

# ecalj note

<https://github.com/tkotani/ecalj>

August 9, 2025

This note becomes rather obsolete. I have to revise things step by step.

## 1 Wannier function

We can generate Wannier functions (maximally localized Wannier Functions or similar) by a script `genMLWF`. It automatically perform cRPA calculation successively. (If it is not necessary, insert 'exit' in `genMLWF`, after it performs `lmaxloc2`).

Try to run examples in `ecalj/MATERIALS/Sample_MLWF/`. Read `README` in it. To run the script `genMLWF`, we need to get `GWinput` by editing `GWinput.tmp`. (`mkGWin_lmf2` contains default Wannier section). In addition, we have some settings (energy windows and so on). This is the example of the initial conditions for Cu case. 5 is the number of Wannier function. The most left one means  $\phi$  index and the right one of it is  $\dot{\phi}$  index. They are written in the `@MNLA_CPHI` file.

Then we can run `genMLWF`. After it finished, we can analyze it results. (if you don't need Wannier function plot, You can skip a line of `wanplot` in `genMLWF`. Then we don't need to set `vis_wan_*` options.)

### 1.1 lwmatK1 and lwmatK2

If you input the following command

```
>grep Wan lwmatK*
```

You will get the following results. (This case : Cu cases)

lwmatK1:	Wannier	1	1	24.644475	0.000000 eV	
lwmatK1:	Wannier	1	2	24.644576	0.000000 eV	
lwmatK1:	Wannier	1	3	25.471361	0.000000 eV	
lwmatK1:	Wannier	1	4	24.644575	0.000000 eV	
lwmatK1:	Wannier	1	5	25.470946	0.000000 eV	
lwmatK2:	Wannier	1	1	0.000000 eV	-21.263759	-0.000000 eV
lwmatK2:	Wannier	1	2	0.000000 eV	-21.263839	0.000000 eV
lwmatK2:	Wannier	1	3	0.000000 eV	-21.931033	-0.000000 eV
lwmatK2:	Wannier	1	4	0.000000 eV	-21.263839	-0.000000 eV
lwmatK2:	Wannier	1	5	0.000000 eV	-21.930702	-0.000000 eV

### Wannier Branch now under developing (imported from T.Miyake's Wannier and H.Kino's).

A. make at ecalj/fpgw/Wannier/ directory, and do make, and make install.

(need to check Makefile first). You first have to install fpgw/exec/ in advance.

B. Samples are at these directories.

MATERIALS/CuMLWFs (small samples),

MATERIALS/CuMLWF/

MATERIALS/CuMLWFs/

MATERIALS/FeMLWF/

MATERIALS/NiOMLWF/

MATERIALS/SrVO3MLWF/

C. With GWinput and ctrl.\*, run

>genMLWF

at these directories.

In GWinput, we supply settings to generate Wannier functions. (Sorry, not documented yet..)

D. After genMLWF, do

>grep Wan lwmatK\*

then compare these with Result.grepWanlwmatK

These are onsite effective interactions (diagonal part only shown).

\*.xsf are for plotting the Maximally localized Wannier functions.

Anyway, documentation on Wannier is on the way half.

Time consuming part (and also the advantage) is for effective interaction in RPA.

Look into the shell script genMLWF; you can skip last part if you don't need the effective interaction.

## 2 How to set local orbitals

As we stated, do "lmfa |grep conf" to check used MTO basis.

We have to set SPEC\_ATOM\_PZ=?,?,?

(they ordered as PZ=s,p,d,f,... ) to set local orbitals.

lmv7 (originally due to ASA in Stuttgart) uses a special terminology

"continuous principle quantum number for each l", which is just

related to the logarithmic derivative of radial functions at MT

boundary. It is defined as

$P = \text{principleQuantumNumber} + 0.5 - 1/\pi \cdot \arctan(r \cdot 1/\phi \cdot d\phi/dr)$ ,

where  $\phi$  is the radial function for each l. For example,

$P = n.5$  for  $l=0$  of free electron (flat potential) because  $\phi=r^0$ ,

$P = n.25$  for  $l=1$  because  $\phi=r^1$ ;

$P = n.147584$  for  $l=2$  because  $\phi=r^2$ ;  $P = n.102416$   $n.077979$  for  $l=3,4$ .  
 (Integer part can be changed). See Logarithmic Derivative Parameters in  
<http://titus.phy.qub.ac.uk/packages/LMT0/lmto.html#section2>

Its fractional part  $0.5 - \arctan(1/\phi \, d\phi/dr)$  is closer to unity for  
 core like orbital, but closer to zero for extended orbitals.

