

Mechanisms for Generating Power-Law Size Distributions, Part 1

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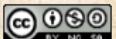
Principles of Complex Systems | @pocsvox
CSYS/MATH 300, Fall, 2018

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Vermont Advanced Computing Core | University of Vermont



Random Walks
The First Return Problem
Examples
Variable transformation
Basics
Holtsmark's Distribution
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References



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Mechanisms, Pt. 1

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On Instagram at [pratchett_the_cat](https://www.instagram.com/pratchett_the_cat/)



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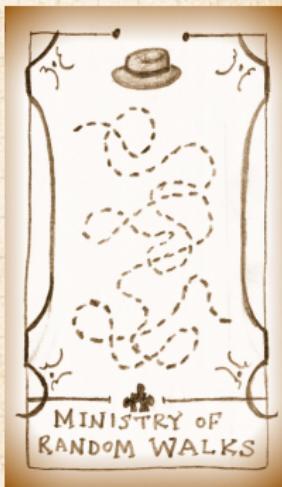
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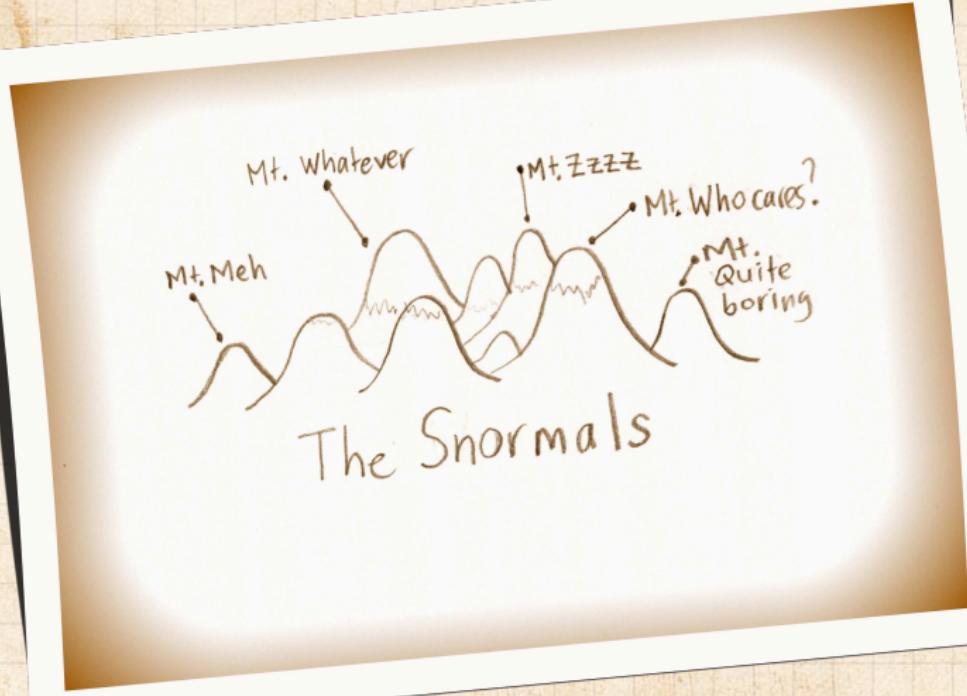
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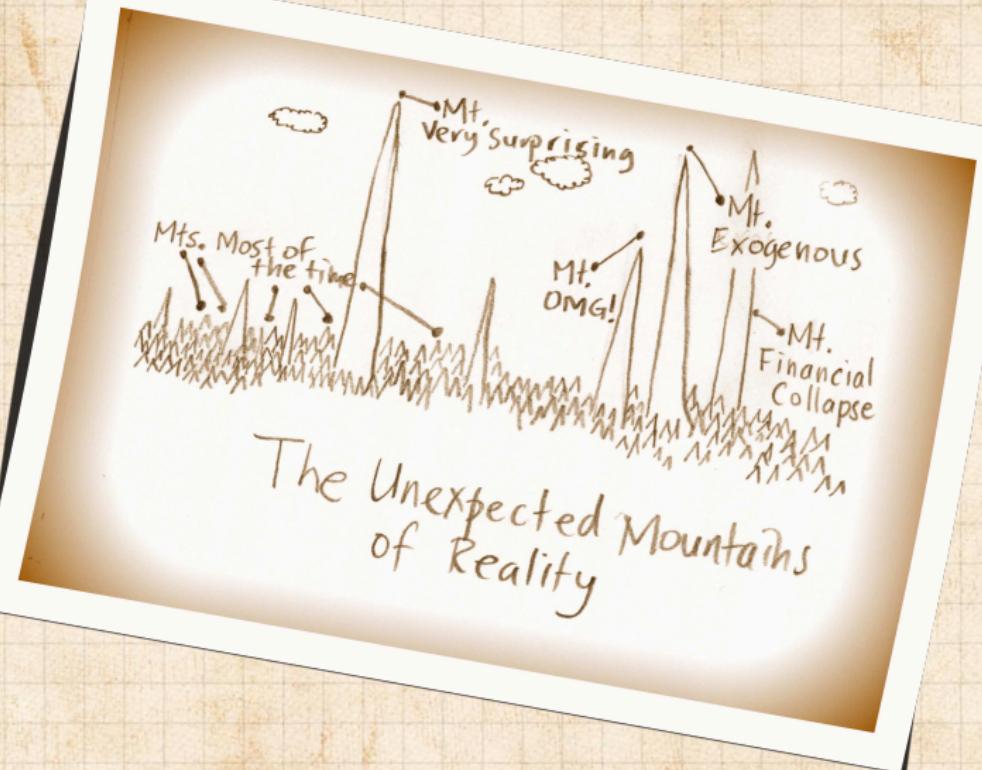
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Mechanisms:

A powerful story in the rise of complexity:

- 🎲 structure arises out of randomness.
- 🎲 Exhibit A: Random walks. ↗

The essential random walk:

- 🎲 One spatial dimension.
- 🎲 Time and space are discrete
- 🎲 Random walker (e.g., a drunk) starts at origin $x = 0$.
- 🎲 Step at time t is ϵ_t :

