

# Vector Calculus

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Based on Lecture Notes by David Tong and Dexter Chua

### **Abstract**

These are a set of notes I compiled when teaching myself vector calculus. Mainly using David Tong's amazing notes supporting myself with Dexter Chua's notes.

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# 1 Curves

## 1.1 Coordinates

We will look at maps of the form

$$f : \mathbb{R} \rightarrow \mathbb{R}^n$$

which are called *parametric curves* and are often denoted as the curve  $C$ .

As the image of the curve is  $\mathbb{R}^n$  we will need an orthonormal basis  $\{\mathbf{e}_i\}_{i=1}^n$  for  $\mathbb{R}^n$ . A vector  $\mathbf{x} \in \mathbb{R}^n$  will be represented as

$$\mathbf{x} = (x^1, x^2, \dots, x^n) = \sum_{i=1}^n x^i \mathbf{e}_i = x^i \mathbf{e}_i \quad (1)$$

where we have used the convention of implicitly summing over  $i = 1, 2, \dots, n$  and each component of  $\mathbf{x}$  is  $x^i$  instead of  $x_i$  which you may have seen before.

## 1.2 Differentiating a Curve

The vector function  $\mathbf{x}(t)$  is differentiable if as  $\delta t \rightarrow 0$ , we can write

$$\mathbf{x}(t + \delta t) - \mathbf{x}(t) = \dot{\mathbf{x}}(t)\delta t + \mathcal{O}(\delta t^2) \quad (2)$$

where  $\mathcal{O}(\delta t^2)$  is the “big-O” notation meaning terms that scale  $\delta t^2$  or smaller as  $\delta t \rightarrow 0$ . The term  $\dot{\mathbf{x}}(t)$  is the *derivative* of  $\mathbf{x}(t)$ ; a vector function that is differentiable everywhere is called *smooth*.

We can write

$$\delta \mathbf{x}(t) = \mathbf{x}(t + \delta t) - \mathbf{x}(t)$$

to have the derivative of the vector function  $\mathbf{x}(t)$  as

$$\frac{d\mathbf{x}}{dt} = \dot{\mathbf{x}}(t) = \lim_{\delta t \rightarrow 0} \frac{\delta \mathbf{x}}{\delta t}$$

To differentiate a vector  $\mathbf{x}(t)$  we can differentiate component wise

$$\mathbf{x}(t) = x^i \mathbf{e}_i \Rightarrow \dot{\mathbf{x}}(t) = \dot{x}^i \mathbf{e}_i$$

where our choice of basis  $\{\mathbf{e}_i\}$  is independent of  $t$ . We will look later at the case when  $\{\mathbf{e}_i\}$  are functions of  $t$ .

In addition, we have the following rules for differentiating a function  $f(t)$  and vector functions  $\mathbf{g}(t)$  and  $\mathbf{h}(t)$ .

$$\begin{aligned} \frac{d}{dt}(f\mathbf{g}) &= \frac{df}{dt}\mathbf{g} + f\frac{d\mathbf{g}}{dt} \\ \frac{d}{dt}(\mathbf{g} \cdot \mathbf{h}) &= \frac{d\mathbf{g}}{dt} \cdot \mathbf{h} + \mathbf{g} \cdot \frac{d\mathbf{h}}{dt} \\ \frac{d}{dt}(\mathbf{g} \times \mathbf{h}) &= \frac{d\mathbf{g}}{dt} \times \mathbf{h} + \mathbf{g} \times \frac{d\mathbf{h}}{dt} \end{aligned}$$

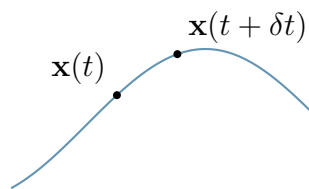
### 1.3 Tangent Vectors

The vector  $\dot{\mathbf{x}}$  is referred to as the *tangent vector*.

The direction of the tangent vector is independent of the choice of parametrisation of the curve  $C$ . In contrast, the magnitude  $|\dot{\mathbf{x}}(t)|$  of the tangent vector does depend of the choice of parametrisation. A curve has a *regular* parametrisation if  $\dot{\mathbf{x}}(t) \neq \mathbf{0}$  for all  $t$ .

### 1.4 Arc Length

Suppose that we have a parametrised curve with a parametrisation  $t$ .