Examples of choice:

Ga p: in this case, choice 0 or choice 2 is recommended.

We usually use lo for semi-core, or virtually unoccupied level.

(0)no lo (4p as valence is default treatment without lo.)

3p core, 4p valence, no lo: default.

Then we have choice that lo is set to be for 3p,4p,5p.

(1)3p lo ---> 4p val (when 3p is treated as valence)

3d semi core, 4d valence

Set PZ=0,3.9

(P is not required to set. \*.9 for core like state. It is just an initial condition.)

(2)5p lo ---> 4p val (PZ>P)

Set PZ=0,5.5

5.5 is just simply given by a guess (no method have yet  
 implemented for

If 5.2 or something, it may fails

because of poorness in linear-dependency. We may need to observe  
 results should not change so much on the value of PZ.

(3xxx)4p lo ---> 5p val (we don't use this usually. this is for test purpose)

4p lo, 5p valence

Set PZ=0,4.5 P=0,5.5 (In this case, set P= simultaneously).

(NOTE: zero for s channel is to use default numbers for s)

Ga d: (in this case, choice 0 or choice 1 is recommended).

(0)no lo (3d core, 4d valence, no lo: default.)

Then we have choice that lo is set to be for 3d,4d,5d.

(1) 3d lo ---> 4d val (when 3d is treated as valence)

Set PZ=0,0,3.9 (P is not required to set)

(2) 5d lo ---> 4d val (PZ>P)

Set PZ=0,0,5.5

(3xxx) 4d lo ---> 5d val (this is for test purpose)

Set PZ=0,0,5.5 P=0,0,4.5

(NOTE: zero for s,p are to use default numbers )

```
% If you like to read from atm.ga file instead of rst file(if exist).  
% You have to do lmf --rs=1,1,0,0,1, for example. See lmf --help  
% Because rst file keeps the setting of MT0, thus change in ctrl is not  
% reflected without the above option to lmf.
```

=====

### 3 Linear response calculations

With these scripts for linear response calculations, **eps\***, we can calculate **q**-dependent dielectric function  $\epsilon(\omega, \mathbf{q})$  (and  $v$ ,  $W$ ) (and  $\chi$  for spin fluctuation). But (because of numerical reason), we can not use  $\mathbf{q} = 0$  limit. (if  $|\mathbf{q}|$  is too small, we have numerical problem, zero divided by zero, because we have not implemented the version to use  $\mathbf{q} = 0$ .)

- **eps\_lmfh**

Dielectric function epsilon with local field correction. Expensive calculation (we may need to reduce number of wing parts in future...).

- **epsPP\_lmfh**

epsilon without local field correction.  $1 - \langle e^{i\mathbf{qr}} | v | e^{i\mathbf{qr}} \rangle \langle e^{i\mathbf{qr}} | (\chi^0) | e^{i\mathbf{qr}} \rangle$

- **epsPP\_lmfh\_chipm**

For spin susceptibility. This essentially calculate non-interacting spin susceptibility. Then it is used for the calculation of full spin susceptibility with **util/calj\_\*.F** programs (small quick programs). See spin wave paper. See spin susceptibility section Sec.??.

- (not maintained now; we will recover this) **eps\_lmfh\_chipm**

This gives full non-interacting spin susceptibility. Testing. We have to determine  $U$  (Stoner  $I$ ) for the determination of full spin susceptibility. TDLDA? or so?

- (This is old mode --- not maintained) **epsPP\_lmfh\_chipm\_q**

For spin susceptibility, spin susceptibility  $\langle e^{i\mathbf{qr}} | \chi(q, \omega) | e^{i\mathbf{qr}} \rangle$  In this script, You have to assign that isp=1 is majority, isp=2 is minority. This is with long wave approximation.

- 
- We use the histogram method (the Hilbert transformation method); we first calculate its imaginary parts with the tetrahedron technique for dielectric functions. Then we get its real part by the Hilbert transformation.

You need to choose `HisBin_dw, HisBin_ratio`. The width of histogram bins are getting larger when omega gets larger. dw is the size of histogram-bin width at omega=0. At omega=omg\_c, its width gets twiced.