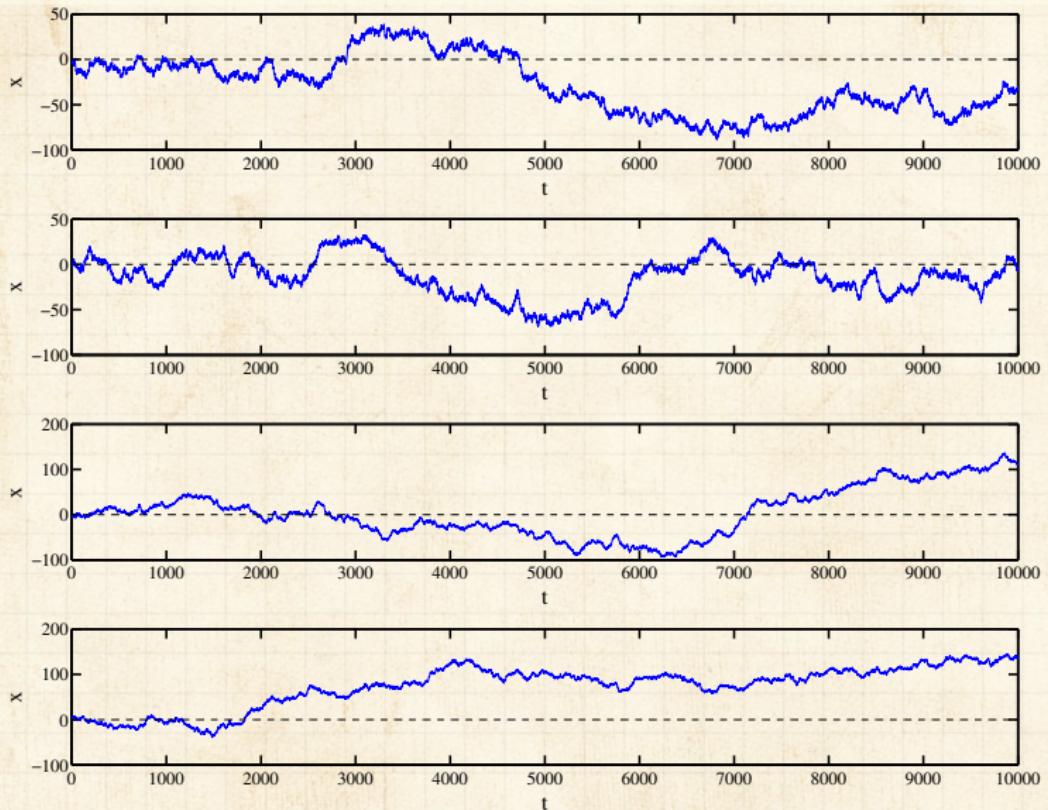
$$\epsilon_t = \begin{cases} +1 & \text{with probability } 1/2 \\ -1 & \text{with probability } 1/2 \end{cases}$$



A few random random walks:

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Random walks:

Displacement after t steps:

$$x_t = \sum_{i=1}^t \epsilon_i$$

Expected displacement:

$$\langle x_t \rangle = \left\langle \sum_{i=1}^t \epsilon_i \right\rangle = \sum_{i=1}^t \langle \epsilon_i \rangle = 0$$

- ⬢ At any time step, we 'expect' our drunkard to be back at the pub.
- ⬢ Obviously fails for odd number of steps...
- ⬢ But as time goes on, the chance of our drunkard lurching back to the pub must diminish, right?

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Variances sum: ↗*

$$\text{Var}(x_t) = \text{Var} \left(\sum_{i=1}^t \epsilon_i \right)$$

$$= \sum_{i=1}^t \text{Var}(\epsilon_i) = \sum_{i=1}^t 1 = t$$

* Sum rule = a good reason for using the variance to measure spread; only works for independent distributions.

So typical displacement from the origin scales as:

$$\sigma = t^{1/2}$$

- ⬢ A non-trivial scaling law arises out of additive aggregation or accumulation.



Stock Market randomness:

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Also known as the bean machine ↗, the quincunx (simulation) ↗, and the Galton box.



Great moments in Televised Random Walks:

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Plinko! ↗ from the Price is Right.



Random walk basics:

Counting random walks:

- Each **specific** random walk of length t appears with a chance $1/2^t$.
- We'll be more interested in how many random walks end up at the same place.
- Define $N(i, j, t)$ as # distinct walks that start at $x = i$ and end at $x = j$ after t time steps.
- Random walk must displace by $+(j - i)$ after t steps.
- Insert question from assignment 3 ↗

$$N(i, j, t) = \binom{t}{(t + j - i)/2}$$

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How does $P(x_t)$ behave for large t ?

- ⬢ Take time $t = 2n$ to help ourselves.
- ⬢ $x_{2n} \in \{0, \pm 2, \pm 4, \dots, \pm 2n\}$
- ⬢ x_{2n} is even so set $x_{2n} = 2k$.
- ⬢ Using our expression $N(i, j, t)$ with $i = 0$, $j = 2k$, and $t = 2n$, we have

$$\Pr(x_{2n} \equiv 2k) \propto \binom{2n}{n+k}$$

- ⬢ For large n , the binomial deliciously approaches the Normal Distribution of Snoredom:

$$\Pr(x_t \equiv x) \simeq \frac{1}{\sqrt{2\pi t}} e^{-\frac{x^2}{2t}}.$$

Insert question from assignment 3 ↗

- ⬢ The whole is different from the parts. #nutritious
- ⬢ See also: Stable Distributions ↗

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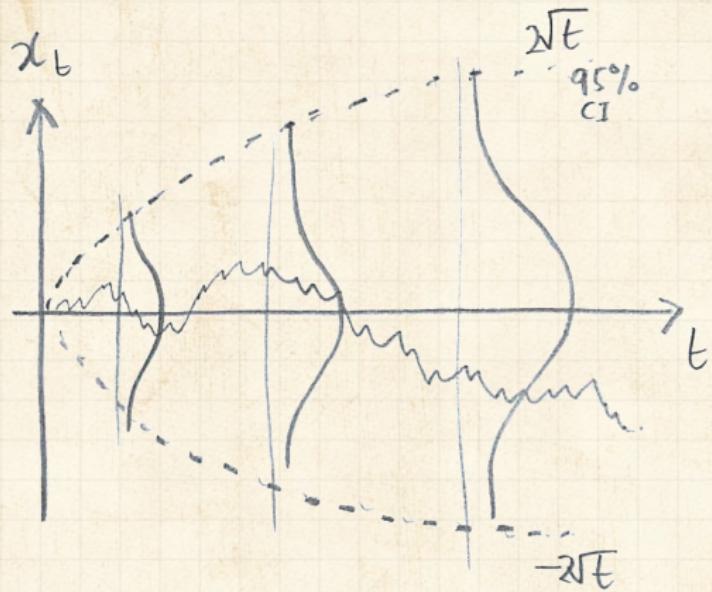
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Universality ↗ is also not left-handed:

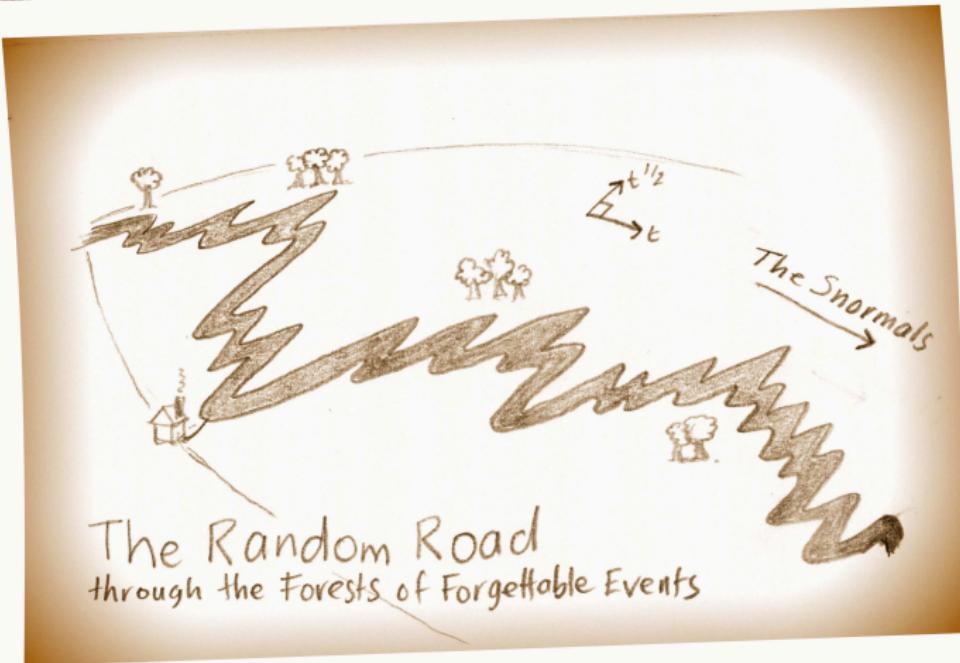
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- ⬢ This is Diffusion ↗: the most essential kind of spreading (more later).
- ⬢ View as Random Additive Growth Mechanism.





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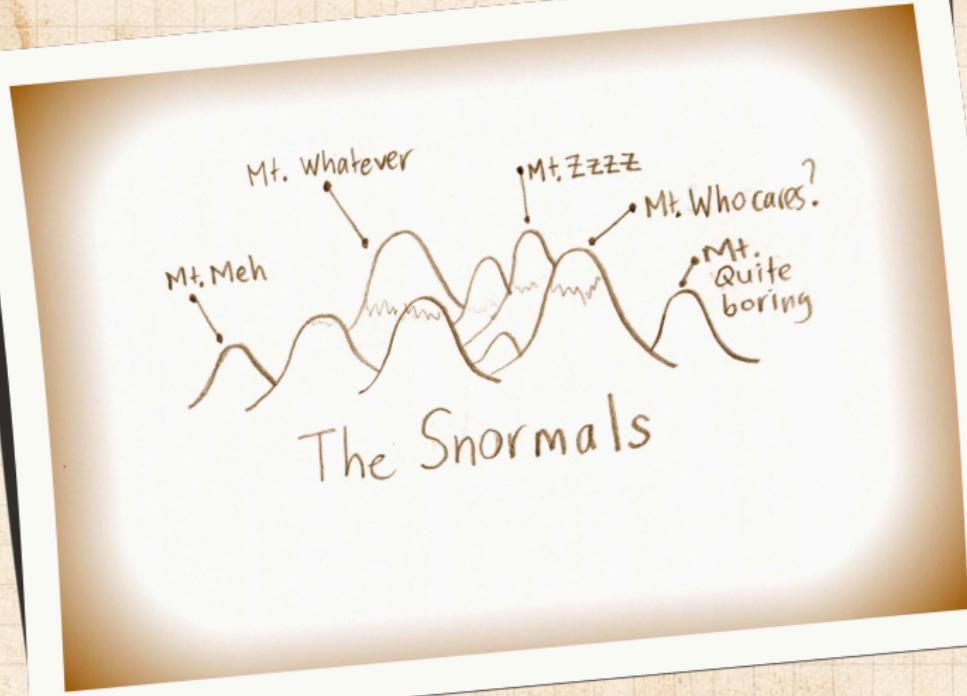
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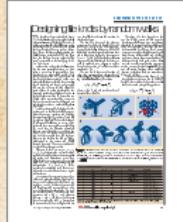
Random walks are even weirder than you might think...

- 🎲 $\xi_{r,t}$ = the probability that by time step t , a random walk has crossed the origin r times.
- 🎲 Think of a coin flip game with ten thousand tosses.
- 🎲 If you are behind early on, what are the chances you will make a comeback?
- 🎲 The most likely number of lead changes is... 0.
- 🎲 In fact: $\xi_{0,t} > \xi_{1,t} > \xi_{2,t} > \dots$
- 🎲 Even crazier:
The expected time between tied scores = ∞

See Feller, Intro to Probability Theory, Volume I [3]



Applied knot theory:



"Designing tie knots by random walks" ↗

Fink and Mao,
Nature, 398, 31–32, 1999. [4]

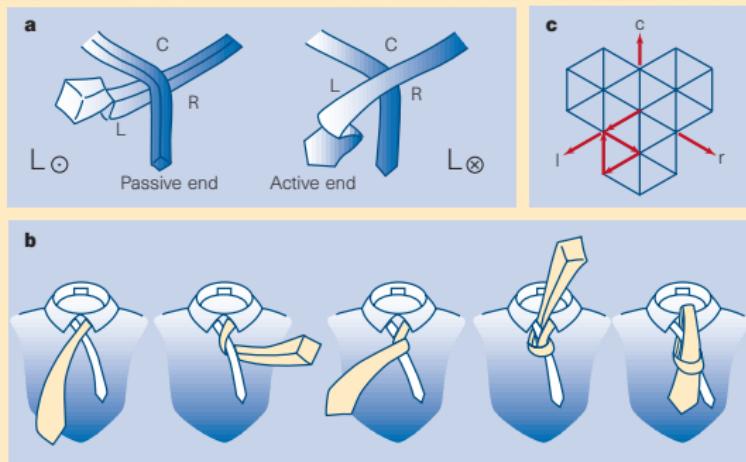


Figure 1 All diagrams are drawn in the frame of reference of the mirror image of the actual tie.
a. The two ways of beginning a knot, L_0 and \bar{L}_0 . For knots beginning with L_0 , the tie must begin inside-out. **b.** The four-in-hand, denoted by the sequence $L_0 \text{ } R_0 \text{ } L_0 \text{ } C_0 \text{ } T$. **c.** A knot may be represented by a persistent random walk on a triangular lattice. The example shown is the four-in-hand, indicated by the walk $1\bar{1}1\bar{1}\bar{1}$.