The distance  $\delta s$  between two nearby points  $\mathbf{x}(t)$  and  $\mathbf{x}(t + \delta t)$  is

$$\delta s = |\delta \mathbf{x}| + \mathcal{O}(|\delta \mathbf{x}|^2) = |\dot{\mathbf{x}}(t)\delta t| + \mathcal{O}(\delta t^2)$$

where the  $\mathcal{O}(|\delta \mathbf{x}|^2)$  comes from the fact that the distance between the two points is not exactly a straight line along the curve and contains higher order terms in  $\delta \mathbf{x}^2$ .

We then have

$$\frac{ds}{dt} = \pm \left| \frac{d\mathbf{x}}{dt} \right|$$

where we get the plus sign for increasing  $t$  and the minus sign for decreasing  $t$ .

Hence, if we choose a starting point  $t_0$  and an end point  $t$  (with  $t_0 < t$ ) then the distance along the curve from  $t_0$  to  $t$  is

$$s = \int_{t_0}^t dt' |\dot{\mathbf{x}}(t')| \quad (3)$$

This is called the *arc length* of the curve.

If we choose a different parametrisation  $\tau(t)$  for the curve (which we take to be invertible and smooth with  $\dot{\tau} > 0$  so that they measure “increasing time” in the same direction) then

$$\frac{d\mathbf{x}}{dt} = \frac{d\mathbf{x}}{d\tau} \cdot \frac{d\tau}{dt}$$

We can then compute the arc length with the parametrisation  $\tau$ ,

$$s = \int_{t_0}^t dt' |\dot{\mathbf{x}}(t')| = \int_{\tau_0}^{\tau} d\tau' \left| \frac{d\mathbf{x}}{d\tau'} \cdot \frac{d\tau'}{dt'} \right| \frac{dt'}{d\tau'} = \int_{\tau_0}^{\tau} d\tau' \left| \frac{d\mathbf{x}}{d\tau'} \right|$$

This means that the arc length is independent of the choice of parametrisation  $t$ .

As the arc length of a curve is independent of the choice of parametrisation, then we can parametrise the curve by its arc length. We can think of  $\mathbf{x}(s)$  with its tangent vector  $d\mathbf{x}/ds$  which we also denote as  $\mathbf{x}'(s)$ .

In addition, as we are viewing  $\mathbf{x}(s)$

$$\frac{ds}{dt} = \pm \left| \frac{d\mathbf{x}}{dt} \right| \Rightarrow \left| \frac{d\mathbf{x}}{ds} \right| = 1$$

This means that our tangent vector is a unit vector when working in terms of the arc length.

### Example

Consider the curve  $\mathbf{x} = (\cos t, \sin t, t)$  for  $t \in [0, 4\pi]$ .

We then have

$$\dot{\mathbf{x}}(t) = (-\sin t, \cos t, 1)$$

Calculating the arc length from  $t = 0$  to  $t$

$$s = \int_0^t dt' \sqrt{2} = t\sqrt{2}$$

We then parametrise our curve by  $s$

$$\mathbf{x}(s) = \left( \cos \frac{s}{\sqrt{2}}, \sin \frac{s}{\sqrt{2}}, \frac{s}{\sqrt{2}} \right)$$

With our tangent vector

$$\mathbf{x}'(s) = \left( -\frac{1}{\sqrt{2}} \sin \frac{s}{\sqrt{2}}, \frac{1}{\sqrt{2}} \cos \frac{s}{\sqrt{2}}, \frac{1}{\sqrt{2}} \right)$$

Which has magnitude

$$|\mathbf{x}'(s)| = \frac{1}{2} \left( \sin^2 \frac{s}{\sqrt{2}} + \cos^2 \frac{s}{\sqrt{2}} + 1 \right) = 1$$

as expected.

