

# Time-Optimal Attitude Maneuvers

**The goal of this class project is to investigate the application of optimal control theory in attitude dynamics. In particular, we will focus on looking up time-optimal trajectories. Since there is no interest in any particular system, we use the dynamics for simple symmetric rigid-body. The problem of interest is to study the optimal control when the input consists of three controls acting on three orthogonal axis that where some norm of the control is bounded. The expected solution should be of bang-bang type.**

## I. Introduction

Attitude maneuvers are a core problem to the control of aerospace vehicles. As such, understanding and optimizing the problem is a key aspect to the design of better performing vehicles. The performance index to be optimized will depend on the specific application and vehicle considered. For high performance aircraft and satellites, a possible performance index is the time used for a rest-to-rest maneuver. This has direct applications to satellite re-targeting or high-acceleration maneuvers for aircraft allowing for faster redirecting of the aircraft path.

A survey on time-optimal maneuvers is provided by Scrivener [1]. Optimal maneuvers for different configurations: number of angular velocity controls, rigid and flexible bodies, and rotating around the eigenaxis are discussed.

Here, since we aim to achieve insight on maneuvers for both spacecraft and aircraft we will focus on rigid-body maneuvers under bounded controls. For large angle reorientations, introducing quaternions will simplify the problem of computing the difference between attitudes, besides avoiding the singularities of Euler angles, as proposed by Wie and Weiss [2]. In this case they introduce a regulator that follows the eigenaxis rotation between initial and final state which is the minimum action path. Chowdhry [3] solves the time-optimization of the attitude stabilization (i.e. drive angular velocity to 0) for different configurations where one angular rate control is not available for its application to super-maneuverable aircraft.

The problem is interesting from a theoretical point of view since it involves a number of different tools. The attitude equations of motion are highly nonlinear, and without any assumptions result in complex motion. Attitude itself calls for the use of quaternions or similar formulations since Euler angle strategies may become badly-defined and little intuitive. Finally, optimal control theory is necessary in order to achieve the necessary conditions for the time optimality of the control as stated by Longuski [4]. Besides, the final control is not readily intuitive: eigenaxis maneuvers provide the minimum length path, however aircraft tend to combine rolling and pitching motion.

In particular, we will follow the approach used by Bilimoria [5] and Fleming [6] but using a slightly different notation.

An axis-symmetric rigid-body will be considered where all three angular velocities can be controlled independently and where the control is norm by some norm. It would also be interesting to follow some of the approaches used by Byers [7] to look for real-time solving capabilities for the control so that it could be flyable.

## II. Analysis

### A. Problem formulation

As covered by Bilimoria the equations of motion for the rotation around principal axes of inertia are:

$$I_x \dot{\Omega}_x + (I_z - I_y) \Omega_z \Omega_y = \tau_x \quad (1)$$

$$I_y \dot{\Omega}_y + (I_x - I_z) \Omega_x \Omega_z = \tau_y \quad (2)$$

$$I_z \dot{\Omega}_z + (I_y - I_x) \Omega_y \Omega_x = \tau_z \quad (3)$$

For simplicity, we will assume that the vehicle is symmetric and therefore all three moments of inertia are identical and equal to  $I$ . Under those conditions the equations of motion decouple into:

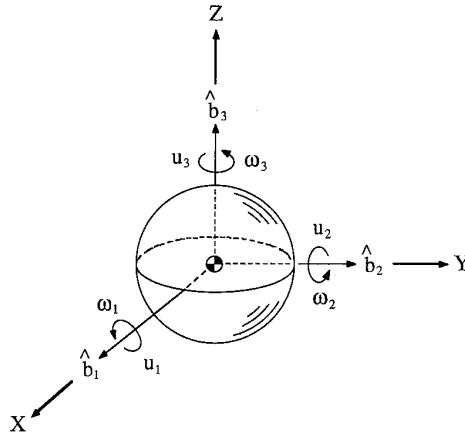
$$I \dot{\Omega}_x = \tau_x \quad (4)$$

$$I \dot{\Omega}_y = \tau_y \quad (5)$$

$$I \dot{\Omega}_z = \tau_z \quad (6)$$

Figure 1 offers a depiction of such a system. The equations of motion above can be expressed in vector notation in a more compact manner:

$$I \vec{\Omega} = \vec{\tau} \quad (7)$$



**Fig. 1 Symmetric Body with controls along three axes. Original Source: Bilimoria [5]**

And the control is bounded by:

$$|\tau_x|^q + |\tau_y|^q + |\tau_z|^q \leq \tau_{max}^q \quad (8)$$

Where  $q$  is the norm of interest; this is a general way of bounding the control. The most interesting norms will be the 2-norm and the infinity-norm ( $\infty$ -norm) corresponding to the maximum applied torque being independent of the direction, and to three independently bounded controls by  $\tau_{max}$  respectively.

We can consider all three controls to be equally bounded by  $\tau_{max}$  we can non-dimensionalize the equations of motion using  $u_i = \tau_i / \tau_{max}$  and a characteristic time  $t = \sqrt{I / \tau_{max}}$  so the characteristic angular velocity  $\omega_i = \Omega_i \sqrt{I / \tau_{max}}$ . The vector equation of motion now simplifies to its final form:

$$\vec{\omega} = \vec{u} \quad (9)$$

Beside the equations relating velocity to the control torques we need the equations of motion for the attitude quaternions. A description on how quaternions relate to euler angles and angular velocities can be found in Parwana and Kothari[8] and useful relations for quaternion derivatives are provided by Graf [9]. For more details on the quaternion convention, or the derivation and partial differentiation we refer to the Appendix A. Our attitude quaternion is  $q$  and we introduce an angular velocity quaternion  $\omega = [0 \quad \vec{\omega}]^T$

$$\dot{q} = 1/2 \, q \circ \omega \quad (10)$$

Now we are ready to state our optimization problem. Our performance index is the final time, which we will minimize, in other words:

$$\min J = t_f \quad (11)$$

Subject to the vector process equations 9 and 10. Our control, by the definition of the non-dimensionalization, is subject to  $|u_i| \leq 1$ . The initial and final states are fixed and defined by the maneuver to be optimized:

$$q(0) = [1 \quad 0 \quad 0 \quad 0]^T \quad (12)$$

$$q(t_f) = q_f \quad (13)$$

$$\vec{\omega}(0) = \vec{\omega}(t_f) = 0 \quad (14)$$

The only free boundary is the final time  $t_f$  which is the performance index to be minimized.

## B. Hamiltonian and co-state equations

With the formulation for the equations of motion we define the Hamiltonian for the motion:

$$H(x, u, \lambda) = \frac{1}{2} \lambda_q^T (q \circ \omega) + \lambda_\omega^T \vec{u} \quad (15)$$

Where  $\lambda_q$  and  $\lambda_\omega$  are the co-state vectors with a size of 4 elements and 3 elements respectively. We realize that, since there is no explicit time dependence, the Hamiltonian is constant throughout the motion.

In order to get the co-state equations we get apply the Euler-Lagrange equations and find the appropriate derivatives. In order to simplify the resulting expressions Matrix calculus is used to come up with more compact expressions, see Appendix A for details on the derivation. The expression for the attitude co-states is:

$$\dot{\lambda}_q = -\frac{\partial H}{\partial q} = \frac{1}{2} \lambda_q \circ \omega \quad (16)$$

The partial derivative for  $\dot{\lambda}_\omega$  depends on omega which appears only in the quaternion product, but it cannot be obtained straight forward using the expressions for quaternion derivatives since  $\dot{\lambda}_\omega$  is not a 4 element vector. However, if we extend it to a 4 element vector where the first element is forced to be 0, we can find the partial derivative in terms of the imaginary part of the quaternion derivative.

$$\dot{\lambda}_\omega = -\frac{\partial H}{\partial \omega} = -\frac{1}{2} \text{Im}(q' \circ \lambda_q) \quad (17)$$

We now look at the boundary conditions. In the problem statement we have stated that the final and initial states are fixed, also the initial time is fixed. Since the only free state is the final time  $t_f$  we look for the associated condition in the transversality condition:

$$H_f dt_f - \lambda^T dx_f + dg = H_f dt_f - \lambda^T dx_f + dt_f = (H_f + 1) dt_f \implies H_f + 1 = 0 \quad (18)$$

Also, since the Hamiltonian is constant we have that throughout the motion:

$$H = -1 \quad (19)$$

## C. Derivation of the optimal control

The control logic is that of a bang-bang controller, which can be expected for a minimum-time type strategy with some modifications. Because of this, we have the possibility of having singular sub-arcs. Later we will have to study the possible consequences on the control.

