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## ABSTRACT

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## 1 COSMIC SHEAR

## 1.1 Power spectra

We compute the shear cross correlation between 2 tomographic bins as

$$C_{\ell}^{ij} = \left\langle \gamma_{i} \gamma_{j}^{*} \right\rangle = \frac{1}{N} \int dz_{p_{i}} p(z_{p_{i}}) \int dz_{s_{i}} p(z_{s_{i}} | z_{p_{i}}) \mathcal{W}(z_{p_{i}}, z_{s_{i}}) \int dz_{p_{j}} p(z_{p_{j}}) \int dz_{s_{j}} p(z_{s_{j}} | z_{p_{j}}) \mathcal{W}(z_{p_{j}}, z_{s_{j}})$$

$$\int dz_{l} \frac{c}{H(z_{l})} \frac{\overline{\rho}_{m}}{\sum_{c} (z_{l}, z_{s_{i}})} \frac{\overline{\rho}_{m}}{\sum_{c} (z_{l}, z_{s_{j}})} \frac{1}{f_{k}(\chi_{l})^{2}} P_{mm}(z_{l})$$

$$N = \int dz_{p_{i}} p(z_{p_{i}}) \int dz_{s_{i}} p(z_{s_{i}} | z_{p_{i}}) \mathcal{W}(z_{p_{i}}, z_{s_{i}}) \int dz_{p_{j}} p(z_{p_{j}}) \int dz_{s_{j}} p(z_{s_{j}} | z_{p_{j}}) \mathcal{W}(z_{p_{j}}, z_{s_{j}})$$

$$(2)$$

Our notation is slightly different from many lensing papers.  $z_{p_i}$  denotes the photo-z distribution for sample  $i, z_{s_i}$  denotes the true redshift for these source galaxies.  $p(z_{p_i})$  is the photometric redshift distribution for these galaxies and  $p(z_{s_i}|z_{p_i})$  is the distribution of true redshift for galaxies with photo-z  $z_{p_i}$ . We will use subscript l to denote quantities related to the matter (such as  $z_l$ ) that is lensing the source galaxies. We use  $d\chi_l = dz_l \frac{c}{H(z_l)}$ , lensing weight  $W_L = \frac{\bar{\rho}_m}{\sum_{c}(z_l, z_{s_2}) f_k(\chi_l)}$  where  $f_k(\chi_l)$  is the transverse separation to redshift  $z_l$ .  $P_{mm}(z_l)$  is the matter power spectrum at redshift  $z_l$ . We use the normalization factor N to correctly normalize the computed power spectra.  $\mathcal{W}(z_{p_i}, z_{s_i})$  are the weights that are applied to the source galaxies.

SS: In the code, we assume that  $z_{p_i}$  is the true redshift for now and hence there is not integral over  $dz_{s_i}$