

1 The Bose-Hubbard Model

The Bose-Hubbard (BH) model describes the ground state of spinless itinerant bosons on a 1D lattice, subject to periodic boundary conditions, and on-site particle interactions. The Hamiltonian matrix representing the BH model is:

$$H = - \sum_{i,j} t_{i,j} b_i^\dagger b_j + \frac{U}{2} \sum_i n_i(n_i - 1) - \sum_i \mu_i n_i \quad (1)$$

where $t_{i,j}$ is the hopping amplitude between sites i and j , U is the on-site interaction strength, μ_i is an on-site chemical potentials, b_i^\dagger (b_i) creates (annihilates) a boson on site i , and n_i counts the number of bosons on site i .

In the Fock configuration basis, the interaction potential and the chemical potential terms are diagonal operators, while the first term, the kinetic energy operator, is not. The Hamiltonian can be rewritten as:

$$H = H_1 + H_0 \quad (2)$$

where $H_1 = - \sum_{i,j} t_{i,j} b_i^\dagger b_j$ is the off-diagonal term, while $H_0 = \frac{U}{2} \sum_i n_i(n_i - 1) - \sum_i \mu_i n_i$ is the diagonal term. Due to H_0 being a diagonal operator and the basis states being orthonormal, the diagonal matrix elements are rather simply obtained:

$$\langle \alpha_p | H_0 | \alpha_q \rangle = \langle \alpha_p | \epsilon_{\alpha_q} | \alpha_q \rangle = \epsilon_{\alpha_q} \langle \alpha_p | \alpha_q \rangle = \epsilon_{\alpha_q} \delta_{p,q} \quad (3)$$

where ϵ_{α_q} is the off-diagonal energy of Fock configuration α_q , and $\delta_{p,q}$ is the Kronecker-Delta function. In other words, the matrix elements vanish when the configurations α_p and α_q are different, and become ϵ_{α_q} otherwise.

The off-diagonal elements of the Hamiltonian are given by:

$$\langle \alpha_p | H_1 | \alpha_q \rangle = - \sum_{i,j} \langle \alpha_p | t_{i,j} b_i^\dagger b_j | \alpha_q \rangle \quad (4)$$

Recall the definition of the bosonic creation and annihilation operators:

$$b^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle \quad (5)$$

$$b |n\rangle = \sqrt{n} |n-1\rangle \quad (6)$$

The basis states can be expressed as:

$$|\alpha_q\rangle = |n_1, n_2, \dots, n_L\rangle \quad (7)$$

where $n_i = b_i^\dagger b_i$ is the number of particles on site i , and L is the number of lattice sites. In this representation, acting first with the annihilation operator on site j , then acting with the creation operator on site i produces the following result:

$$b_i^\dagger b_j |n_1, n_2, \dots, n_L\rangle = \sqrt{n_j(n_j+1)} |\dots, n_i+1, \dots, n_j-1, \dots\rangle \quad (8)$$

Only sites i and j are modified. The matrix element in Eq. (4) becomes:

$$\langle \alpha_p | H_1 | \alpha_q \rangle = - \sum_{i,j} t_{i,j} \sqrt{n_j(n_j+1)} \langle m_1, m_2, \dots, m_L | \dots, n_i+1, \dots, n_j-1, \dots \rangle \quad (9)$$

Due to orthonormality of the basis states:

$$\langle m_1, m_2, \dots, m_L | \dots, n_i+1, \dots, n_j-1, \dots \rangle = \delta_{m_i, n_i+1} \delta_{m_j, n_j-1} \prod_{k \neq i,j} \delta_{m_k, n_k} \quad (10)$$

where m_i is the number of particles on site i of the state $|\alpha_p\rangle$ (i.e, the "bra" state) and the product is carried over all sites k , excluding i and j . The elements of the kinetic operator then become:

$$\langle \alpha_p | H_1 | \alpha_q \rangle = - \sum_{i,j} t_{i,j} \sqrt{n_j(n_j+1)} \delta_{m_i, n_i+1} \delta_{m_j, n_j-1} \prod_{k \neq i,j} \delta_{m_k, n_k} \quad (11)$$

Eqs. (3) and (11) give the diagonal and off-diagonal elements of the Bose-Hubbard Hamiltonian (1). These elements will be necessary in order to sample the partition function of a Bose-Hubbard configuration via Path Integral Monte Carlo. In the next section, the path integral formulation of the partition will be discussed.