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Applied knot theory:

Table 1 **Aesthetic tie knots**

h	γ	γ/h	$K(h, \gamma)$	s	b	Name	Sequence
3	1	0.33	1	0	0		$L \circ R \circ C \circ T$
4	1	0.25	1	-1	1	Four-in-hand	$L \circ R \circ L \circ C \circ T$
5	2	0.40	2	-1	0	Pratt knot	$L \circ C \circ R \circ L \circ C \circ T$
6	2	0.33	4	0	0	Half-Windsor	$L \circ R \circ C \circ L \circ R \circ C \circ T$
7	2	0.29	6	-1	1		$L \circ R \circ L \circ C \circ R \circ L \circ C \circ T$
7	3	0.43	4	0	1		$L \circ C \circ R \circ C \circ L \circ R \circ C \circ T$
8	2	0.25	8	0	2		$L \circ R \circ L \circ C \circ R \circ L \circ R \circ C \circ T$
8	3	0.38	12	-1	0	Windsor	$L \circ C \circ R \circ L \circ C \circ R \circ L \circ C \circ T$
9	3	0.33	24	0	0		$L \circ R \circ C \circ L \circ R \circ C \circ L \circ R \circ C \circ T$
9	4	0.44	8	-1	2		$L \circ C \circ R \circ C \circ L \circ C \circ R \circ L \circ C \circ T$

Knots are characterized by half-winding number h , centre number γ , centre fraction γ/h , knots per class $K(h, \gamma)$, symmetry s , balance b , name and sequence.

➊ $h = \text{number of moves}$

➋ $s = \sum_{i=1}^h x_i$ where $x = -1$ for L and $+1$ for R .

➌ $\gamma = \text{number of center moves}$

➍ $b = \frac{1}{2} \sum_{i=2}^{h-1} |\omega_i + \omega_{i-1}|$ where $\omega = \pm 1$ represents winding direction.

➎ $K(h, \gamma) = 2^{\gamma-1} \binom{h-\gamma-2}{\gamma-1}$



Random walks #crazytownbananapants

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The problem of first return:

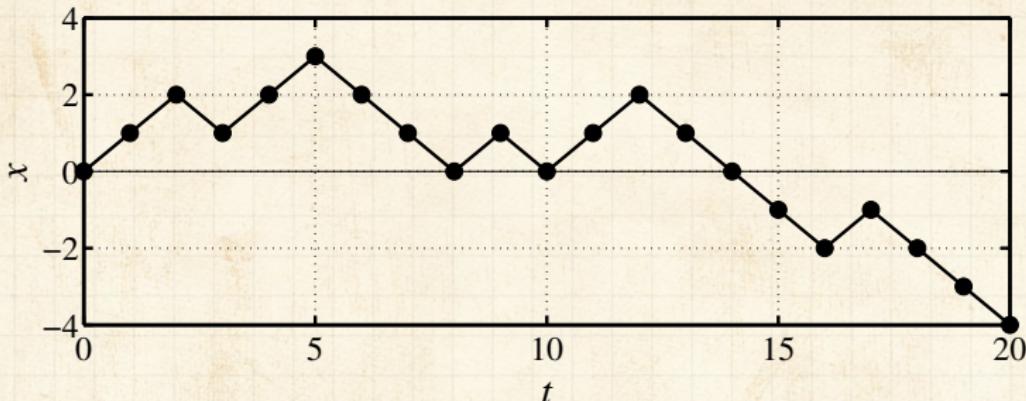
- What is the probability that a random walker in one dimension returns to the origin for the first time after t steps?
- Will our drunkard always return to the origin?
- What about higher dimensions?

Reasons for caring:

- We will find a power-law size distribution with an interesting exponent.
- Some physical structures may result from random walks.
- We'll start to see how different scalings relate to each other.

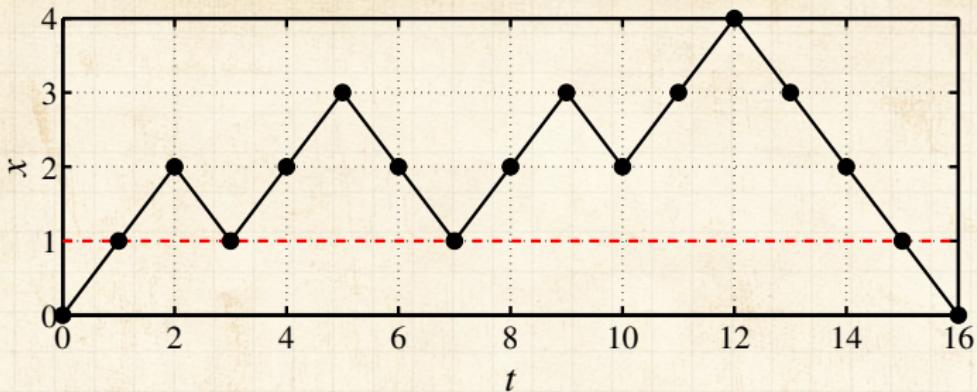


For random walks in 1-d:



- cube A **return** to origin can only happen when $t = 2n$.
- cube In example above, returns occur at $t = 8, 10$, and 14 .
- cube Call $P_{\text{fr}(2n)}$ the probability of **first return** at $t = 2n$.
- cube Probability calculation \equiv Counting problem (combinatorics/statistical mechanics).
- cube Idea: Transform first return problem into an easier return problem.





- 🎲 Can assume drunkard first lurches to $x = 1$.
- 🎲 Observe walk first returning at $t = 16$ stays at or above $x = 1$ for $1 \leq t \leq 15$ (dashed red line).
- 🎲 Now want walks that can return many times to $x = 1$.
- 🎲 $P_{\text{fr}}(2n) = 2 \cdot \frac{1}{2} Pr(x_t \geq 1, 1 \leq t \leq 2n - 1, \text{ and } x_1 = x_{2n-1} = 1)$
- 🎲 The $\frac{1}{2}$ accounts for $x_{2n} = 2$ instead of 0.
- 🎲 The 2 accounts for drunkards that first lurch to $x = -1$.



