

## Uma Antena Linear Centralmente Alimentada

Vamos considerar uma antena bem simples, consistindo de duas barras condutoras idênticas, retas e finas, ambas alinhadas ao longo de um eixo que vamos chamar de eixo  $z$ , mas com uma pequena separação entre elas, desprezível para nossa finalidade, onde o sinal elétrico é inserido. Vamos adotar a origem do sistema de coordenadas exatamente no centro dessa pequena separação entre as barras. Seja  $d/2$  o comprimento de cada uma das barras condutoras. A ideia é supormos que a densidade de corrente elétrica estabelecida ao longo da antena seja simétrica com relação à origem de coordenadas e dependente senoidalmente de  $z$  e harmonicamente de  $t$ . Assim, vamos tomar a densidade de corrente complexa como

$$\mathbf{J}_c(\mathbf{r}) = \hat{\mathbf{z}} I \delta(x) \delta(y) \sin\left(\frac{kd}{2} - k|z|\right), \text{ para } |z| < \frac{d}{2},$$

sendo nula para  $|z| \geq d/2$  e

$$k = \frac{\omega}{c}.$$

As funções de Dirac indicam que a antena é estreita sobre o eixo  $z$ . Na zona de radiação,

$$\begin{aligned} \mathbf{A}_c^{\text{rad}}(\mathbf{r}) &= \frac{\exp(ikr)}{rc} \int_V d^3r' \mathbf{J}_c(\mathbf{r}') \exp(-ik\hat{\mathbf{r}} \cdot \mathbf{r}') \\ &= \frac{\exp(ikr)}{rc} \int_V d^3r' \hat{\mathbf{z}} I \delta(x') \delta(y') \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ik\hat{\mathbf{r}} \cdot \mathbf{r}'). \end{aligned}$$

Integrando em  $x'$  e  $y'$  fornece

$$\mathbf{A}_c^{\text{rad}}(\mathbf{r}) = \hat{\mathbf{z}} \frac{I \exp(ikr)}{rc} \int_{-d/2}^{d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta),$$

onde  $\theta$  é o ângulo entre o eixo  $z$  e o vetor posição do ponto de observação,  $\mathbf{r}$ . Calculemos a integral:

$$\begin{aligned} \int_{-d/2}^{+d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta) &= \int_{-d/2}^0 dz' \sin\left(\frac{kd}{2} + kz'\right) \exp(-ikz' \cos \theta) \\ &+ \int_0^{d/2} dz' \sin\left(\frac{kd}{2} - kz'\right) \exp(-ikz' \cos \theta), \end{aligned}$$

que, com a mudança de variável

$$z' \rightarrow -z'$$

na primeira integral do membro direito, resulta em

$$\begin{aligned} \int_{-d/2}^{+d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta) &= \int_0^{d/2} dz' \sin\left(\frac{kd}{2} - kz'\right) \exp(ikz' \cos \theta) \\ &+ \int_0^{d/2} dz' \sin\left(\frac{kd}{2} - kz'\right) \exp(-ikz' \cos \theta) \end{aligned}$$

e, portanto,

$$\int_{-d/2}^{+d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta) = 2 \int_0^{d/2} dz' \sin\left(\frac{kd}{2} - kz'\right) \cos(kz' \cos \theta).$$

Mas,

$$\begin{aligned}\sin\left(\frac{kd}{2} - kz'\right) \cos(kz' \cos \theta) &= \frac{1}{2} \sin\left(\frac{kd}{2} - kz' + kz' \cos \theta\right) + \frac{1}{2} \sin\left(\frac{kd}{2} - kz' - kz' \cos \theta\right) \\ &= -\frac{1}{2} \sin\left[k(1 + \cos \theta)z' - \frac{kd}{2}\right] - \frac{1}{2} \sin\left[k(1 - \cos \theta)z' - \frac{kd}{2}\right]\end{aligned}$$

