimeter observations of this star. In Table A.1 we list the majority of these observations and plot their flux densities as a function of frequency in Figure 6.2 (red diamonds). In comparison to other single red giants,  $\alpha$  Boo had been relatively well observed at radio continuum wavelengths before this study, including detections in four VLA bands (i.e., Q, K, Ku, and C). No Ku band receivers were available during the commissioning phase of the VLA in early 2011 so we can compare three of our detections with previous ones.

Previous detections of  $\alpha$  Boo at 6 cm ranged from a  $3\sigma$  upper limit of 0.18 mJy to a  $3\sigma$  detection at 0.39 mJy. Our 6 cm value agrees to within  $\sim 10\%$  of the highest S/N ( $5\sigma$ ) value of [Drake & Linsky \(1986\)](#). There is no significant difference between our 1.3 cm value and that of [Dehaes \*et al.\* \(2011\)](#). There is however a notable difference in flux density values at 0.7 cm where [Dehaes \*et al.\* \(2011\)](#) report values that are lower than ours by over 40%. Although we do not rule out such a level of chromospheric radio variability, it is not expected based on the small level of UV variability observed from such supposedly inactive stars ([Harper \*et al.\*, 2013](#)). Another possibility for the difference in values is that the

### 6.3 Results vs Previous Observations

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**Table 6.2:** Compilation of Previous Radio Observations ( $\nu \leq 250$  GHz).

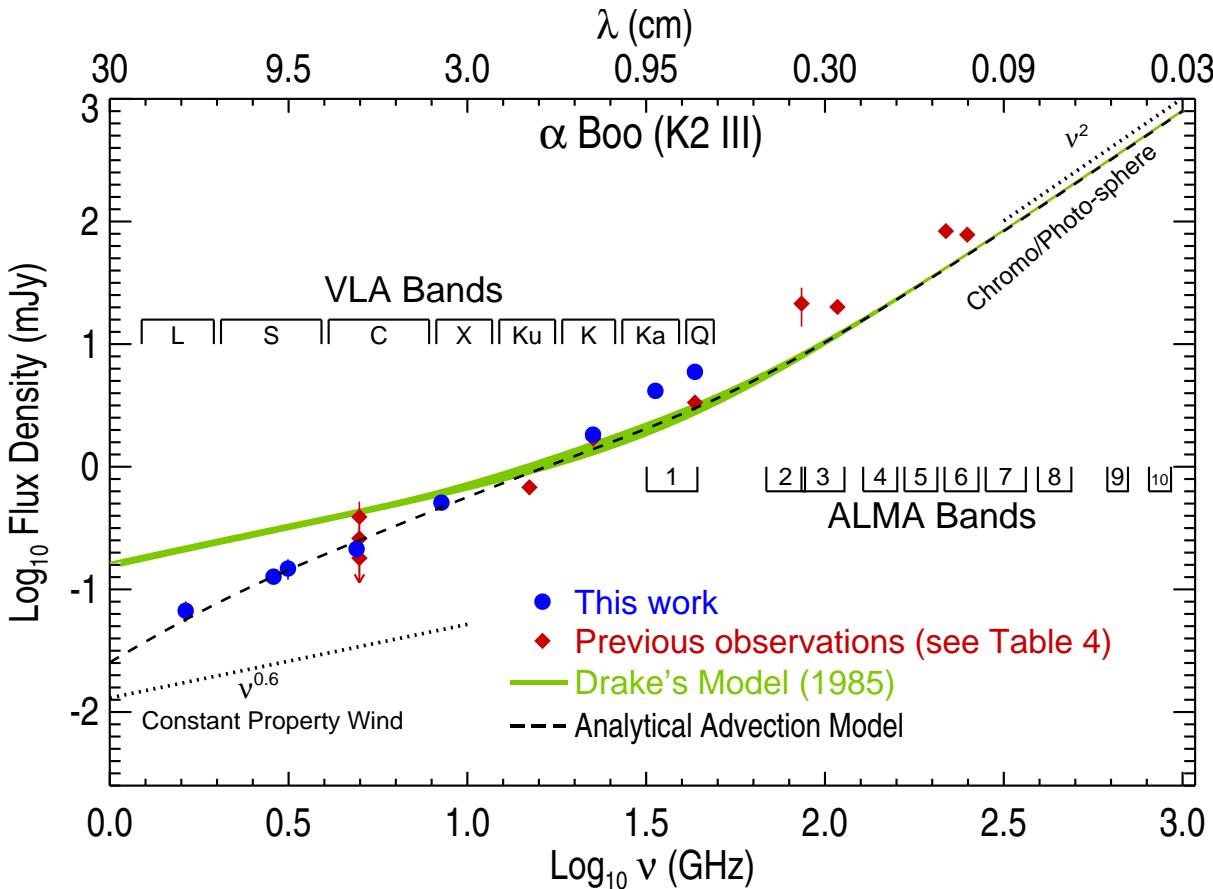
Source	$\nu$ (GHz)	Date	$F_\nu$ (mJy)	S/N	Reference
$\alpha$ Boo	4.9	1983 Jan 21	0.39	3.0	1
	4.9	1983 May 20	0.26	3.3	1
	4.9	1983 Dec 26	$\leq 0.18(3\sigma)$	-	1
	4.9	1984 Mar 17	0.24	4.8	1
	15.0	1984 Nov 6	0.68	7.6	1
	22.5	1999 Jan 06	1.7	8.5	2
	43.3	1999 Jan 06	3.3	8.3	2
	43.3	2004 Jan 25	3.34	41.8	2
	86.0	1985 Nov	21.4	3.0	3
	108.4	1997 Nov - 2000 Jun	20.1	29.1	4
$\alpha$ Tau	217.8	1997 Nov - 2000 Jun	83.5	48.8	4
	250.0	1986 Dec - 1989 Mar	78.0	9.8	5
	4.9	1983 Jan 21	$\leq 0.27(3\sigma)$	-	1
$\alpha$ Tau	4.9	1984 Nov 6	$\leq 0.22(3\sigma)$	-	1
	5.0	1997 Sep 27	$\leq 0.07(3\sigma)$	-	6
	8.5	1997 Sep 27	0.28	9.3	6
	14.9	1997 Sep 27	0.95	11.9	6
	15.0	1984 Nov 6	0.60	6.0	1
	108.4	1997 Nov - 2000 Dec	14.0	9.6	4
	217.8	1999 Sep - 2000 Dec	25.8	4.6	4
	250.0	1986 Dec - 1987 Jan	51.0	8.5	5

References.-(1) [Drake & Linsky \(1986\)](#); (2) [Dehaes \*et al.\* \(2011\)](#); (3) [Altenhoff \*et al.\* \(1986\)](#); (4) [Cohen \*et al.\* \(2005\)](#); (5) [Altenhoff \*et al.\* \(1994\)](#); (6) [Wood \*et al.\* \(2007\)](#);

longer cycle time used by [Dehaes \*et al.\* \(2011\)](#), which was over double our value, may lead to larger phase errors and thus lower final flux density values. Future high frequency VLA observations of  $\alpha$  Boo will clarify this discrepancy at 0.7 cm but past detections at longer wavelengths appear to be in good agreement with our data.

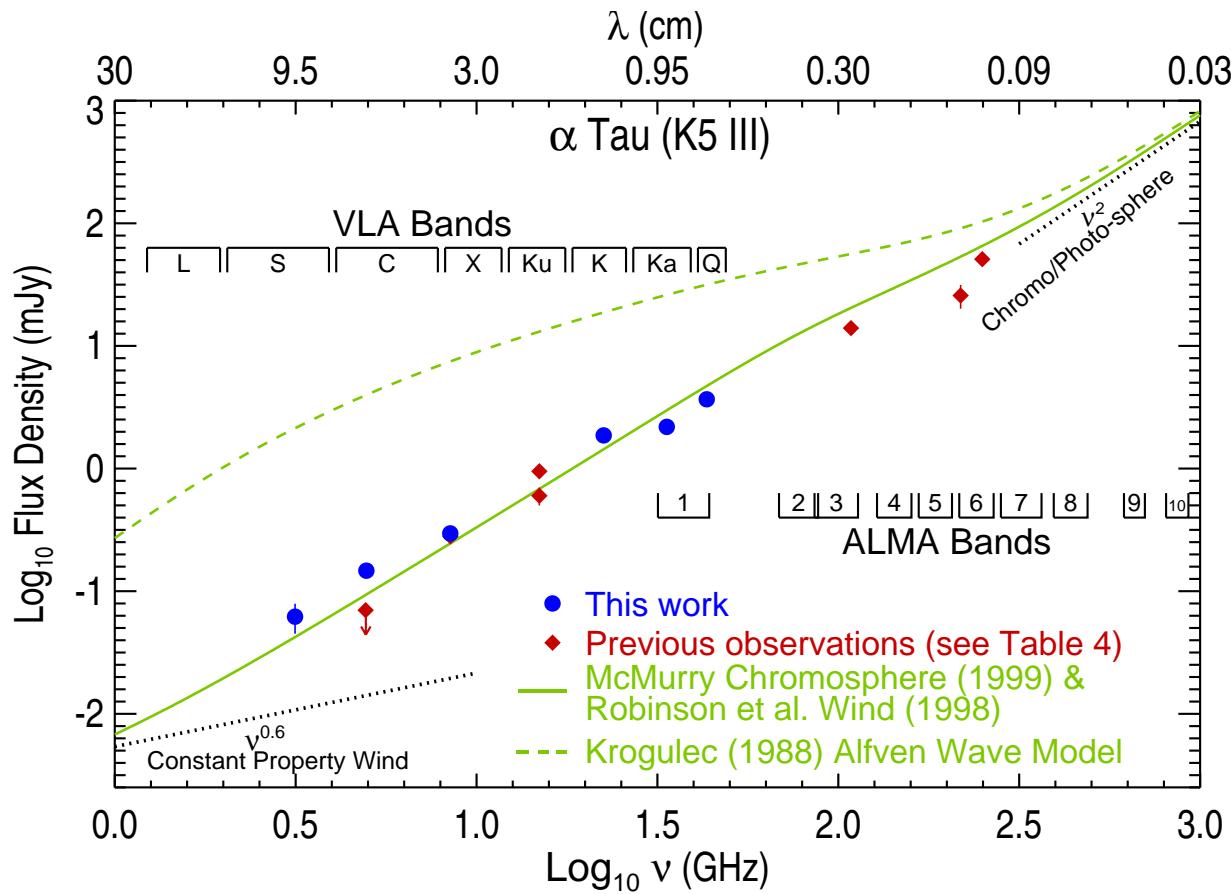
In Figure 6.3 we plot the previous radio measurements of  $\alpha$  Tau at all frequencies below 250 GHz (i.e.  $> 0.12$  cm). Prior to this study,  $\alpha$  Tau had only been detected at two VLA bands (i.e., X and Ku) and had never been detected at wavelengths longer than 3 cm due to its relatively low mass-loss rate. Our

