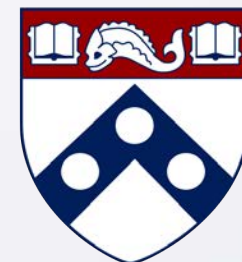


ENM 360: Introduction to Data-driven Modeling

Lecture #3: Approximation of functions and data

Paris Perdikaris
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Penn
UNIVERSITY of PENNSYLVANIA

Numerical differentiation with finite differences

$$A = \begin{bmatrix} -2 & 1 & & & & \\ 1 & -2 & 1 & & & \\ & 1 & -2 & 1 & & \\ & & 1 & -2 & 1 & \\ & & & 1 & -2 & 1 \\ & & & & \ddots & \ddots & \ddots \\ & & & & & 1 & -2 \end{bmatrix}$$

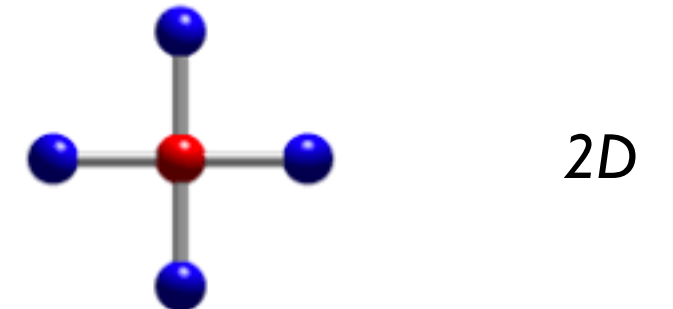
Central difference stencil for second derivative approximation in 1D

$$\begin{bmatrix} 4 & -1 & & & & \\ -1 & 4 & -1 & & & \\ & -1 & 4 & -1 & & \\ & & -1 & 4 & -1 & \\ -1 & & & 4 & -1 & -1 \\ -1 & -1 & -1 & -1 & 4 & -1 \\ & -1 & -1 & -1 & -1 & 4 & -1 \\ & & -1 & -1 & -1 & -1 & 4 & -1 \\ & & & -1 & -1 & -1 & -1 & 4 & -1 \\ & & & & -1 & -1 & -1 & -1 & 4 \end{bmatrix}$$

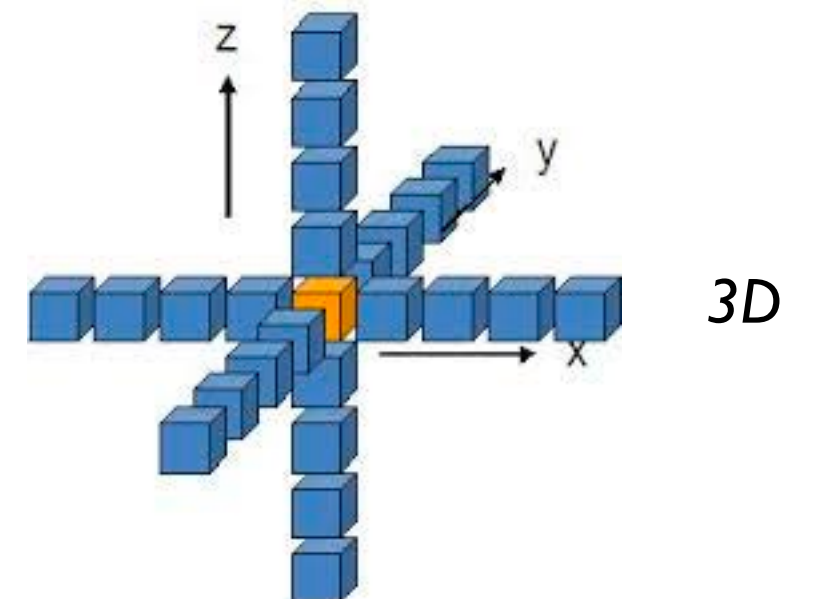
Central difference stencil for second derivative approximation in 2D