In order to find the optimal control we recognize that we have an inequality constraint in the control of the following form:

$$1 = |u_x|^q + |u_y|^q + |u_z|^q \quad (20)$$

This means that  $H_u^* = 0$  cannot be directly satisfied. This bound on the control does is not a switching function neither, so we have to apply Pontryagin's Minimum Principle in its most basic form. We have to look for the optimal control such that:

$$H(x^*, u^*, \lambda^*) \leq H(X^*, u, \lambda^*) \quad (21)$$

Or equivalently for our Hamiltonian:

$$\lambda_{\omega}^{*T} \vec{u}^* \leq \lambda_{\omega}^{*T} \vec{u} \quad (22)$$

Since the expression above is linear in terms of  $\vec{u}$  the minimum will necessarily lie on the inequality constraint of the norm of  $\vec{u}$ . We introduce the constraint in the expression above by solving for  $u_z$  in terms of  $u_x$  and  $u_y$  and the sign of  $u_z$  which we don't know yet:

$$u_z = \text{sgn}(u_z) (1 - |u_x|^q - |u_y|^q)^{\frac{1}{q}} \quad (23)$$

Now we have to look for a minimum of:

$$\mathcal{L} = \lambda_{\omega_x} u_x + \lambda_{\omega_y} u_y + \text{sgn}(u_z) \lambda_{\omega_z} (1 - |u_x|^q - |u_y|^q)^{\frac{1}{q}} \quad (24)$$

In order to find the minimum we take partial derivatives with respect  $u_x$  and  $u_y$ . In the derivatives we identify the expression for  $|u_z|$  so we substitute it back:

$$\frac{\partial \mathcal{L}}{\partial u_x} = 0 = \lambda_{\omega_x} - \text{sgn}(u_x u_z) \lambda_{\omega_z} (1 - |u_x|^q - |u_y|^q)^{\frac{1-q}{q}} |u_x|^{q-1} = \lambda_{\omega_x} - \text{sgn}(u_x u_z) \lambda_{\omega_z} |u_z|^{1-q} |u_x|^{q-1} \quad (25)$$

$$\frac{\partial \mathcal{L}}{\partial u_y} = 0 = \lambda_{\omega_y} - \text{sgn}(u_y u_z) \lambda_{\omega_z} (1 - |u_x|^q - |u_y|^q)^{\frac{1-q}{q}} |u_y|^{q-1} = \lambda_{\omega_y} - \text{sgn}(u_y u_z) \lambda_{\omega_z} |u_z|^{1-q} |u_y|^{q-1} \quad (26)$$

And on the expressions above we solve for  $u_x$  or  $u_y$ :

$$\lambda_{\omega_{x,y}} - \text{sgn}(u_{x,y} u_z) \lambda_{\omega_z} |u_z|^{1-q} |u_{x,y}|^{q-1} = 0 \quad (27)$$

$$\implies \text{sgn}(u_{x,y}) |u_{x,y}|^{q-1} = \frac{\lambda_{\omega_{x,y}}}{\lambda_{\omega_z}} \text{sgn}(u_z) |u_z|^{q-1} \quad (28)$$

$$\implies \text{sgn}(u_{x,y}) |u_{x,y}| = \text{sgn} \left( \frac{\lambda_{\omega_{x,y}}}{\lambda_{\omega_z}} \right) \left| \frac{\lambda_{\omega_{x,y}}}{\lambda_{\omega_z}} \right|^{\frac{1}{q-1}} \text{sgn}(u_z) |u_z| \quad (29)$$

$$= u_{x,y} = \text{sgn} \left( \frac{\lambda_{\omega_{x,y}}}{\lambda_{\omega_z}} \right) \left| \frac{\lambda_{\omega_{x,y}}}{\lambda_{\omega_z}} \right|^{\frac{1}{q-1}} u_z \quad (30)$$

And substituting these expressions in the expression we obtained in 23 we obtain:

$$|u_z| = \left( \frac{|u_z|^q}{|u_z|^q} \left| \frac{\lambda_{\omega_z}}{\lambda_{\omega_z}} \right|^{\frac{q}{q-1}} |u_z|^q - \left| \frac{\lambda_{\omega_x}}{\lambda_{\omega_z}} \right|^{\frac{q}{q-1}} |u_z|^q - \left| \frac{\lambda_{\omega_y}}{\lambda_{\omega_z}} \right|^{\frac{q}{q-1}} |u_z|^q \right)^{\frac{1}{q}} \quad (31)$$

$$\Rightarrow |\lambda_{\omega_z}|^{\frac{q}{q-1}} = \frac{1}{|u_z|^q} |\lambda_{\omega_z}|^{\frac{q}{q-1}} - |\lambda_{\omega_x}|^{\frac{q}{q-1}} - |\lambda_{\omega_y}|^{\frac{q}{q-1}} \quad (32)$$

$$\Rightarrow \frac{1}{|\lambda_{\omega_z}|^{\frac{q}{q-1}}} |\lambda_{\omega_z}|^{\frac{q}{q-1}} + |\lambda_{\omega_x}|^{\frac{q}{q-1}} + |\lambda_{\omega_y}|^{\frac{q}{q-1}} = \frac{1}{|u_z|^q} \quad (33)$$

$$\Rightarrow |u_z| = \frac{|\lambda_{\omega_z}|^{\frac{1}{q-1}}}{\left[ |\lambda_{\omega_x}|^{\frac{q}{q-1}} + |\lambda_{\omega_y}|^{\frac{q}{q-1}} + |\lambda_{\omega_z}|^{\frac{q}{q-1}} \right]^{\frac{1}{q}}} \quad (34)$$

These last expressions define the magnitude of each control input as function of the magnitude of the co-states, but we still don't know how to find the sign of  $u_z$ . However to pick the right sign we just need to recall that our goal is to minimize the Hamiltonian, therefore we have to pick the sign with the opposite sign to  $\lambda_{\omega_z}$ .

Once we have found the right  $u_z$  we can solve for  $u_{x,y}$ , using 30. The expression for all three controls becomes similar, as it could be expected:

$$u_i = \frac{-\text{sgn}(\lambda_{\omega_i}) |\lambda_{\omega_i}|^{\frac{1}{q-1}}}{\left[ |\lambda_{\omega_x}|^{\frac{q}{q-1}} + |\lambda_{\omega_y}|^{\frac{q}{q-1}} + |\lambda_{\omega_z}|^{\frac{q}{q-1}} \right]^{\frac{1}{q}}} \quad (35)$$

Finally, we have to consider the possibility of singular control. In order to rule it out, we recall the  $H = -1$  condition obtained from the transversality condition. This allows us to guarantee that there will be non singular sub-arcs. Singular subarcs will only happen if  $\lambda_{\omega} = 0$  for a finite interval, this implies that  $\dot{\lambda}_{\omega} = 0$ , but since  $q$  is an attitude quaternion,  $\dot{\lambda}_{\omega} = 0$  will only happen if  $\lambda_q = 0$ . This cannot happen at the same time that  $\lambda_{\omega} = 0$  since that would make the Hamiltonian vanish, contradicting the transversality condition. Since there are no singular sub-arcs the controller above satisfies the necessary conditions for an optimal control.