### 6.3 Results vs Previous Observations



**Figure 6.2:** Spectral energy distribution of α Boo for  $1 \text{ GHz} \leq \nu \leq 1 \text{ THz}$ . Our new multi-frequency VLA observations which were mainly acquired over a few days in February 2011 are the blue circles and disagree with the existing chromospheric and wind models of Drake (1985). The overlap between the two models is represented by the green shaded area. The red diamonds are previous observations which were acquired sporadically over the last three decades with the ‘old’ VLA, IRAM and BIMA. The black dashed line is the expected radio emission from the Drake model which undergoes rapid wind cooling beyond  $\sim 2.3 R_\star$  (see Section 6.5 and 6.6).

### 6.3 Results vs Previous Observations



**Figure 6.3:** Spectral energy distribution of  $\alpha$  Tau for  $1 \text{ GHz} \leq \nu \leq 1 \text{ THz}$ . Our new multi-frequency VLA observations of  $\alpha$  Tau (blue circles) were acquired in just two days in February 2011. The red diamonds are the previous radio observations of the star which were acquired over many years (see Table 6.1). The green line is the expected radio emission from the existing hybrid chromosphere and wind model, while the dashed green line is the expected radio emission from a theoretical Alfvén wave driven model atmosphere.

## 6.4 Results vs Existing Models

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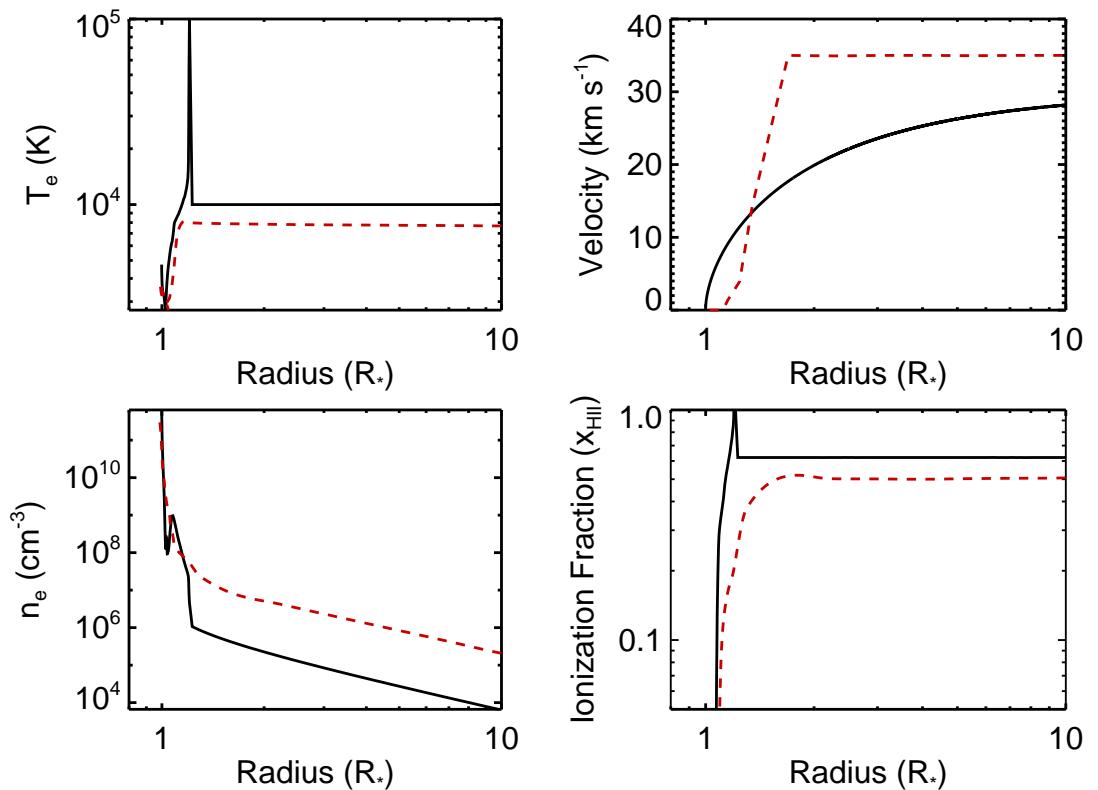
lack of a Ku-band measurement means that we can only compare the previous 3 cm detection reported in [Wood \*et al.\* \(2007\)](#) to ours. We find that there is no significant difference between the two. Interestingly, [Wood \*et al.\* \(2007\)](#) report a non-detection of  $\alpha$  Tau at 6 cm and placed a  $3\sigma$  upper limit of 0.07 mJy on its emission. In stark contrast to this, we were able to detect the star at 6 cm with a flux density over two times greater than this value. This hint of variability at long wavelengths would be consistent with the predictions of the broadband nonlinear Alfvén wave model of [Airapetian \*et al.\* \(2010\)](#) but can again only be confirmed with future high S/N observations.

## 6.4 Results vs Existing Models

One of the most important diagnostic features indicating mass outflows in late-type evolved stars are the blue shifted absorption components present in the Ca II H and K and Mg II h and k resonance lines. Figure 6.4 shows one of the two chromosphere and wind models of  $\alpha$  Boo ([Drake, 1985](#), ‘model A’) which is based on the Mg II k 2796 Å emission line observed with the International Ultraviolet Explorer. The line was modeled by solving the radiative transfer equation in a spherical co-moving frame and the effects of partial redistribution (e.g., [Drake & Linsky, 1983b](#)) were taken into account. Both of Drakes atmospheric models are semi-empirical and contain no assumptions about the wind driving mechanism. They contain the photospheric model of [Ayres & Linsky \(1975\)](#), predict the wind to reach a terminal velocity of 35 - 40 km s<sup>-1</sup> by  $2 R_\star$ , and reach a maximum microturbulence of 5 km s<sup>-1</sup>. They contain a broad temperature plateau with  $T_e \approx 8,000$  K between 1.2 and  $\sim 20 R_\star$  with a cooler region further out, and hydrogen is 50% ionized throughout. We compute the radio spectrum from these models assuming spherical 1-D geometry ([Harper, 1994](#)) with the free-free Gaunt factors from [Hummer \(1988\)](#). The radiative transfer equation is solved using the Feautrier technique ([Mihalas, 1978](#)) and the boundary condition is determined by ensuring the atmosphere is optically thick at the deepest layers. [Drake \(1985\)](#) predicts that their atmospheric model would produce a flux density value of 0.4 mJy at 6 cm and encouragingly, our radio spectrum reproduces this value. Departures from spherical symmetry are to be expected in magnetic stellar atmospheres. For

## 6.4 Results vs Existing Models

example,  $\alpha$  Boo has an inclination axis of  $58^\circ \pm 25^\circ$  (Gray *et al.*, 2006) and a global magnetic dipole could cause density variations between the equator and the polar regions. Despite this fact, the study of a spherically symmetric atmosphere forms the basis of understanding the more complex environments in real stellar atmospheres.



**Figure 6.4:** 1-D semi-empirical model atmospheres for  $\alpha$  Boo and  $\alpha$  Tau. The dashed red line is the model atmosphere for  $\alpha$  Boo (Drake, 1985) while the solid black line is the model atmosphere of  $\alpha$  Tau which consists of the chromosphere and transition region model of McMurry (1999) and the wind model of Robinson *et al.* (1998).