Counting first returns:

Approach:

- cube Move to counting numbers of walks.
- cube Return to probability at end.
- cube Again, $N(i, j, t)$ is the # of possible walks between $x = i$ and $x = j$ taking t steps.
- cube Consider all paths starting at $x = 1$ and ending at $x = 1$ after $t = 2n - 2$ steps.
- cube Idea: If we can compute the number of walks that hit $x = 0$ at least once, then we can subtract this from the total number to find the ones that maintain $x \geq 1$.
- cube Call walks that drop below $x = 1$ excluded walks.
- cube We'll use a method of images to identify these excluded walks.

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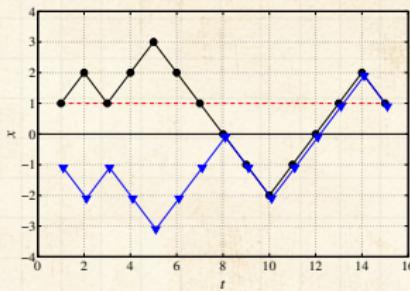
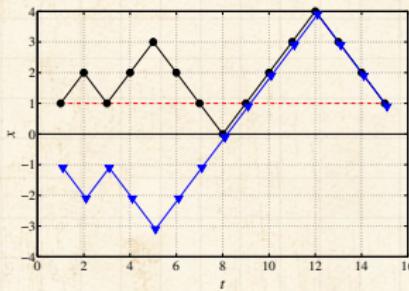
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Examples of excluded walks:



Key observation for excluded walks:

- ⬢ For any path starting at $x=1$ that hits 0, there is a unique matching path starting at $x=-1$.
- ⬢ Matching path first mirrors and then tracks after first reaching $x=0$.
- ⬢ # of t -step paths starting and ending at $x=1$ and hitting $x=0$ at least once
 = # of t -step paths starting at $x=-1$ and ending at $x=1$ = $N(-1, 1, t)$
- ⬢ So $N_{\text{first return}}(2n) = N(1, 1, 2n-2) - N(-1, 1, 2n-2)$



Probability of first return:

Insert question from assignment 3 ↗ :

Find

$$N_{\text{fr}}(2n) \sim \frac{2^{2n-3/2}}{\sqrt{2\pi n^{3/2}}}.$$

Normalized number of paths gives probability.

Total number of possible paths = 2^{2n} .



$$P_{\text{fr}}(2n) = \frac{1}{2^{2n}} N_{\text{fr}}(2n)$$

$$\simeq \frac{1}{2^{2n}} \frac{2^{2n-3/2}}{\sqrt{2\pi n^{3/2}}}$$

$$= \frac{1}{\sqrt{2\pi}} (2n)^{-3/2} \propto t^{-3/2}.$$



- 🎲 We have $P(t) \propto t^{-3/2}$, $\gamma = 3/2$.
- 🎲 Same scaling holds for continuous space/time walks.
- 🎲 $P(t)$ is normalizable.
- 🎲 Recurrence: Random walker always returns to origin
- 🎲 But mean, variance, and all higher moments are infinite. **#totalmadness**
- 🎲 Even though walker must return, expect a long wait...
- 🎲 One moral: Repeated gambling against an infinitely wealthy opponent must lead to ruin.

Higher dimensions ↗:

- 🎲 Walker in $d = 2$ dimensions must also return
- 🎲 Walker may not return in $d \geq 3$ dimensions
- 🎲 Associated genius: George Pólya ↗



Random walks

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Mechanisms, Pt. 1

On finite spaces:

- ⬢ In any finite homogeneous space, a random walker will visit every site with equal probability
- ⬢ Call this probability the **Invariant Density** of a dynamical system
- ⬢ Non-trivial Invariant Densities arise in chaotic systems.

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On networks:

- ⬢ On networks, a random walker visits each node with frequency \propto node degree #groovy
- ⬢ Equal probability still present: walkers traverse **edges** with equal frequency.



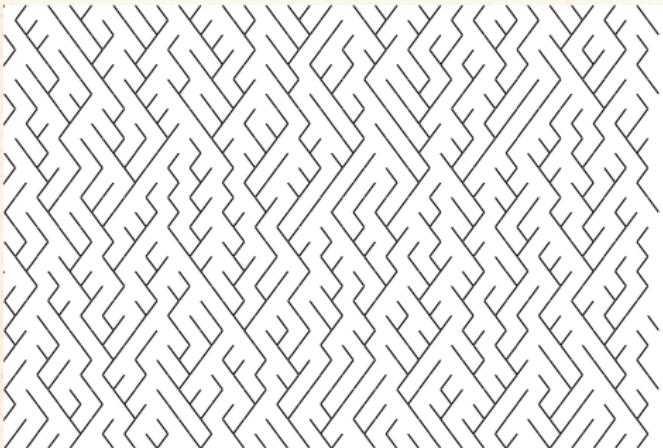
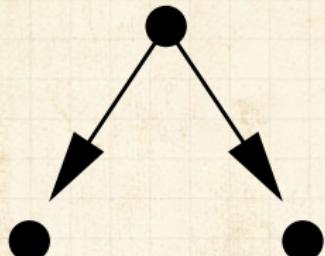
#totallygroovy



Scheidegger Networks [9, 2]

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- _RANDOM DIRECTIONALITY_ Random directed network on triangular lattice.
- _RANDOM DIRECTIONALITY_ Toy model of real networks.
- _RANDOM DIRECTIONALITY_ 'Flow' is southeast or southwest with equal probability.



Scheidegger networks

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- Creates basins with random walk boundaries.
- Observe that subtracting one random walk from another gives random walk with increments:

$$\epsilon_t = \begin{cases} +1 & \text{with probability } 1/4 \\ 0 & \text{with probability } 1/2 \\ -1 & \text{with probability } 1/4 \end{cases}$$

- Random walk with probabilistic pauses.
- Basin termination = first return random walk problem.
- Basin length ℓ distribution: $P(\ell) \propto \ell^{-3/2}$
- For real river networks, generalize to $P(\ell) \propto \ell^{-\gamma}$.

