# ME136/236: Cheat sheet

## **AERODYNAMICS**

Reynold's number,  $Re := \frac{v_{\infty}l\rho}{\mu}$ , Mach number,  $M := \frac{v_{\infty}}{v_{\text{sound}}}$ , with free-stream air speed  $v_{\infty}$ .

## Aerofoils:

- $f_l = \frac{1}{2}\rho v_\infty^2 cC_l$
- α: aerofoil angle of attack [rad]
- Thin aerofoil model:
  - Lift curve slope:  $\frac{\partial C_l}{\partial \alpha} = 2\pi$
  - $-C_l = C_{l_0} + \frac{\partial C_l}{\partial \alpha} \alpha = C_{l_0} + 2\pi\alpha$

# Finite wings:

- Lift:  $f_L = \frac{1}{2}\rho v_{\infty}^2 SC_L$ :
  - S: Surface area
  - -c: chord (wing width, varies along wing)
  - b: span (wing tip-to-tip distance)

  - AR: aspect ratio,  $AR = \frac{b^2}{S}$  Analytical model for a finite wing, flat plate aerofoil:  $C_L = \frac{AR}{AR+2}C_l$
  - Induced drag:  $f_{D_i} := f_L \sin \alpha_i = \frac{1}{2} \rho v_{\infty}^2 SC_{D_i}$ Analytical model:  $C_{D_i} \approx \frac{C_L^2}{\pi A B}$

**Drag:** Form drag (pressure drag):

$$f_D = \frac{1}{2}\rho v_\infty^2 SC_D$$

**Propellers** Thrust  $f_T$  [N] from scalar speed  $\Omega$  [rads<sup>-1</sup>]  $f_T = C_T \Omega^2$ , Torque:  $\tau = \gamma f_T$ 

Momentum theory: propeller radius  $r_p$ , power consumed P:

$$P = \frac{f_T^{3/2}}{r_p \sqrt{2\pi\rho}}$$

## TENSORS & DYNAMICS

- Displacement of point A w.r.t. point B:  $\underline{s_{AB}}$
- Velocity of point A w.r.t. frame E ( $\overline{B}$  is fixed in E),  $v_A^E := D^E \underline{s_{AB}} : v_A^E$
- Rotation of frame B w.r.t. frame A:  $R^{BA}$
- Ang. velocity of frame B w.r.t. frame  $\overline{A}$ :  $\omega^{BA}$
- (capital letter) Skew symmetric form of  $\omega^{\overline{BA}}$  is  $\Omega^{BA}$ :
- Identity tensor (identity matrix in any coord. sys:) I
- Overbar is 'transpose':  $\overline{x}$

Skew-symmetric form:

if 
$$[\underline{x}]^A = \begin{bmatrix} x_1 \\ x_2 \\ x_3 \end{bmatrix}$$
 then  $[\underline{X}]^A = \begin{bmatrix} 0 & -x_3 & x_2 \\ x_3 & 0 & -x_1 \\ -x_2 & x_1 & 0 \end{bmatrix}$ 

Changing coordinate systems:  $[T]^{AB}$  transforms to A w.r.t. B. **Definition of tensors:** for any coordinate systems A, B:

- First order tensor  $x: [x]^A = [T]^{AB} [x]^B$
- Second order tensor  $\underline{Y}$ :  $[\underline{Y}]^A = [T]^{AB} [\underline{Y}]^B \overline{[T]}^{AB}$

Contraction of indices:

- Displacements:  $\underline{s_{AC}} = \underline{s_{AB}} + \underline{s_{BC}}$
- Transformations:  $[T]^{A\overline{C}} = [T]^{AB}[T]^{BC}$

Length of a tensor:

•  $\|\underline{x}\| = \sqrt{\overline{x}} = \sqrt{\left[\underline{x}\right]^A \left[\underline{x}\right]^A} = \sqrt{x_1^2 + x_2^2 + x_3^2}$  for any coordinate system A, and wherein  $(x_1,x_2,x_3)=[x]^A$ .