## 1.5 Curvature & Torsion

Given a curve  $C$  which is parametrised by its arc length  $s$ , we have shown that the tangent vector

$$\mathbf{t} = \frac{d\mathbf{x}}{ds}$$

has unit length. We further define the *curvature* as the magnitude of the derivative of the tangent vector with respect to its arc length

$$\kappa(s) = \left| \frac{d^2\mathbf{x}}{ds^2} \right| \quad (4)$$

If we normalise the derivative of the tangent vector with respect to its arc length (provided that  $\kappa \neq 0$ ) we arrive at the *principle normal*

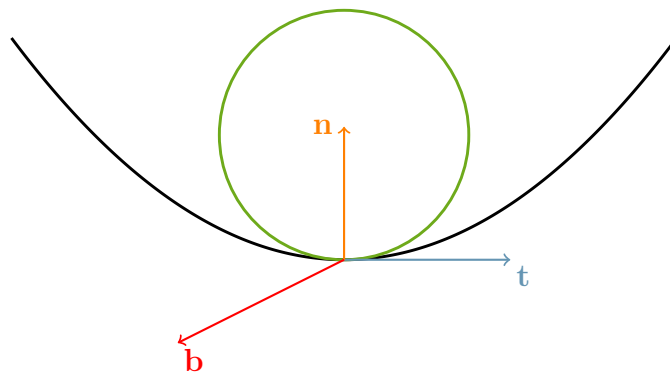
$$\mathbf{n} = \frac{1}{\kappa} \frac{d^2 \mathbf{x}}{ds^2} = \frac{1}{\kappa} \frac{d\mathbf{t}}{ds} \quad (5)$$

Which has a unit length,  $|\mathbf{n}| = 1$ . Using the fact that  $\mathbf{t}$  is a unit vector with  $\mathbf{t} \cdot \mathbf{t} = 1$

$$\frac{d}{ds}(\mathbf{t} \cdot \mathbf{t}) = 2\kappa \mathbf{n} \cdot \mathbf{t} = 0$$

As  $\kappa \neq 0$  then  $\mathbf{n} \cdot \mathbf{t} = 0$  and the principle normal is orthogonal to the tangent vector. Hence, the vectors  $\{\mathbf{t}, \mathbf{n}\}$  define a plane called the *osculating plane*.

For each point  $s$  on the curve there is an associated osculating plane, draw a circle in each plane that touches the curve at the point  $s$  whose curvature is  $\kappa(s)$ ; this is called the *osculating circle*.



Working in  $\mathbb{R}^3$ , we define the *binormal* as the normal to the osculating plane

$$\mathbf{b} = \mathbf{t} \times \mathbf{n} \quad (6)$$

As both  $\mathbf{t}$  and  $\mathbf{n}$  are unit vectors, then  $\mathbf{b}$  is a unit vector. Using the same argument as before, as  $\mathbf{b}$  is a unit vector with  $\mathbf{b} \cdot \mathbf{b} = 1$

$$\frac{d}{ds}(\mathbf{b} \cdot \mathbf{b}) = 2\mathbf{b} \cdot \frac{d\mathbf{b}}{ds} = 0 \Rightarrow \mathbf{b} \cdot \frac{d\mathbf{b}}{ds} = 0$$



We also have that  $\mathbf{t} \cdot \mathbf{b} = 0$  and  $\mathbf{n} \cdot \mathbf{b} = 0$  from the definition of the cross product, which gives

$$\begin{aligned} \frac{d}{ds}(\mathbf{t} \cdot \mathbf{b}) &= \frac{d\mathbf{t}}{ds} \cdot \mathbf{b} + \mathbf{t} \cdot \frac{d\mathbf{b}}{ds} = \kappa(\mathbf{n} \cdot \mathbf{b}) + \mathbf{t} \cdot \frac{d\mathbf{b}}{ds} \\ &\Rightarrow \mathbf{t} \cdot \frac{d\mathbf{b}}{ds} = 0 \end{aligned}$$

Which means that  $d\mathbf{b}/ds$  is orthogonal to both  $\mathbf{t}$  and  $\mathbf{b}$  and thus must lie parallel to  $\mathbf{n}$ . We then define the *torsion* as how much the binormal changes

$$\frac{d\mathbf{b}}{ds} = -\tau(s)\mathbf{n} \quad (7)$$

Intuitively, the curvature tells us how much the curve fails to be a straight line and the torsion how much the curve fails to be a plane.

