Polytropes and Models of White Dwarf Stars

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Abstract

I. Introduction

Understanding stellar mechanics requires the use of mathematical models of the internal structure of stars. We understand stars to be nearly spherical collections of hot gas held together by self-gravitation and holding themselves up against gravitational collapse by fluid and radiation pressure. Equations from Newtonian gravitational and fluid mechanics allow us to create models for the how the pressure, density, and temperature of stars varies with their mass and size, but they are complicated and highly coupled.[2] Additionally, as mass and density increase, quantum and relativistic effects become important.

Prior to the development of advanced electronic computers, solutions to these equations were difficult to impossible to compute. In 1870, American physicist Jonathan Homer Lane proposed a simplified model[3] by assuming the gas pressure depends only on the density of the gas (a polytropic fluid), eliminating any explicit dependence on temperature and decoupling the equations for pressure, temperature, and density.

$$P(r) = K\rho^{1+1/n} \tag{1}$$

Later Swiss physicist Robert Emden formalized the model in the dimensionless differential equation that bears their names, the Lane-Emden equation, whose derivation can be found in the appendix.

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) + \theta^n = 0 \tag{2}$$

where ξ is a dimensionless function of the radius, θ is a dimensionless function relating

density and pressure, and n is the *polytropic* index of the fluid. Solutions to this equation are called *polytropes*.

While Lane's polytropic simplification seemed unrealistic for most situations conceivable at the time, the vast simplification of its solutions over more realistic models proved a generous return on the trade with results that were still close enough to observed data to be very useful. Closed-form solutions[4] can be found readily found for polytropic indices n=0 and n=1, and with some algebraic substitution another one can be found for n=5 (which has an infinite radius and will not be discussed further), and numerical solutions can be found using techniques known at the time of Lane's original publication.

In the 20th century new utility was found for this model with the development of quantum mechanics and the discovery of white dwarfs, stellar remnants composed of extremely dense gas. The electron density in white dwarfs approaches the density of available quantum energy states; and since the Pauli exclusion principle disallows multiple electrons from occupying the same quantum state[1], this results in what is called *electron degeneracy pressure* which is the main force opposing the dwarf's own gravity since white dwarfs are no longer actively undergoing fusion. White dwarfs are therefore composed of nearly fully degenerate matter. Serendipitously, fully degenerate matter has the following equation of state: [2, pp.163-166]

$$P(r) = K\rho(r)^{\gamma} \tag{3}$$

identical in form to eq. 1, *i.e.*, it is a polytropic gas!

II. METHODS

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III. Results

We obtained the following parameters for non-relativistic (n = 1.5) and relativistic (n = 3) polytropes using a Runge-Kutta scheme in Matlab:

Table 1: Solutions obtained with Runge-Kutta

n	ξ_f	$\theta'(\xi_f)$	$ ho_c$
1.5	3.6838	-0.2033	5.9907
	6.8968	-0.0424	54.1825

Compared with the parameters outlined in *Stellar Interiors*[2] any discrepancies can be attributed to rounding differences.

Table 2: Parameters for n = 1.5 and n = 3 polytropes[2]

\overline{n}	ξ_f	$\theta'(\xi_f)$	$ ho_c$
1.5	3.6538	-0.20330	5.991
3	6.8969	-0.20330 -0.04243	54.1825

IV. DISCUSSION

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