2 Partition function (weights)

The ground state can be exactly obtained by:

$$|\psi_0\rangle = \lim_{\beta \rightarrow \infty} e^{-\frac{\beta}{2}H} |\psi_T\rangle \quad (12)$$

where $\beta = \frac{1}{k_B T}$, H is the ground state Hamiltonian, and $|\psi_T\rangle$ is a trial wavefunction. Under this definition, the ground state weight Z_0 can then be approximated as the expectation value of the density operator $e^{\beta H}$

$$Z_0 = \langle \psi_0 | \hat{I} | \psi_0 \rangle = \lim_{\beta \rightarrow \infty} \langle \psi_T | e^{-\beta H} | \psi_T \rangle \quad (13)$$

To save some typing effort in what follows, define a "trial weight" \mathcal{Z}_T , such that:

$$\mathcal{Z}_T \equiv \langle \psi_T | e^{-\beta H} | \psi_T \rangle \quad (14)$$

The trial weight, and hence the ground state weight can be sampled via Path Integral Monte Carlo (PIMC). To do this, \mathcal{Z}_T has to be rewritten in the path integral formulation.

Due to completeness of the wavefunction, $|\psi_T\rangle$ can be expanded as:

$$|\psi_T\rangle = \sum_{\alpha} C_{\alpha} |\alpha\rangle \quad (15)$$

where the C_{α} 's are complex expansion coefficients, and the sum is carried over

all possible states $|\alpha\rangle$ of a conveniently chosen basis. For the Bose-Hubbard model, the basis that we will "conveniently choose" is the set of all configurations of N particles on a lattice of L sites (i.e, the Fock states). The trial weight becomes:

$$\mathcal{Z}_{\mathcal{T}} = \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* \langle \alpha_0 | e^{-\beta H} | \alpha_1 \rangle \quad (16)$$

where each of the sums is carried over all Fock states. The state labels could be chosen arbitrarily, but with the power of hindsight, numerical subscripts have been used to distinguish the bra-ket pair, as it will ease the notation in the following steps. Nevertheless, keep in mind that both α_0 and α_1 will span the space of all Fock configurations.

In the interaction picture of quantum mechanics, the density operator can be expressed as:

$$e^{-\beta H} = e^{-\beta H_0} T_{\tau} e^{-\int_0^{\beta} d\tau_1 H_1(\tau_1)} \quad (17)$$

where T_{τ} is the imaginary time-ordering operator and $\tau = it/\hbar$. The integrand is the interaction picture representation of the off-diagonal part of the Hamiltonian (kinetic energy operator):

$$H_1(\tau_1) = e^{\tau_1 H_0} H_1 e^{-\tau_1 H_0} \quad (18)$$

where H_1 is the off-diagonal part and H_0 , the diagonal. The time-ordered exponential can be Taylor expanded as:

$$e^{-\int_0^{\beta} d\tau_1 H_1(\tau_1)} = \mathbb{1} - \int_0^{\beta} H_1(\tau_1) d\tau_1 + \int_0^{\beta} d\tau_1 \int_0^{\tau_1} d\tau_2 H_1(\tau_1) H_1(\tau_2) - \dots \quad (19)$$

With this Taylor expansion, the trial weight \mathcal{Z}_T becomes:

$$\mathcal{Z}_T = \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* e^{-\epsilon_{\alpha_0} \beta} \langle \alpha_0 | (\mathbb{1} - \int_0^\beta H_1(\tau_1) d\tau_1 + \int_0^\beta d\tau_1 \int_0^{\tau_1} d\tau_2 H_1(\tau_1) H_1(\tau_2) - \dots) | \alpha_1 \rangle (20)$$

By evaluating the first few terms of the expansion, a generating formula for the weight \mathcal{Z}_T can be obtained.

The zeroth-order term is:

$$\begin{aligned} \mathcal{Z}_T^{(0)} &= \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* e^{-\epsilon_{\alpha_0} \beta} \langle \alpha_0 | \mathbb{1} | \alpha_1 \rangle \\ &= \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* e^{-\epsilon_{\alpha_0} \beta} \langle \alpha_0 | \alpha_1 \rangle \\ &= \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* e^{-\epsilon_{\alpha_0} \beta} \delta_{\alpha_0, \alpha_1} \\ \mathcal{Z}_T^{(0)} &= \sum_{\alpha_0} |C_{\alpha_0}|^2 e^{-\epsilon_{\alpha_0} \beta} \end{aligned} \quad (21)$$