To plot dielectric function with reasonable resolution, it might be better to set **dw** 0.001 and **omg\_c** 0.1 for example. You may have to choose small enough omega for spin wave mode as 0.001 Ry (Or smaller). omg\_c is given like 0.05 Ry or so. But sometimes it can be like 1Ry.

- **epsPP\_lmfh** only calculation an matrix element of dielectric function for  $\exp(i\mathbf{qr})$ . Thus very faster than **eps\_lmfh** mode. It uses a a special product basis set for cases without inversion (problem is in how to expand  $\exp(i\mathbf{qr})$  in the MPB; the product basis is not from phi and phidot, but from spherical Bessel functions).

In `*_lmfh_* modes( I now use little for *_lmf_* modes)`, you can use small enough delta. Use small enough delta ( $=-1e-8$  a.u.) for spin wave modes (also you can use it for dielectric function and GW). This is necessary because pole is too smeared if you use larger delta.

### 3.1 `eps_lmfh`, `epsPP_lmfh`: the dielectric functions

You can invoke the script, e.g. as `"eps_lmfh si"`.

---

Specify  $\mathbf{q}$  point in `<QforEPS>` or so. Mesh for  $\omega$  is specified by `dw, omg_c`.

The obtained data are in `EPS*.dat` and `EPS*.nlfc.dat`. `EPS*.nlfc.dat` contains the result without local-field correction `EPS*.dat` contains the result with local-field correction (this is generated only for `eps_lmfh`. Both of them contains

$\mathbf{q}(1:3)$ ,  $\omega$ ,  $\text{Re}(\epsilon)$   $\text{Im}(\epsilon)$ ,  $\text{Re}(1/\epsilon)$ ,  $\text{Im}(1/\epsilon)$

in each line.

### 3.2 `epsPP_lmfh`: the dielectric function(No LFC— faster)

You can calculate  $\epsilon$  without LFC by `epsPP_lmfh`. It is very faster than `eps_lmfh`.

To calculate  $\epsilon(\mathbf{q}, \omega)$  without LFC accurately, the best basis set for the expansion of the Coulomb matrix within MT is apparently not the product basis, but the Bessel functions corresponding to the plane waves  $\exp(i\mathbf{q}\mathbf{r})$ . We use such a basis in this mode. However, our experience shows that the changes are little even with the usual product basis (we don't describe this here).

### 3.3 How to calculate correct dielectric function?

(this subsection is essentially OK... but need to clean it up. dec2014)

There are problems to calculate correct epsilon.

At first, we talk about `epsPP_lmfh`, which is No LFC. Main problem are

---

1. Convergence for number of k point(specified by `n1n2n3`).

Roughly speaking, `20x20x20` is required for not-so-bad results for Fe and Ni.

It is better to do `30x30x30` to see convergence check.

However, in the case of ZB-MnAs (maybe because of simple structure around Ef), it requires less q points.

figs are for GaAs.

fig001: `n1n2n3` convergence for `Chi_RegQbz = on` case.

fig002: `n1n2n3` convergence for `Chi_RegQbz = off` case.

(`Chi_RegQbz` is explained in General section in this manual).

As you see,  $k$  points convergence looks a little better in `Chi_RegQbz=off` (mesh not including  $\gamma$ ). However a little problem is that its threshold around 0.5eV is too high and slowly changing.

fig003: Alouanis'(from Arnaud) vs. ```Chi_RegQbz = on''` vs. ```Chi_RegQbz = off''`  
 As you see, the threshold of the Red line (20x20x20 `Chi_RegQbz=on`) and Alouani's are almost the same, but the red line is too oscillating at the low energy part. On the other hand, ```Chi_RegQbz = off''` in Green broken line is not so satisfactory at the low energy part.

fig.gas\_eps\_kconf.pdf shows the convergence behavior of epsilon for

2.  $q \rightarrow 0$  convergence (this is related to whether `Chi_RegQbz=on` or `off`).

If you use very small  $q$  like  $q=0.001$  in GaAs, it can cause a problem.

Use  $q=0.01$  or larger (maybe  $q=0.02$  or more is safer).

Very small  $q$  can give numerical error for high-energy region.

