

1 Description of algorithm

1. Evolve(Configuration * config)

- Iterate over all partons in the configuration - if they are switched off ignore them.
Evolve(Parton * parton), see for details below, returns a boolean (true in the case of finding a splitting with evolution parameter t above the cut-off t_0 , $t > t_0$, false if no allowed splitting has been found). It also keeps track of the winning splitting Splitting * p_winner, i.e. the splitting with the largest t to date. If such a splitting exists, (p_winner!=NULL), the respective parton will be split according to the parameters stored in it. Irrespective of whether the splitting is viable in PerformSplitting (it will always be possible, unless there is a bug in how the offsprings of the splitting are inserted into the configuration), the evolution parameter of the configuration will be set to the t of the splitting.
- Add a final weight to the overall weight, see below.

2. Evolve(Parton * parton)

- Start with p_winner = NULL, i.e. the default is that we do not find a successful splitting.
- Iterate over all spectators (one in the case of quarks, two in the case of gluons).
For each spectator, produce Kernels, a vector of Kernel's, which encapsulate the splitting function and the coupling part of the splitting kernel. They depend on the splitter and on the initial/final-state nature of both splitter and spectator. Create a trial t for this combination by calling GenerateTestSplitting(Kernels * kernels, Splitting & split), where a new Splitting * split is initiated for each splitter-spectator pair. If successful the Splitting split will be filled, and if its t is the largest so far it will be kept as the winning splitting, p_winner. If not, split will be deleted.
- The method returns true or false dependent on whether we found a winner or not.

3. GenerateTestSplitting(Kernels * kernels, Splitting & split)

- kernels->Integral(split, p_massselector) iterates over all allowed splitting kernels, and extracts the sum of the integrals of their over-

estimators, $\sum_K \bar{I}_K$. The integrals are given by

$$\bar{I}_K = \frac{C_K \alpha_S(t_0)}{2\pi} \int_0^1 dz \bar{P}_{a \rightarrow bc}(z) = \frac{C_K \alpha_S(t_0)}{2\pi} \bar{I}_{a \rightarrow bc}, \quad (1)$$

where $C_K = C_F = 4/3$ for quarks, $C_K = C_A/2 = 3/2$ for gluons splitting into gluons, and $C_K = 1/2$ for gluons splitting into quarks. $\alpha_S(t_0)$ obviously is the maximum of the strong coupling, since we assume t to be the relevant scale for all splittings. The overestimated splitting functions are given by

$$\begin{aligned} \bar{P}_{q \rightarrow qg} = \bar{P}_{g \rightarrow gg}^{(1)} &= \frac{2(1-z)}{(1-z)^2 + \kappa_0^2} \cdot (1 + \bar{K}), \\ \bar{P}_{q \rightarrow gq} = \bar{P}_{g \rightarrow gg}^{(2)} &= \frac{2z}{z^2 + \kappa_0^2} \cdot (1 + \bar{K}), \\ \bar{P}_{g \rightarrow gg} &= 1. \end{aligned} \quad (2)$$

We have two gluon splitting functions into gluons, which will need to be added. κ_0^2 is the overestimator of the Curci-Furmanski-Petronzio regulator, t_0/Q^2 , and \bar{K} is the overestimator of the soft enhancement function,

$$\bar{K} = K|_{t=t_0} = \frac{\alpha_S(t_0)}{2\pi} \left[C_A \left(\frac{67}{18} - \frac{\pi^2}{6} \right) - \frac{10}{9} T_{Rn_f}(t_0) \right]. \quad (3)$$

The z -integrals of the splitting functions are given by

$$\begin{aligned} \bar{I}_{q \rightarrow qg} = \bar{I}_{q \rightarrow gq} = \bar{I}_{g \rightarrow gg}^{(1)} = \bar{I}_{g \rightarrow gg}^{(2)} &= \log \left(1 + \frac{Q^2}{t_0} \right) \cdot (1 + \bar{K}) \\ \bar{P}_{g \rightarrow gg} &= 1. \end{aligned} \quad (4)$$

Starting from a given upper scale for the evolution parameter, T , a new trial evolution parameter t emerges as solution of the equation

$$\# \exp \left[- \int_{t_0}^T \frac{dt'}{t'} \sum_K \bar{I}_K \right] = \exp \left[- \int_{t_0}^t \frac{dt'}{t'} \sum_K \bar{I}_K \right], \quad (5)$$

with $\#$ a random number:

$$t = T \cdot \exp \left[\frac{\#}{\sum_K \bar{I}_K} \right]. \quad (6)$$