The controller structure is fundamentally of a bang-bang nature as anticipated. There is a component controlling the amplitude of the control to the different actuators, but on top of this amplitude allocation, the sign is controlled in the same fashion as a usual bang-bang control. Choosing a norm will determine the behavior of the control amplitude allocation, allowing for more independence between the controls that will lead to more and more complex control profiles and possible faster maneuvers since the set of possible controls increases with the norm.

In the limit, where the norm goes to infinity the controls become completely independent with the form of switching functions. Each control component must satisfy  $|u_i| \leq 1$ . Then from the limit of the control in 35, or following the usual reasoning for switching functions minimizing the Hamiltonian we get:

$$u_i = -\text{sgn}(\lambda_{\omega_i}) \quad (36)$$

### III. Numerical Solution

Once the TPBVP has been defined, the next step is to solve it. The discussion of the numerical solution will be split in three different topics, the first one will discuss the general strategy used for the problem and the statement of the boundary conditions, the second one will cover the initial guesses used for the solver, the third one considers the possible methods used for the solution and the advantages and disadvantages of each one them.

#### A. Statement of the constraints

The problem to be solved right now can be thought of a feasibility problem, where we satisfy the constraints of the TPBVP. As such, we will resort to numerical methods for its solution, but we should try to reduce the dimensionality of the problem and improve its conditioning as much as possible.

The first guess for the numerical solver will be obtained from the numerical integration of some initial conditions. Since we know the initial values for the states, these will be fixed and will not be allowed to change, this helps to reduce the dimensionality of the problem by 7 variables. In the most basic form of the solution, we have reduced the problem to finding the initial values for the 7 co-states.

Also, we know that the Hamiltonian is constant throughout the motion and  $H = -1$ . The information that this constraint provides is merely in the terms of scaling of the co-states. However, for our particular problem the scaling of the co-states is not critical since we use them in three ways: modulating amplitude (where the relative values is the used information), in the switching function (where only the sign matters) and in taking quaternion derivatives (which are preserved with scale changes). This means that in order to avoid ill-conditioned numerical problems we can reformulate the final Hamiltonian constraint as:

$$|\lambda_{q_f}| = 1 \quad (37)$$

But recalling the costate equation for  $\dot{\lambda}_q$  we can observe that it is just the rotation in time of  $\lambda_q$  under the angular velocity  $\vec{\omega}$ . This means that the magnitude of  $\lambda_q$  stays constant in time, and therefore we check the magnitude at the initial conditions, where it is more convenient to check for our initial guess. Therefore the transversality condition is substituted by:

$$|\lambda_{q_0}| = 1 \quad (38)$$

Other than this condition, we will also check for state constraints at the final time. Since the final time is variable, it is also another variable to be found in our constraints feasibility problem solution. We have a total of 8 constraints, 7 states and the just found condition on the module of  $\lambda_q$ , and 8 free variables, the initial states and the final time, so the problem is well defined.

## B. Initial guesses and continuation strategy

After defining the constraints to be satisfied, the next step is how to come up with initial guesses for the solver. All solvers will work by solving an approximated problem, therefore we must come up with an initial guess close enough to the final solution in order for the solver to converge.

A possible strategy is to use a continuation on the norm. If the solution for a specific norm is known, its initial co-states can be used as a first guess for the problem of a neighboring norm. However, in any case a initial solution for a particular norm must be known. In the next paragraph we come up with an initial guess for the 2-norm controller, so for solving higher norm controllers we will continue the norm up to our desired final norm.

In the case of the 2-norm control, the initial solution can be found analytically with the following thought process. This solution is not meant to be formal since it is only used as an initial guess. The 2-norm control has no preferred direction of control since all controls are embedded in a sphere, therefore the optimal control will be the one that uses the shortest path. This shortest path will correspond to the rotation about the eigenaxis that connects the initial and final states. Since the initial quaternion is fixed to  $q_0 = [1 \ 0 \ 0 \ 0]^T$ , the final quaternion provides the angle  $\theta$  to be rotated and the eigenaxis  $\vec{n}$  information in the angle-axis representation of the quaternion:

$$q(t) = \left( \cos \frac{\theta}{2}, \sin \frac{\theta}{2} \vec{n} \right) \quad (39)$$

In order to consider the time evolution of the rotation, we can consider the problem as a 1-d rotation. Here, the rest-to-rest maneuver will follow a triangular angular velocity profile with one switch of the control exactly at the half of the path. The time to the control switch since the beginning of the maneuver can be found as:

$$\frac{1}{2} t_{sw}^2 = \frac{\theta_f}{2} \implies t_f = 2t_{sw} = 2\sqrt{\theta_f} \quad (40)$$

The angular velocity must be parallel to the eigenaxis at all moments, therefore  $\vec{u}$  must also be parallel. The switching must happen in all three components at the same time so they must be equal to zero at  $t_f/2$ . We previously said that  $\lambda_q$  rotates in time with  $\omega$  so it must follow also follow an eigenaxis rotation. We can express it as the composition of the rotation of  $q$  and the initial offset:

$$\lambda_q = q(t)\lambda_{q_0} \quad (41)$$

And then for  $\lambda_\omega$  we get:

$$\dot{\lambda}_\omega = -\frac{1}{2} \text{Im}(q(t)' \circ q(t) \circ \lambda_{q_0}) = -\frac{1}{2} \text{Im}(\lambda_{q_0}) \quad (42)$$

So from the expression above,  $\lambda_\omega(t)$  is linear with time and knowing that  $\lambda_\omega$  must be parallel to the rotation axis we can write the expression below. For that we use the fact that  $\lambda_\omega(0)$  should oppose the rotation axis for the input to be



aligned with it and positive (it is a switching function):

$$\lambda_\omega(0) = \frac{1}{2}\text{Im}(\lambda_{q_0})(t_f/2) = -k\vec{n}$$

So from this expression we get that:

$$\lambda_{q_0} = -[0 \quad \vec{n}]^T \quad (43)$$

$$\lambda_{\omega_0} = -\frac{t_f}{4}\vec{n} \quad (44)$$

### C. Solving methods

Different solvers for the feasibility problem were considered . The options are summarized below with a discussion, with their advantages and disadvantages, on their applicability to our problem of interest

#### 1. Collocation methods

They are based on discretization of the motion into several intervals where the motion is integrated usually in an implicit way [10]. These methods are characterized by their efficiency but lack the accuracy of other methods. MATLAB's `bvp4c` was the first option to be tested due to its familiarity.

Its first problem comes from the collocation of the switching and signum functions which cause the Jacobian to be singular. This was solved by approximating the signum function by a hyperbolic tangent so the continuation strategy would increase the norm and also the accuracy of the hyperbolic tangent approximation.

They are able to converge to valid solutions for low order norms but they were not able to converge with the desired accuracy as the continuation increased.

#### 2. Multiple shooting method with Newton Iterations

A shooting method consists on running solving the feasibility problem by obtaining a first order approximation of how changes in the initial state around some reference guess change the final conditions. This leads to a Newton-Raphson iteration scheme where the errors and their Jacobian are used to compute the new initial conditions until the error vanishes.

The Jacobian approximation for changes in the initial conditions is only valid for an interval around the initial guess depending on the sensitivity of the problem. In order to minimize this sensitivity multiple shooting is used[11]. It decreases the sensitivity by splitting the time interval into several intervals. On each of these intervals a shooting is performed. Since the number of states has increased, you have a initial state guess at the beginning of each shooting, extra constraints have to be added. These constraints ensure the continuity between shootings so that all different

shootings are consistent. This decreases the sensitivity and increases the robustness of the method at the expense of more design constraints and increased computational cost. The preferred number of shooting intervals was set to 3.