In fig004, we show the high energy tail part of  $\text{Im } \epsilon(\omega)$  for GaAs case. At  $q=0.01$  (this means  $q = 2\pi/\text{alat} * (0.001)$ ), the imaginary part is a little too large. Less than 80eV,  $q=0.02$  gives good results when compared with other high  $q$  results, though it still has noise above 80eV.

In fig005, I showed the same results compared with Alouani's (his is up to 40eV). Both gives rather good agreements. As you see,  $q=0.06$  or above might be necessary to get reasonable convergence for high energy part above 40eV.

We have to be careful for this pooriness in high energy part--- it may effect low-energy  $\text{Re } \epsilon$  through KK relation. However this can be very small enough.

In fig.gas\_eps\_qconv.jpg, we checked the convergence of  $\epsilon(\omega=0, q)$  for  $q \rightarrow 0$ . As you see, it gives convergence, however,  $q=0.01$  is a little out of curve---this should be because of the pooriness in the high energy part. so  $q=0.02$  or  $q=0.03$  is safer, and you can get  $\epsilon$  within 1 percent accuracy.

3. Including Core for dielectric constant is dangerous.

It can cause very poor results if you include core part in `GWinput`.

You need to include core just as valence (with local orbital).

In fig008, we showed core effects. It starts from  $\approx 16\text{eV}$

(this is core to conduction transition).

fig007 showd the check about the q point dependence---even with large q, it would not change.

These shows that the core excitation can have larger energy range.

This is in contrast to the valence case

(then the most of excitaion is limited to less than 10eV).

We have to be careful for such high-energy exciation... The LMT0 basis might be not so good for high energy.

#### 4. basis set.

Use QpGcut\_psi \approx 3.0 a.u. or so (as same as GW calculation).

In the case of epsPP\* mode,

QpGcut\_cou can be very small--- In our codes now,

ngc>=1 should be for all q vector shown in lqg4gw02 (output of echo 2|qg4gw).

[In principle, it should be only for the q vector for which we calculate epsilon.

But there is a technical poorness in our code---

(maybe) a problem here; the plane-wave part of the eigenfunction generated in lmf2gw is not correctly passed to lmf2gw when ngc=0].

-- eps\_lmfh: including LFC -----

To include eps with LFC, do eps\_lmfh.

But lcutmx=2 seems to be good enough to get 0.5 percent error (maybe better than this).

Test it 10x10x10 or so. (I need to repeat if necessary).

Further you can use smaller QpGcut\_cou like 2.2 or so,

with rather smaller product basis (up to p timed d, not including f).

Note: epsPP\_lmfh is designed to use good basis to calculate eps

without LFC. This is usually in agreement with what you obtained by eps\_lmfh;

however it can give slight difference when you use small product basis.

---Summary -----

So in conclusion, I think a best way to do is

1. set q=0.02 [q=2pi/alat(0 0 0.02)] or so for GaAs case.

If you want to check, do q=0.03 and q=0.06 also.

``Chi\_RegQbz = off'' is better for matrials like GaAs with direct gap.



2. You can use small QpGcut\_cou but all ngc should be one or more.

3. As for the Product basis setting in epsPP\* scripts, only

lcutmx and tolerance (this can be like 0.001 or so) are relevant.to determine eps(omega=0, c

E.g. set lcutmx=4 or so.

5. To get eps with LFC, set QpGcut\_cut as xxx, and set lcutmx=2 where

4. Do nk=20 18 16 and take interpolariion

(occupied sp) \timex (unoccupied spd) are included.

But correct EPS\*.nolfc.d is rather from epsPP\_lmfh script.