First-order term:

$$\begin{aligned} \mathcal{Z}_T^{(1)} &= - \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* \int_0^\beta d\tau_1 e^{-\epsilon_{\alpha_0} \beta} \langle \alpha_0 | H_1(\tau_1) | \alpha_1 \rangle \\ &= - \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* \int_0^\beta d\tau_1 e^{-\epsilon_{\alpha_0} \beta} \langle \alpha_0 | e^{\tau_1 H_0} H_1 e^{-\tau_1 H_0} | \alpha_1 \rangle \\ &= - \sum_{\alpha_0, \alpha_1} C_{\alpha_0} C_{\alpha_1}^* \int_0^\beta d\tau_1 e^{-\epsilon_{\alpha_0} (\beta - \tau_1)} e^{-\epsilon_1 \tau_1} \langle \alpha_0 | H_1 | \alpha_1 \rangle \end{aligned} \quad (22)$$

Second-order term:

$$\begin{aligned}
\mathcal{Z}_T^{(2)} &= \sum_{\alpha_0, \alpha_2} C_{\alpha_0} C_{\alpha_2}^* \int_0^\beta d\tau_1 \int_0^{\tau_1} d\tau_2 e^{-\epsilon_0 \beta} \langle \alpha_0 | H_1(\tau_1) H_1(\tau_2) | \alpha_2 \rangle \\
&= \sum_{\alpha_0, \alpha_1, \alpha_2} C_{\alpha_0} C_{\alpha_2}^* \int_0^\beta d\tau_1 \int_0^{\tau_1} d\tau_2 e^{-\epsilon_0 \beta} \langle \alpha_0 | H_1(\tau_1) | \alpha_1 \rangle \langle \alpha_1 | H_1(\tau_2) | \alpha_2 \rangle \\
&= \sum_{\alpha_0, \alpha_1, \alpha_2} C_{\alpha_0} C_{\alpha_2}^* \int_0^\beta d\tau_1 \int_0^{\tau_1} d\tau_2 e^{-\epsilon_0 \beta} \langle \alpha_0 | e^{\tau_1 H_0} H_1 e^{-\tau_1 H_0} | \alpha_1 \rangle \langle \alpha_1 | e^{\tau_2 H_0} H_1 e^{-\tau_2 H_0} | \alpha_2 \rangle \\
\mathcal{Z}_T^{(2)} &= \sum_{\alpha_0, \alpha_1, \alpha_2} C_{\alpha_0} C_{\alpha_2}^* \int_0^\beta d\tau_1 \int_0^{\tau_1} d\tau_2 e^{-\epsilon_{\alpha_0}(\beta - \tau_1)} e^{-\epsilon_1(\tau_1 - \tau_2)} e^{-\epsilon_2 \tau_2} \langle \alpha_0 | H_1 | \alpha_1 \rangle \langle \alpha_1 | H_1 | \alpha_2 \rangle
\end{aligned} \tag{23}$$

Proceeding in similar fashion, the Q -th term of \mathcal{Z}_T can be generated:

$$\mathcal{Z}_T^{(Q)} = (-1)^Q \sum_{\alpha_0, \alpha_1, \dots, \alpha_Q} C_{\alpha_0} C_{\alpha_Q}^* e^{-\epsilon_{\alpha_Q} \tau_Q} \prod_{q=1}^Q \int_0^{\tau_{q-1}} d\tau_q e^{-\epsilon_{\alpha_{q-1}}(\tau_{q-1} - \tau_q)} \langle \alpha_{q-1} | H_1 | \alpha_q \rangle \tag{24}$$

where $Q \geq 1$ and $\tau_0 = \beta$. Summing $\mathcal{Z}_T^{(Q)}$ over all Q from 0 to ∞ , the full expression for \mathcal{Z}_T is obtained:

$$\mathcal{Z}_T = \mathcal{Z}_T^{(0)} + \sum_{Q=0}^{\infty} \mathcal{Z}_T^{(Q)} \tag{25}$$

Recall from Eq. (13) that:

$$\mathcal{Z}_0 = \langle \psi_0 | \hat{I} | \psi_0 \rangle = \lim_{\beta \rightarrow \infty} \langle \psi_T | e^{-\beta H} | \psi_T \rangle = \lim_{\beta \rightarrow \infty} \mathcal{Z}_T \tag{26}$$

