

Analytic framework

The analytic framework includes equations from the atmospheric fluid dynamics and wind farm fluid dynamics. The equations from atmospheric fluid dynamics concern steady-state, horizontal, large-scale flows under neutral conditions above an infinite homogeneous surface with roughness z_0 [1][2][3]. First, we recall the governing equations of the geostrophic flow, i.e., the momentum equations [4]:

$$-fV_g = -\frac{1}{\rho} \frac{\partial p}{\partial x}, \quad fU_g = -\frac{1}{\rho} \frac{\partial p}{\partial y}, \quad (1)$$

where U_g and V_g are the geostrophic wind components (G is the modulus), p the pressure, ρ the density, and f the Coriolis parameter. The Coriolis parameter is given by $2\Omega \sin \varphi$, where Ω is the rotation rate of the Earth ($7.2921 \cdot 10^{-5}$ rad/s), and φ the latitude. In the Ekman layer, mechanical turbulence becomes important, and the resulting momentum equations then become [5]:

$$-fv = -\frac{1}{\rho} \frac{\partial p}{\partial x} - \frac{\partial(\overline{u'w'})}{\partial z}, \quad fu = -\frac{1}{\rho} \frac{\partial p}{\partial y} - \frac{\partial(\overline{v'w'})}{\partial z}, \quad (2)$$

where u and v are the mean wind components, and $\overline{u'w'}$ and $\overline{v'w'}$ are the turbulent fluxes (also called Reynolds stresses). Substituting the pressure-gradient terms of Eq. 2 with the Coriolis force terms from Eq. 1, the equations of motion can be written as [5]:

$$f(v - V_g) = \frac{\partial(\overline{u'w'})}{\partial z}, \quad f(u - U_g) = -\frac{\partial(\overline{v'w'})}{\partial z}. \quad (3)$$

From these governing equations, Blackadar and Tennekes [1] derived, for the case of large Rossby numbers, $Ro = G/(fz_0)$, two kinds of self-similar solutions, one valid only in the Ekman layer well outside the surface layer and another valid inside the surface layer [3]. By matching those solutions in a region of overlap, they derived the following expressions:

$$\frac{U_g}{u_*} = \frac{1}{\kappa} \ln \left(\frac{u_*}{f z_0} \right) - 4, \quad \frac{V_g}{u_*} = -12, \quad (4)$$

$$\frac{G}{u_*} = \sqrt{\left(\frac{U_g}{u_*} \right)^2 + \left(\frac{V_g}{u_*} \right)^2} = \sqrt{\left(\frac{1}{\kappa} \ln \left(\frac{u_*}{f z_0} \right) - 4 \right)^2 + 12^2}, \quad (5)$$

where u_* is the surface friction velocity and κ the Von Kármán constant (≈ 0.4). The resistant constants (4 and 12) on the right-hand side of Eq. 4 are derived in Refs. [1][3]. These equations relate the large-scale flow (geostrophic flow and Ekman layer) with the small-scale flow (surface layer, where turbines are located).

The equations from wind farm fluid dynamics concern instead steady-state, horizontal flows under neutral conditions within and over an infinite wind farm [6][7]. In this case, the derivation neglects any effect of the Coriolis force and focuses therefore on smaller spatial scales. This analysis provides an equivalent surface roughness, $z_{0,wf}$, for the effect of a large wind farm on the overlying atmospheric boundary layer [6]:

$$z_{0,wf} = z_H \left(1 + \frac{D}{2z_H} \right)^{v_w^*/(1+v_w^*)} \exp \left\{ - \left\{ \frac{c_{ft}}{2\kappa^2} + \left\{ \ln \left[\frac{z_H}{z_0} \left(1 - \frac{D}{2z_H} \right)^{v_w^*/(1+v_w^*)} \right] \right\}^{-2} \right\}^{-0.5} \right\}, \quad (6)$$

where D and z_H are the turbine diameter and hub height, respectively, z_0 the actual surface roughness of the underlying terrain, and

$$v_w^* = \frac{\sqrt{0.5c_{ft}} U_H D}{\kappa u_{*,wf} z_H}, \quad (7)$$

$$c_{ft} = \frac{\pi C_T}{4s_x s_y}, \quad (8)$$

where U_H is the hub-height wind speed, $u_{*,wf}$ the equivalent surface friction velocity generated by the wind farm, C_T is the thrust coefficient provided by the manufacturer (see Supplementary

Fig. 1), and s_x and s_y are the horizontal nondimensional spacings. For the case of aligned layout and installed capacity density of 9 W/m^2 , $s_x = s_y = 1000/D = 6.1$, whereas for the case of staggered layout and installed capacity density of 4.5 W/m^2 , $s_x = s_y = \sqrt{2} \cdot 1000/D = 8.6$.

The hub-height wind speed is calculated according to [6]:

$$U_H = \frac{u_{*,wf}}{\kappa} \ln \left[\frac{z_H}{z_{0,wf}} \left(1 + \frac{D}{2z_H} \right)^{v_w^*/(1+v_w^*)} \right]. \quad (9)$$

Here, we combine Eq. 5 with Eqs. 6-9 to provide a closed-form system that relates the influence of geostrophic wind and Coriolis parameter (latitude-dependent) with the wind farm power density. z_0 and u_* in Eq. 5 are replaced with $z_{0,wf}$ and $u_{*,wf}$ to account for the effect of the wind farm on the Ekman layer. The system represented by Eqs. 5-9, with unknowns $u_{*,wf}$, $z_{0,wf}$, v_w^* , c_{ft} and U_H , can be solved iteratively and provides a value for U_H , with which we calculate the power produced according to the power curve given by the manufacturer.

References

- [1] A. K. Blackadar and H. Tennekes, “Asymptotic Similarity in Neutral Barotropic Planetary Boundary Layers,” *J. Atmos. Sci.*, vol. 25, no. 6, pp. 1015–1020, 1968.
- [2] H. Tennekes and J. L. Lumley, *A first course in turbulence*. MIT press, 1972.
- [3] H. Tennekes, “The Logarithmic Wind Profile,” *J. Atmos. Sci.*, vol. 30, no. 2, pp. 234–238, 1973.
- [4] M. Mak, *Atmospheric Dynamics*. Cambridge University Press, 2011.
- [5] M. Z. Jacobson, *Fundamentals of Atmospheric Modeling*. Cambridge University Press, 2005.
- [6] M. Calaf, C. Meneveau, and J. Meyers, “Large eddy simulation study of fully developed

wind-turbine array boundary layers,” *Phys. Fluids*, vol. 22, no. 1, pp. 1–16, 2010.

- [7] C. Meneveau, “The top-down model of wind farm boundary layers and its applications,” *J. Turbul.*, vol. 13, pp. 1–12, 2012.