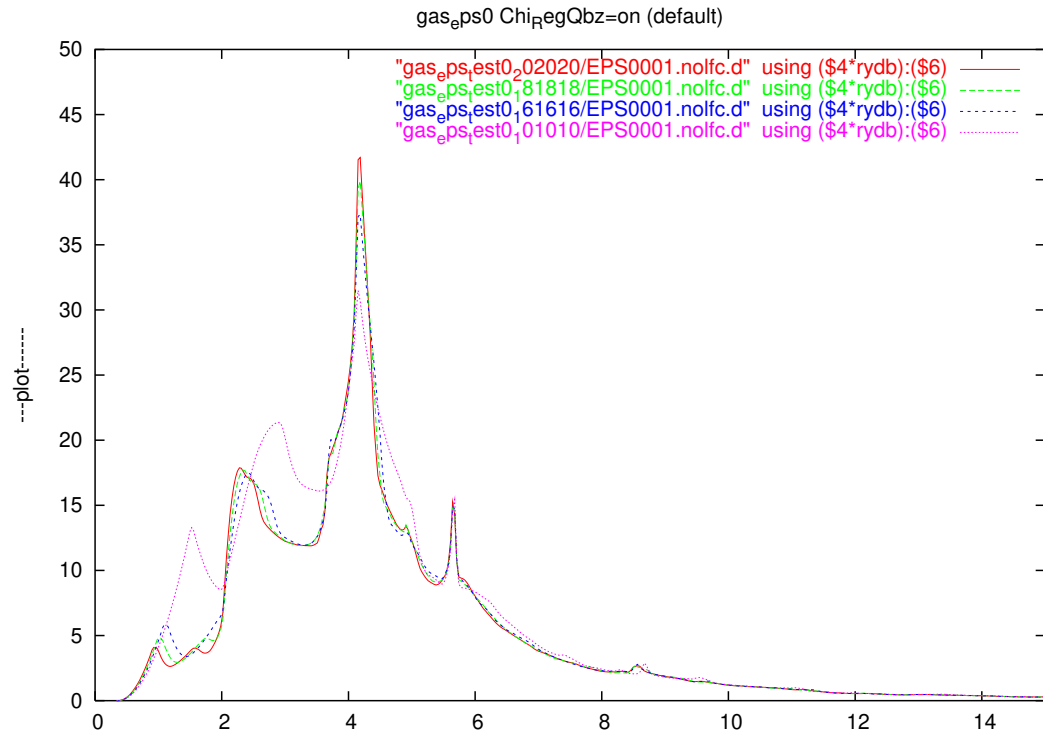


fig001

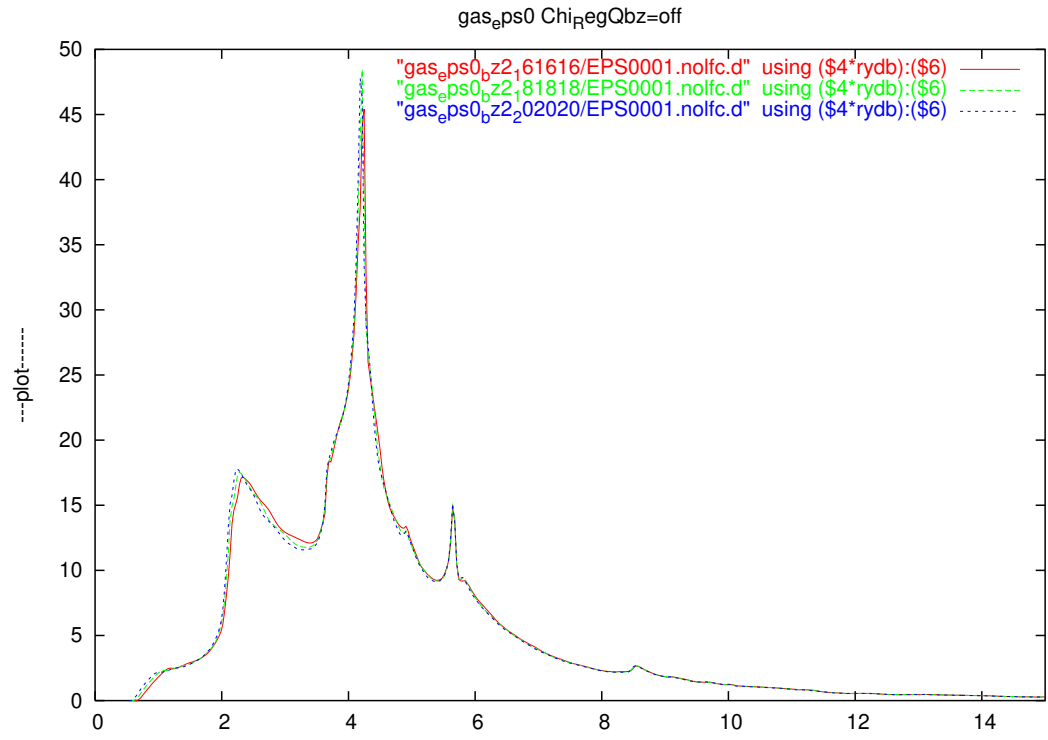


