Computational Physics Homework 4

Yi-Hsuan Hsu

10/16/2014

1 Problem 1

Newtonian mechanic describes planetary orbit in second order differential equation.

$$\frac{d^2u}{d\phi^2} + u = \frac{GM}{l^2} \tag{1}$$

Solving the equation analytically and numerically. Euler method, two stages Runge-Kutta and four stages Runge-Kutta method are applied. Compare behaviors such as error, convergence rate between each method.

1.1 Theory

(a) prove exact solution of equation (1) is

$$u = \frac{GM}{l^2}(1 + \epsilon \cos\phi) \tag{2}$$

Explicitly, we have homogenious and inhomogenious solution

$$u_0 = \frac{GM}{l^2} \tag{3}$$

$$u_H = A\cos(\phi + \phi_0) \tag{4}$$

where A and ϕ_0 are arbitrary parameter depends on initial condition. We choose $\phi_0 = 0$, and $A = \epsilon \frac{GM}{l^2}$ where

$$\epsilon = \sqrt{1 + 2EL^2/(GMm)^2} \tag{5}$$

hence, the exact solution is u_H plus u_0

$$u(\phi) = u_0 + u_H$$

$$= \frac{GM}{l^2} (1 + \epsilon \cos(\phi))$$
(6)

An ellipse with eccentricity ϵ are obtained.

1.2 Algorithms

1.2.1 Euler Method

Suppose we we want to solve a first order ODE with initial value

$$y'(t) = f(t, y(t)), y(t_0) = y_0$$
(7)

Set step sizeh, one step of Euler method from t_n to t_{n+1} is

$$y_{n+1} = y_n + hf(t_n, y_n) \tag{8}$$

Now derive local truncation error, using Taylor expansion in equation (8),

$$y(t_{n+1} + h) = y(t_n) + hy'(t_n) + \frac{1}{2}y''(t_n) + O(h^3)$$
(9)

Therefore,

$$\delta = y(t_0 + h) - y_1 = \frac{1}{2}h^2y''(t_0) + O(h^3)$$
(10)

which shows that truncation error is proportional to h^2 .

In this homework, we can reduce second order ODE into two first order ODE problem.

$$\frac{du}{d\phi} = v = f_1(u, n, t) \qquad , u(0) = 1 + \epsilon, u'(0) = 0$$

$$\frac{dv}{d\phi} = -u + 1 = f_2(v, n, t) \qquad , v(0) = 0, v'(0) = 0$$
(11)

where we let $\frac{GM}{l^2} = 1$ with proper unit. Therefore, Euler method become

$$u_{n+1} = u_n + h f_1(u_n, v_n, t_n)$$

$$v_{n+1} = v_n + h f_2(u_n, v_n, t_n)$$
(12)

1.2.2 2-order Runge-Kutta Method

2-order Runge-Kutta method are given by formula

$$y_{n+1} = y_n + hf(t_n + \frac{h}{2}), y_n + \frac{h}{2}f(t_n, y_n)$$
(13)

truncation error is proportional to h^2 .

In our practice, the formula become

$$l_{1} = f_{1}(u_{n} + \frac{h}{2}f_{1}(u_{n}, v_{n}, t_{n}), v_{n}, t_{n} + \frac{h}{2})$$

$$k_{1} = f_{2}(u_{n} + \frac{h}{2}f_{1}(u_{n}, v_{n}, t_{n}), v_{n} + \frac{h}{2}f_{2}(u_{n}, v_{n}, t_{n}), t_{n} + \frac{h}{2})$$

$$u_{n+1} = u_{n} + hl_{1}$$

$$v_{n+1} = v_{n} + hk_{1}$$
(14)

1.2.3 4-order Runge-Kutta Method

4-order Runge-Kutta method are given by formula

$$y_{n+1} = y_n + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$

$$k_1 = f(t_n, y_n)$$

$$k_2 = f(t_n + \frac{h}{2}, y_n + \frac{h}{2}k_1)$$

$$k_3 = f(t_n + \frac{h}{2}, y_n + \frac{h}{2}k_2)$$

$$k_4 = f(t_n + h, y_n + hk_3)$$
(15)

truncation error is proportional to h^4 .