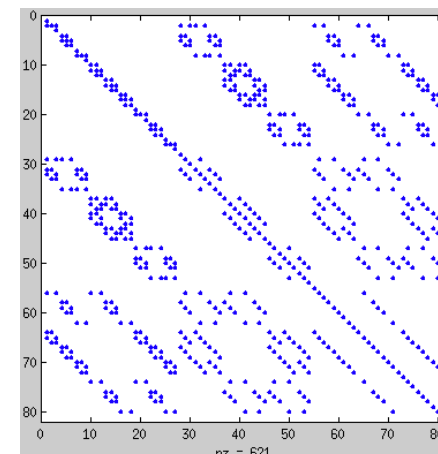
1D



2D



3D



Sparsity pattern of a realistic FD stencil matrix

Numerical differentiation with spectral methods

For *continuous* periodic function $f(x)$, $f(x + 2\pi) = f(x)$, represented by a Fourier series:

$$f(x) = \sum_{n=-\infty}^{\infty} \hat{f}_n e^{inx}$$

The differentiation of $f(x)$ can then be evaluated by:

$$\frac{df(x)}{dx} = \sum_{n=-\infty}^{\infty} \underbrace{(in\hat{f}_n)}_{\text{Fourier coefficient of } f'(x)} e^{inx}$$

∴ Once the coefficients of the Fourier series \hat{f}_n is obtained, the differentiation can be evaluated by summing the Fourier series with new coefficients $(in\hat{f}_n)$

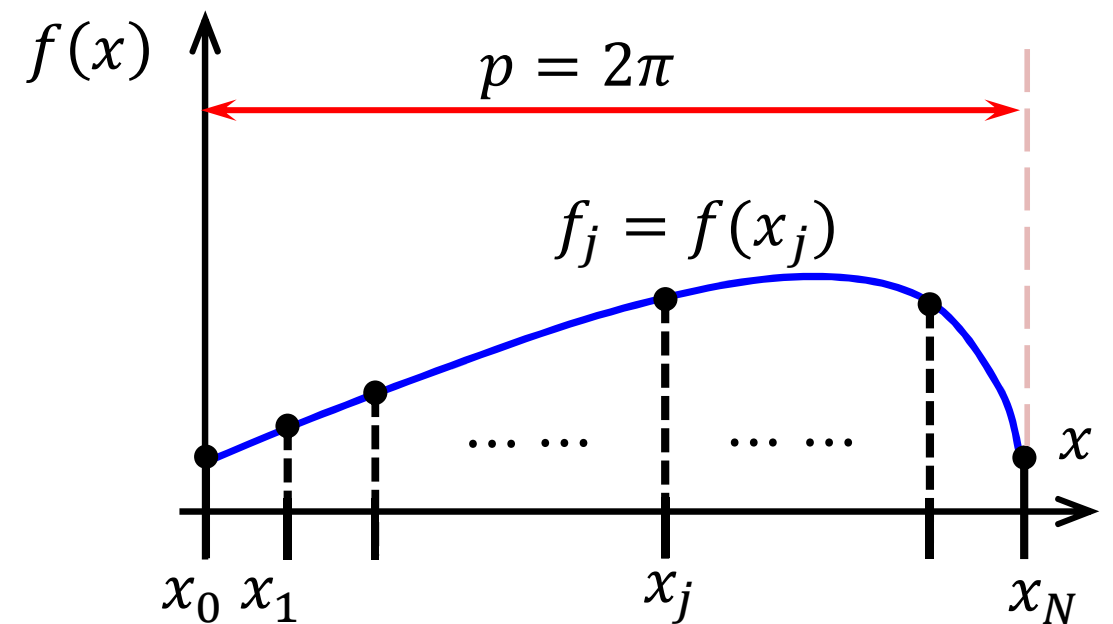
Numerical differentiation with spectral methods