Observe that using  $\mathbf{b} = \mathbf{t} \times \mathbf{n}$ ,

$$\mathbf{b} \times \mathbf{t} = (\mathbf{t} \times \mathbf{n}) \times \mathbf{t} = (\mathbf{t} \cdot \mathbf{t})\mathbf{n} - (\mathbf{n} \cdot \mathbf{t})\mathbf{t} = \mathbf{n}$$

Which we can then use to get

$$\mathbf{n}' = (\mathbf{b} \times \mathbf{t})' = \mathbf{b}' \times \mathbf{t} + \mathbf{b} \times \mathbf{t}' = (-\tau\mathbf{n}) \times \mathbf{t} + \mathbf{b} \times (\kappa\mathbf{n}) = \tau\mathbf{b} - \kappa\mathbf{t}$$

using our previous relationships. We then have three sets of equations

$$\frac{d\mathbf{t}}{ds} = \kappa\mathbf{n} \quad (8)$$

$$\frac{d\mathbf{b}}{ds} = -\tau\mathbf{n} \quad (9)$$

$$\frac{d\mathbf{n}}{ds} = \tau\mathbf{b} - \kappa\mathbf{t} \quad (10)$$

which are referred to as the *Frenet-Serret* equations. If we are given  $\kappa(s)$  and  $\tau(s)$  with initial conditions on  $\mathbf{b}$  and  $\mathbf{t}$ , we can then solve for  $\mathbf{b}$ ,  $\mathbf{t}$  and the curve  $\mathbf{x}(s)$ .

## 1.6 Line Integral

### Scalar Fields

We will first look at the map

$$\begin{aligned}\phi : \mathbb{R}^n &\rightarrow \mathbb{R} \\ \mathbf{x} &\rightarrow \phi(\mathbf{x})\end{aligned}$$

We could define the function  $\phi(\mathbf{x}(t))$  for a parametrised curve  $C$  in  $\mathbb{R}^n$  and then integrate over  $t$  in the usual way. Yet this will depend on the parametrisation.

Instead, we will work with the arc length  $s$ . Integrating from the point  $\mathbf{x}(s_a) = \mathbf{a}$  to  $\mathbf{x}(s_b) = \mathbf{b}$  on a curve  $C$ , we define the line integral as

$$\int_C \phi \, ds = \int_{s_a}^{s_b} \phi(\mathbf{x}(s)) \, ds$$

If instead we are given a curve  $C$  with a different parametrisation  $\mathbf{x}(t)$  with  $\mathbf{x}(t_a) = \mathbf{a}$  and  $\mathbf{x}(t_b) = \mathbf{b}$ , then

$$\int_C \phi \, ds = \int_{t_a}^{t_b} \phi(\mathbf{x}(t)) \frac{ds}{dt} \, dt \quad \text{with} \quad \frac{ds}{dt} = \pm |\dot{\mathbf{x}}(t)|$$

where the sign depend if  $t_a < t_b$  or  $t_b < t_a$ .

### Vector Fields

Consider the map

$$\mathbf{F} : \mathbb{R}^n \rightarrow \mathbb{R}^n$$

We could integrate component wise, treating each component as a scalar field to give a vector as an answer.

On the other hand, we can integrate the component of the vector field that lies tangent to the curve to give a scalar as an answer.

Suppose our curve  $C$  has a parametrisation  $\mathbf{x}(t)$  and we wish to integrate from  $\mathbf{x}(t_a) = \mathbf{a}$  to  $\mathbf{x}(t_b) = \mathbf{b}$ . We define the line integral as

$$\int_C \mathbf{F}(\mathbf{x}) \cdot d\mathbf{x} = \int_{t_a}^{t_b} \mathbf{F}(\mathbf{x}(t)) \cdot \dot{\mathbf{x}}(t) dt$$

The orientation is in the direction along the curve, or the direction of the tangent vector  $\dot{\mathbf{x}}$ . The orientation is from  $\mathbf{a}$  to  $\mathbf{b}$  regardless if  $t_a < t_b$  or  $t_b < t_a$ .

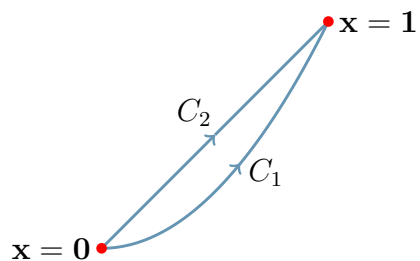
### Example

Let  $\phi(\mathbf{x}) = x + y + z$  and  $\mathbf{F}(\mathbf{x}) = (xe^y, z^2, xy)$ . We define the curves

$$C_1 : \mathbf{x}(t) = (t, t^2, t^3)$$

$$C_2 : \mathbf{x}(t) = (t, t, t)$$

and wish to integrate each of the fields along curves  $C_1$  and  $C_2$  from  $\mathbf{x} = \mathbf{0}$  to  $\mathbf{x} = \mathbf{1}$ .