e, assim,

$$\begin{aligned}\int_{-d/2}^{+d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta) &= -\int_0^{d/2} dz' \sin\left[k(1 + \cos \theta)z' - \frac{kd}{2}\right] \\ &\quad - \int_0^{d/2} dz' \sin\left[k(1 - \cos \theta)z' - \frac{kd}{2}\right] \\ &= \frac{\cos\left[k(1 + \cos \theta)\frac{d}{2} - \frac{kd}{2}\right] - \cos\left(\frac{kd}{2}\right)}{k(1 + \cos \theta)} \\ &\quad + \frac{\cos\left[k(1 - \cos \theta)\frac{d}{2} - \frac{kd}{2}\right] - \cos\left(\frac{kd}{2}\right)}{k(1 - \cos \theta)}.\end{aligned}$$

Simplificando, obtemos

$$\begin{aligned}\int_{-d/2}^{+d/2} dz' \sin\left(\frac{kd}{2} - k|z'|\right) \exp(-ikz' \cos \theta) &= \frac{1}{k} \left[ \cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right) \right] \\ &\quad \times \left[ \frac{1}{1 + \cos \theta} + \frac{1}{1 - \cos \theta} \right] \\ &= \frac{2}{k} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right].\end{aligned}$$

Logo,

$$\mathbf{A}_c^{\text{rad}}(\mathbf{r}) = \hat{\mathbf{z}} \frac{2I \exp(ikr)}{rck} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right].$$

Calculemos o campo indução magnética de radiação:

$$\begin{aligned}\mathbf{B}_c(\mathbf{r}) &= \nabla \times \mathbf{A}_c^{\text{rad}}(\mathbf{r}) \\ &= \nabla \times \left\{ \hat{\mathbf{z}} \frac{2I \exp(ikr)}{rck} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right] \right\} \\ &= -\hat{\mathbf{z}} \times \nabla \left\{ \frac{2I \exp(ikr)}{rck} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right] \right\}.\end{aligned}$$

Como

$$\nabla = \hat{\mathbf{r}} \frac{\partial}{\partial r} + \frac{\hat{\boldsymbol{\theta}}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\boldsymbol{\varphi}}}{r \sin \theta} \frac{\partial}{\partial \varphi},$$

segue

$$\begin{aligned}\mathbf{B}_c(\mathbf{r}) &\approx -\hat{\mathbf{z}} \times \hat{\mathbf{r}} \frac{\partial}{\partial r} \left\{ \frac{2I \exp(ikr)}{rck} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right] \right\} \\ &\approx -\hat{\mathbf{z}} \times \hat{\mathbf{r}} ik \frac{2I \exp(ikr)}{rck} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin^2 \theta} \right]\end{aligned}$$

e, como

$$\hat{\mathbf{r}} = \hat{\mathbf{x}} \sin \theta \cos \varphi + \hat{\mathbf{y}} \sin \theta \sin \varphi + \hat{\mathbf{z}} \cos \theta,$$

então

$$\begin{aligned}\hat{\mathbf{z}} \times \hat{\mathbf{r}} &= \hat{\mathbf{z}} \times \hat{\mathbf{x}} \sin \theta \cos \varphi + \hat{\mathbf{z}} \times \hat{\mathbf{y}} \sin \theta \sin \varphi + \hat{\mathbf{z}} \times \hat{\mathbf{z}} \cos \theta \\ &= \hat{\mathbf{y}} \sin \theta \cos \varphi - \hat{\mathbf{x}} \sin \theta \sin \varphi \\ &= \sin \theta (\hat{\mathbf{y}} \cos \varphi - \hat{\mathbf{x}} \sin \varphi) \\ &= \sin \theta \hat{\boldsymbol{\varphi}}.\end{aligned}$$

Portanto,

$$\mathbf{B}_c^{\text{rad}}(\mathbf{r}) = -\hat{\boldsymbol{\varphi}} i \frac{2I \exp(ikr)}{rc} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin \theta} \right].$$