Again, in our problem the formula become

$$l_1 = f_1(t_n, u_n, v_n)$$

$$k_1 = f(t_n, u_n + \frac{h}{2}l_1, v_n)$$

$$l_2 = f(t_n + \frac{h}{2}, u_n + \frac{h}{2}l_1, v_n + \frac{h}{2}k_1)$$

$$k_2 = f(t_n + \frac{h}{2}, u_n + \frac{h}{2}l_2, v_n + \frac{h}{2}k_1)$$

$$l_3 = f(t_n + \frac{h}{2}, u_n + \frac{h}{2}l_2, v_n + \frac{h}{2}k_2)$$

$$k_{3} = f(t_{n} + \frac{h}{2}, u_{n} + \frac{h}{2}l_{3}, v_{n} + \frac{h}{2}k_{2})$$

$$l_{4} = f(t_{n} + h, u_{n} + hl_{3}, v_{n} + 2k_{3})$$

$$k_{4} = f(t_{n} + h, u_{n} + hl_{4}, v_{n} + 2k_{3})$$

$$u_{n+1} = u_{n} + \frac{h}{6}(l_{1} + 2l_{2} + 2l_{3} + l_{4})$$

$$v_{n+1} = v_{n} + \frac{h}{6}(k_{1} + 2k_{2} + 2k_{3} + k_{4})$$
(16)

1.3 Sample output

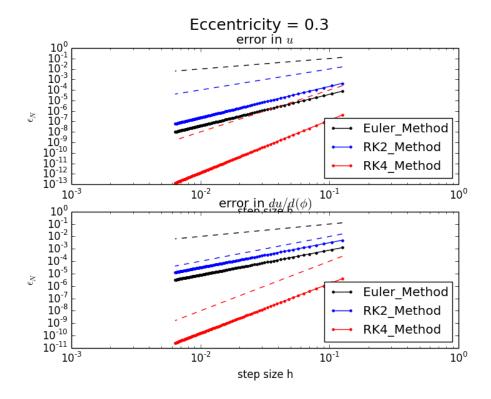


Figure 1: Sample output eccentricity=0.3

From sample output, RK2 and RK4 converge roughly follow the theoritical predict. But, Euler method go faster then it should be.

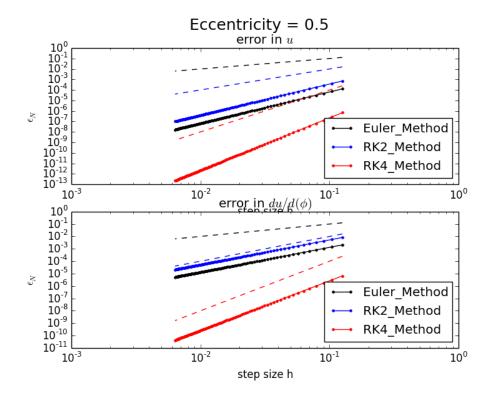


Figure 2: Sample output eccentricity=0.5

2 Problem 2

2.1 Theory

General relativity corrects equation (1) to

$$\frac{d^2u}{d\phi^2} + u = \frac{GM}{l^2} + \frac{3GM}{c^2}u^2 \tag{17}$$

or rewrite into

$$u'' + u = 1 + 3\lambda u^2 \tag{18}$$

from perturbation theory, we assume the solution is homogenious exact solution plus perturbation term