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Connections between exponents:

- 3 For a basin of length ℓ , width $\propto \ell^{1/2}$
- 3 Basin area $a \propto \ell \cdot \ell^{1/2} = \ell^{3/2}$
- 3 Invert: $\ell \propto a^{2/3}$
- 3 $d\ell \propto d(a^{2/3}) = 2/3a^{-1/3}da$
- 3 $\Pr(\text{basin area} = a)da$
 $= \Pr(\text{basin length} = \ell)d\ell$
 $\propto \ell^{-3/2}d\ell$
 $\propto (a^{2/3})^{-3/2}a^{-1/3}da$
 $= a^{-4/3}da$
 $= a^{-\tau}da$

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Connections between exponents:

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- _both basin area and length obey power law distributions
- Observed for real river networks
- Reportedly: $1.3 < \tau < 1.5$ and $1.5 < \gamma < 2$

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Generalize relationship between area and length:

- Hack's law^[5]:

$$\ell \propto a^h.$$

- For real, large networks $h \simeq 0.5$
- Smaller basins possibly $h > 1/2$ (see: allometry).
- Models exist with interesting values of h .
- Plan: Redo calc with γ , τ , and h .



Connections between exponents:

Given

$$\ell \propto a^h, P(a) \propto a^{-\tau}, \text{ and } P(\ell) \propto \ell^{-\gamma}$$

 $d\ell \propto d(a^h) = ha^{h-1}da$

 Find τ in terms of γ and h .

 $\Pr(\text{basin area} = a)da$
 $= \Pr(\text{basin length} = \ell)d\ell$
 $\propto \ell^{-\gamma}d\ell$
 $\propto (a^h)^{-\gamma}a^{h-1}da$
 $= a^{-(1+h(\gamma-1))}da$



$$\boxed{\tau = 1 + h(\gamma - 1)}$$

 Excellent example of the **Scaling Relations** found between exponents describing power laws for many systems.

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Connections between exponents:

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With more detailed description of network structure, $\tau = 1 + h(\gamma - 1)$ simplifies to:^[1]

$$\tau = 2 - h$$

and

$$\gamma = 1/h$$

- Only one exponent is independent (take h).
- Simplifies system description.
- Expect Scaling Relations where power laws are found.
- Need only characterize Universality class with independent exponents.

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Other First Returns or First Passage Times:

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Mechanisms, Pt. 1

Failure:

- ⬢ A very simple model of failure/death: [11]
- ⬢ x_t = entity's 'health' at time t
- ⬢ Start with $x_0 > 0$.
- ⬢ Entity fails when x hits 0.

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Streams

- ⬢ Dispersion of suspended sediments in streams.
- ⬢ Long times for clearing.



More than randomness

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Power-Law
Mechanisms, Pt. 1

- Can generalize to Fractional Random Walks [7, 8, 6]
- Levy flights, Fractional Brownian Motion
- See Montroll and Shlesinger for example: [6]
“On $1/f$ noise and other distributions with long tails.”
Proc. Natl. Acad. Sci., 1982.
- In 1-d, standard deviation σ scales as

$$\sigma \sim t^\alpha$$

$\alpha = 1/2$ — diffusive

$\alpha > 1/2$ — superdiffusive

$\alpha < 1/2$ — subdiffusive

- Extensive memory of path now matters...