Figure 6.2 shows the resulting predicted radio spectrum between 1 GHz and 1 THz for  $\alpha$  Boo from these chromosphere and wind models (green line). At high frequencies the radio spectra produced by these models have a blackbody-like slope (i.e.,  $\sim \nu^2$ ) as a result of the small ion density scale heights close to the star where the temperature is changing slowly. At low frequencies however,

## 6.4 Results vs Existing Models

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where the Drake models predict the wind to have constant velocity, ionization fraction, and temperature, the slopes approaches the well known  $\sim \nu^{0.6}$  limit (Olnon, 1975; Panagia & Felli, 1975; Wright & Barlow, 1975). The paucity and, in some cases low S/N, of previous observations made it difficult to discern the validity of this model prior to our multi-frequency study of  $\alpha$  Boo. Our new data reveal significant deviations from the semi-empirical model at both low and high frequencies (in this case below  $\sim 8$  GHz and above  $\sim 25$  GHz). At high frequencies our VLA data indicates a flux excess which is in agreement with previous mm-observations. This may be due to larger chromospheric ion densities or to the possible presence of transition region plasma not accounted for in the Drake model. The discrepancy at low frequencies may be due to a lower ionization fraction in the wind, or a lower mass loss rate than that used in the Drake model and this will be discussed further in Section 6.6.

In Figure 6.3 we plot the expected radio spectrum of  $\alpha$  Tau based on the semi-empirical 1-D chromosphere and transition region model of McMurry (1999) embedded in the 1-D wind model of Robinson *et al.* (1998). The semi-empirical McMurry model was created by using the radiative transfer code MULTI (Carlssson, 1986) to reproduce the fluxes of collisionally excited C I, C II, Si III, Mg II, and C IV lines in a plane-parallel, hydrostatic, one-component atmosphere. It contains the photospheric model of Johnson (1973) and reaches a maximum temperature of  $10^5$  K at  $1.2 R_\star$ . As it does not contain a wind outflow, we use Robinson et al.'s wind characteristics beyond  $1.2 R_\star$  to describe the outflow velocity. In this wind model, the wind reaches  $\sim 80\%$  of its terminal value of  $30 \text{ km s}^{-1}$  by  $3 R_\star$ . The Robinson et al. wind characteristics are based on matching the Fe II 2755 Å line and the O I triplet near 1304 Å with a simplified wind model using the SEI computer code (Lamers *et al.*, 1987). We assume the wind to have a constant temperature of 10,000 K and have a constant ionization fraction of  $x_e = 0.6$  throughout, based on the ionization fraction at the corresponding temperature in the McMurry model. This idealized *hybrid* model atmosphere for  $\alpha$  Tau is plotted in Figure 6.4.

The radio flux densities at high frequencies (i.e.,  $\nu > 30$  GHz) are overestimated by the combination of both atmospheric models, although this approach

## 6.4 Results vs Existing Models

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does well in reproducing the VLA flux densities below 30 GHz. The VLA, Institut de Radioastronomie Millimétrique (IRAM) 30 m-telescope and Berkeley Illinois Maryland Association (BIMA) continuum flux densities confirm that this model predicts a flux excess at even higher frequencies. One possible explanation for this is that the inner atmosphere contains extensive amounts of cooler gas than that predicted by the 1-D static chromospheric model of McMurry. This scenario agrees with the findings of [Wiedemann \*et al.\* \(1994\)](#) who conclude that cool regions exist close to the stellar surface with large ( $> 99\%$ ) filling factors i.e., a thermally bifurcated CO-mosphere ([Ayres, 1996](#)). The wind which we have overlain on top of the McMurry chromosphere and transition region is found to be optically thin at nearly all VLA wavelengths, and only contributes a very small flux density at the longest wavelengths. As our model matches the data reasonably well below 30 GHz we conclude that  $\alpha$  Tau's wind is optically thin and the VLA radio emission at all wavelengths emanates from the inner atmosphere.

We also include the predicted radio spectrum from the theoretical Alfvén wave-driven outflow model for  $\alpha$  Tau ([Krogulec, 1989](#)) in Figure 6.3 to demonstrate how radio observations can empirically challenge theoretical models. This model has a fully-ionized outflow inside  $10 R_\star$ , and has a mass loss rate of  $6.3 \times 10^{-9} M_\odot \text{ yr}^{-1}$ , more than two orders of magnitude higher than the more recent estimate given in Table 6.1. As the radio opacity is proportional to  $n_e n_{\text{ion}}$ , where  $n_e$  and  $n_{\text{ion}}$  are the electron and ion number densities respectively, this model greatly overestimates the actual radio flux density at all VLA wavelengths. The linear Alfvén wave models for  $\alpha$  Boo ([Krogulec, 1988](#)) also assume full ionization and have higher mass loss rates than the value given in Table 6.1, predicting higher flux densities than observed. The lack of agreement between the Alfvén wave-driven wind models of [Krogulec \(1988, 1989\)](#) and our observed radio fluxes may not necessarily be due to an incorrect wind driving mechanism and instead may be due to the simplifications and uncertainties in these models, such as wind densities, magnetic field strengths, damping lengths, and flow geometries close to the star. For example, the mass loss rate is very sensitive to the radial surface magnetic field strength (i.e.,  $\dot{M} \propto B^4$ ) in these Alfvén wave models ([Holzer \*et al.\*, 1983](#)) so a small uncertainty in the mass loss rate can lead to a large uncertainty in the magnetic field strength. Relaxing some of these simplifications such as

purely radial flows or non-assumption of the WKB approximation ([Charbonneau & MacGregor, 1995](#)) may also also lead to better agreement with our radio data.

## 6.5 Spectral Indices

Long wavelength radio emission from non-dusty K spectral-type red giants is due to thermal free-free emission in their partially ionized winds while shorter wavelength radio emission emanates from the near static and more ionized lower atmospheric layers. The radio flux density-frequency relationship for these stars is usually found to be intermediate between that expected from the isothermal stellar disk emission, where  $\alpha$  follows the Rayleigh-Jeans tail of the Planck function (i.e.,  $\alpha = +2$ ), and that from an optically thin plasma (i.e.,  $\alpha = -0.1$ ). We have shown in Chapter 1 that the expected radio spectrum from a spherically symmetric isothermal outflow with a constant velocity and ionization fraction varies as  $\nu^{0.6}$  ([Olonon, 1975](#); [Panagia & Felli, 1975](#); [Wright & Barlow, 1975](#)). In reality however, thermal gradients will exist in the wind when the heating mechanisms become insufficient to counteract adiabatic and line cooling, so one would expect a temperature decrease in the wind at some point. Also, if the radio emission emanates from the wind acceleration zone then the electron density will not follow  $n_e \propto r^{-2}$ .

We therefore relax some of the constant property wind model assumptions and assume that the electron density and temperature vary as a function of distance from the star  $r$ , and have the power-law form  $n_e \propto r^{-p}$  and  $T_e \propto r^{-n}$ , respectively (e.g., [Seaquist & Taylor, 1987](#)). Finding the spectral index for an outflow with these conditions is non-trivial so we highlight the main steps required to do so here. We assume the same geometry and notation used for the constant property wind model in Chapter 1, and again start by calculating the total optical depth for a ray with an impact parameter  $b$ , through the atmosphere:

$$\tau_\nu = \frac{0.212Z^2 n_e^2(r_0) r_0^{2p}}{T^{1.35} \nu^{2.1} T_e^{1.35}(r_0) r_0^{1.35n}} \int_{-\infty}^{\infty} \frac{1}{(b^2 + z^2)^{(2p-1.35n)/2}} dz. \quad (6.1)$$

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This integral can be solved using the relationship

$$\int_{-\infty}^{\infty} \frac{1}{(b^2 + z^2)^{t/2}} dz = b^{1-t} \sqrt{\pi} \left[ \frac{\Gamma(0.5t - 0.5)}{\Gamma(0.5t)} \right] \quad (6.2)$$

and setting  $t = (2p - 1.35n)$ . Here  $\Gamma$  is the gamma function i.e.,  $\Gamma(y) = \int_0^\infty u^{y-1} e^{-u} du$ . The total optical depth along a ray is then

$$\tau_\nu(b) = G b^{1-2p+1.35n} \quad (6.3)$$

where  $G$  is a constant that incorporates  $\nu$ . As the total flux density is

$$F_\nu = \frac{2\pi}{D^2} \int_0^\infty B_\nu [1 - e^{-\tau_\nu(b)}] b db, \quad (6.4)$$

we can now substitute in the Rayleigh Jeans function for  $B_\nu$  (remembering that  $T_e$  now depends on distance from the star) to get

$$F_\nu = \frac{4\pi k \nu^2 T_e(r_0) r_0^n}{D^2 c^2} \int_0^\infty (1 - e^{-\tau_\nu}) (b^2 + z^2)^{-\frac{n}{2}} b db. \quad (6.5)$$

Expansion of the second term inside the integral gives

$$F_\nu \simeq \frac{4\pi k \nu^2 T_e(r_0) r_0^n}{D^2 c^2} \int_0^\infty (1 - e^{-\tau_\nu}) b^{1-n} db. \quad (6.6)$$