Starting with  $\phi(\mathbf{x})$ , we see that

$$\int_{C_1} \phi ds = \int_0^1 dt (t + t^2 + t^3) \cdot |(1, 2t, 3t^2)| = \int_0^1 dt (t + t^2 + t^3) \sqrt{1 + 4t^2 + 9t^4} \approx 2.715$$

We can also re-parametrise  $C_2$  in terms of its arc length

$$\mathbf{x}(s) = \left( \frac{s}{\sqrt{3}}, \frac{s}{\sqrt{3}}, \frac{s}{\sqrt{3}} \right) \Rightarrow \int_{C_2} \phi ds = \int_0^{\sqrt{3}} \frac{3s}{\sqrt{3}} ds = \frac{3\sqrt{3}}{2}$$

Secondly, integrating  $\mathbf{F}$  along  $C_1$

$$\int_{C_1} \mathbf{F} \cdot d\mathbf{x} = \int_0^1 dt (te^{t^2}, t^6, t^3) \cdot (1, 2t, 3t^2) = \int_0^1 dt (te^{t^2} + 2t^7 + 3t^5) = \frac{1}{4}(1 + 2e)$$

Finally, along  $C_2$

$$\int_{C_2} \mathbf{F} \cdot d\mathbf{x} = \int_0^1 dt (te^t, t^2, t^2) \cdot (1, 1, 1) = \int_0^1 dt (te^t + 2t^2) = \frac{5}{3}$$

If we wish to integrate along a closed curve  $C$  who has the same start and end point, then we denote it as

$$\oint_C \mathbf{F} \cdot d\mathbf{x}$$

In addition, if we have a closed curve  $C$  such that  $C = \sum_i C_i$  with each  $C_i$  smooth and joined at the end points, then

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = \sum_i \int_{C_i} \mathbf{F} \cdot d\mathbf{x}$$

We will also define the curve  $-C$  for a given curve  $C$  which is in the opposite direction,

$$\int_{-C} \mathbf{F} \cdot d\mathbf{x} = - \int_C \mathbf{F} \cdot d\mathbf{x}$$

### Example

Define the curve  $C = C_1 - C_2$  with  $\mathbf{F} = (xe^y, z^2, xy)$ ,  $C_1$  and  $C_2$  from the previous example. As  $C$  is a closed curve then

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = \int_{C_1} \mathbf{F} \cdot d\mathbf{x} - \int_{C_2} \mathbf{F} \cdot d\mathbf{x} = \frac{1}{4}(1 + 2e) - \frac{5}{3}$$

## 1.7 The Gradient

Take a scalar field  $\phi : \mathbb{R}^n \rightarrow \mathbb{R}$ , we have a partial derivative of  $\phi(x^1, x^2, \dots, x^n)$  as

$$\frac{\partial \phi}{\partial x^i} = \lim_{h \rightarrow 0} \frac{\phi(x^1, \dots, x^i + h, \dots, x^n) - \phi(x^1, \dots, x^n)}{h}$$

for  $i = 1, \dots, n$ . This is a derivative in the  $x^i$  direction, treating all other variables as constant.

We can write each of these partial derivatives as

$$\nabla\phi = \frac{\partial\phi}{\partial x^i} \mathbf{e}_i$$

where  $\{\mathbf{e}_i\}$  is an orthonormal basis for  $\mathbb{R}^n$ ; this is called the *gradient* of  $\phi$ . This  $\nabla\phi$  is a vector field yet it is convention not to write  $\nabla\phi$  and understand that  $\nabla\phi$  is a vector.

Take a normal unit vector  $\hat{\mathbf{n}}$ , the rate of change of  $\phi$  in the direction of  $\hat{\mathbf{n}}$  is

$$D_{\hat{\mathbf{n}}}\phi = \hat{\mathbf{n}} \cdot \nabla\phi$$

which is called the *directional derivative*. We see that this is maximised when  $\hat{\mathbf{n}}$  is parallel to the gradient  $\nabla\phi$ , which tells us that  $\nabla\phi$  points towards where  $\phi(\mathbf{x})$  changes the quickest.