fig002

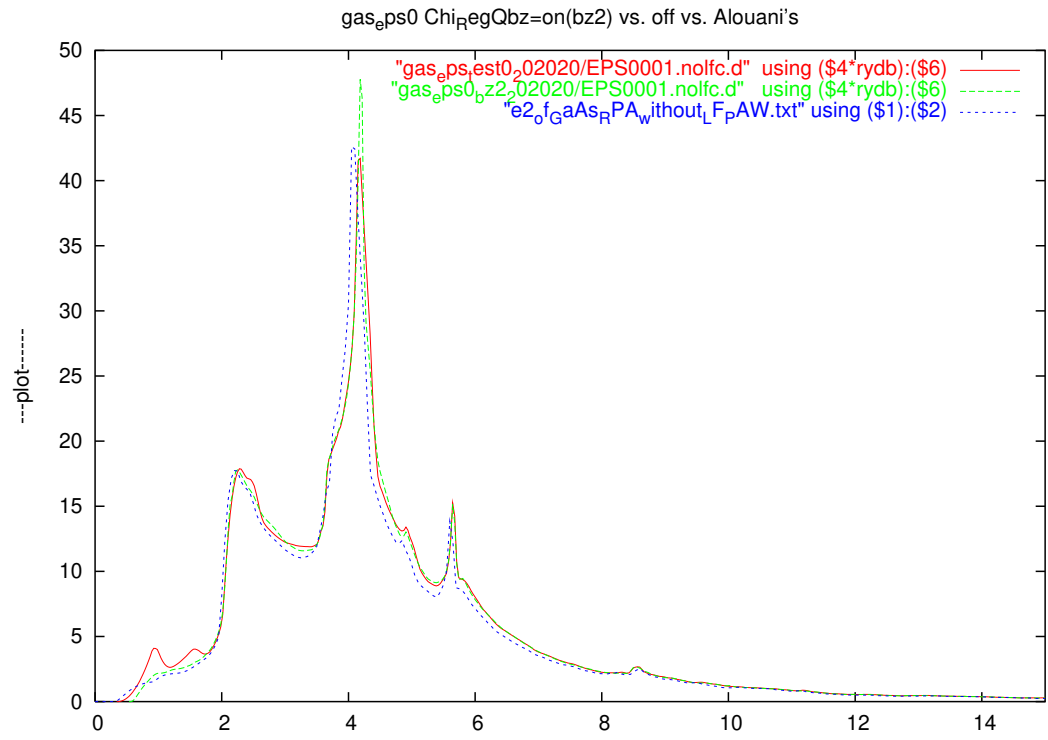


fig003

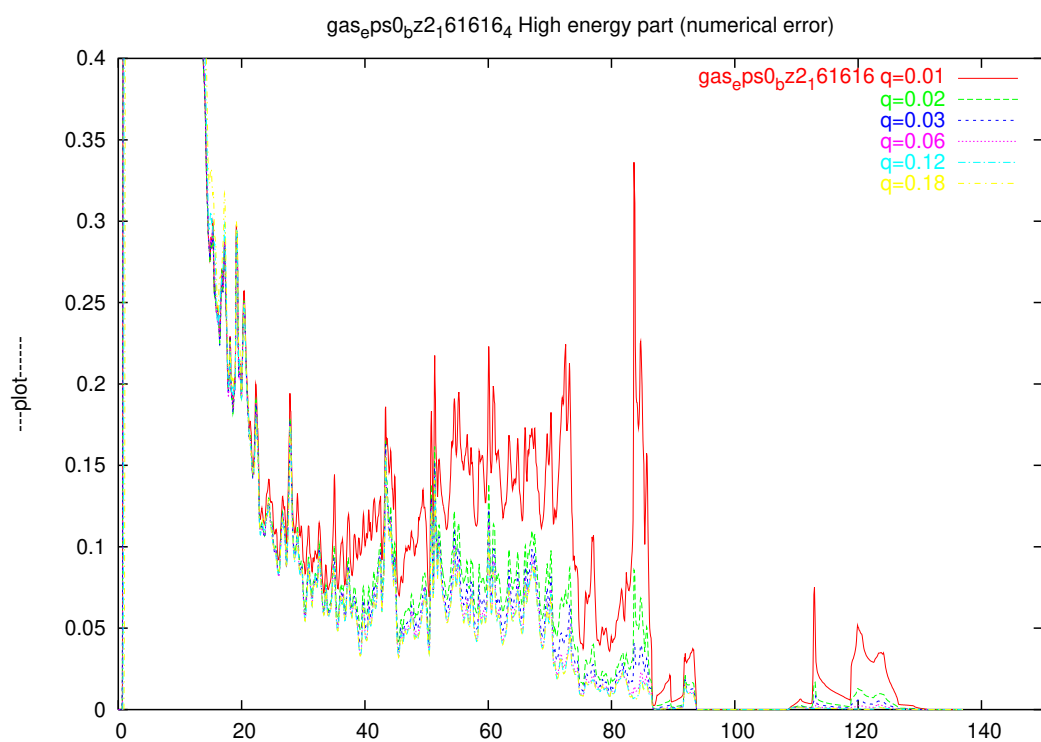