To progress further, Equation 6.3 can be rearranged to find  $b$  and  $db$  in terms of  $\tau_\nu$ , and can then be inserted into Equation 6.6 to give

$$F_\nu \simeq \frac{4\pi k \nu^2 T_e(r_0) r_0^n}{D^2 c^2} \int_0^\infty (1 - e^{-\tau_\nu}) \left( \frac{\tau_\nu}{G} \right)^{\frac{1-2.35n+2p}{1-2p+1.35n}} \frac{d\tau_\nu}{G(1-2p+1.35n)}. \quad (6.7)$$

This allows  $\nu$  to be separated out to give

$$F_\nu \propto \nu^2 G^{\frac{2.35n-2p-1}{1-2p+1.35n}} G^{-1} \quad (6.8)$$

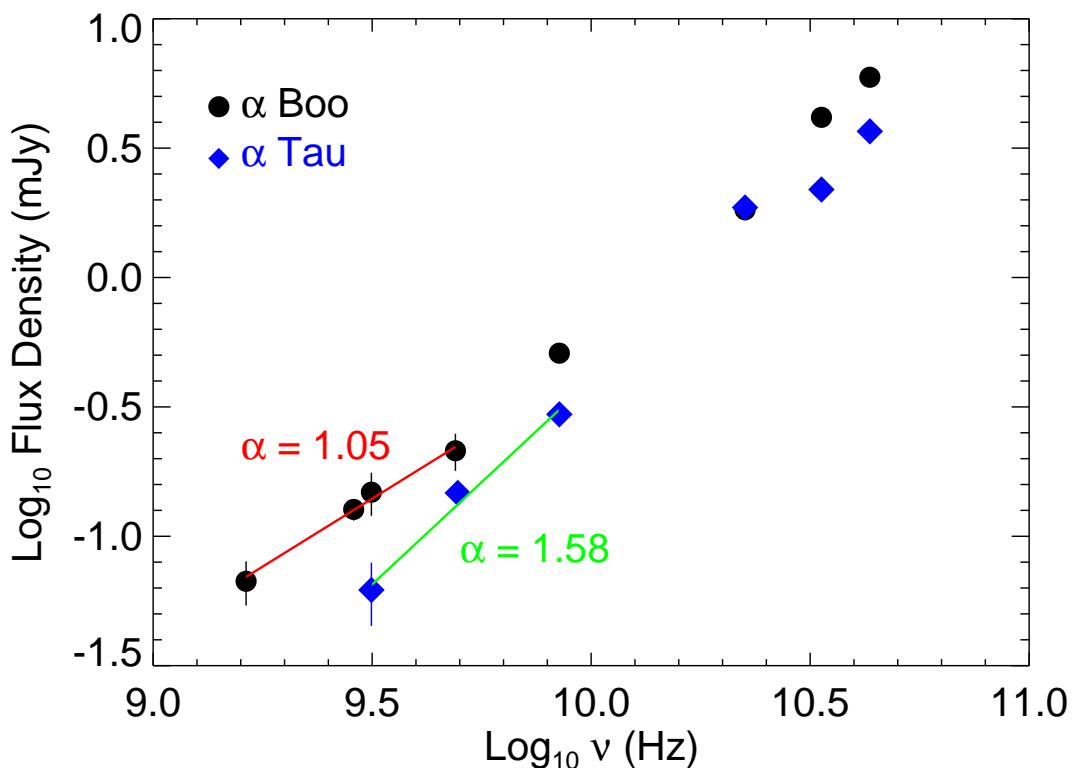
and as  $G \propto \nu^{-2.1}$  we get

$$F_\nu \propto \nu^2 \nu^{\frac{4.2-2.1n}{1-2p+1.35n}}, \quad (6.9)$$

i.e.,

$$\alpha = \frac{4p - 6.2 - 0.6n}{2p - 1 - 1.35n}. \quad (6.10)$$

Therefore, if the spectral index of a stellar outflow can be measured, and if we make an assumption about the property of either the thermal or electron density profile of the wind, then Equation 6.10 provides us with information on how the other value varies.



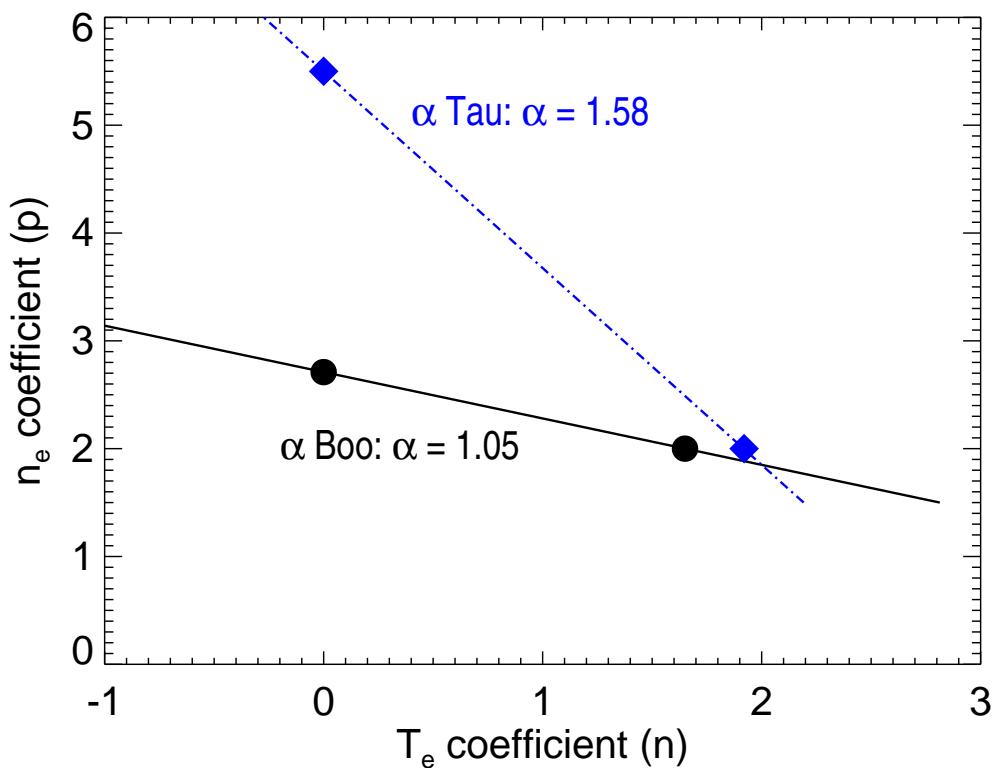
**Figure 6.5:** Radio spectra for  $\alpha$  Boo and  $\alpha$  Tau, together with the best fit power law to their long wavelength flux densities and the resulting spectral indices. The spectral indices for  $\alpha$  Boo and  $\alpha$  Tau are found to be 1.05 and 1.58, respectively, which are both larger than the 0.6 value expected for a constant property wind model.

The radio spectra for both stars are shown in Figure 6.5, together with the power laws that were fitted to the long wavelength flux densities by minimizing the chi-square error statistic. For  $\alpha$  Boo, a power law with  $F_\nu \propto \nu^{1.05 \pm 0.05}$  fits the four longest wavelength data points well. This spectral index is larger than the 0.8 value obtained by [Drake & Linsky \(1986\)](#) whose value was based on a shorter wavelength (2 cm) value and a mean value of four low S/N measurements

## 6.5 Spectral Indices

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at 6 cm.  $\alpha$  Tau was found to have a larger spectral index and a power law with  $S_\nu \propto \nu^{1.58 \pm 0.25}$  best fitted the three longest wavelength data points. This value is in agreement with [Drake & Linsky \(1986\)](#) who report a value  $\geq 0.84$  and is lower than the value of 2.18 that can be derived from the shorter wavelength data given in [Wood \*et al.\* \(2007\)](#). It should be emphasized that the spectral index for both stars is a lot steeper than that expected from the idealized constant property wind model.



**Figure 6.6:** The variation of density and temperature coefficients for the empirically derived spectral indices. The density coefficients for an isothermal flow ( $n=0$ ) along with the temperature coefficients for a constant outflow velocity ( $p=2$ ) are also shown for both stars.

