

## 8.311 Recitation Notes

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## I. INTRODUCTION

Today I will review some of the material that could be on your third exam. The exam will cover the same material covered in problem sets eight through ten.

## II. MAXWELL'S EQUATIONS IN MATERIALS

In materials, Maxwell's equations are:

$$\nabla \cdot \mathbf{D} = \rho_{\text{f}}, \quad (1)$$

$$\nabla \cdot \mathbf{B} = 0, \quad (2)$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad (3)$$

$$\nabla \times \mathbf{H} = \mathbf{J}_{\text{f}} + \frac{\partial \mathbf{D}}{\partial t}, \quad (4)$$

where:

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}, \quad (5)$$

$$\mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M}, \quad (6)$$

and the subscript f denotes free charges and currents. Considering the bound charges and currents:

$$\rho_{\text{b}} = -\nabla \cdot \mathbf{P}, \quad (7)$$

$$\mathbf{J}_{\text{b}} = \nabla \times \mathbf{M} + \frac{\partial \mathbf{P}}{\partial t}, \quad (8)$$

these macroscopic Maxwell's equations reproduce the traditional Maxwell's equations, with:

$$\rho = \rho_{\text{f}} + \rho_{\text{b}}, \quad (9)$$

$$\mathbf{J} = \mathbf{J}_{\text{f}} + \mathbf{J}_{\text{b}}. \quad (10)$$

Often, we consider linear media, where we take:

$$\mathbf{D} = \epsilon \mathbf{E}, \quad (11)$$

$$\mathbf{H} = \frac{1}{\mu} \mathbf{B}. \quad (12)$$

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### III. DISPERSION RELATIONS

As proven in the eleventh recitation notes, the general dispersion relation for a medium is:

$$\det \left( \mathbf{k} \otimes \mathbf{k} - k^2 + \frac{\omega^2}{\epsilon_0 c} \boldsymbol{\epsilon} \right) = 0. \quad (13)$$

As:

$$\boldsymbol{\epsilon} = \epsilon_0 + \frac{i\boldsymbol{\sigma}}{\omega}, \quad (14)$$

this is equivalent to the dispersion relation:

$$\det \left( \mathbf{k} \otimes \mathbf{k} - k^2 + \frac{\omega^2}{c^2} + i\mu_0\omega\boldsymbol{\sigma} \right) = 0. \quad (15)$$

### IV. SNELL'S LAW

Snell's law is the statement that, for a ray of light leaving a medium with an index of refraction  $n_1$  and entering a medium with an index of refraction  $n_2$ , the angles of refraction when taken to the normal are related by:

$$n_1 \sin(\theta_1) = n_2 \sin(\theta_2). \quad (16)$$

### V. BREWSTER'S ANGLE

Brewster's angle  $\theta_B$  is a critical angle such that unpolarized light is reflected back perfectly polarized; it is given by:

$$\theta_B = \arctan \left( \frac{n_2}{n_1} \right). \quad (17)$$

### VI. REFLECTION AND TRANSMISSION

For an electric field with normal polarization to the plane of incidence leaving a medium with an index of refraction  $n_1$  and entering a medium with an index of refraction  $n_2$  at zero angle of incidence, the Fresnel equations give that the amplitude reflection and transmission coefficients are:

$$r = \frac{n_1 - n_2}{n_1 + n_2}, \quad (18)$$

$$t = \frac{2n_1}{n_1 + n_2}. \quad (19)$$

## VII. TIME AVERAGES OF COMPLEX FIELDS

When evaluating a time average of a product—say,  $\langle \mathbf{J} \cdot \mathbf{B} \rangle$ —what one should actually be evaluating is  $\langle \text{Re}(\mathbf{J}) \cdot \text{Re}(\mathbf{B}) \rangle$ , *not*  $\langle \text{Re}(\mathbf{J} \cdot \mathbf{B}) \rangle$ . If one wants to take the real part at the end of the calculation, one must consider the fact that:

$$\langle \text{Re}(\mathbf{J}) \cdot \text{Re}(\mathbf{B}) \rangle = \frac{1}{2} \text{Re}(\mathbf{J} \cdot \mathbf{B}^*). \quad (20)$$

## VIII. APPLICATIONS OF FARADAY'S LAW

The integral form of Faraday's law states that the electromotive force  $\mathcal{E}$  around a loop is given by:

$$\mathcal{E} = -\frac{d\Phi_B}{dt}, \quad (21)$$

where  $\Phi_B$  is the magnetic flux through the surface defined by the loop. For a battery with voltage  $V$  in series with a resistor of resistance  $R$ , the magnitude of the electromotive force between anode of the resistor and the anode of the battery is given by:

$$|\mathcal{E}| = |V - IR|, \quad (22)$$

where  $I$  is the current through the loop.