$$u = u_0 + \lambda * f_1(\phi) \tag{19}$$

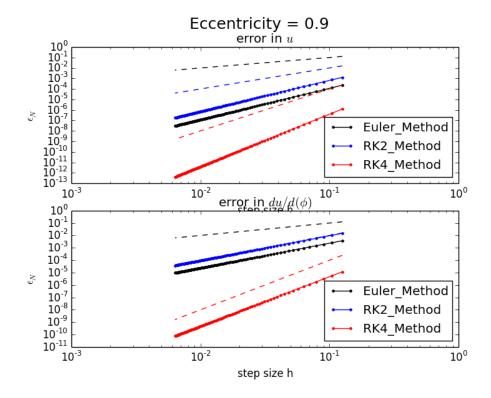


Figure 3: Sample output eccentricity=0.9

where $u_0 = 1 + \epsilon \cos(\phi)$ as we shown before, and λ is a small number. Substitute equation (19) into equation (18), we have

$$\lambda f'' + \lambda f = 3\lambda (1 + \epsilon \cos(\phi))^{2}$$

$$f'' + f = 3(1 + 2\epsilon \cos(\phi)) + O(\epsilon^{2})$$
 (20)

Solving $f(\phi)$, we have

$$f(\phi) = 3(1 + \epsilon \phi \sin(\phi)) \tag{21}$$

Therefore, equation (19) becomes

$$u = 1 + \epsilon \cos(\phi) + 3\lambda(1 + \epsilon\phi\sin(\phi)) \tag{22}$$

for small λ , one can write

$$cos(\phi(1-3\alpha)) = cos\phi cos 3\lambda\phi + sin\phi sin 3\lambda\phi$$

$$\approx cos\phi + 3\lambda\phi sin\phi$$
(23)

$$u \approx 1 + \epsilon \cos[\phi(1 - 3\lambda)] \tag{24}$$

finally,

$$\Delta \phi^{shift} = \frac{2\pi}{1 - 3\lambda} - 2\pi = 6\pi\lambda + O(\lambda^2)$$
 (25)

shows that $\Delta \phi^{shift}$ is proportional to small number of λ , the slope is 6π .

2.2 Algorithm

4-order Runge-Kutta Method were described before. Here we just modified $f_2(u, v, t)$ into

$$\frac{dv}{d\phi} = f_2(u, v, t) = -u + 1 + 3\lambda u^2$$
 (26)

2.3 Sample output

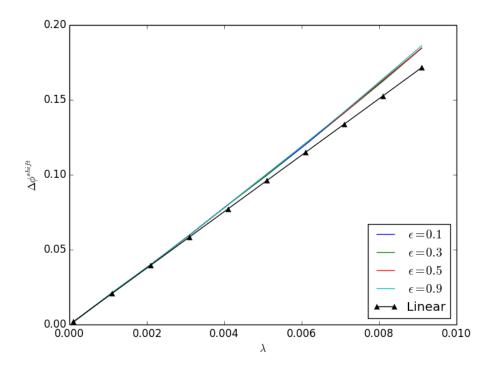


Figure 4: Relationship between λ and $\Delta \phi^{shift}$ follow linear approximation when λ is small. When λ goes larger, the numerical approach behave like quadratic due to higher order term

2.4 problem 2 (c)

From NASA database, the key fact of Mercury

$$G = 6.67 \times 10^{11} m^3 kg^{-1} s^{-2}$$

$$M_{sun} = 1.988 \times 10^{30} kg$$

$$c = 3.0 \times 10^8 ms^{-1}$$

$$R_{orbit} = 46 \times 10^9 m$$

$$v = 58.98 \times 10^3 ms^{-1}$$

$$T_{period} = 87.968 days$$
(27)

Hence

$$\lambda = \frac{GM^2}{lc} = 2.6541 \times 10^{-8} \tag{28}$$

Therefore, $\Delta \phi^{shift}$ perperiod

$$\Delta \phi^{shift} = 6\pi \lambda = 0.103191$$
 (29)

and shift per century are given by

$$\Delta \phi_{century}^{shift} = \Delta \phi^{shift} \frac{100}{T_{period}} = 42.8458$$
 (30)