Thus, the weight \mathcal{Z}_0 approximately becomes:

$$\begin{aligned}
\mathcal{Z}_0 &\approx \sum_{\alpha_0} |C_{\alpha_0}|^2 e^{-\epsilon_0 \beta} \\
&+ \sum_{Q=1}^{\infty} (-1)^Q \sum_{\alpha_0, \alpha_1, \dots, \alpha_Q} C_{\alpha_0} C_{\alpha_Q}^* e^{-\epsilon_{\alpha_Q} \tau_Q} \prod_{q=1}^Q \int_0^{\tau_{q-1}} d\tau_q e^{-\epsilon_{\alpha_{q-1}}(\tau_{q-1} - \tau_q)} \langle \alpha_{q-1} | H_1 | \alpha_q \rangle
\end{aligned} \tag{27}$$

where the off-diagonal elements $\langle \alpha_{q-1} | H_1 | \alpha_q \rangle$ are given by Eq. (11) and the

diagonal elements ϵ_α are given by Eq. (3) and the approximation becomes exact in the limit $\beta \rightarrow \infty$.

NEED TO INCLUDE DERIVATION

"Worm kinetic energy"

$$H_1^{(worm)} = \sum_i (\eta_i b_i^\dagger + \eta_i^* b_i) \quad (28)$$

"Worm kinetic energy" matrix elements

$$\langle \alpha_p | H_1^{(worm)} | \alpha_q \rangle = \sum_i \eta_i \sqrt{n_i + 1} \delta_{m_i, n_i + 1} + \eta_i^* \sqrt{n_i} \delta_{m_i, n_i - 1} \prod_{k \neq i} \delta_{m_k, n_k} \quad (29)$$

3 Metropolis sampling

In this section, a brief walkthrough of the Metropolis-Hastings algorithm [1] will be given. First, the process of obtaining acceptance ratios that satisfy the Principle of Detailed Balance will be given. Then, the Metropolis-Hastings algorithm will be stated.

3.1 Deriving Metropolis acceptance ratios

The Principle of Detailed Balance states that in an equilibrium system, the number of transitions from a configuration x to a configuration x' , must be equal to the number of transitions from x' to x . This principle can be expressed as:

$$p(x)P(x \rightarrow x') = p(x')P(x' \rightarrow x) \quad (30)$$

where $p(x)$ is the steady state probability distribution, and $P(x \rightarrow x')$ is the probability of transitioning to x' conditional to original state x . Solving for the ratio of transition probabilities:

$$\frac{P(x \rightarrow x')}{P(x' \rightarrow x)} = \frac{p(x')}{p(x)} \quad (31)$$

The transition probability can be rewritten as a factor of proposal and acceptance probabilities:

$$P(x \rightarrow x') = g(x \rightarrow x')A(x \rightarrow x') \quad (32)$$

where $g(x \rightarrow x')$ is the probability of proposing an update that takes the state of the system from x to x' , and $A(x \rightarrow x')$ is the probability of accepting such update. Substituting the factored transition probabilities into Eq. (31):

$$\frac{g(x \rightarrow x')A(x \rightarrow x')}{g(x' \rightarrow x)A(x' \rightarrow x)} = \frac{p(x')}{p(x)} \quad (33)$$

The ratio of acceptance probabilities becomes:

$$\frac{A(x \rightarrow x')}{A(x' \rightarrow x)} = \frac{p(x')g(x' \rightarrow x)}{p(x)g(x \rightarrow x')} \equiv R \quad (34)$$

Depending on the value of R , there are various ways in which the acceptance probabilities can be chosen, such that they satisfy Eq. (34).

If $R \leq 1$:

$$A(x \rightarrow x') = R \text{ and } A(x' \rightarrow x) = 1$$

If $R > 1$:

$$A(x \rightarrow x') = 1 \text{ and } A(x' \rightarrow x) = 1/R$$

The above selections can be summarized as:

$$A(x \rightarrow x') = \min\{1, R\} \text{ and } A(x' \rightarrow x) = \min\{1, 1/R\} \quad (35)$$

As part of Metropolis sampling, proposed configurations will be accepted with the probabilities in Eq. (35).

3.2 The Metropolis-Hastings algorithm

$$\pi \approx 3.14 \dots \quad (36)$$

Essentially, Eq. (36) is an approximation. In the incredible paper by Prokockokokok [2], the authors play some Pokemon. In the Metropolis paper [1]. On the Brian Keng's website [3]

4 Lattice Worm Algorithm (WA) Updates

In this section, a set of ergodic lattice worm updates is introduced. First, an explanation of how each update changes the system is given.. Then, a walkthrough of the decisions that comprise each update will be given, including the probabilities for each desicion's outcome. Finally, the Metropolis conditions are derived or the direct sampling explained, depending on the update.