## 1.8 Conservative Fields

For a vector field  $\mathbf{F}$ , we say it is *conservative* if

$$\mathbf{F} = \nabla\phi$$

for some scalar field  $\phi$ ; this scalar field is called the *potential*.

**Claim:** The line integral around any closed curve vanishes if and only if the vector field  $\mathbf{F}$  is conservative.

**Proof:** Suppose that  $\mathbf{F} = \nabla\phi$ . Integrating along a closed curve  $C$  from  $\mathbf{x}(t_a) = \mathbf{a}$  to  $\mathbf{x}(t_b) = \mathbf{b}$

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = \int_{t_a}^{t_b} dt \mathbf{F}(\mathbf{x}(t)) \cdot \frac{d\mathbf{x}}{dt} = \int_{t_a}^{t_b} dt \frac{\partial\phi(\mathbf{x}(t))}{\partial x^i} \cdot \frac{dx^i}{dt} = \int_{t_a}^{t_b} dt \frac{d}{dt} \phi(\mathbf{x}(t))$$

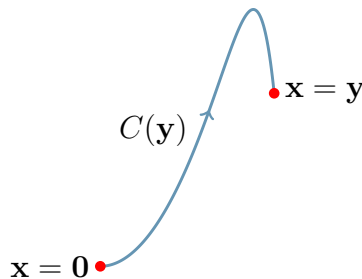
by the chain rule. Now we have the integral of a total derivative, which gives

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = \phi(\mathbf{b}) - \phi(\mathbf{a}) = 0$$

Suppose that  $\oint_C \mathbf{F} \cdot d\mathbf{x} = 0$  for any curve  $C$ . Define a scalar curve  $\phi : \mathbb{R}^n \rightarrow \mathbb{R}$  with  $\phi(\mathbf{0}) = 0$ . At any point  $\mathbf{x} = \mathbf{y}$ , we define

$$\phi(\mathbf{y}) = \int_{C(\mathbf{y})} \mathbf{F} \cdot d\mathbf{x}$$

where  $C(\mathbf{y})$  is the curve that starts at  $\mathbf{x} = \mathbf{0}$  and ends at  $\mathbf{x} = \mathbf{y}$ .



By our assumption,

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = 0$$

for any closed curve  $C$ , so it does not matter which closed curve  $C$  we take, they all give the same answer.

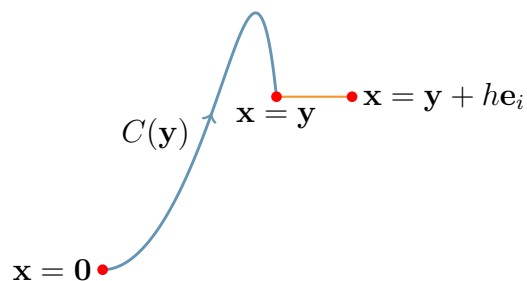
Consider,

$$\frac{\partial \phi}{\partial x^i}(\mathbf{y}) = \lim_{h \rightarrow 0} \frac{1}{h} \left\{ \int_{C(\mathbf{y} + h\mathbf{e}_i)} \mathbf{F} \cdot d\mathbf{x} - \int_{C(\mathbf{y})} \mathbf{F} \cdot d\mathbf{x} \right\}$$

The first integral goes along the line  $C(\mathbf{y})$  and then along the red line in the  $\mathbf{e}_i$  direction as shown in the figure below. The second integral goes back along the curve  $C(\mathbf{y})$ .

As a result the difference in the integrals is simply the integral along the orange line

$$\frac{\partial \phi}{\partial x^i}(\mathbf{y}) = \lim_{h \rightarrow 0} \frac{1}{h} \int_{\text{orange line}} \mathbf{F} \cdot d\mathbf{x}$$



As we are integrating along a straight line in the  $\mathbf{e}_i$  direction, we are projecting the  $F_i$  component of  $\mathbf{F}$ . As  $h$  is small we get

$$\int_{\text{orange line}} \mathbf{F} \cdot d\mathbf{x} \approx F_i h$$

and after taking the limit as  $h \rightarrow 0$ ,

$$\frac{\partial \phi}{\partial x^i}(\mathbf{y}) = F_i(\mathbf{y}) \Rightarrow \mathbf{F} = \nabla \phi \quad \square$$

In order to check if there exists a scalar field  $\phi$  such that  $\mathbf{F} = \nabla \phi$  we have

$$F_i = \frac{\partial \phi}{\partial x^i}$$

by definition, differentiating again gives

$$\frac{\partial F_i}{\partial x^j} = \frac{\partial^2 \phi}{\partial x^i \partial x^j} = \frac{\partial F_j}{\partial x^i}$$

as for “nice”  $\mathbf{F}$  we have  $\partial^2 F / \partial x \partial y = \partial^2 F / \partial y \partial x$ . Hence, for  $\mathbf{F}$  to be conservative we need  $\partial_i F_j = \partial_j F_i$  for all  $i, j$ .