Equation 6.10 can be used in conjunction with our new spectral index for each star to calculate the density and temperature coefficients that may describe their outflows. The combinations of the electron temperature and density coefficients are shown for each star in Figure 6.6 ( $\alpha$  Boo is represented by the solid line and

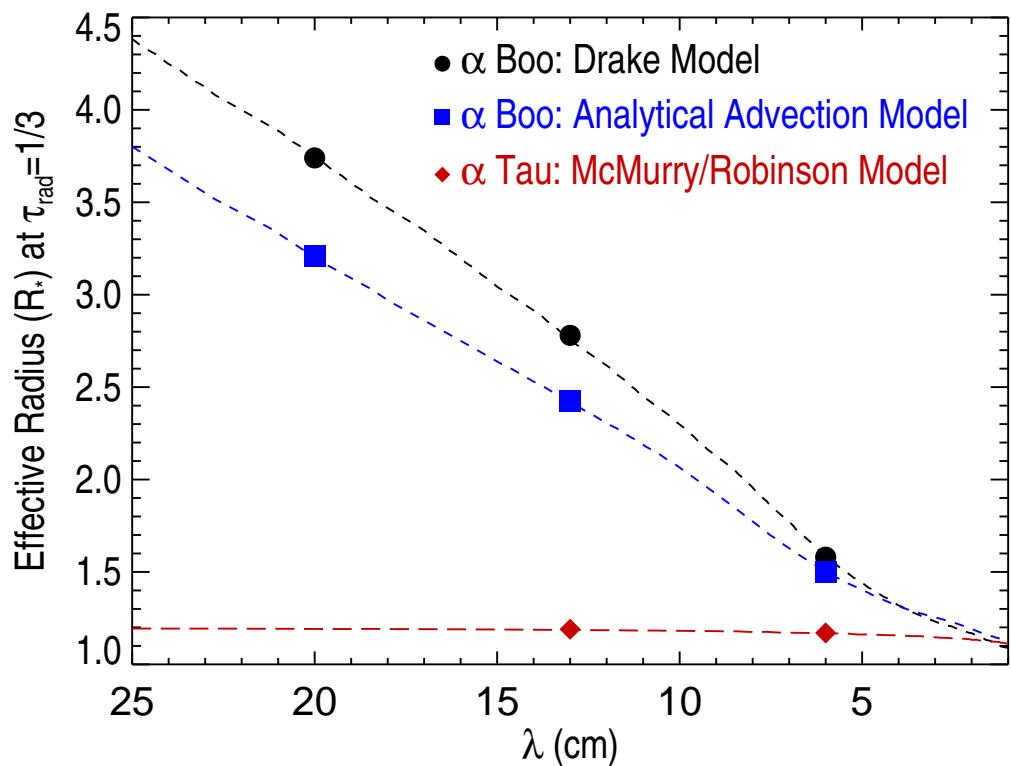
## 6.5 Spectral Indices

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$\alpha$  Tau by the dash-dotted line) along with the coefficients obtained by assuming either an isothermal flow ( $n = 0$ ) or a constant velocity flow ( $p = 2$ ). One potential explanation for spectral indices of stellar outflows being larger than 0.6 is that the wind is still accelerating in the radio emitting region, if the thermal gradients are assumed to be small. For an isothermal flow, the density coefficients are  $p = 2.7$  and  $5.5$  for  $\alpha$  Boo and  $\alpha$  Tau, respectively. From mass conservation and assuming a steady flow, the power law coefficients for the velocity profiles of each star can be found, ie.,  $v(r) \propto r^{p-2}$ . For  $\alpha$  Boo we find  $v(r) \propto r^{0.7}$  while for  $\alpha$  Tau we find  $v(r) \propto r^{3.5}$ . This suggests that our long VLA wavelengths may probe a steep acceleration region for  $\alpha$  Tau's outflow but for  $\alpha$  Boo, may probe a region where the wind is close to its terminal velocity.

The assumption of shallow thermal gradients in a stellar outflow is probably unreliable however. It is likely that some form of Alfvèn waves are required to lift the material out of the gravitational potential as suggested by [Hartmann & MacGregor \(1980\)](#). These waves need damping lengths which are much larger than the chromospheric density scale height  $H$  (where  $H \sim 0.01 R_\star$  for our targets), in order to lift the material out of the gravitational potential but  $\lesssim 1 R_\star$  in order to avoid wind terminal velocities greater than those observed (e.g. [Holzer et al., 1983](#)). It has also been shown that these waves are expected to produce substantial heating near the base of the wind (e.g. [Hartmann et al., 1982](#)). If the long wavelength radio emission from  $\alpha$  Tau is indeed emanating from the wind acceleration region, then the dissipation of these Alfvèn waves may introduce thermal gradients in this region and its velocity profile will not be described by  $v(r) \propto r^{3.5}$ . Furthermore, diverging flow geometries have been invoked as a more realistic representation of stellar atmospheres ([Hartmann et al., 1982](#); [Jatenco-Pereira & Opher, 1989](#); [Vidotto et al., 2006](#)) and so one could write the area of a flux tube (normalized to its value at  $r_1$ ) as  $(r/r_1)^2 f(r)$  where  $f(r)$  is a function describing the divergence from a purely radial flow [i.e., when  $f(r) = 1$ ]. If the non-radial expansion term can be described by a power law  $f(r) \propto r^s$ , where  $s > 1$  indicates super-radial expansion, then  $v(r) \propto r^{3.5-s}$  would describe the velocity profile of  $\alpha$  Tau (assuming an isothermal flow with a constant ionization fraction). Therefore, including diverging geometries reduces the magnitude of the acceleration. Further out in the wind, where it has reached its terminal

velocity, one would also expect a thermal gradient (but now of opposite sign) due to adiabatic expansion and line cooling. If the long wavelength radio emission emanates from this region of the wind then Equation ?? provides us with a direct estimate of the temperature coefficient as we can assume the density coefficient is  $p = 2$ . This may be the case for our long VLA wavelength measurements of  $\alpha$  Boo in which case  $T_e(r) \propto r^{-1.65}$ .



**Figure 6.7:** Predicted effective radius (dashed lines) as a function of wavelength derived from the existing atmospheric models of  $\alpha$  Boo and  $\alpha$  Tau. Also plotted is the predicted effective radius derived from our analytical advection model for  $\alpha$  Boo (discussed in Section 6.6). Points corresponding to our long wavelength VLA measurements are also shown. At the same radio wavelengths the lower mass loss rate of  $\alpha$  Tau's outflow results in a smaller effective radius than that for  $\alpha$  Boo.

To investigate this matter further, we estimate the effective radius of the radio emitting region as a function of wavelength based on the Drake model for  $\alpha$  Boo and the hybrid McMurry and Robinson et al. model for  $\alpha$  Tau. We follow

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

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the approach used by [Cassinelli & Hartmann \(1977\)](#) and assume that the radio emission at each wavelength is characterized by emission from a radial optical depth  $\tau_{\text{rad}} \sim 1/3$ . This is a modification of the Eddington-Barbier relation for an extended atmosphere where emission from smaller optical depths has added weight. Since the radio free-free opacity increases at longer wavelengths the optical depth along a line of sight into the stellar outflow also increases at longer wavelengths. This implies that the effective radius (i.e., the radius where  $\tau_\lambda = \tau_{\text{rad}}$ ) will increase with longer wavelengths and will be greater for outflows with higher densities of ionized material as  $\tau_\lambda(r) \propto \lambda^{2.1} \int n_{\text{ion}}(r)n_e(r)dr$ .

The larger mass loss rate of  $\alpha$  Boo in comparison to  $\alpha$  Tau means that the latter has a substantially smaller effective radius at longer wavelengths, as seen in Figure 6.7. At 6, 13, and 20 cm the effective radius of  $\alpha$  Boo at  $\tau_{\text{rad}} = 1/3$  is predicted to be 1.6, 2.8, and 3.7  $R_\star$  but is only  $\sim 1.2 R_\star$  at 6 and 13 cm for  $\alpha$  Tau. [Robinson et al. \(1998\)](#) predict that  $\alpha$  Tau's wind reaches  $\sim 80\%$  of its terminal velocity by  $3 R_\star$ , but even our longest-wavelength observations are highly unlikely to sample the wind outside the lower velocity layers closer to the star. For  $\alpha$  Boo however, [Drake \(1985\)](#) predicts that the wind has reached its terminal velocity by  $\sim 2 R_\star$  so based on this model our longest-wavelength measurements are of the region where the wind has reached a steady terminal velocity. From Figure 6.6, this implies that the  $n_e$  coefficient is  $p = 2$  and thus the  $T_e$  coefficient is  $n = 1.65$ . Pure adiabatic spherical expansion cooling with no heat source has  $n = 1.33$  so additional cooling routes must be operating, possibly due to line cooling. Finally, the wind ionization balance may not have become *frozen-in* in the region of  $\alpha$  Boo's wind where the radio emission emanates from. If this is true, then the excess slope of the spectral index could be due to a combination of both cooling and changing ionization fraction. In this scenario the temperature coefficient  $n$ , would be smaller than our derived value because Equation 6.10 assumes a constant ionization fraction.

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

A failure of the Drake model for  $\alpha$  Boo is that it overestimates the long wavelength VLA radio flux densities which sample the outer atmosphere, as clearly shown

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

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in Figure 6.2. If these wavelengths are indeed sampling the wind at its terminal velocity then one reason for this overestimation is that the wind is cooling closer in than predicted by the existing model, which assumes a constant temperature of 8,000 K out to  $\sim 20 R_\star$ . The main mechanism for such cooling would be adiabatic expansion (O'Gorman & Harper, 2011) and would cause lower electron densities than those predicted by the existing model due to larger recombination rates. In the next two sections we derive the fraction of ionized hydrogen in a stellar outflow with a temperature gradient based on the work of Glassgold & Huggins (1986), and use this method to insert a temperature gradient into the outer atmosphere of the Drake model to see if such an atmosphere could better reproduce the new long wavelength VLA flux densities.