- Now, for discrete periodic function f_j defined at $x_j, j = 0, 1, \dots, N - 1$:

$$f_j = \sum_{n=-\frac{N}{2}}^{\frac{N}{2}-1} \hat{f}_n e^{inx_j} = \sum_{n=0}^{N-1} \hat{f}_n e^{inx_j}$$

where \hat{f}_n is the discrete Fourier transform:

$$\hat{f}_n = \frac{1}{N} \sum_{j=0}^{N-1} f(x_j) e^{-inx_j}$$



- Can the differentiation of f defined at x_j , i.e., $\left. \frac{df}{dx} \right|_j$ be evaluated by:

$$\left. \frac{df}{dx} \right|_j = \sum_{n=-\frac{N}{2}}^{\frac{N}{2}-1} (in) \hat{f}_n e^{inx_j}$$

Similarly,

$$\left. \frac{d^2 f}{dx^2} \right|_j = \sum_{n=-\frac{N}{2}}^{\frac{N}{2}-1} (in)^2 \hat{f}_n e^{inx_j} = - \sum_{n=-\frac{N}{2}}^{\frac{N}{2}-1} n^2 \hat{f}_n e^{inx_j}$$

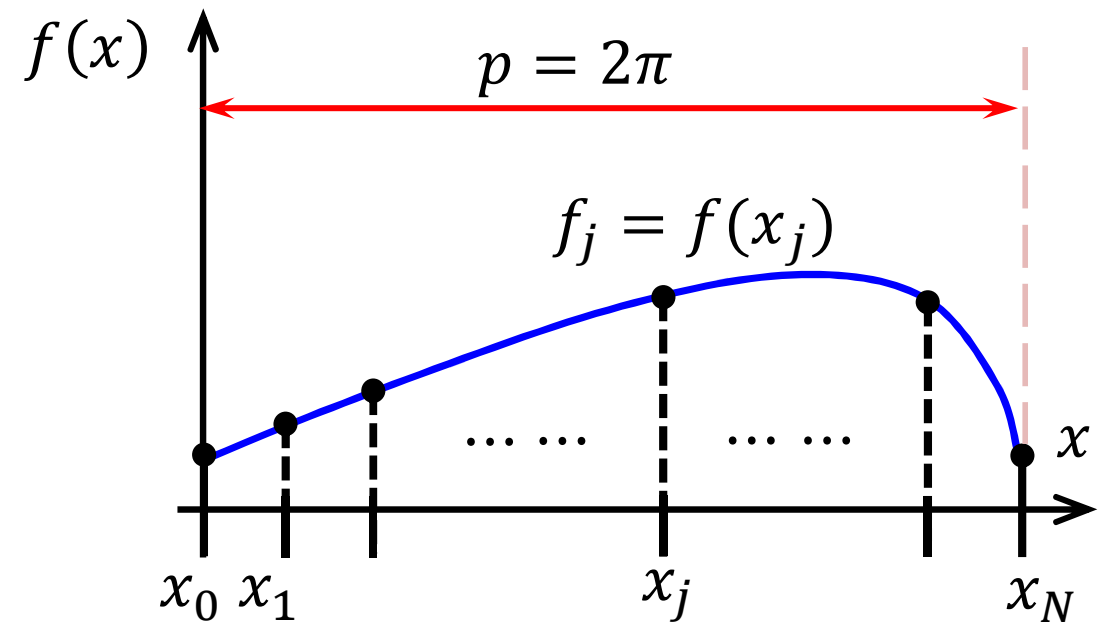
Numerical differentiation with spectral methods

- Now, for discrete periodic function f_j defined at $x_j, j = 0, 1, \dots, N - 1$:

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Numerical differentiation with spectral methods

- The **spectral** derivative is much more accurate than any **finite-difference** schemes for *periodic functions*.
- The major cost involved is the use of fast Fourier transform.
- However, it is inaccurate and does not converge when the derivative is discontinuous.

