

**Exercício 4: 04/09** Calcule, detalhadamente, a partir dos potenciais de Liénard & Wiechert, o campo elétrico de uma partícula de carga  $q$  e massa  $m$  percorrendo uma trajetória arbitrária.

Queremos demonstrar o campo  $\vec{E}$ :

$$\vec{E}(\vec{r}, t) = q \left[ \frac{(\hat{R} - \vec{\beta})(1 - \beta^2)}{R^2(1 - \hat{R} \cdot \vec{\beta})^3} + \frac{\hat{R} \times [(\hat{R} - \vec{\beta}) \times \vec{a}]}{Rc^2(1 - \hat{R} \cdot \vec{\beta})^3} \right] \quad (4.1)$$

partindo de,

$$\phi(\vec{r}, t) = \frac{q}{R - \vec{R} \cdot \vec{\beta}} \quad (4.2)$$

e

$$\vec{A}(\vec{r}, t) = \frac{q\vec{\beta}}{R - \vec{R} \cdot \vec{\beta}} \quad (4.3)$$

uma vez que,

$$\vec{E}(\vec{r}, t) = -\vec{\nabla}\phi - \frac{1}{c} \frac{\partial}{\partial t} \vec{A}. \quad (4.4)$$

Vamos calcular primeiro  $-\vec{\nabla}\phi$ :

$$\begin{aligned} -\vec{\nabla}\phi &= \vec{\nabla} \left[ \frac{q}{R - \vec{R} \cdot \vec{\beta}} \right] = -q\vec{\nabla} [(R - \vec{R} \cdot \vec{\beta})^{-1}] \\ &= q [(R - \vec{R} \cdot \vec{\beta})^{-2}] \vec{\nabla}(R - \vec{R} \cdot \vec{\beta}) \\ &= \frac{q}{(R - \vec{R} \cdot \vec{\beta})^2} (\partial_j R - \vec{\beta} \cdot (\partial_j \vec{R}) - \vec{R} \cdot (\partial_j \vec{\beta})) \end{aligned} \quad (4.5)$$

Vamos agora calcular as derivadas diretamente:

$$\partial_j R = \frac{1}{2R} \partial_j R^2 = \frac{1}{2R} \partial_j (\vec{R} \cdot \vec{R}) = \hat{R} \cdot \partial_j \vec{R} \quad (4.6)$$

$$\partial_j \vec{R} = \partial_j (\vec{r} - \vec{r}_0(t_R)) = \hat{e}_j - \frac{dr_0(t_R)}{dt_R} \partial_j t_R = \hat{e}_j - \vec{v} \partial_j t_R \quad (4.7)$$

$$\begin{aligned} \partial_j t_R &= \partial_j \left( t - \frac{R}{c} \right) = \partial_j t - \frac{\partial_j R}{c} = -\frac{\partial_j R}{c} = \frac{\hat{R} \cdot \partial_j \vec{R}}{c} \\ &= \frac{\hat{R} \cdot \hat{e}_j - \hat{R} \cdot \vec{v} \partial_j t_R}{c} = \frac{-\hat{R} \cdot \hat{e}_j}{c} + \vec{\beta} \cdot \hat{R} \partial_j t_R \end{aligned} \quad (4.8)$$

$$\partial_j t_R - \vec{\beta} \cdot \hat{R} \partial_j t_R = \partial_j t_R (1 - \vec{\beta} \cdot \hat{R}) = \frac{-\hat{R} \cdot \hat{e}_j}{c} \implies \partial_j t_R = \frac{-\hat{R} \cdot \hat{e}_j}{c(1 - \vec{\beta} \cdot \hat{R})} \quad (4.9)$$

Finalmente,

$$\partial_j t_R = \frac{\hat{R} \cdot \hat{e}_j}{c(1 - \vec{\beta} \cdot \hat{R})} \quad (4.10)$$

$$\partial_j \vec{R} = \hat{e}_j + \vec{\beta} \frac{\hat{R} \cdot \hat{e}_j}{c(1 - \vec{\beta} \cdot \hat{R})} \quad (4.11)$$

$$\partial_j R = \hat{R} \cdot \hat{e}_j + \frac{(\hat{R} \cdot \vec{\beta})(\hat{R} \cdot \hat{e}_j)}{c(1 - \vec{\beta} \cdot \hat{R})} \quad (4.12)$$

substituído em 4.5 temos,

$$\begin{aligned} -\vec{\nabla} \phi &= \frac{q}{(R - \vec{R} \cdot \vec{\beta})^2} (\partial_j R - \vec{\beta} \cdot (\partial_j \vec{R}) - \vec{R} \cdot (\partial_j \vec{\beta})) \\ &= \frac{q}{(R - \vec{R} \cdot \vec{\beta})^2} \left[ \hat{R} \cdot \hat{e}_j + \frac{(\hat{R} \cdot \vec{\beta})(\hat{R} \cdot \hat{e}_j)}{c(1 - \vec{\beta} \cdot \hat{R})} - \vec{\beta} \cdot \hat{e}_j - \frac{(\vec{\beta} \cdot \vec{\beta})(\hat{R} \cdot \hat{e}_j)}{c(1 - \vec{\beta} \cdot \hat{R})} + \frac{\vec{a}}{c^2} \frac{(\hat{R} \cdot \hat{e}_j)}{(1 - \vec{\beta} \cdot \hat{R})} \right] \end{aligned} \quad (4.13)$$