4.1 Insert/Delete worm

The insert worm update creates a particle on a flat region of a worldline, then destroys it after a certain time, also inside the same flat region. Formally, the particle is created by acting on the state at that imaginary time with the bosonic creation operator (worm tail) and annihilated by acting with the bosonic annihilation operator (worm head). An antiworm can instead be inserted by first inserting the worm head and then the tail. In other words, an antiworm will first annihilate a particle and create one at a later time inside the flat region. The update proceeds as follows:

0. Do with probability p_{iw} **What is this actually? There's already an attempt and acceptance prob.**
1. Randomly choose an integer $i \in [0, L - 1]$, where L is the number of sites on the lattice. The i^{th} site will be chosen wih probability $p_i = 1/L$.
2. Randomly choose an integer $f \in [0, F - 1]$, where F is the number of flat regions on site i . The f^{th} flat region will be chosen with probability $p_f = 1/F$

3. Count the number of particles n_{flat} on the flat region and check if inserting an antiworm is possible:

(a) If $n_{flat} = 0$: Only a worm can be inserted with probability $p_{type} = 1$

(b) Else: A worm or antiworm can be inserted with probability $p_{type} = 1/2$

4. Randomly choose a real number $\Delta\tau_{worm} = rand() * \Delta\tau_{flat}$, where $\Delta\tau_{flat}$ is the length of the flat region and $rand()$ is a random number from the uniform distribution in the interval $[0, 1)$. The probability of the worm being of length $\Delta\tau_{worm}$ is $p_{len} = 1/\Delta\tau_{flat}$

5. Randomly select a real number $\tau = \tau_{min} + rand() * (\Delta\tau_{flat} - \Delta\tau_{worm})$, where τ_{min} is the lower bound of the flat region. The probability of inserting the worm (antiworm) tail (head) at τ is $p_{\tau} = 1/(\Delta\tau_{flat} - \Delta\tau_{worm})$.

6. Calculate the ratio of weights of configurations pre and post worm insertion:

$$\frac{W_+}{W_-} = \eta^2 e^{-\Delta\tau_{worm} \Delta V} \quad (37)$$

where W_+ and W_- are the weights of the configuration with a worm and no worm, respectively, η is the worm fugacity and ΔV is the change in potential energy pre and post worm insertion.

-Derive ΔV here for insert worm and antiworm

-Give a similar walkthrough of the delete worm update

-Derive the Metropolis condition from Detailed Balance

$$\pi \approx 3.14159... \quad (38)$$

4.2 Insert/Delete ground state (gs) worm

These updates will be very similar to the insert/delete worm update. The main difference is that these are particular to the case of zero temperature. Remember that at zero temperature, imaginary time is not subject to periodic boundary conditions, as is the case at finite temperature. Instead, open boundary conditions are imposed in the imaginary time direction, while keeping the space direction periodic.

The open boundary conditions will allow now for the insertion of worms that will have one of its ends go past either the $\tau = 0$ or the $\tau = \beta$ boundaries. These worms that go past the imaginary time boundaries will be called ground state (gs) worms. In practice, inserting a gs-worm will look like inserting only one worm end to the worldline configuration at a time, and analogously for deletion.

4.2.1 Insert gs-worm

At $T = 0$, a ground state (gs) worm or antiworm can be inserted. One of the ends of this worm will lie either past the $\tau = 0$ or the $\tau = \beta$ boundaries. The other end, will lie in the flat region preceding this boundary (i.e, the first or the last flat region). In practice, this will look like inserting a single worm end to the worldline configuration at these flat regions. This worm end can only be inserted if there are no worm ends or one worm end present. The update proceeds as follows:

0. Do with probability p_{gsiw}
1. Choose site with probability $p = 1/L$
2. Choose first or last flat:
 - (a) if $n_{flats} = 1$: treat the flat as first or last (depending on the implementation) with $p = 1$

(b) else: $p = 1/2$

3. Choose gs-worm or gs-antiworm: if $n_i = 0$: can only insert worm with prob. $p = 1$ else: choose either with $p = 1/2$
4. Choose the insertion time of the worm end $\tau = \tau_{min} + rand()(\tau_{max} - \tau_{min})$ with probability $p = 1/(\tau_{max} - \tau_{min})$

4.2.2 Delete gs-worm

The Delete gs-worm update can only be done if there's at least one worm end on the worldline configuration. A random worm end is chosen and then deleted. The update proceeds as follows:

0. Do with probability p_{gsdw}
1. Choose the worm end to be deleted with probability $p = 1/2$

4.2.3 Acceptance probabilities of gs-worm insert/delete

From detailed balance:

$$w_- p_{gsiw}^{att} p_{gsiw}^{acc} = W_+ p_{gsdw}^{att} p_{gsdw}^{acc} \quad (39)$$

Solving for the ratio of acceptance probabilities:

$$\frac{p_{gsiw}^{acc}}{p_{gsdw}^{acc}} = \frac{p_{gsdw}^{att} W_+}{p_{gsiw}^{att} W_-} \quad (40)$$

Inserting the attempt probabilities:

$$\frac{p_{gsiw}^{att}}{p_{gsdw}^{att}} = \frac{p_{gsdw}^{att}}{p_{gsiw}^{att}} \frac{L(\tau_{max} - \tau_{min})}{p_{flat}} \frac{W_+}{W_-} \quad (41)$$

where p_{flat} is the probability of choosing either the first or last flat region of the worldline. If there are no kinks in the worldline, $p_{flat} = 1$, and if there

are kinks, $p_{flat}=1/2$.

The ratio of weights for configurations after a gs-worm insertion and after it is deleted is:

$$\frac{W_+}{W_-} = \eta e^{-\Delta\tau\Delta V} \quad (42)$$

4.3 Timeshift

4.3.1 Forward

The timeshift update consists of moving either a worm head or tail either forward or backwards in imaginary time. The update proceeds as follow:

0. Do with probability p_{fw} :
1. Choose which worm end will move, head or tail, with probability $p = 1/2$
2. Randomly choose a real number $\tau_{new} = \tau + rand() * (\tau_{max} - \tau)$, where τ_{max} is the upper bound of the flat region delimited by the moving end and the next kink or worm end and τ is the original time of the moving worm end. The probability that the worm end will move to this time is $p = 1/(\tau_{max} - \tau)$
3. Calculate the ratio of weights of configurations pre and post timeshifting forward:

$$\frac{W_+}{W_-} = e^{-\Delta\tau_{worm}\Delta V} = 1 \quad (43)$$

where W_+ and W_- are the weights of the configuration post and pre moving forward, respectively, and ΔV is the change in potential energy. Since the number of particles on each site remains constant before and after the update, the change in potential energy is $dV = 0$ and thus the exponential becomes unity.

4.3.2 Backward

0. Do with probability p_{bw} :
1. Choose which worm end will move, head or tail, with probability $p = 1/2$
2. Randomly choose a real number $\tau_{new} = \tau_{min} + rand() * (\tau - \tau_{min})$, where τ_{min} is the lower bound of the flat region delimited by the moving end and the preceding kink or worm end and τ is the original time of the moving worm end. The probability that the worm end will move to this time is $p = 1/(\tau - \tau_{min})$
3. Calculate the ratio of weights of configurations pre and post timeshifting backward:

$$\frac{W_-}{W_+} = e^{\Delta\tau_{worm}\Delta V} = 1 \quad (44)$$

4.3.3 Detailed balance

$$p_{fw}^{att} p_{fw}^{acc} W_- = p_{fw}^{att} p_{fw}^{acc} W_+ \quad (45)$$

The attempt probabilities for the forward and backward timeshift can be read off from the description of the update. Substituting them in and solving for $p_{fw}^{acc}/p_{bw}^{acc}$:

$$\frac{p_{fw}^{acc}}{p_{bw}^{acc}} = \frac{\tau_{max} - \tau}{\tau - \tau_{min}} = R \quad (46)$$

4.4 Spaceshift before

4.5 Spaceshift after

References

- [1] N. Metropolis, A. W. Rosenbluth, M. N. Rosenbluth, A. H. Teller, and E. Teller. “Equation of state calculations by fast computing machines.” The Journal of Chemical Physics **21**, 1087 (1953).
<https://doi.org/10.1063/1.1699114>

- [2] N. V. Prokof'ev, B. V. Svistunov, and I. S. Tupitsyn. "Exact, complete, and universal continuous-time worldline monte carlo approach to the statistics of discrete quantum systems." *Journal of Experimental and Theoretical Physics* **87**, 310 (1998).
<https://doi.org/10.1134%2F1.558661>
- [3] B. Keng. "Markov chain monte carlo methods, rejection sampling and the metropolis-hastings algorithm."
<http://bjlkeng.github.io/posts/markov-chain-monte-carlo-mcmc-and-the-metropolis-hastings-algorithm/>