In order to apply this method an integration scheme has to be used. Here MATLAB's ode 45 Runge-Kutta integrator was used with tight integration constraints to prevent numerical error from interfering in the computation of the Jacobian matrix. Specifically the relative tolerance was set to  $10^{-11}$  and the absolute tolerance to  $10^{-14}$ .

This method works as long as the Jacobian approximation works. It seemed to have convergence problems as the norm increased. It may have been solved by decreasing the step in the continuation procedure of the norm.

### 3. Multiple shooting method with *fmincon*

This method is based on the previous method. The integration and constraint formulation to ensure continuity between different shootings is formulated in the same way. The difference lies in the solver used to find the new initial guess after each iteration.

We used MATLAB's *fmincon*, a convex optimizer that allows for constraint handling directly. Therefore, instead of trying to minimize the constraint functions in the way of an augmented Lagrangian objective function, we let *fmincon* handle the constraints. This solver introduces methods for satisfying the constraint that go beyond the Newton's method we were using before. This leads to a greater radius of convergence.

The objective function to be minimized was time, and the constraints were the ones already stated. However, we are not particularly interested in minimizing time: we already know that the solution will be time optimal if the constraints are satisfied. Optimization parameters are chosen according to this idea. The tolerance for the constraints is set to  $10^{-4}$  and for the objective function optimality the tolerance is set to 1 since we are not really concerned about it.

This approach was the one that allowed for successful computation of the optimal control. The parameters for the integration are the same as those above. Results discussed below were found using this method.

## IV. Results

Results were obtained for an arbitrary maneuver. The initial state is neutral attitude: roll, pitch and yaw all equal to zero. The final state are the following roll, pitch and yaw angles (in that order):

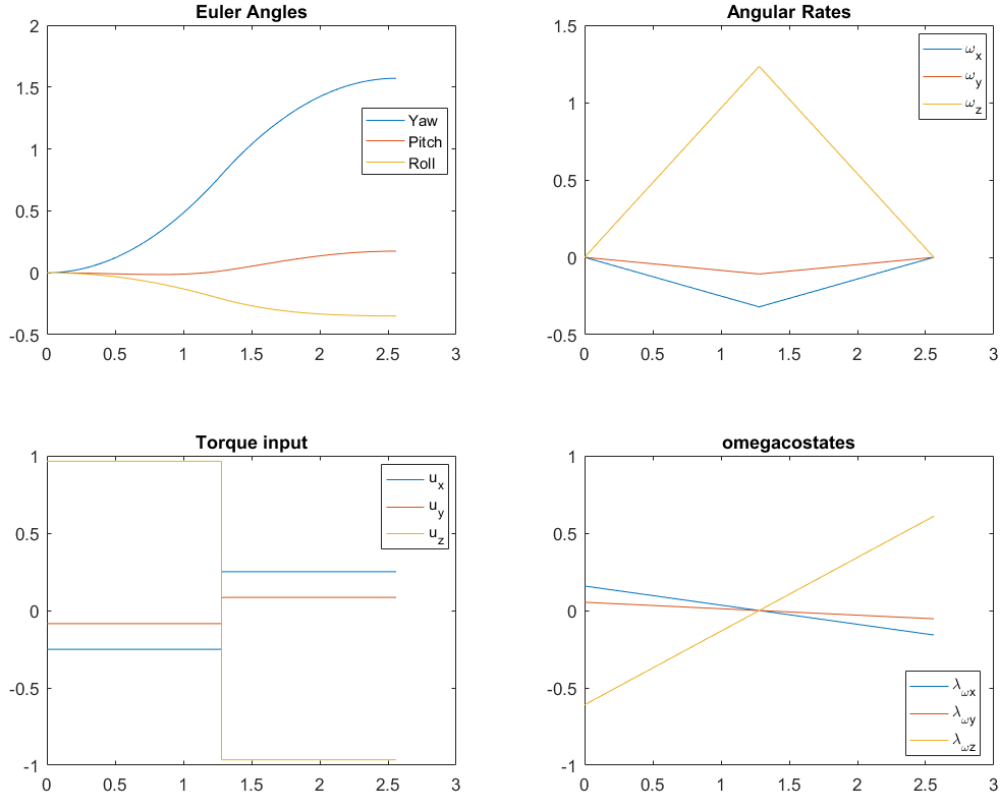
$$\phi = -20^\circ \quad \theta = 10^\circ \quad \psi = 90^\circ \quad (45)$$

We are particularly interested in two of the norms. The 2-norm and  $\infty$ -norm. The 2-norm shows the maneuver for the shortest path, it can be computed in real time using the approximation provided in section IIIc, and works as a reference solution. The  $\infty$ -norm uses three independent bounds, resembles more closely the limits of an attitude control system, and offers the possibility for unexplored more optimal solutions than those found by the 2-norm.

We start by discussing in detail the solution for the 2-norm solution. Figure 2 shows the main characteristics of the maneuver. Time-history is provided for the main parameters of interest: the states, attitude (in Euler angles) and angular velocity, the applied torques and the co-states that determine the input. The co-states of the attitude are not represented since they do not offer any additional insight. The initial co-states for the maneuver are:

$$\lambda_{q_0} = [-0.0006 \quad 0.2474 \quad 0.0833 \quad -0.9532]^T \quad \lambda_{\omega_0} = [0.1583 \quad 0.0533 \quad -0.6093] \quad (46)$$

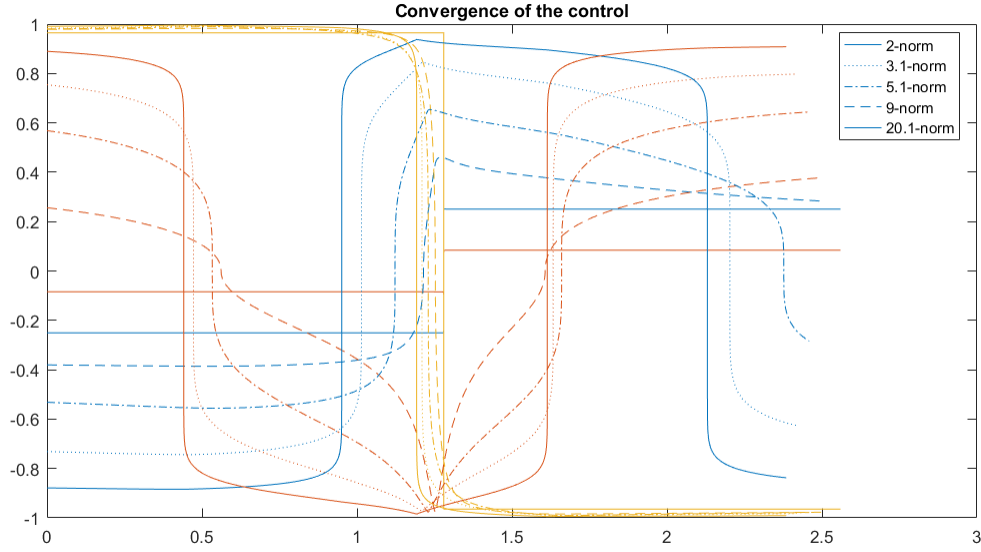
The total time used for the maneuver is 2.5596 non-dimensional time units and the amplitude of total rotated angle is 1.6378 radians. A closer look to the graphics reveals that our initial guess for the solution is indeed correct. The  $\lambda_{\omega}$  behave linearly with time and being equal to 0 at exactly the half of the motion where the inputs reverse. The amplitude of these is set so that they are aligned with the eigenaxis of the rotation.