### 6.6.1 H II recombination in a stellar outflow

The time dependent non-static rate/statistical equations can be written as

$$\frac{\partial n_i}{\partial t} + \nabla \cdot (n_i v) = \sum_{i \neq j}^n n_j P_{ji} - n_i \sum_{i \neq j}^n P_{ij} \quad (6.11)$$

where  $v$  is the flow velocity,  $n_{i,j}$  are the population densities of the energy levels  $i, j$  which are functions of radial distance  $r$ , and  $P_{ij} = C_{ij} + R_{ij}$  are the total transition probabilities ( $s^{-1}$ ) from energy levels  $i \rightarrow j$ .  $C_{ij}$  are the total collision rates (electrons, protons, and ions), and  $R_{ij}$  are the problematic radiative rates. In 1-D spherical geometry

$$\nabla \cdot \mathbf{A} = \frac{1}{r^2} \frac{d}{dr} (r^2 A), \quad (6.12)$$

where  $\mathbf{A}$  is some vector quantity and equals  $n_i v$  in our case. Thus, for a steady flow the rate equations become

$$\frac{1}{r^2} \frac{d}{dr} (r^2 v n_{\text{tot}} \frac{n_i}{n_{\text{tot}}}) = \sum_{i \neq j}^n n_j P_{ji} - n_i \sum_{i \neq j}^n P_{ij} \quad (6.13)$$

where  $n_{\text{tot}}$  is the total hydrogen number density (i.e.  $n_{\text{tot}} = n_{\text{HI}} + n_{\text{HII}}$ ). We note that  $r^2 v n_{\text{tot}}$  is some constant of the flow defined by the mass continuity equation,

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

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and if we define the relative populations as  $x_i = n_i/n_{\text{tot}}$  we get

$$v \frac{d}{dr}(x_i) = \sum_{i \neq j}^n x_j P_{ji} - x_i \sum_{i \neq j}^n P_{ij}. \quad (6.14)$$

To simplify Equation 6.14 further, we follow the approach of Glassgold & Huggins (1986) and assume:

1. A constant velocity mass outflow, i.e.,  $n_{\text{tot}} = C/r^2$  where  $C$  is a constant proportional to the ratio of the mass loss rate divided by the terminal velocity.
2. All hydrogen ionization processes cease beyond some distance  $r_1$  (i.e.  $P_{ij} = 0$ ). The ionization of hydrogen in the chromosphere and wind is a two stage process: the  $n = 2$  level is excited by electron collisions and Lyman-alpha scattering, followed by photoionization by the optically thin Balmer continuum. When the temperature begins to decrease in the wind the collisional excitation rate and thus ionization rate decrease rapidly.
3. Only radiative recombination of hydrogen is considered [i.e.  $R_{ji} = \alpha_{\text{rr}}(T_e)n_e$  where  $\alpha_{\text{rr}}$  is the total radiative recombination coefficient].
4. A fixed ion contribution from metals with a low first ionization potential,  $x_{\text{ion}} = n_{\text{ion}}/n_{\text{tot}}$ , as these are easily ionized in the outflow. The total electron density is then  $n_e = n_{\text{tot}}x_{\text{HII}} + n_{\text{tot}}x_{\text{ion}}$  where  $x_{\text{HII}} = n_{\text{HII}}/(n_{\text{HI}} + n_{\text{HII}})$ .

Using these assumptions Equation 6.14 becomes

$$\frac{dx_{\text{HII}}}{dr} = \frac{\alpha_{\text{rr}}C}{vr^2} (x_{\text{HII}}^2 + x_{\text{ion}}x_{\text{HII}}) \quad (6.15)$$

which can be rearranged and integrated between  $r_1$  and  $r$  to give

$$\int_{x_{\text{HII}(r_1)}}^{x_{\text{HII}(r)}} \frac{dx_{\text{HII}}}{x_{\text{HII}}(x_{\text{HII}} + x_{\text{ion}})} = \int_{r_1}^r \frac{\alpha_{\text{rr}}C}{vr^2} dr. \quad (6.16)$$

The ionization fraction beyond  $r_1$  is then given by

$$x_{\text{HII}}(r) = \frac{x_{\text{HII}}(r_1)x_{\text{ion}}e^{-I(r)}}{x_{\text{ion}} + x_{\text{HII}}(r_1)[1 - e^{-I(r)}]} \quad (6.17)$$

where

$$I(r) = \int_{r_1}^r \frac{x_{\text{ion}}\alpha_{\text{rr}}C}{vr^2} dr \quad (6.18)$$

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

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If we further assume a constant velocity,  $T_e \propto r^{-n}$ , and  $\alpha_{rr} \propto T_e^{-\beta}$  giving  $\alpha_{rr} \propto r^{-n\beta}$ , then

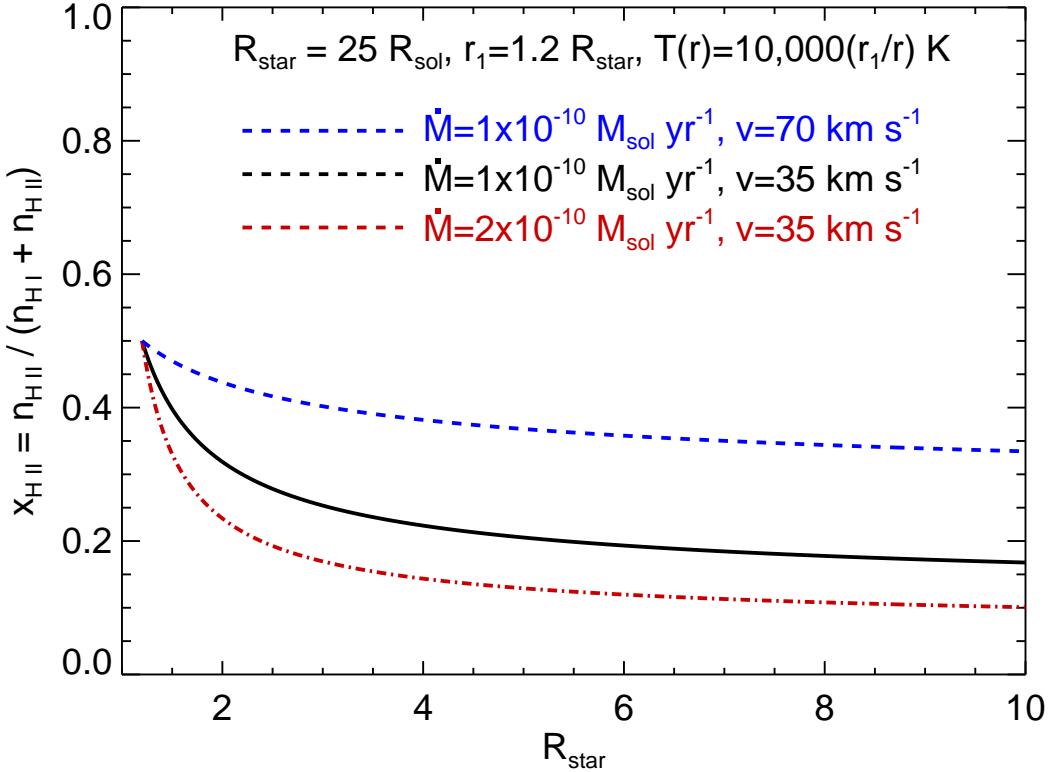
$$x_{\text{HII}}(r) = \frac{x_{\text{HII}}(r_1)x_{\text{ion}}e^{-I(r)}}{x_{\text{ion}} + x_{\text{HII}}(r_1)(1 - e^{-I(r)})} \quad (6.19)$$

where

$$I(r) = w \left[ 1 - \left( \frac{r_1}{r} \right)^{1-n\beta} \right] \quad (6.20)$$

and

$$w = \frac{x_{\text{ion}}\alpha_{rr}(r_1)C}{vr_1(1 - n\beta)}. \quad (6.21)$$



**Figure 6.8:** Illustration of how the ionization fraction *freezes-in* in a stellar wind for a hypothetical  $25 R_{\star}$  star. The stars outflow has a maximum temperature of 10,000 K at  $1.2 R_{\star}$  which then falls off linearly with distance. Blue dashed line:  $\dot{M} = 1 \times 10^{-10} M_{\odot} \text{ yr}^{-1}$  and  $v = 70 \text{ km s}^{-1}$ . Black smooth line:  $\dot{M} = 1 \times 10^{-10} M_{\odot} \text{ yr}^{-1}$  and  $v = 35 \text{ km s}^{-1}$ . Red dot-dashed line:  $\dot{M} = 2 \times 10^{-10} M_{\odot} \text{ yr}^{-1}$  and  $v = 35 \text{ km s}^{-1}$ .

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

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As  $w$  is just a constant,  $I(r)$  approaches a constant value for large  $r$ . This means that  $x_{\text{HII}}(r)$  also approaches a constant value for large  $r$ , and the ionization fraction gets *frozen-in*. This is a universal property of all stellar winds. It occurs when the recombination timescale of the ions in the wind  $\tau_{\text{rec}} = (\alpha_{\text{rec}} n_e)^{-1}$ , becomes larger than the wind expansion timescale,  $\tau_{\text{exp}} = |\rho(d\rho/dt)^{-1}| \simeq 0.5r/v$ . Figure 6.5 shows that the ionization fraction in the outflow of a *typical* red giant *freezes-in* to higher values for winds with higher velocities and lower mass loss rates.