### Exact Differentials

Given a scalar field  $\phi(\mathbf{x})$  on  $\mathbb{R}^n$ , we define the *differential* as

$$d\phi = \frac{\partial \phi}{\partial x^i} dx^i = \nabla \phi \cdot d\mathbf{x}$$

This is a function of  $\mathbf{x}$  and captures how much the function  $\phi$  changes as we move in any direction around  $\mathbf{x}$ .

Consider a field  $\mathbf{F}(\mathbf{x})$  on  $\mathbb{R}^n$ . Taking an inner product with an infinitesimal vector to get

$$\mathbf{F} \cdot d\mathbf{x}$$

is a *differential form*. If a differential form can be written as

$$\mathbf{F} \cdot d\mathbf{x} = d\phi = \nabla\phi \cdot d\mathbf{x}$$

for some  $\phi$  it is said to be an *exact* differential form. This is a re-formalisation of our idea: a differential is exact if and only if the vector field is conservative.

### An Application: Work & Potential Energy

Consider a particle with trajectory  $\mathbf{x}(t)$ . By Newton's second law

$$\mathbf{F}(\mathbf{x}) = m\ddot{\mathbf{x}}$$

where  $\mathbf{F}(\mathbf{x})$  is the force field. Recall that a particle has kinetic energy given by  $K = \frac{1}{2}m\dot{\mathbf{x}}^2$ ; this kinetic energy changes in time as

$$K(t_2) - K(t_1) = \int_{t_1}^{t_2} dt \frac{dK}{dt} = \int_{t_1}^{t_2} dt (m\dot{\mathbf{x}} \cdot \ddot{\mathbf{x}}) = \int_{t_1}^{t_2} dt \dot{\mathbf{x}} \cdot \mathbf{F} = \int_C \mathbf{F} \cdot d\mathbf{x}$$

where  $C$  is the trajectory of the curve. The line integral of the force  $\mathbf{F}$  along the trajectory  $C$  of the particle is known as the *work done*. If our force field  $\mathbf{F}$  is conservative, then

$$\mathbf{F} = -\nabla V$$

for some  $V$ . From the above, for a conservative force the work done only depends on the end points of our trajectory

$$\begin{aligned} K(t_2) - K(t_1) &= \int_C \mathbf{F} \cdot d\mathbf{x} = -V(t_2) + V(t_1) \\ \Rightarrow K(t) + V(t) &= \text{constant} \end{aligned}$$



**A Subtlety**

Consider the vector field given by

$$\mathbf{F} = \left( \frac{-y}{x^2 + y^2}, \frac{x}{x^2 + y^2} \right)$$

This is a conservative field with

$$\mathbf{F} = \nabla \phi \quad \text{and} \quad \phi = \arctan \left( \frac{y}{x} \right)$$

Integrating  $\mathbf{F}$  along a closed curve  $C$  given by a circle of radius  $R$  about the origin  $\mathbf{x}(t) = (R \cos t, R \sin t)$ :

$$\oint_C \mathbf{F} \cdot d\mathbf{x} = \int_0^{2\pi} dt \mathbf{F} \cdot \frac{d\mathbf{x}}{dt} = \int_0^{2\pi} dt \frac{1}{R^2} (R^2 \sin^2 t + R^2 \cos^2 t) = 2\pi$$

Yet we have  $\oint_C \mathbf{F} \cdot d\mathbf{x} \neq 0$  when  $\mathbf{F}$  is a conservative field!? This does not vanish as our function  $\phi$  is not continuous along the  $y$ -axis. In fact, we have implicitly assumed that our function  $\phi$  is continuous. We should have  $\mathbf{F} = \nabla \phi$  with  $\phi$  continuous to guarantee that  $\oint_C \mathbf{F} \cdot d\mathbf{x} = 0$ .