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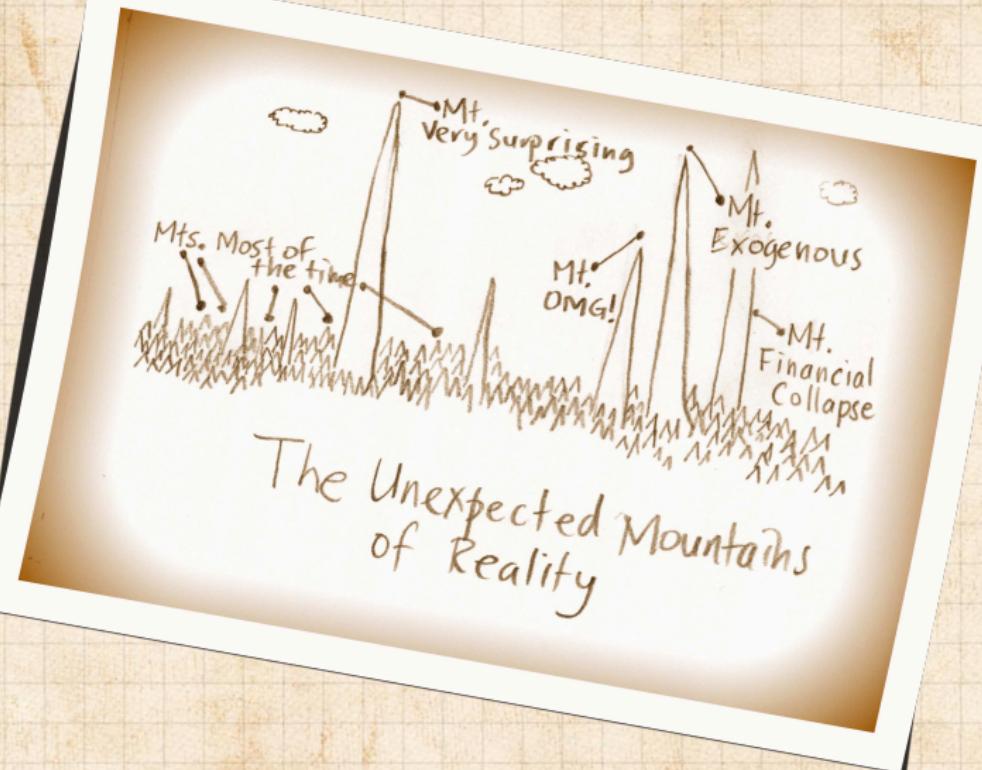
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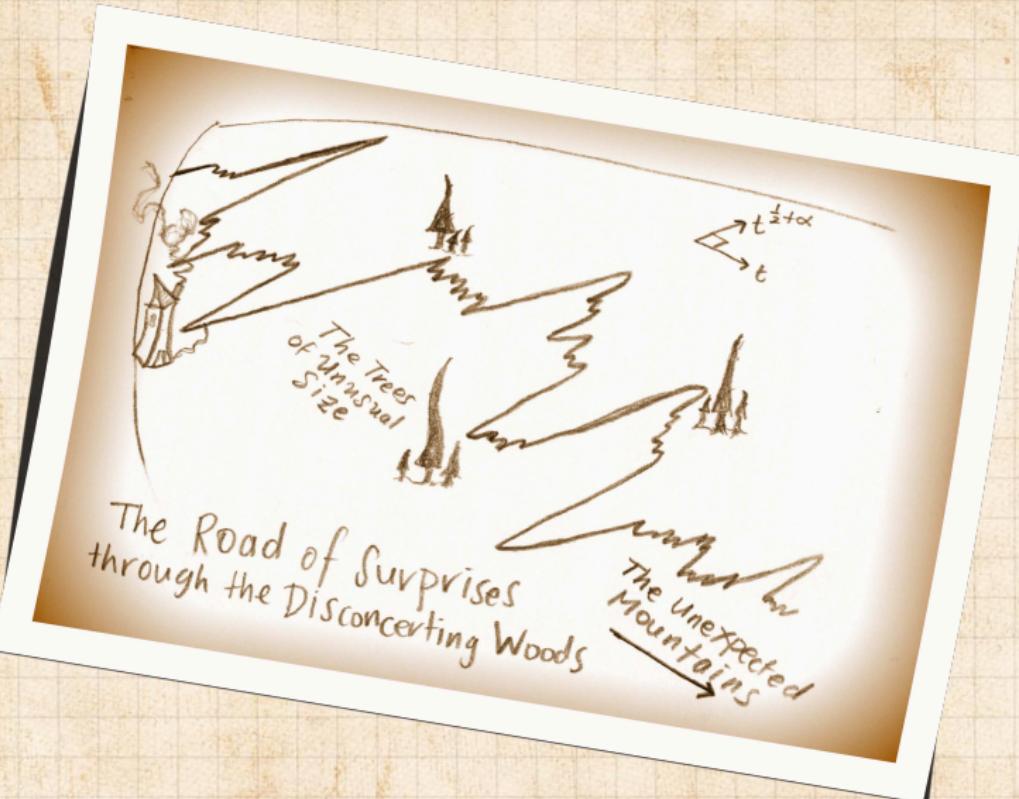
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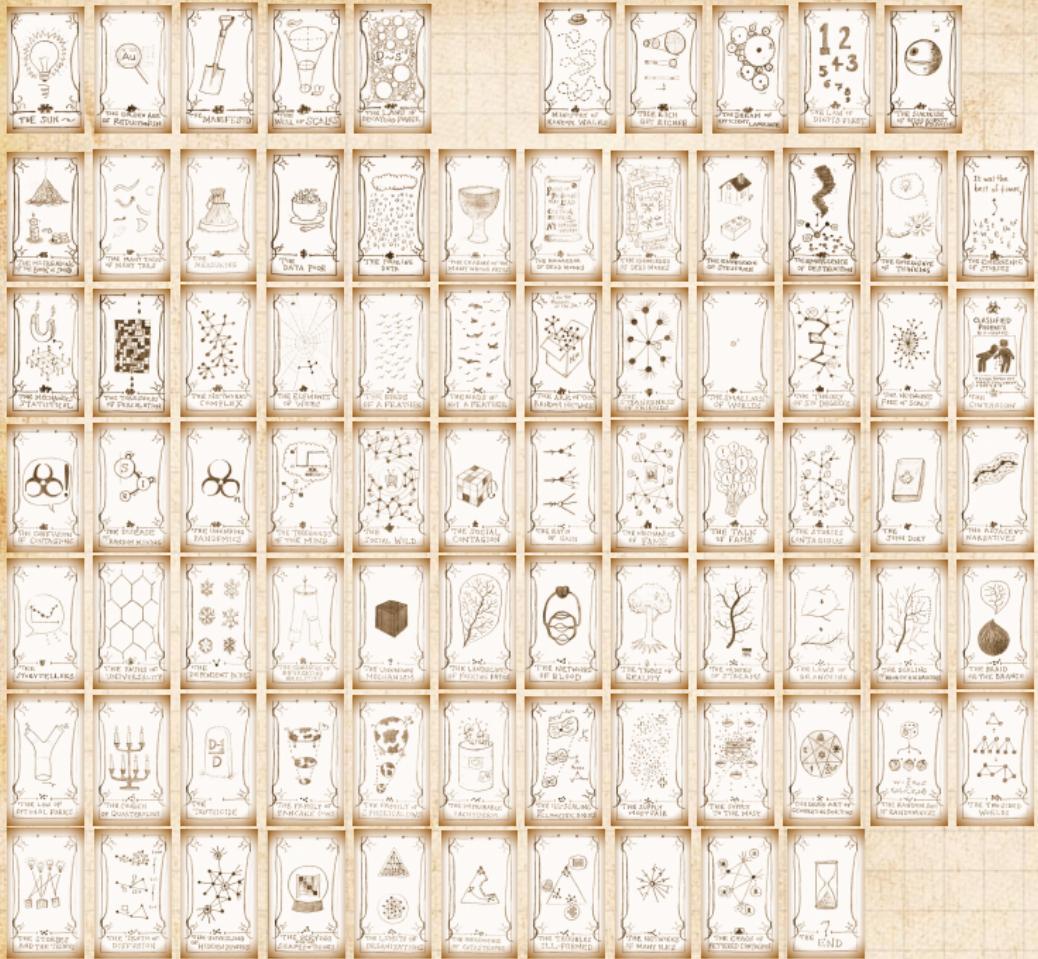
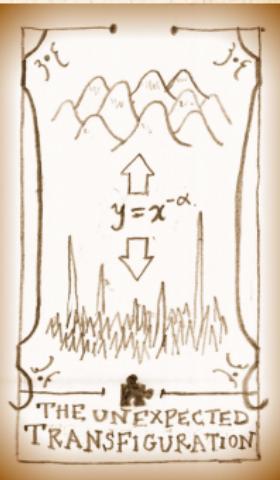
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Variable Transformation

Understand power laws as arising from

1. Elementary distributions (e.g., exponentials).
2. Variables connected by power relationships.

- Random variable X with known distribution P_x
- Second random variable Y with $y = f(x)$.

$$\begin{aligned} P_Y(y)dy &= \\ \sum_{x|f(x)=y} P_X(x)dx &= \\ \sum_{y|f(x)=y} P_X(f^{-1}(y)) \frac{dy}{|f'(f^{-1}(y))|} \end{aligned}$$

- Often easier to do by hand...

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General Example

Assume relationship between x and y is 1-1.