### 6.6.2 Application to $\alpha$ Boo's Wind

To investigate the possibility that  $\alpha$  Boo's wind may be undergoing more rapid cooling closer in to the star than originally predicted by the Drake Model, we adjust one of the existing models [referred to as 'Model A' in [Drake \(1985\)](#)] to include a temperature power-law falloff of the form

$$T_e(r) = T_e(r_1) \left(\frac{r_1}{r}\right)^n, \quad (6.22)$$

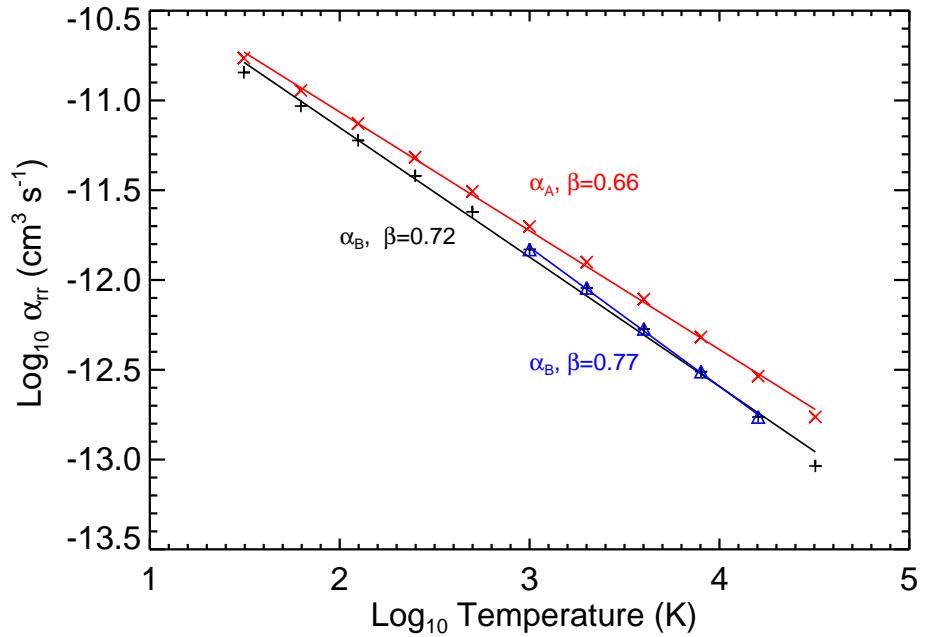
at some distance  $r_1$  from the star. We set the temperature coefficient to the value obtained from our new VLA data which was derived assuming a constant velocity flow, i.e.  $n = 1.65$  (see Figure 6.6). We introduce the distance  $r_1$  as the outer limit to ionization processes; at  $r > r_1$ , the ionization fraction is only determined by recombination. To calculate the new electron densities in the wind regime where this temperature falloff occurs, we use Equations 6.19, 6.20, and 6.21 to calculate the hydrogen ionization fraction,  $x_{\text{HII}} = n_{\text{HII}}/n_{\text{H}}$ . We assume a fixed ion contribution from metals with a low first ionization potential of  $x_{\text{ion}} = 1 \times 10^{-4}$ . So, knowing that the total hydrogen density does not change, the electron density at any point in the wind post  $r_1$  can be found by the following formula

$$n_e(r) = n_{\text{tot}}(r)[x_{\text{HII}}(r) + 1 \times 10^{-4}]. \quad (6.23)$$

Using the same wind terminal velocity as that used in 'Model A' (i.e. 35 km s<sup>-1</sup>) and the mass loss rate defined in Table ??, we find the constant proportional to the ratio of the mass loss rate divided by the terminal velocity to be  $C = 1.7 \times 10^{32}$  cm<sup>-1</sup> for  $\alpha$  Boo. To calculate the temperature dependent radiative

## 6.6 Analytical Advection Model for $\alpha$ Boo's Wind

recombination coefficient  $\alpha_{rr}$  and its power law coefficient  $\beta$  used in Equation 6.21, we consider only radiative recombination of hydrogen and exclude captures to the n=1 level. Spitzer (1978) has tabulated values for the variation of this recombination coefficient with temperature and it is often referred to as  $\alpha_B$  in the literature when captures to the n=1 level are excluded. To find the power law form of  $\alpha_B$  we found the best fit to the tabulated data between the expected range of temperatures in the outflow of  $\alpha$  Boo. This approach is shown in Figure 6.9, where  $\beta$  was found to be 0.77 by fitting a power law to values of  $\alpha_B$  between 1,000 and 16,000 K. For completeness, we also show in Figure 6.9 the slight variation in  $\beta$  if a wider range of temperatures is used and also if captures to the n=1 level are included. It can then be easily shown that the ionization fraction beyond  $r_1$



**Figure 6.9:** The temperature dependent recombination rates for hydrogen. The variation of the recombination rate excluding captures to the n=1 level  $\alpha_B$ , with temperature is derived for the tabulated temperatures in Spitzer (1978) (black) and for the more realistic temperatures existing in  $\alpha$  Boo's outflow (blue). We also show the how the recombination coefficient varies with temperature if captures to the n=1 level are included (red).

## 6.7 Constraining $\alpha$ Tau's Molsphere

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is given by

$$x_{\text{HII}}(r) = \frac{x_{\text{HII}}(r_1)x_{\text{ion}}e^{-I(r)}}{x_{\text{ion}} + x_{\text{HII}}(r_1)[1 - e^{-I(r)}]} \quad (6.24)$$

where

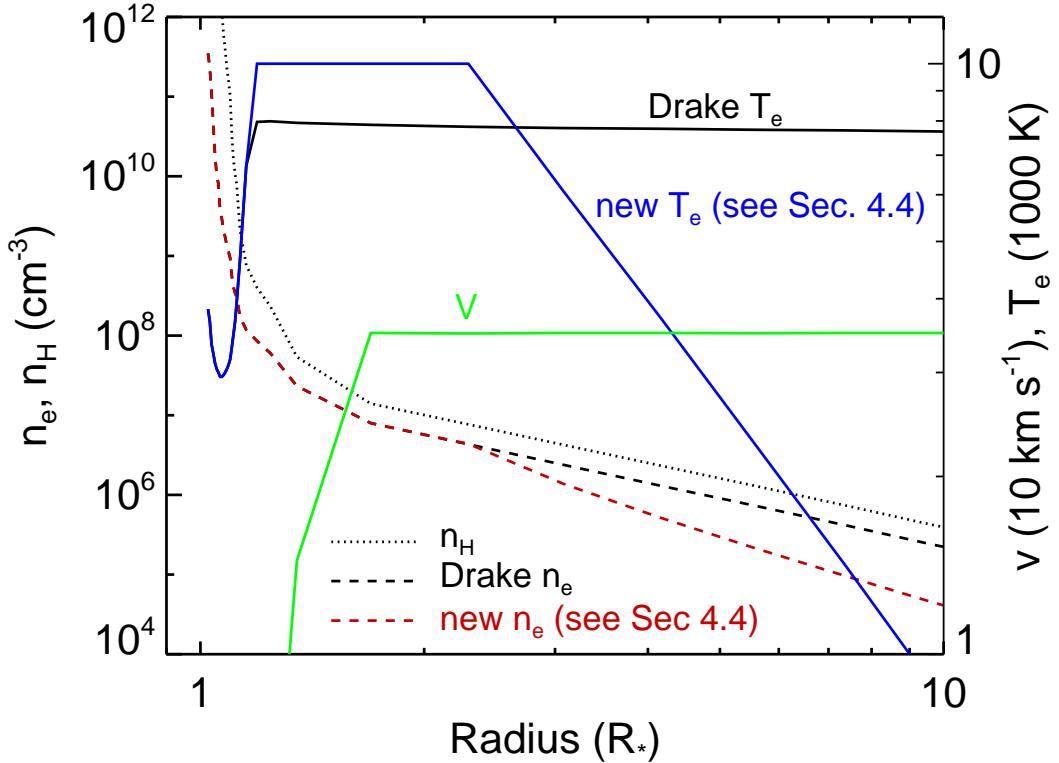
$$I(r) = \frac{4.7 \times 10^9}{r_1} \left[ \left( \frac{r_1}{r} \right)^{-0.27} - 1 \right], \text{ and } r \geq r_1. \quad (6.25)$$

We then adjusted the value of  $r_1$  to obtain the best fit to our long wavelength observations and found this happened when  $r_1 = 2.3 R_\star$ . To get this best fit, the existing atmospheric model needed to be adjusted so that it now has a narrower and slightly larger temperature plateau of  $T_e = 10,000$  K between 1.2 and  $2.3 R_\star$ , and a temperature profile and a density profile governed by Equation 6.22 and Equation 6.24 beyond  $r_1 = 2.3 R_\star$ , respectively. This gives much better agreement with our new long wavelength VLA data as shown in Figure 6.2. This new *hybrid* atmospheric model which is plotted along with the original Drake model in Figure 6.10, still has the original ionization fraction of  $x_{\text{HII}} \approx 0.5$  inside  $2.3 R_\star$  but now contains an initial rapid decrease in  $x_{\text{HII}}$  beyond  $2.3 R_\star$  which then *freezes-in* to a constant value of  $\sim 0.04$  beyond  $\sim 10 R_\star$ .