Numerical integration: The midpoint rule

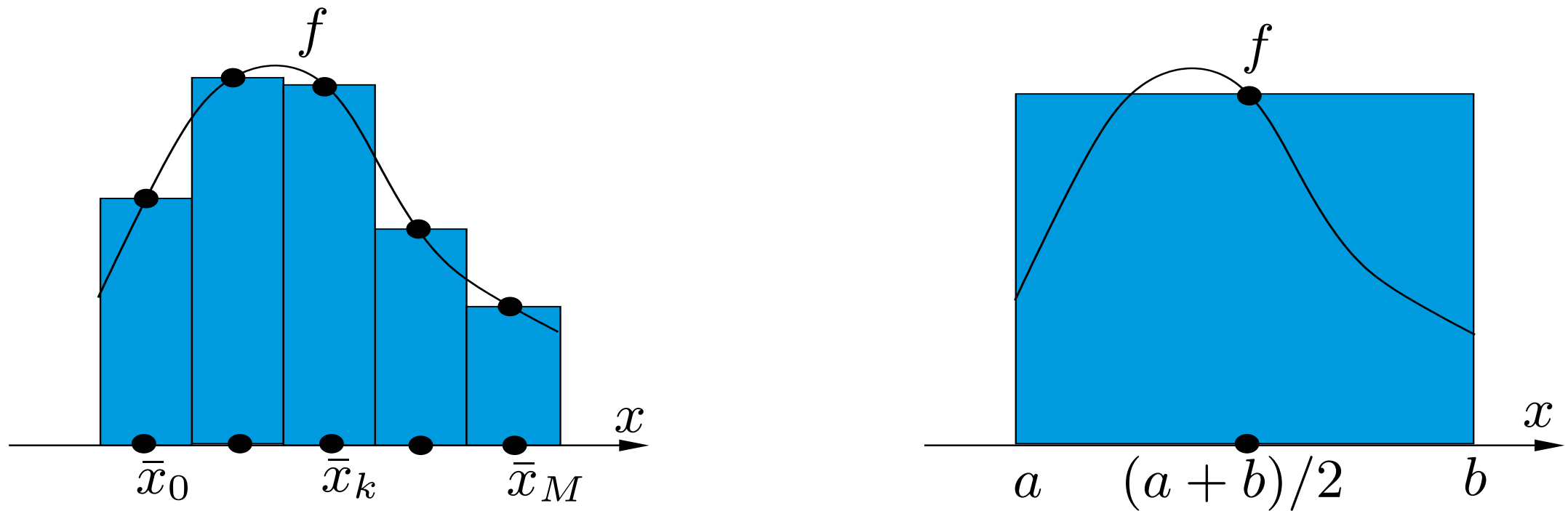


Fig. 4.3. The composite midpoint formula (*left*); the midpoint formula (*right*)