**Fig. 2** Solution, inputs and co-states for the switching function for the 2-norm optimal maneuver.

We now proceed to discuss the continuation process to the infinity norm. This was done by increasing the norm of the controller a about a 10 percent each iteration up to a norm of 20. Figure 3 shows some selected norms in the continuation process that show the evolution of the controller. As the norm of the controls increase the coupling between

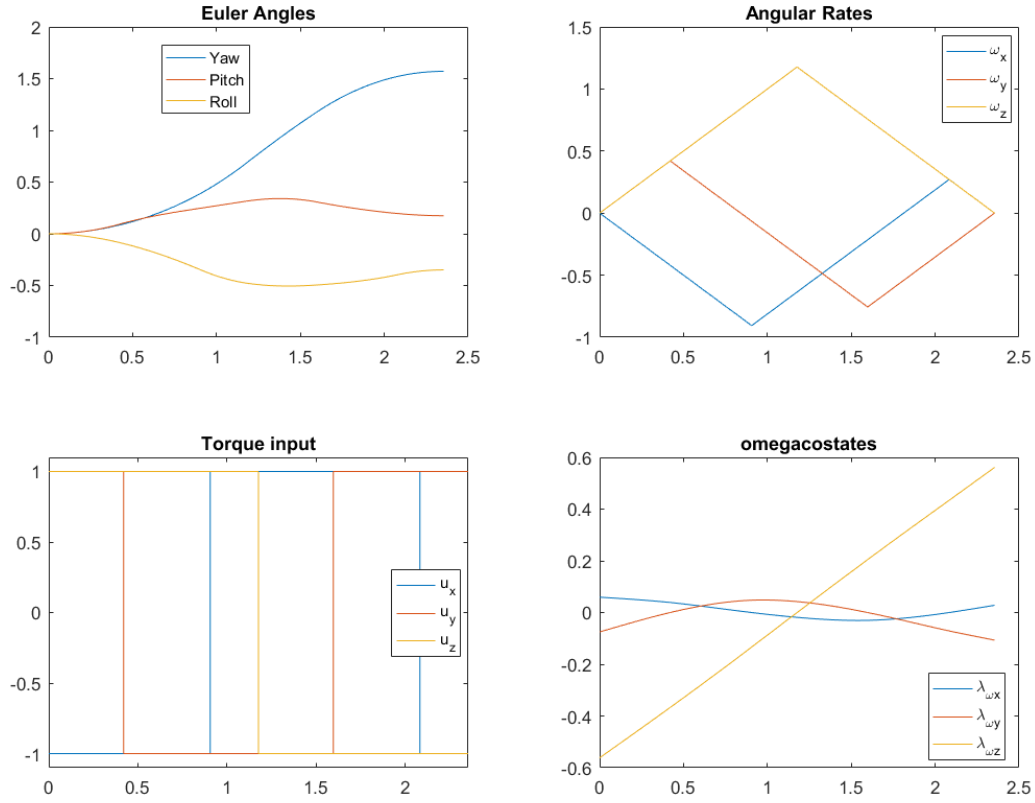
axes becomes more and more relaxed, for the  $\infty$ -norm they should be three independent controls bounds so the all three controls should lie on the  $\pm 1$  bound.



**Fig. 3** Three axes controls for different norms ranging from 2 to 20. As the norm increases the three controls become more and more decoupled. Reproduced from Bilimoria [5].

Once the 20-norm controller is found, it is used as the starting solution for the  $\infty$ -norm controller implemented as in equation 36. The three controls become completely independent, the characteristics for the new control are plotted in figure 4 in a similar fashion to those of the 2-norm control. The initial value for the co-states is:

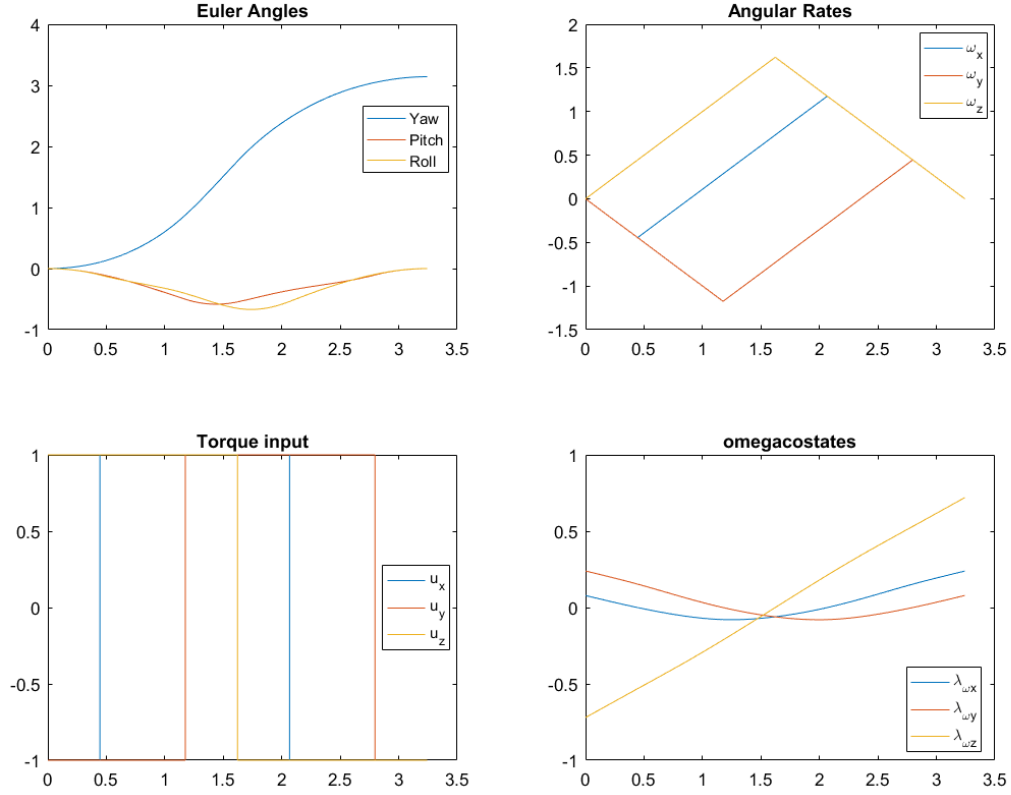
$$\lambda_{q_0} = [-0.0405 \quad 0.0809 \quad -0.3818 \quad -0.9179]^T \quad \lambda_{\omega_0} = [0.0594 \quad -0.0751 \quad -.5617] \quad (47)$$



**Fig. 4** Solution, inputs and costates for the switching function for the  $\infty$ -norm optimal maneuver.

The maneuver time has decreased to 2.3540. A decrease above 8% in the maneuver time. The controls have the form of independent switching functions as it could be expected. This means that the angular velocity is also composed of linear intervals, this implies that the trajectory will be made up of analytically solvable intervals. A phase-space approach to the problem may offer additional insight on how to choose the switches and define a global strategy, but it will definitely be complex due the high number of states. The shape for the co-states of the switching function has become more complex. The results are coherent with those found by Bilimoria[5] in terms of the shape of the switching function and the observed reduction in time.

Finally, in order to attempt to reproduce the results from Bilimoria [5], a 180 degree in yaw maneuver was attempted. While this maneuver is the largest angle maneuver, and possibly will show the most benefit from our control strategy, the control to achieve it is not unique by the rotational symmetry of the problem. In this last piece of the results we attempted to obtain a control similar to the one he provides in Figure 6 in [5]. The angular velocity history obtained here is provided in figure 5. The results are identical, other than a mirrored  $\omega_y$  due to the mentioned non-uniqueness of optimal control, the optimal time for the trajectory is 3.243 non-dimensional units same as the paper states.



**Fig. 5 180 yaw degree maneuver. Reproduced from Bilimoria[5]**

It is interesting to note that the maneuver time has decreased with respect to the eigenaxis maneuver even though we are following a longer path. This is because the controller is reorienting the body so that a torque greater than any individual limit is applied in the desired rotation of motion from an inertial perspective as pointed out by Bilimoria[5].