Power-law relationship between variables:

$$y = cx^{-\alpha}, \alpha > 0$$

Look at y large and x small



$$dy = d(cx^{-\alpha})$$

$$= c(-\alpha)x^{-\alpha-1}dx$$

invert: $dx = \frac{-1}{c\alpha}x^{\alpha+1}dy$

$$dx = \frac{-1}{c\alpha} \left(\frac{y}{c}\right)^{-(\alpha+1)/\alpha} dy$$

$$dx = \frac{-c^{1/\alpha}}{\alpha} y^{-1-1/\alpha} dy$$



Now make transformation:

$$P_y(y)dy = P_x(x)dx$$

$$P_y(y)dy = P_x\left(\left(\frac{y}{c}\right)^{-1/\alpha}\right)\overbrace{\frac{c^{1/\alpha}}{\alpha}y^{-1-1/\alpha}dy}^{dx}$$

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🎲 If $P_x(x) \rightarrow$ non-zero constant as $x \rightarrow 0$ then

$$P_x(y) \propto y^{-1-1/\alpha} \text{ as } y \rightarrow \infty.$$

🎲 If $P_x(x) \rightarrow x^\beta$ as $x \rightarrow 0$ then

$$P_y(y) \propto y^{-1-1/\alpha-\beta/\alpha} \text{ as } y \rightarrow \infty.$$



Example

Exponential distribution

Given $P_x(x) = \frac{1}{\lambda} e^{-x/\lambda}$ and $y = cx^{-\alpha}$, then

$$P(y) \propto y^{-1-1/\alpha} + O(y^{-1-2/\alpha})$$

- ➊ Exponentials arise from randomness (easy)...
- ➋ More later when we cover robustness.

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Gravity

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- Select a random point in the universe \vec{x}
- Measure the force of gravity $F(\vec{x})$
- Observe that $P_F(F) \sim F^{-5/2}$.



Matter is concentrated in stars: [10]

- ⬢ F is distributed unevenly
- ⬢ Probability of being a distance r from a single star at $\vec{x} = \vec{0}$:

$$P_r(r)dr \propto r^2 dr$$

- ⬢ Assume stars are distributed randomly in space (oops?)
- ⬢ Assume only one star has significant effect at \vec{x} .
- ⬢ Law of gravity:

$$F \propto r^{-2}$$

- ⬢ invert:

$$r \propto F^{-\frac{1}{2}}$$

- ⬢ Connect differentials: $dr \propto dF^{-\frac{1}{2}} \propto F^{-\frac{3}{2}} dF$

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Transformation:

Using $r \propto F^{-1/2}$, $dr \propto F^{-3/2} dF$, and $P_r(r) \propto r^2$



$$P_F(F) dF = P_r(r) dr$$



$$\propto P_r(\text{const} \times F^{-1/2}) F^{-3/2} dF$$



$$\propto (F^{-1/2})^2 F^{-3/2} dF$$



$$= F^{-1-3/2} dF$$



$$= F^{-5/2} dF.$$

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- Mean is finite.
- Variance = ∞ .
- A **wild** distribution.
- Upshot:** Random sampling of space usually safe but can end badly...



Doctorin' the Tardis

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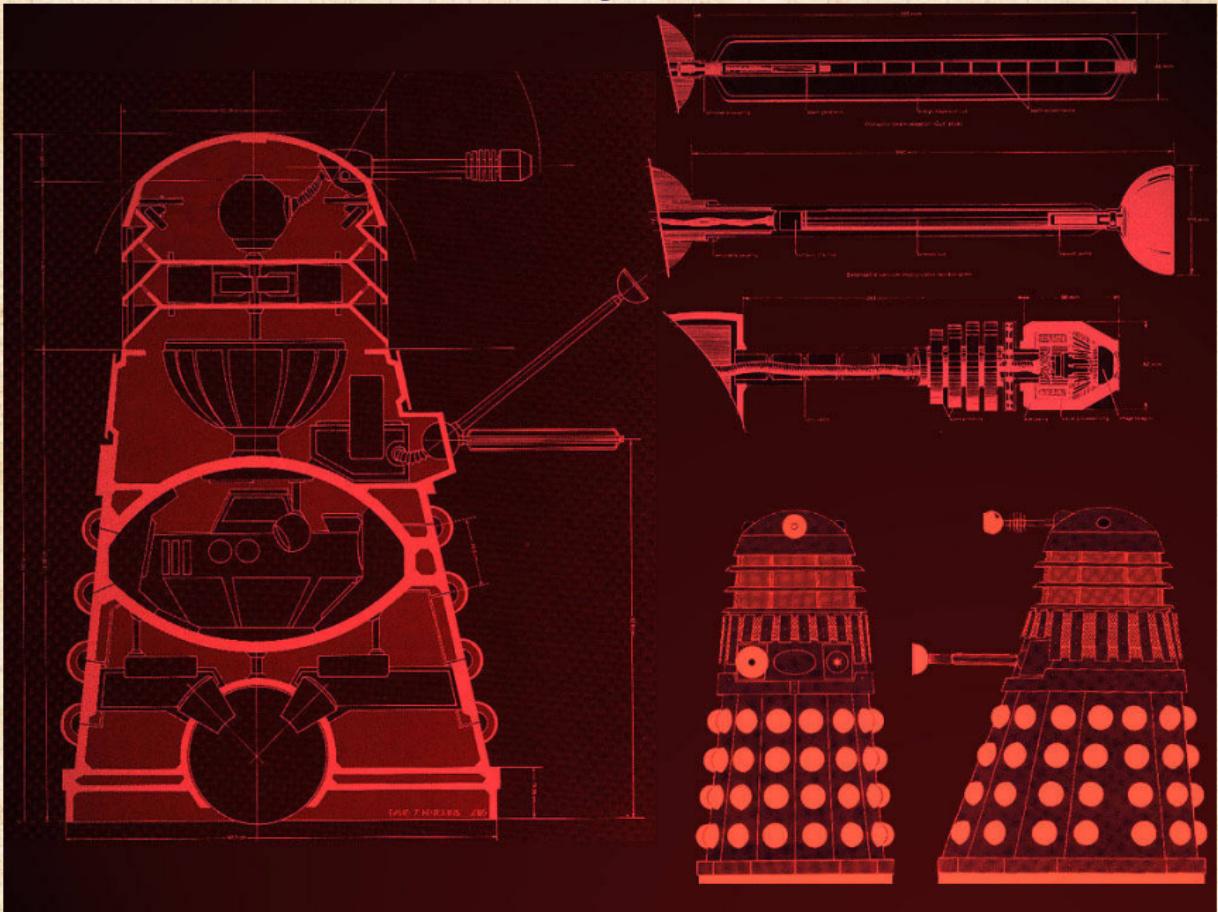
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Todo: Build Dalek army.



Extreme Caution!

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- ⬢ PLIPLO = Power law in, power law out
- ⬢ Explain a power law as resulting from another unexplained power law.
- ⬢ Yet another homunculus argument ↗ ...
- ⬢ Don't do this!!! (slap, slap)
- ⬢ MIWO = Mild in, Wild out is the stuff.
- ⬢ In general: We need mechanisms!



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