$$I_{mp}^c(f) = H \sum_{k=1}^M f(\bar{x}_k)$$

$$I_{mp}(f) = (b - a) f[(a + b)/2]$$

Numerical integration: The trapezoidal rule

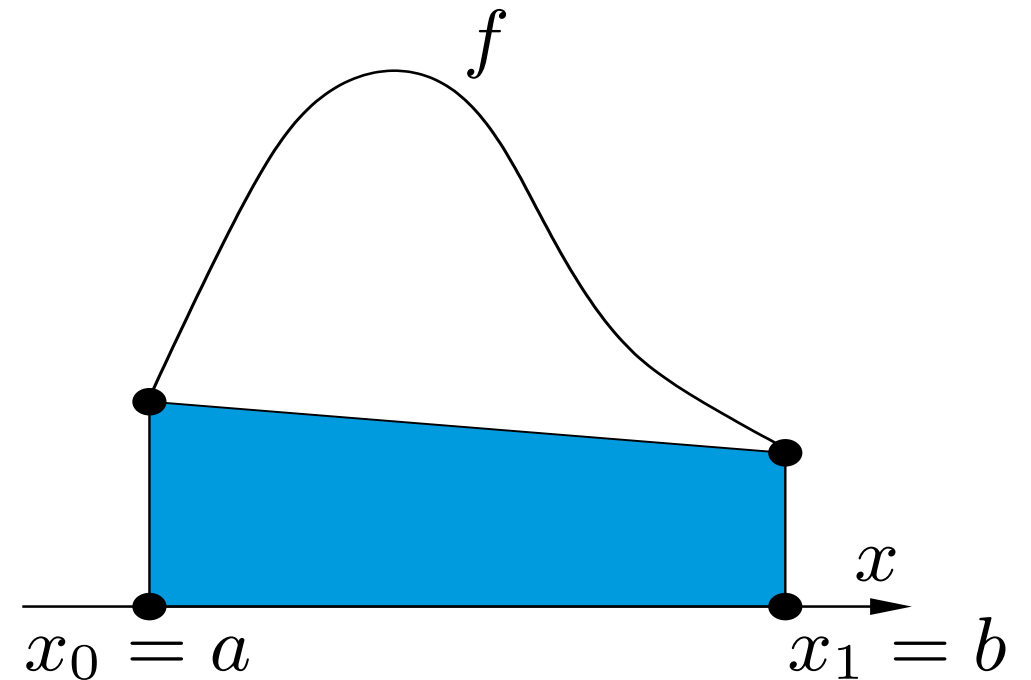
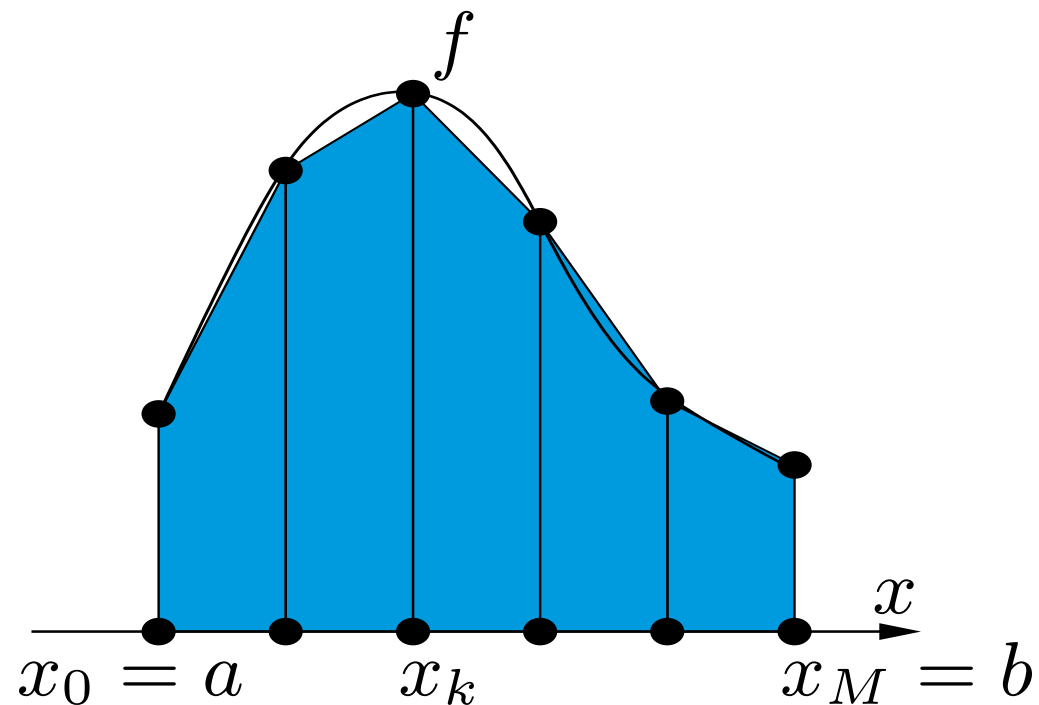