## V. Conclusions

The time-optimal maneuver has been studied for the two most relevant cases: the 2-norm and the  $\infty$ -norm bounded controls. The conclusions from this work can be grouped in three main categories: optimality of the control, those regarding the solution, and finally the applicability and insight obtained from the results.

About the control we can state that:

- 1) First order necessary conditions for an arbitrary norm control have been found using Pontryagin's Minimum Principle. The resulting control has the form of a modified switching function.
- 2) For the  $\infty$ -norm controller, the resulting controls are independent switching functions.
- 3) The existence of singular sub-arcs for the optimal solution has been ruled out.

- 4) The trajectory between switches in the control can be solved analytically. A phase-space like approach could offer new insight.

When it comes to solution of the system, the following can be concluded:

- 1) The condition arising from the transversality condition can be rewritten to improve numerical conditioning in this particular problem.
- 2) An analytical solution was found for the 2-norm controller or eigenaxis maneuver.
- 3) A continuation strategy based on the norm can be used to find any specific norm. The  $\infty$ -norm was found following this procedure starting in the 2-norm.
- 4) Multiple shooting methods can be used for increased robustness and accuracy in the solution at the expense of increased computation cost.

Finally, about the method used and the results obtained we can state:

- 1) Solutions of a norm higher than 2 improve the time used for a rest-to-rest attitude maneuver. In our particular example an improvement of 8% was found.
- 2) An arbitrary maneuver was tested. It has no special characteristics so this method should be applicable to any maneuver.
- 3) The computation expense of the  $\infty$ -norm maneuver prevents it from its application to attitude control. The computation time used to find the initial co-states is not reasonable to be performed in real time. This can be done for 2-norm control since an analytical solution is known.

This last finding is of critical importance to the applicability of this kind of maneuvers in real systems. Suboptimal approaches could be used to allow fast attitude maneuvers such as the 2-norm control. Other authors such as Byers[7] have worked on this paper with intention of leveraging the known structure of the control to simplify the computations. As mentioned in the last point in the control related conclusions, a phase-space analysis may provide additional insight on this approach, and could lead into a global closed loop control, although the problem is likely to suffer from the curse of dimensionality.

Advances on the computational cost would allow the implementation of this kind of optimal maneuvers in the AOCS/FCS systems of spacecraft and aircraft. If the topic was to be researched more deeply the next steps would be in this direction.

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## Appendices

### A. Mathematical notation

#### A. Quaternion and Vector notation and derivatives

Here we use the so called Hamilton quaternion convention. A unitary quaternion  $q$  with components  $q_0, q_x, q_y, q_z$  is defined as:

$$q = \begin{bmatrix} q_0 & q_x & q_y & q_z \end{bmatrix}^T = \left( \cos \frac{\theta}{2}, \sin \frac{\theta}{2} \vec{n} \right) \quad (48)$$

Where  $q_0$  represents the real part of the quaternion associated with the rotation angle  $\theta$  and  $q_x, q_y, q_z$  represent the imaginary part associated with the rotation axis  $\vec{n}$ .

The quaternion product of the quaternions  $q$  and  $p$  is defined as usual and denoted by  $q \circ p$ . The quaternion conjugate is denoted by  $q'$ . We highlight the quaternion product as a matrix product as shown in Parwana [8] and its relations with the conjugate quaternion:

$$q \circ p = (q' \circ p')' = Q(q)p = \bar{Q}(p)q = \begin{bmatrix} q_0 & -q_x & -q_y & -q_z \\ q_x & q_0 & -q_z & q_y \\ q_y & q_z & q_0 & -q_x \\ q_z & -q_y & q_x & q_0 \end{bmatrix} \begin{bmatrix} p_0 \\ p_x \\ p_y \\ p_z \end{bmatrix} = \begin{bmatrix} p_0 & -p_x & -p_y & -p_z \\ p_x & p_0 & p_z & -p_y \\ p_y & -p_z & p_0 & p_x \\ p_z & p_y & -p_x & p_0 \end{bmatrix} \begin{bmatrix} q_0 \\ q_x \\ q_y \\ q_z \end{bmatrix} \quad (49)$$

The time derivative of  $q$  as a function of the angular velocity in the rotating body frame  $\vec{\omega}$  is

$$\dot{q} = 1/2 q \circ [0 \quad \vec{\omega}]^T \quad (50)$$

In the Hamiltonian we get an expression of the form  $\lambda^T(q \circ p)$ , where  $\lambda$  is a 4 element column vector. This expression is scalar. Expressing it as a matrix product we can find its derivative with respect to  $q$  and  $p$ , using matrix calculus to obtain more compact expressions for the costates ODEs, we will use a column gradient convention. We start by finding the derivative of the quaternion product:

$$\frac{\partial}{\partial p} \lambda^T(q \circ p) = \frac{\partial}{\partial p} \lambda^T Q(q)p = (\lambda^T Q(q))^T = Q(q)^T \lambda = Q(q') \lambda = q' \circ \lambda \quad (51)$$

$$\frac{\partial}{\partial q} \lambda^T(q \circ p) = \frac{\partial}{\partial q} \lambda^T \bar{Q}(p)q = (\lambda^T \bar{Q}(p))^T = \bar{Q}(p') \lambda = \lambda \circ p' \quad (52)$$

Where the derivative of the quaternion is obtained directly from Matrix calculus and the second one is obtained by inspection of the resulting matrix. We now obtain the expressions for the partial of the Hamiltonian taylorred to the

dynamics obtained for our problem:

$$H(x, u, \lambda) = \frac{1}{2} \lambda_q^T (q \circ \omega) + \lambda_\omega^T \vec{u} \quad (53)$$

$$-\frac{\partial H}{\partial q} = -\frac{1}{2} \frac{\partial}{\partial q} \lambda_q^T (q \circ \omega) = -\frac{1}{2} \lambda_q \circ \omega' = \frac{1}{2} \lambda_q \circ \omega \quad (54)$$

$$-\frac{\partial H}{\partial \omega} = -\frac{1}{2} \frac{\partial}{\partial \omega} \lambda_q^T (q \circ \omega) = -\frac{1}{2} q' \circ \lambda_q \quad (55)$$

## B. Derivative of a norm

The norm used in this report is defined in terms of an absolute value for the greatest generality. We prove here the derivative:

$$\frac{d}{dx} |x|^q = \begin{cases} \frac{d}{dx} x^q = q x^{q-1} = q |x|^{q-1} & \text{for } x > 0 \\ \frac{d}{dx} (-x)^q = q (-x)^{q-1} = q |x|^{q-1} & \text{for } x < 0 \end{cases} \quad (56)$$

Or in a more compact way using the signum function:

$$\frac{d}{dx} |x|^q = \text{sgn}(x) q |x|^{q-1} \quad (57)$$

We will also use extensively the following property of the signum function:

$$\text{sgn}(x \cdot y) = \text{sgn}(x) \text{sgn}(y) \quad (58)$$

## B. MATLAB code

This appendix includes all relevant code developed in MATLAB. For the sake of simplicity only the code regarding Multiple Shooting with `fmincon` is included. All other numerical methods can be found on GitHub: <https://github.com/ebabio/attitudeTimeOptimalManeuvers.git> (including this very same report)