E o campo elétrico de radiação fica

$$\begin{aligned}\mathbf{E}_c(\mathbf{r}) &= \frac{i}{k} \boldsymbol{\nabla} \times \mathbf{B}_c^{\text{rad}}(\mathbf{r}) \\ &= \frac{i}{k} \boldsymbol{\nabla} \times \left\{ -\hat{\boldsymbol{\varphi}} i \frac{2I \exp(ikr)}{rc} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin \theta} \right] \right\}.\end{aligned}$$

Como

$$\begin{aligned}\boldsymbol{\nabla} \times \mathbf{F}(\mathbf{r}) &= \frac{\hat{\mathbf{r}}}{r \sin \theta} \left[ \frac{\partial}{\partial \theta} (\sin \theta F_\varphi) - \frac{\partial F_\theta}{\partial \varphi} \right] + \hat{\boldsymbol{\theta}} \left[ \frac{1}{r \sin \theta} \frac{\partial F_r}{\partial \varphi} - \frac{1}{r} \frac{\partial}{\partial r} (r F_\varphi) \right] \\ &+ \frac{\hat{\boldsymbol{\varphi}}}{r} \left[ \frac{\partial}{\partial r} (r F_\theta) - \frac{\partial F_r}{\partial \varphi} \right],\end{aligned}$$

para um campo vetorial  $\mathbf{F}(\mathbf{r})$  arbitrário, obtemos

$$\begin{aligned}\mathbf{E}_c(\mathbf{r}) &= \frac{i}{k} \boldsymbol{\nabla} \times \left\{ -\hat{\boldsymbol{\varphi}} i \frac{2I \exp(ikr)}{rc} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin \theta} \right] \right\} \\ &= -\frac{i}{k} \hat{\mathbf{r}} i \frac{2I \exp(ikr)}{r^2 c \sin \theta} \frac{\partial}{\partial \theta} \left[ \cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right) \right] \\ &+ \frac{i}{k} \hat{\boldsymbol{\theta}} \frac{1}{r} i \frac{2I k \exp(ikr)}{c} \left[ \frac{\cos\left(\frac{kd}{2} \cos \theta\right) - \cos\left(\frac{kd}{2}\right)}{\sin \theta} \right],\end{aligned}$$

isto é,

$$\mathbf{E}_c^{\text{rad}}(\mathbf{r}) = -\frac{i}{k}\hat{\boldsymbol{\theta}}\frac{2Ik\exp(ikr)}{rc}\left[\frac{\cos\left(\frac{kd}{2}\cos\theta\right)-\cos\left(\frac{kd}{2}\right)}{\text{sen}\theta}\right].$$

A média temporal do vetor de Poynting de radiação é dada por

$$\begin{aligned}\langle\mathbf{S}^{\text{rad}}\rangle &= \frac{c}{8\pi}\text{Re}\left\{\mathbf{E}_c^{\text{rad}}(\mathbf{r})\times[\mathbf{B}_c^{\text{rad}}(\mathbf{r})]^*\right\} \\ &= \hat{\boldsymbol{\theta}}\times\hat{\boldsymbol{\varphi}}\frac{I^2}{2\pi r^2c}\left[\frac{\cos\left(\frac{kd}{2}\cos\theta\right)-\cos\left(\frac{kd}{2}\right)}{\text{sen}\theta}\right]^2.\end{aligned}$$

Logo,

$$\langle\mathbf{S}^{\text{rad}}\rangle = \hat{\mathbf{r}}\frac{I^2}{2\pi r^2c}\left[\frac{\cos\left(\frac{kd}{2}\cos\theta\right)-\cos\left(\frac{kd}{2}\right)}{\text{sen}\theta}\right]^2$$

e a distribuição angular da média temporal da potência irradiada por unidade de área é dada por

$$\begin{aligned}\frac{dP}{d\Omega} &= \hat{\mathbf{r}}\cdot\langle\mathbf{S}^{\text{rad}}\rangle r^2 \\ &= \frac{I^2}{2\pi c}\left[\frac{\cos\left(\frac{kd}{2}\cos\theta\right)-\cos\left(\frac{kd}{2}\right)}{\text{sen}\theta}\right]^2.\end{aligned}$$