Fig. 4.4. Composite trapezoidal formula (*left*); trapezoidal formula (*right*)

$$\begin{aligned} I_t^c(f) &= \frac{H}{2} \sum_{k=1}^M [f(x_k) + f(x_{k-1})] \\ &= \frac{H}{2} [f(a) + f(b)] + H \sum_{k=1}^{M-1} f(x_k) \end{aligned}$$

$$I_t(f) = \frac{b-a}{2} [f(a) + f(b)]$$

Numerical integration: Simpson's rule

$$I_s(f) = \frac{b-a}{6} [f(a) + 4f((a+b)/2) + f(b)]$$

Simpson's formula

$$I_s^c(f) = \frac{H}{6} \sum_{k=1}^M [f(x_{k-1}) + 4f(\bar{x}_k) + f(x_k)]$$

The composite Simpson's rule

Gauss-Legendre quadrature

$$I_s(f) = \sum_{j=1}^n w_j f(x_j)$$

w_j weights
 x_j nodes

| n | x_j | w_j |
|-----|---|---|
| 1 | $\{\pm 1/\sqrt{3}\}$ | $\{1\}$ |
| 2 | $\{\pm\sqrt{15}/5, 0\}$ | $\{5/9, 8/9\}$ |
| 3 | $\left\{ \pm(1/35)\sqrt{525 - 70\sqrt{30}}, \right.$ $\left. \pm(1/35)\sqrt{525 + 70\sqrt{30}} \right\}$ | $\{(1/36)(18 + \sqrt{30}),$ $(1/36)(18 - \sqrt{30})\}$ |
| 4 | $\left\{ 0, \pm(1/21)\sqrt{245 - 14\sqrt{70}} \right.$ $\left. \pm(1/21)\sqrt{245 + 14\sqrt{70}} \right\}$ | $\{128/225, (1/900)(322 + 13\sqrt{70})$ $(1/900)(322 - 13\sqrt{70})\}$ |

Table 4.1. Nodes and weights for some quadrature formulae of Gauss-Legendre on the interval $(-1, 1)$. Weights corresponding to symmetric couples of nodes are reported only once

Monte Carlo approximation

$$\mathbb{E}_{x \sim p(x)} [f(x)] = \int f(x) p(x) dx \approx \frac{1}{n} \sum_{i=1}^n f(x_i),$$

where x_i are drawn iid from $p(x)$

Monte Carlo approximation

Example: estimating π by Monte Carlo integration

MC approximation can be used for many applications, not just statistical ones. Suppose we want to estimate π . We know that the area of a circle with radius r is πr^2 , but it is also equal to the following definite integral:

$$I = \int_{-r}^r \int_{-r}^r \mathbb{I}(x^2 + y^2 \leq r^2) dx dy \quad (2.99)$$

Hence $\pi = I/(r^2)$. Let us approximate this by Monte Carlo integration. Let $f(x, y) = \mathbb{I}(x^2 + y^2 \leq r^2)$ be an indicator function that is 1 for points inside the circle, and 0 outside, and let $p(x)$ and $p(y)$ be uniform distributions on $[-r, r]$, so $p(x) = p(y) = 1/(2r)$. Then

$$I = (2r)(2r) \int \int f(x, y) p(x) p(y) dx dy \quad (2.100)$$

$$= 4r^2 \int \int f(x, y) p(x) p(y) dx dy \quad (2.101)$$

$$\approx 4r^2 \frac{1}{S} \sum_{s=1}^S f(x_s, y_s) \quad (2.102)$$

We find $\hat{\pi} = 3.1416$ with standard error 0.09 (see Section 2.7.3 for a discussion of standard errors). We can plot the points that are accepted/ rejected as in Figure 2.19.