### A. attitudeManeuver

Main script, checks if in continuation or not, solves the problem and displays results.

```
%% 3d Attitude Test

%% Script mode

reuse = 1

%% Setup workspace
addpath(genpath(pwd))
if(reuse == 1 && exist('orbitFinal', 'var'))
    clearvars -except orbitFinal reuse continuation parameters fileName
else
    clearvars -except continuation fileName
end
clc

%% Define boundary conditions

% Initial / final conditions
roll = deg2rad(0.01);
pitch = deg2rad(-0.01);
yaw = deg2rad(179.99);

q0 = [1 0 0 0]';
qf = eul2quat([yaw pitch roll])';

omegaBCs = zeros(3,1);
```

```

x0BCs = [q0; omegaBCs];
xfBCs = [qf; omegaBCs];

% Control parameters
if(exist('continuation', 'var') && exist('parameters', 'var') && continuation)
    continuationCoefficient = 1.08;
    parameters.lNorm = floor(10 * continuationCoefficient * parameters.lNorm) /
        10; % round to first decimal
    parameters.k = continuationCoefficient * parameters.k;
else
    parameters.lNorm = 2;
    parameters.k = 5e1;
end
display(['Looking for optimal control with norm ' num2str(parameters.lNorm) ':
    '])

%% Setup

% 1. Trajectory Integration routine
control = @(t,x) omegaControl(t, x, parameters);
shootingOde = @(t, x) (attitudeODEs(t, x, control));

orbit = Orbit(zeros(14,1), shootingOde);
orbit.odeOptions = odeset('RelTol',1e-11,'AbsTol',1e-14);

% 2. Boundary Conditions
shootingBCs = @(x0, xf) (attitudeBCs(x0BCs, x0, xfBCs, xf, shootingOde));

% 3. Optimizer setup
fObjective = @(x) x(end);

```

```

% 3.a: fmincon
struct.nIntervals = 3;
struct.setupOrbit = orbit;
gConstraints = @(x) ( nonlconWrapper([x0BCs; x], shootingBCs, struct) );
fminconoptions = optimoptions('fmincon', 'Display', 'iter', '
    FiniteDifferenceStepSize', 1e-4, ...
    'ConstraintTolerance', 1e-4, 'OptimalityTolerance', 1, ...
    'MaxIterations', 1e3, 'MaxFunctionEvaluations', 1e6, ...
    'UseParallel', true);

% 3.b: fminunc
fLagrangian = @(x) (sqrt(gConstraints(x))*gConstraints(x));
fminuncoptions = optimoptions('fminunc', 'Display', 'iter', '
    FiniteDifferenceStepSize', 1e-4, 'UseParallel', true, 'OptimalityTolerance',
    1e-4);

%% Initial solution

if(~exist('orbitFinal', 'var'))
    % First guess
    thetaf = 2*acos(qf(1));
    n = qf(2:4)/norm(qf(2:4));
    tf = 2*sqrt(thetaf); % * (.98 + 0.02*randn(1));

    lambda_q = -[0; n];% + .2*randn(4,1);
    lambda_omega = -tf/4*n; % + .2*randn(3,1);
    lambda0 = [lambda_q; lambda_omega];

    % First actual integration
    orbit.x0 = [ x0BCs; lambda0];
    orbit.integrateX0(1*tf);
else
    % Reuse previous solution as final guess

```

```

    orbit.x0 = orbitFinal.x0; % + 0.01*randn(size(orbit.x0));
    orbit.integrateX0(1*orbitFinal.tf);
end

% Setup independent shootings
dim = numel(orbit.x0);
tfVector = orbit.tf * (1:struct.nIntervals)./struct.nIntervals;
for i=1:struct.nIntervals
    if(i==1)
        x0Opt = orbit.x0(1:dim);
    else
        index0 = (i-1)*dim;
        x0Opt(index0+1:index0+dim) = deval(orbit.odeSol, tfVector(i-1));
    end
end
x0Opt(end+1) = orbit.tf;

%% Solve
% 1. Test for initial guess (initial point BCs should be satisfied)
f0 = shootingBCs(orbit.x0, orbit.xf);
f0(8:end)
orbit.x0(8:14)

% 2. Iterate shooting
tic
[xFinal, tf, exitflaf, output] = fmincon(f0objective, x0Opt(dim/2+1:end), [],
    [], [], [], [], gConstraints, fminconoptions);
%[xf, tf, exitflaf, output] = fminunc(fLagrangian, x0Opt(dim/2+1:end),
    fminuncoptions);
toc
orbitFinal = copy(orbit);

```

```

orbitFinal.x0 = [x0BCs; xFinal(1:dim/2)];
orbitFinal.integrateX0(tf);

% 3. Test for final guess (all BCs should be satisfied)
f = shootingBCs(orbitFinal.x0, orbitFinal.xf)
orbitFinal.x0(8:14)

%% Evaluate the initial solution at all times

maneuver0.t = orbit.t;
maneuver0.x = orbit.x;

maneuver0.q = maneuver0.x(1:4,:);
maneuver0.eul = quat2eul(maneuver0.q)';
maneuver0.omega = maneuver0.x(5:7,:);
maneuver0.lambda_q = maneuver0.x(8:11,:);
maneuver0.lambda_omega = maneuver0.x(12:14,:);

maneuver0.u = zeros(3, numel(maneuver0.t));
for i=1:numel(maneuver0.t)
    maneuver0.u(:,i) = control(maneuver0.t, maneuver0.x(:,i));
end

% plot Data
figure(1)
clf reset

subplot(2,2,1)
plot(maneuver0.t, maneuver0.eul)
title('Euler Angles')
legend('Yaw', 'Pitch', 'Roll', 'Location', 'best')
subplot(2,2,2)

```

```

plot(maneuver0.t, maneuver0.omega)
title('Angular Rates')
legend('\omega_x', '\omega_y', '\omega_z', 'Location', 'best')
subplot(2,2,3)
plot(maneuver0.t, maneuver0.u)
title('Torque input')
legend('u_x', 'u_y', 'u_z')
subplot(2,2,4)
plot(maneuver0.t, maneuver0.lambda_omega)
title('omegacostates')
legend('\lambda_\omega_x', '\lambda_\omega_y', '\lambda_\omega_z', 'Location',
      'best')

%% Evaluate the final solution at all times
maneuver.t = orbitFinal.t;
maneuver.x = orbitFinal.x;

maneuver.q = maneuver.x(1:4,:);
maneuver.eul = quat2eul(maneuver.q)';
maneuver.omega = maneuver.x(5:7,:);
maneuver.lambda_q = maneuver.x(8:11,:);
maneuver.lambda_omega = maneuver.x(12:14,:);

maneuver.u = zeros(3, numel(maneuver.t));
for i=1:numel(maneuver.t)
    maneuver.u(:,i) = control(maneuver.t, maneuver.x(:,i));
end

%plot data
figure(2)
clf reset

```



```

subplot(2,2,1)
plot(maneuver.t, maneuver.eul)
title('Euler Angles')
legend('Yaw', 'Pitch', 'Roll', 'Location', 'best')
subplot(2,2,2)
plot(maneuver.t, maneuver.omega)
title('Angular Rates')
legend('\omega_x', '\omega_y', '\omega_z', 'Location', 'best')
subplot(2,2,3)
plot(maneuver.t, maneuver.u)
title('Torque input')
legend('u_x', 'u_y', 'u_z', 'Location', 'best')
subplot(2,2,4)
plot(maneuver.t, maneuver.lambda_omega)
title('omegacostates')
legend('\lambda_\omega_x', '\lambda_\omega_y', '\lambda_\omega_z', 'Location',
      'best')

```

**B. attitudeManeuverLoop**

Initializes and runs the continuation scheme up to norm 20 while logging all intermediate information.  $\infty$ -norm must be solved manually.

```

%% Run attitudeManeuver until hitting stopping condition

%% Setup workspace
clear all

fileName = ['results/' datestr(datetime, 'mmdd-HHMM' )];
diary ([fileName ' Diary.txt'])

%% Iterate
continuation = 1;
parameters.lNorm = 2;

while(parameters.lNorm < 20)
    attitudeManeuver
    save([fileName 'lNorm' num2str(parameters.lNorm) ' Workspace'])
    pause(.1);
end

display('continuation successfully finished')
diary off