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Transformation matrix:

$$[T]^{BA} = \begin{bmatrix} \underline{1}^A \end{bmatrix}^B \quad [\underline{2}^A]^B \quad [\underline{3}^A]^B \end{bmatrix} = \begin{bmatrix} \underline{\overline{[1}^B]}^A \\ \underline{\overline{[2}^B]}^A \end{bmatrix}$$

**Euler angles:** Encode transformation matrix  $[T]^{BE}$  with three parameters:  $[T]^{BE} = [T]^{BY}[T]^{YX}[T]^{XE}$ :

- Yaw (ψ): rotate E about 3<sup>E</sup> = 3<sup>X</sup> to generate X system
  Pitch (θ): rotate X about 2<sup>X</sup> = 2<sup>Y</sup> to generate Y system
  Roll (φ): rotate Y about 1<sup>Y</sup> = 1<sup>B</sup> to generate B system

$$[T]^{BE} = \begin{bmatrix} c\psi c\theta & s\psi c\theta & -s\theta \\ c\psi s\theta s\phi - s\psi c\phi & s\psi s\theta s\phi + c\psi c\phi & c\theta s\phi \\ c\psi s\theta c\phi + s\psi s\phi & s\psi s\theta c\phi - c\psi s\phi & c\theta c\phi \end{bmatrix}$$

**Rotation tensor** For frames A, B:  $1^A = R^{AB}1^B$  (same for  $2^A$  and  $3^A$ ).

$$\left[\underline{R^{BA}}\right]^B\!=\!\left[\underline{R^{BA}}\right]^A\!=\!\overline{\left[T\right]}^{BA}\!=\!\left[T\right]^{AB}$$

Rotational time derivative: (Note:  $\left[\frac{\mathrm{d}}{\mathrm{d}t}x\right]^A \neq \left[T\right]^{AB} \left[\frac{\mathrm{d}}{\mathrm{d}t}x\right]^B$ in general.)

$$\left[\underline{D}^{A}\underline{x}\right]^{B} := \left[\underline{\frac{\mathrm{d}}{\mathrm{d}t}}\underline{x}\right]^{B} + \left[T\right]^{BA} \overline{\left[\underline{\frac{\mathrm{d}}{\mathrm{d}t}}T\right]}^{BA} \left[\underline{\underline{x}}\right]^{B}$$

Preserves  $[\underline{D^C x}]^A = [T]^{AB} [\underline{D^C x}]^B$  for any A, B, C.

When superscripts match:

$$\left[\underline{D}^{A}\underline{x}\right]^{A} = \left[\frac{\mathrm{d}}{\mathrm{d}t}\underline{x}\right]^{A} = \begin{bmatrix} \dot{x}_{1} \\ \dot{x}_{2} \\ \dot{x}_{3} \end{bmatrix} \text{ for } \left[\underline{x}\right]^{A} = \begin{bmatrix} x_{1} \\ x_{2} \\ x_{3} \end{bmatrix}$$

Euler transformation:  $D^A x = D^B x + \Omega^{BA} x$ 

- Translational velocity:  $v_B^A := D^A s_{BA_1}$  where  $A_1$  fixed w.r.t.
- Translational acceleration:  $a_B^A := D^A v_B^A$ .
- Angular velocity:  $\Omega^{BA} := (D^A R^{BA}) \overline{R^{BA}}$  (second order tensor).
  - $-\Omega^{BA} = -\Omega^{AB} = \overline{\Omega^{AB}}$
  - Related first order tensor is  $\omega^{BE}$  (if B is body-fixed, E is earth-fixed):

$$\begin{bmatrix} \underline{\omega}^{BE} \end{bmatrix}^B = \begin{bmatrix} p \\ q \\ r \end{bmatrix}; \quad \begin{bmatrix} \underline{\Omega}^{BE} \end{bmatrix}^B = \begin{bmatrix} 0 & -r & q \\ r & 0 & -p \\ -q & p & 0 \end{bmatrix}$$
 
$$\begin{bmatrix} \frac{\mathrm{d}}{\mathrm{d}t} T \end{bmatrix}^{EB} = [T]^{EB} [\underline{\Omega}^{BE}]^B$$

Time derivative of Euler angles

$$\frac{\mathrm{d}}{\mathrm{d}t} \begin{bmatrix} \phi \\ \theta \\ \psi \end{bmatrix} = \begin{bmatrix} 1 & \sin\phi \tan\theta & \cos\phi \tan\theta \\ 0 & \cos\phi & -\sin\phi \\ 0 & \sin\phi/\cos\theta & \cos\phi/\cos\theta \end{bmatrix} \begin{bmatrix} p \\ q \\ r \end{bmatrix}$$