fig004

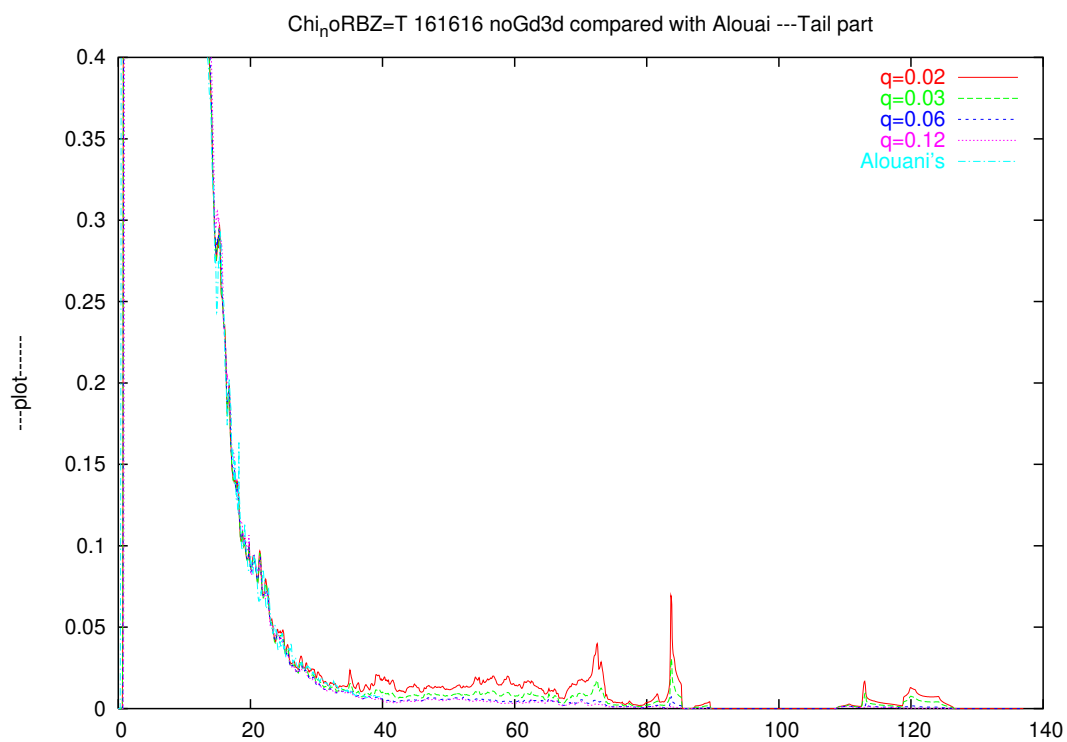


fig005