```

**C. attitudeODEs**

ODEs for the motion other than control

```
function dxdt = attitudeODEs(t, x, control)

% Reassign variables
x = x';           %since quaternion and row vectors in matlab
q = x(1:4);
omega = x(5:7);
lambda_q = x(8:11);
% lambda_omega = x(12:14);

% Synthesize control
u = control(t,x);

% Compute derivatives
% a. Quaternion multiplication way
qDot = 1/2 * quatmultiply(q, [0 omega]); % check derivation
omegaDot = u;
lambdaDot_q = 1/2 * quatmultiply(lambda_q, [0 omega]); % check derivation
lambdaDot_omega = -1/2 * quatmultiply(quatconj(q), lambda_q); % check
    derivation

dxdt = [qDot omegaDot lambdaDot_q lambdaDot_omega(2:4)]';
return
```

**D. attitudeBCs**

BCs for the motion other than control

```
function [BCs, pathCs] = attitudeBCs(x0BCs, x0, xfBCs, xf, xDotFun, tf)
```

```

BCs = zeros(15,1);
pathCs = zeros(2,1);

if(nargin == 6)
    collocation = 1;
else
    collocation = 0;
end

%% Allow for xDotFun collocation parameters
if(collocation)
    xDotWrapper = @(tau, x) (xDotFun(tau, x, tf));
else
    xDotWrapper = @(t,x) (xDotFun(t, x));
end

%% Initial state constraints
% Reassign variables
q0 = x0(1:4);
omega0 = x0(5:7);
lambda_q0 = x0(8:12);
lambda_omega0 = x0(13:14);

% Hamiltonian
xDot = xDotWrapper([],x0);
H0 = [lambda_q0;lambda_omega0]' * xDot(1:7);

BCs(1:7) = [q0-x0BCs(1:4); omega0-x0BCs(5:7)];
pathCs(1) = H0+1;
BCs(15) = norm(lambda_q0)-1;

```

```

%% Final State Constraints
qf = xf(1:4);
omegaf = xf(5:7);
lambda_qf = xf(8:12);
lambda_omegaf = xf(13:14);

% Hamiltonian
xDot = xDotWrapper([],xf);
Hf = [lambda_qf;lambda_omegaf]' * xDot(1:7);

BCs(8:14) = [qf-xfBCs(1:4); omegef-xfBCs(5:7)];
pathCs(2) = Hf+1;
% BCs(15) = pathCs(2);
return

```

### E. omegaControl

Implementation of the control independently of the equation of motion. If the required norm is greater it will directly compute the infinity norm controller.

```

function u = omegaControl(t, x, parameters)

lNorm = parameters.lNorm;
k = parameters.k;

% Reassign variables
q = x(1:4);
omega = x(5:7);
lambda_q = x(8:11);
lambda_omega = x(12:14);

```

```

% Modified control
lambda_omega_exp = @(i) abs(lambda_omega(i))^(1/(lNorm-1));
den = (lambda_omega_exp(1)^lNorm + lambda_omega_exp(2)^lNorm +
      lambda_omega_exp(3)^lNorm)^(1/lNorm);

fun = @(x) sign(x);

if(k<20)
    u(1) = -fun(k*lambda_omega(1)) * lambda_omega_exp(1) / den;
    u(2) = -fun(k*lambda_omega(2)) * lambda_omega_exp(2) / den;
    u(3) = -fun(k*lambda_omega(3)) * lambda_omega_exp(3) / den;
else
    u(1) = -fun(k*lambda_omega(1));
    u(2) = -fun(k*lambda_omega(2));
    u(3) = -fun(k*lambda_omega(3));
end

```

**F. Orbit**

Class for integration of the motion in a convenient manner.

```

classdef Orbit < matlab.mixin.Copyable
    % Important note on handle classes: handle objects behave as handles/
    % pointers
    % To create a different object with the same values use COPY method

    properties
        odefun

        x0

        x
        t

        xf
        tf

        odeSol
        odeOptions
        odeIntegrator
    end

    methods
        function obj = Orbit(x0, odefun)
            if(nargin ==2)
                obj.odefun = odefun;
                obj.x0 = x0;
                obj.x = obj.x0;
                obj.t = 0;
            elseif(nargin ==0)

```

```

        %empty constructor
    else
        warning('number of inputs not recognized');
    end
end

function addDeltaX0(obj, deltaX0)
    if(size(deltaX0,1)>numel(obj.x0))    %update vector may contain
        final time
        obj.x0 = obj.x0 + deltaX0(1:numel(obj.x0));
        obj.tf = obj.tf + deltaX0(numel(obj.x0)+1);
    else
        obj.x0 = obj.x0 + deltaX0;
    end
end

function obj = integrateX0(obj, tEnd)
    if(nargin < 2)
        tEnd = obj.tf;
    end

    obj.odeSol = ode45(obj.odefun, [0 tEnd], obj.x0, obj.odeOptions);
    obj.t = obj.odeSol.x;
    obj.x = obj.odeSol.y;

    obj.tf = obj.t(end);
    obj.xf = obj.x(:,end);
end

function phi = STM(obj, tEnd, dimensions)
    % STM found by variation of the initial conditions
    n = numel(obj.x0);

```



```

if(nargin < 2)
    tEnd = obj.tf;
end
if(nargin < 3)
    activeDimensions = ones(1,n);
    withTime = 0;
else
    % transform from indexes to an indicator vector
    withTime = find(dimensions==n+1);
    activeDimensions(dimensions) = ones(1,numel(dimensions));
    activeDimensions = activeDimensions(1:n);
end

phi = eye(n);
delta = 1e-3;

% Reference solution
obj.odeSol = ode45(obj.odefun, [0 tEnd], obj.x0, obj.odeOptions);
obj.t = obj.odeSol.x;
obj.x = obj.odeSol.y;

obj.tf = obj.t(end);
obj.xf = obj.x(:,end);

% Compute STM
for i=find(activeDimensions==1)
    perturbation = delta * circshift([1; zeros(n-1,1)], i-1);
    xDelta0 = obj.x0 + perturbation;
    [~,xDeltaf] = ode45(obj.odefun, [0 tEnd], xDelta0, obj.
        odeOptions);
    phi(:,i) = (xDeltaf(end,:) - obj.xf)/delta;
end

```

```

end
for i=find(activeDimensions==0)
    phi(:,i) = zeros(n,1);
end

if(withTime)
    phi(:,end+1) = obj.odefun([], obj.xf);
end
end
end

end

```

### G. nonlconWrapper

Wrapper function for the nonlinear constraints. It wraps the integration of the motion and the constraints as MATLAB's optimization toolbox expects. It also handles the number of intervals for multiple shooting method and defects between intervals.

```

function [g, h] = nonlconWrapper(x, shootingBCs, parameters)

%% Assign parameters
nIntervals = parameters.nIntervals;
setupOrbit = parameters.setupOrbit;

dim = (numel(x)-1)/nIntervals;
%% Evaluate orbits

tf = x(end);

iterOrbit(nIntervals) = Orbit();

```

```

tfVector = tf * (1:nIntervals)./nIntervals;
for i=1:nIntervals
    iterOrbit(i) = copy(setupOrbit);
    if(i==1)
        iterOrbit(i).x0 = x(1:dim);
        iterOrbit(i).tf = tfVector(i);
    else
        index0 = (i-1)*dim;
        iterOrbit(i).x0 = x(index0+1:index0+dim);
        iterOrbit(i).tf = tfVector(i) - tfVector(i-1);
    end
    iterOrbit(i).integrateX0();
end

%% Evaluate constraints

% Trajectory constraints
constraintErr = shootingBCs(iterOrbit(1).x0, iterOrbit(end).xf);

% Defects
defectErr = zeros(dim * (nIntervals-1),1);
for i = 2:nIntervals
    index0 = dim*(i-2);
    defectErr(index0+1:index0+dim) = iterOrbit(i).x0 - iterOrbit(i-1).xf;
end

%% Wrap up for returning
g = 0;
h = [constraintErr; defectErr];

```