For small angles only

$$\frac{\mathrm{d}}{\mathrm{d}t}(\phi,\theta,\psi) \approx (p,q,r)$$

Center of mass: That point C for which

$$\sum_{i} m^{i} \underline{s_{P_{i}C}} = 0$$

Translational momentum:

$$\underline{p_C^E} := m^B \underline{v_C^E}$$

Newton's law:

$$D^{I}\underline{p_{C}^{I}} = \underline{f}$$

**Inertia tensor**: Body B, reference point R:  $(S_{P,R})$  is skewsymmetric form of  $s_{P_iR}$ )

$$\underline{J_R^B} := \sum_i m^i \overline{S_{P_i R}} \underline{S_{P_i R}}$$

(See expansion at end of cheat-sheet)

**Angular momentum**: For ref. point equal to center of mass:  $l_B^{BA} = J_B^B \underline{\omega}^{BA}$ 

Moment of force ("torque"):

$$\underline{n_R} = \underline{S_{PR}f}$$

Euler's law: 
$$D^I l_B^{BI} = \underline{n_B}$$

For a rigid body, coordinated in B:

$$\begin{bmatrix} \frac{\mathrm{d}}{\mathrm{d}t} \omega^{BE} \end{bmatrix}^B = \left( \left[ \underline{J}_B^B \right]^B \right)^{-1} \left( \left[ \underline{n}_B \right]^B - \left[ \underline{\Omega}^{BE} \right]^B \left[ \underline{J}_B^B \right]^B \left[ \underline{\omega}^{BE} \right]^B \right)$$

#### SENSORS

**Accelerometer:** Output is  $[\underline{\alpha}]^S$ 

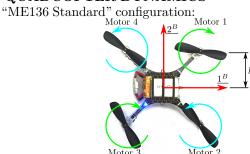
$$\left[\underline{\alpha}\right]^{S} = \left[\underline{a}_{B}^{E}\right]^{S} - \left[\underline{g}\right]^{S}$$

Rate gyroscope: Output is  $\lceil \gamma \rceil^S$ 

$$\left[\underline{\gamma}\right]^S\!=\!\left[\underline{\omega^{BE}}\right]^S$$

$$\begin{bmatrix} \sigma_1 \\ \sigma_2 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix} \left( -\frac{1}{\|s_{FS}\|} \left[ \underline{v_S^E} \right]^S + \left[ \underline{\Omega}^{BE} \right]^S \left[ \underline{3}^S \right]^S \right)$$

### QUADCOPTER DYNAMICS



$$\begin{bmatrix} c_{\Sigma} \\ n_1 \\ n_2 \\ n_3 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 \\ l & -l & -l & l \\ -l & -l & l & l \\ \kappa & -\kappa & \kappa & -\kappa \end{bmatrix} \begin{bmatrix} c_{P_1} \\ c_{P_2} \\ c_{P_3} \\ c_{P_4} \end{bmatrix}$$

$$\begin{bmatrix} c_{P_1} \\ c_{P_2} \\ c_{P_3} \\ c_{P_4} \end{bmatrix} = \frac{1}{4} \begin{bmatrix} 1 & l^{-1} & -l^{-1} & \kappa^{-1} \\ 1 & -l^{-1} & -l^{-1} & \kappa^{-1} \\ 1 & -l^{-1} & l^{-1} & \kappa^{-1} \\ 1 & l^{-1} & l^{-1} & -\kappa^{-1} \end{bmatrix} \begin{bmatrix} c_{\Sigma} \\ n_1 \\ n_2 \\ n_3 \end{bmatrix}$$

## DESCRIPTIONS OF ROTATIONS

Rotation vector, axis-angle: Rotate from frame A about axis  $n^{BA}$  by an angle  $\rho^{BA}$  to generate frame B. Alternatively, rotation vector  $\underline{\lambda}^{BA} := \rho^{BA} \underline{n}^{BA}$ .