Encouraging as it is that such a simple analytical model can reproduce values close to the observed radio fluxes at long wavelengths, it must be stressed that this *hybrid* model is just a first order approximation. It assumes that the excess slope from the radio spectrum is a result of rapid cooling only. It still does not reproduce the radio fluxes at wavelengths shorter than  $\sim 3$  cm and therefore a new atmospheric model is still required that can reproduce all of the observed flux densities. To do so, the non-trivial task of simultaneously solving the radiative transfer equation and non-LTE atomic level populations which include advection will be required.

## 6.7 Constraining $\alpha$ Tau's Molsphere

Recently, Ohnaka (2013) has detected a layer of CO in the outer atmosphere of  $\alpha$  Tau (i.e. a so-called MOLsphere) which extends out to  $\sim 2.5 R_\star$ , has a temperature of  $1500 \pm 200$  K, and a CO column density of  $\sim 1 \times 10^{20} \text{ cm}^{-2}$ . They were unable to constrain the geometrical thickness  $\Delta L$ , of the MOLsphere from



**Figure 6.10:** Existing atmospheric model for  $\alpha$  Boo (Drake, 1985, ‘model A’) along with the same model which undergoes rapid wind cooling beyond  $\sim 2.3 R_\star$ . The original Drake Model’s have a temperature plateau of  $\sim 8,000$  K between  $1.2$  and  $\sim 20 R_\star$  (solid black line), reach a terminal velocity of  $35 - 40 \text{ km s}^{-1}$  within  $2 R_\star$  (solid green line), and have a wind which is 50% ionized (dashed and dotted black lines).

the data however, and arbitrarily set it to  $0.1 R_\star$ . We now try and constrain the thickness of this MOLsphere from our VLA data.

At VLA wavelengths, the opacity corrected for stimulated emission is

$$\kappa_\nu = \frac{0.212 n_e n_{ion} Z_{ion}^2}{T_e^{1.35} \nu^{2.1}} \text{ cm}^{-1} \quad (6.26)$$

and the optical depth of a shell of width  $\Delta L$  is

$$\tau_\nu = \kappa_\nu \Delta L. \quad (6.27)$$

To calculate the electron density within the MOLsphere we have conservatively

assumed that the electrons only come from singly ionized metals and have an abundance of  $n_{\text{ion}} = 1 \times 10^{-5} n_{\text{tot}}$ , where  $n_{\text{tot}}$  is the total hydrogen number density whose value is constant throughout the MOLsphere. We estimate the CO number density by dividing the CO column density by the geometrical thickness of the shell and assume the CO filling factor to be unity. If the MOLsphere extends from the stars surface out to the derived distance of  $2.5 R_{\star}$ , then  $\Delta L = 1.5$  and  $n_{\text{co}} = 2.2 \times 10^7 \text{ cm}^{-3}$ . This equates to a total hydrogen number density of  $n_{\text{tot}} = 1.2 \times 10^{11} \text{ cm}^{-3}$  (Ohnaka, 2013) giving an abundance ratio of  $n_{\text{co}}/n_{\text{tot}} = 3.3 \times 10^{-4}$ . We can then use Equations 6.26 and 6.27 to show that such a MOLsphere would be optically thin at all VLA wavelengths that we observed the star with. If on the other hand, the MOLsphere has only a geometric width of  $\Delta L = 0.1 R_{\star}$  then assuming the same abundances we get  $n_{\text{H}} = 1.76 \times 10^{12} \text{ cm}^{-3}$ , and the MOLsphere would be optically thick at 3.5, 6, and 9.5 cm (i.e. X, C, and S-band). Focusing on our high S/N C-band measurement, it can be shown that such a MOLsphere would have an optical depth of  $\tau_{6 \text{ cm}} = 4.6$  and would produce a corresponding flux density of 0.06 mJy (assuming an optically thick disk) which is considerably lower than our measurement of 0.15 mJy. The disagreement in values between our radio data and the predicted flux for an optically thick disk could mean that the MOLsphere is optically thin at long VLA wavelengths and that the radio emission emanates from the more ionized material closer to the star. It can also be shown that the MOLsphere becomes optically thin at C-band (i.e.  $\tau_{6 \text{ cm}} < 1$ ) for  $\Delta L > 0.45 R_{\star}$ . In the next section we argue that the radio emission from  $\alpha$  Tau even at long VLA wavelengths comes from a region closer in to the star where the wind is undergoing rapid acceleration, which suggests that the MOLsphere either has a geometrical width  $> 0.45 R_{\star}$  or has a filling factor less than unity.

## 6.8 Ionized Mass Loss Rates

The free-free radio emission from the atmospheres of red giants is only sensitive to the ionized component of the outflow. Non-coronal and hybrid atmosphere red giants have only partially ionized atmospheres and so we are only ever capable of finding a lower limit on the mass loss rate from their thermal continuum radio emission. Nevertheless, it is a sufficiently worthwhile goal and can test the validity

## 6.8 Ionized Mass Loss Rates

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of existing values obtained from optically thick emission lines in the UV or optical regime. For a partially ionized outflow such as that expected from  $\alpha$  Boo and  $\alpha$  Tau, the ionized mass loss rate can be estimated from the radio emission assuming a constant velocity, constant temperature, and constant ionization fraction and can be written as

$$\dot{M}_{\text{ion}} \simeq 5.58 \times 10^{-14} \left( \frac{v_{\infty}}{\text{km s}^{-1}} \right) \left( \frac{D}{\text{pc}} \right) \left( \frac{F_{\nu}}{\text{mJy}} \right)^{0.75} \left( \frac{\lambda}{\text{cm}} \right)^{0.45} \left( \frac{T_e}{10^4 \text{K}} \right)^{0.1} \dot{M}_{\odot} \text{ yr}^{-1} \quad (6.28)$$

where  $D$  is the distance and  $v_{\infty}$  is the wind terminal velocity defined previously in Table 3.1. Using this expression, the ionized mass loss rates were derived for each of the long wavelength measurements for both stars and are listed in Table 6.3. These values are based on the assumption of a  $T_e = 10^4$  K wind for both stars.

**Table 6.3:** Ionized mass loss rates for  $\alpha$  Boo and  $\alpha$  Tau.

Star	Wavelength (cm)	$\dot{M}_{\text{ion}} (M_{\odot} \text{ yr}^{-1})$
$\alpha$ Boo	6.0	$5.9 \times 10^{-11}$
	9.5	$5.5 \times 10^{-11}$
	10.0	$5.1 \times 10^{-11}$
	20.0	$4.3 \times 10^{-11}$
$\alpha$ Tau	6.0	$\leq 8.2 \times 10^{-11}$
	9.5	$\leq 5.3 \times 10^{-11}$

We have discussed in the previous sections that for  $\alpha$  Boo, these long wavelengths sample the outer atmosphere of the star where the wind is close to or has reached its terminal velocity and is beginning to cool due to gas expansion. In this case, these  $\dot{M}_{\text{ion}}$  values would be lower than those given in Table 6.3 as the electron temperature is probably lower than 10,000 K. However, as  $\dot{M}_{\text{ion}}$  is only weakly dependent on  $T_e$  this makes only a small difference and amounts to about a 7% increase in the ionized mass loss if the temperature is actually lower by 50%. If the ionization balance is frozen-in in the regions of the outflow where these long wavelengths sample, then the derived  $\dot{M}_{\text{ion}}$  should be the same for these wavelengths. The decrease in  $\dot{M}_{\text{ion}}$  with longer wavelengths is very small as

## **6.8 Ionized Mass Loss Rates**

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seen in Table 6.3 and may be a hint that the ionization fraction may be close to frozen-in at these wavelengths. If we were to assume that the total stellar mass loss rate derived by Drake (1985) is correct, then a comparison with our longest wavelength value for the ionized mass loss rate suggests an ionization fraction of  $\sim 0.2$  in the outer wind of  $\alpha$  Boo. Deriving the exact ionized mass loss rate for  $\alpha$  Tau is more difficult because as we have argued in the previous sections, the longest wavelengths are still sampling the acceleration zone, and thus the velocity in the region where the long wavelength radio emission emanates from will be less than  $30 \text{ km s}^{-1}$ . In Table 6.3 the values are derived assuming the radio emission emanates from the outer atmosphere where the wind has reached its terminal velocity (i.e.,  $30 \text{ km s}^{-1}$ ) and so are just upper limits.

# A

## List of Abbreviations

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**Table A.1:** List of Abbreviations

Abbreviation	Meaning
BIMA	Berkeley Illinois Maryland Association
CARMA	Combined Array for Research in Millimeter-wave Astronomy
CSE	Circumstellar Envelope
e-MERLIN	e-Multi-Element Radio Linked Interferometer Network
FOV	Field of View
HPBW	Half Power Beamwidth
HST	Hubble Space Telescope
IOTA	Infrared Optical Telescope Array
IR	Infrared
IRAM	Institut de Radioastronomie Millimétrique
IUE	International Ultraviolet Explorer
LSR	Local Standard of Rest
OVRO	Owens Valley Radio Observatory
S/N	signal-to-noise
UV	Ultraviolet
VLA	Karl G. Jansky Very Large Array
VLBA	Very Long Baseline Array
VLT	Very Large Telescope

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