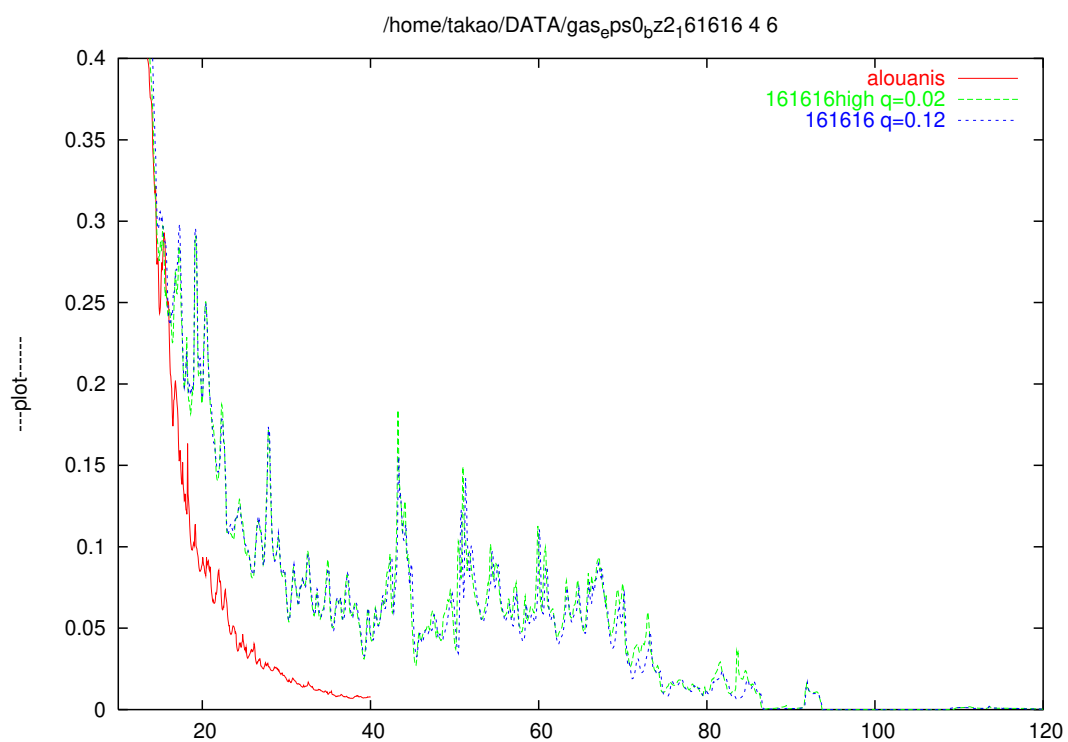


fig007

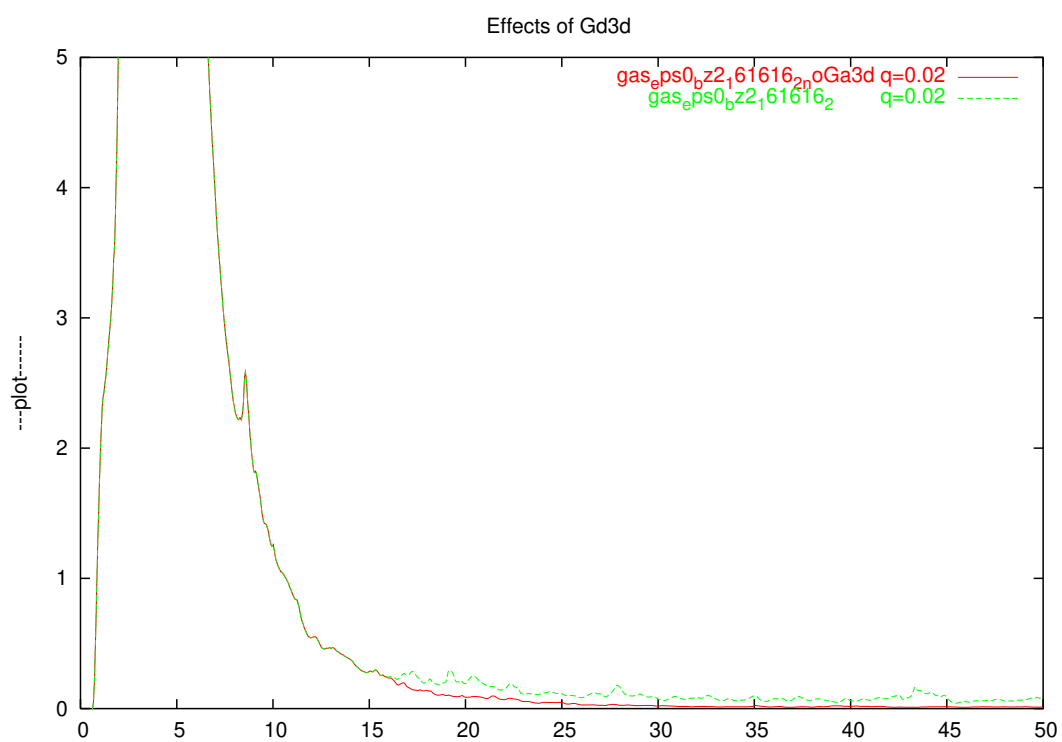


fig008