$$[T]^{BA} = \exp\left(-\left[\underline{\Lambda}^{BA}\right]^{B}\right)$$

$$\begin{split} &= [\underline{I}] + \left( - \left[\underline{\Lambda}^{BA}\right]^B \right) + \frac{1}{2!} \left( - \left[\underline{\Lambda}^{BA}\right]^B \right)^2 + \frac{1}{3!} \left( - \left[\underline{\Lambda}^{BA}\right]^B \right)^3 \dots \\ &= \cos \rho^{BA} [\underline{I}] + \left( 1 - \cos \rho^{BA} \right) \left[\underline{n}^{BA}\right]^B \overline{\left[\underline{n}^{BA}\right]}^B - \sin \rho^{BA} \left[\underline{N}^{BA}\right]^B \end{split}$$

$$\begin{split} & \textbf{Euler Symmetric Parameters:} \\ & \left\{q\right\}^{BA} := \begin{bmatrix} \cos\left(\frac{1}{2}\rho^{BA}\right) \\ \left[\underline{n^{BA}}\right]^{B} \sin\left(\frac{1}{2}\rho^{BA}\right) \end{bmatrix} =: \begin{bmatrix} q_{0} \\ \left[\underline{q_{v}}\right] \end{bmatrix} = (q_{0},q_{1},q_{2},q_{3}) \\ & \left[T\right]^{BA} = q_{0}^{2}[\underline{I}] - \overline{\left[q_{v}\right]} \, \underline{\left[q_{v}\right]} \, \underline{\left[I\right]} + 2 \underline{\left[q_{v}\right]} \, \overline{\left[q_{v}\right]} - 2q_{0}[\underline{Q_{v}}] \\ & = \begin{bmatrix} q_{0}^{2} + q_{1}^{2} - q_{2}^{2} - q_{3}^{2} & 2(q_{1}q_{2} + q_{0}q_{3}) & 2(q_{1}q_{3} - q_{0}q_{2}) \\ 2(q_{1}q_{2} - q_{0}q_{3}) & q_{0}^{2} - q_{1}^{2} + q_{2}^{2} - q_{3}^{2} & 2(q_{2}q_{3} + q_{0}q_{1}) \\ 2(q_{1}q_{3} + q_{0}q_{2}) & 2(q_{2}q_{3} - q_{0}q_{1}) & q_{0}^{2} - q_{1}^{2} - q_{2}^{2} + q_{3}^{2} \end{bmatrix} \end{split}$$

$$\{q\}^{CA} = \overline{\{Q\}}^{CB} \{q\}^{BA}$$

$$\{Q\}^{BA} = q_0\{I\} + \begin{bmatrix} 0 & \overline{[q_v]} \\ -[q_v] & [Q_v] \end{bmatrix}$$

$$\left\{\frac{\mathrm{d}}{\mathrm{d}t}q\right\}^{BA} = \frac{1}{2}\overline{\{\Omega\}}^{BA}\{q\}^{BA}; \text{ with } \{\omega\}^{BA} = \begin{bmatrix}0\\ \left[\underline{\omega}^{BA}\right]^{B}\end{bmatrix}$$

### CONTROL

Canonical model:  $\dot{x} = f(x,u), x \in \mathbb{R}^n, u \in \mathbb{R}^m$ .

Equilibrium:  $x^*, u^*$  such that  $\dot{x} = 0$ 

First order system:  $\dot{x} = -\frac{1}{\tau}x$ 

Second order, damped system:  $\ddot{x}+2\zeta\omega_n\dot{x}+\omega_n^2x=0$ 

$$\left[ \underline{J_R^B} \right]^B \coloneqq \begin{bmatrix} \sum_i m^i \Big( \big( s_{P_i R}^2 \big)^2 + \big( s_{P_i R}^3 \big)^2 \Big) & - \sum_i m^i s_{P_i R}^1 s_{P_i R}^2 & - \sum_i m^i s_{P_i R}^1 s_{P_i R}^3 \\ - \sum_i m^i s_{P_i R}^1 s_{P_i R}^2 & \sum_i m^i \Big( \big( s_{P_i R}^1 \big)^2 + \big( s_{P_i R}^3 \big)^2 \Big) & - \sum_i m^i s_{P_i R}^2 s_{P_i R}^3 \\ - \sum_i m^i s_{P_i R}^1 s_{P_i R}^3 & - \sum_i m^i s_{P_i R}^2 s_{P_i R}^3 & \sum_i m^i \Big( \big( s_{P_i R}^1 \big)^2 + \big( s_{P_i R}^2 \big)^2 \Big) \end{bmatrix}$$