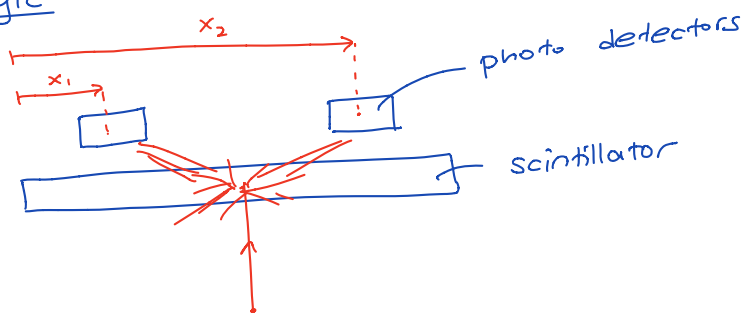


CHAPTER 8 - PLANAR SCINTIGRAPHY

①

Positioning Logic :

in 1D :



$$\bar{x} = \frac{I_1 \cdot x_1 + I_2 \cdot x_2}{I_1 + I_2}$$

Image Formation :

Event position Estimation :

* the height of the response from each photomultiplier tube is related to its distance to the scintillation event.

a_k : amplitude response from photomultiplier tube k
 (x_k, y_k) : positions of " " "

$$Z = \sum_{k=1}^K a_k$$

: total amplitude response
 ("mass" of the light distribution)

Then,

$$X = \frac{1}{Z} \sum_{k=1}^K x_k a_k$$

$$Y = \frac{1}{Z} \sum_{k=1}^K y_k a_k$$

} "Center of mass" equations

So,

(x, y, z)

position amplitude of signal

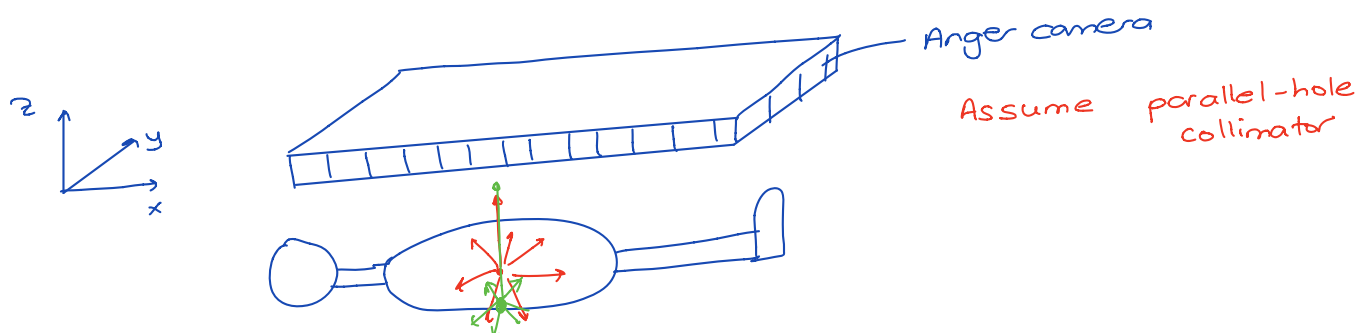
Anger Camera Imaging Equation

(2)

Radioactivity in the body: $A(x, y, z)$

Energy of each photon (γ -ray photon): E

We want to image: photon fluence rate coming from the patient
of photons per unit area per unit time.



* For radioactivity at point (x, y, z) only:

$$\phi_d = \underbrace{\frac{A}{4\pi z^2}}_{\text{photons are emitted in all directions}} \cdot \exp \left\{ \underbrace{- \int_z^0 \mu(x, y, z'; E) dz'}_{\gamma\text{-photons have to pass through the body} \Rightarrow \text{attenuation along the line}} \right\}$$

* Combine all sources in line with the collimation:

$$\phi(x, y) = \int_{-\infty}^{\infty} \frac{A(x, y, z)}{4\pi z^2} \exp \left\{ - \int_z^0 \mu(x, y, z'; E) dz' \right\} dz$$

* This is like projection

- simpler because energy range is restricted
- more complex because two sources of depth-dependent signal loss:

- 1) inverse square law
- 2) object dependent attenuation

\Rightarrow Activity close to the camera contributes more to the image.

Planar source : a simplified case

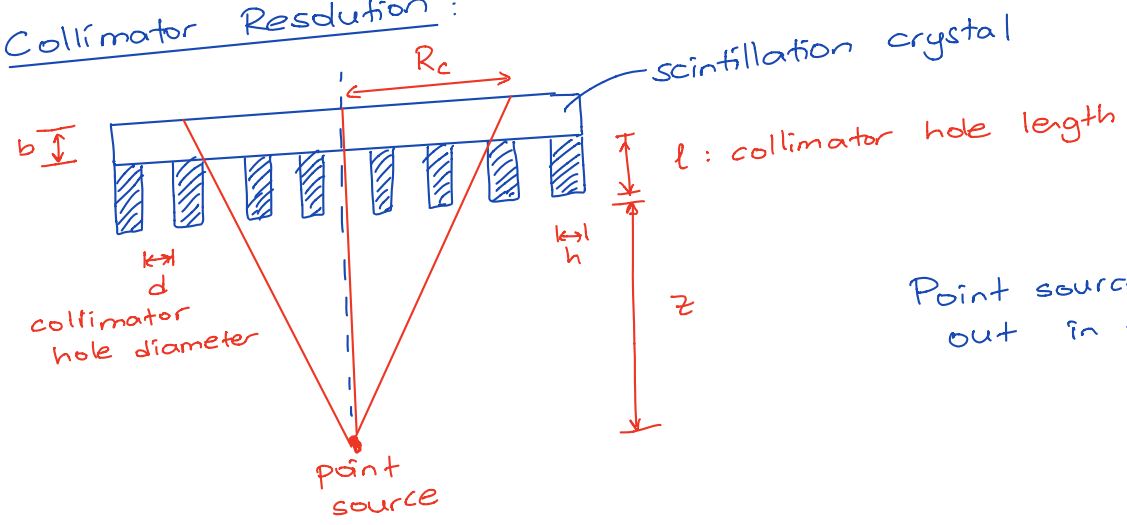
$$A(x,y,z) = A_{z_0}(x,y) \cdot \delta(z-z_0)$$

$$\phi(x,y) = \frac{A_{z_0}(x,y)}{4\pi z_0^2} \exp \left\{ - \int_{z_0}^0 \mu(x,y,z';E) dz' \right\}$$

Even for planar source, intensity depends on μ along the path, which is not uniform.

Image Quality : resolution vs. sensitivity trade-off

Collimator Resolution :



Point source spread out in the image

from geometry: $\frac{d}{l} = \frac{R_c}{l+b+|z|}$

$$R_c = \frac{d}{l} (l+b+|z|)$$

↑ resolution depends on depth $|z|$
blurred more.

So, targets farther away are blurred more.

Solution? make collimator hole longer : $l \rightarrow 2l$

$$R_c = \frac{d}{2l} (2l+b+|z|) = d + \frac{b+|z|}{2l}$$

⇒ reduces sensitivity. fewer activity is detected.

* limit to resolution R_c : $l \rightarrow \infty$, $R_c \rightarrow d$

* making d smaller reduces sensitivity : less events detected.

Trade-off between resolution and sensitivity.

(4)

* Collimator PSF \sim Gaussian with $\text{FWHM} = R_c$

Remember: for a Gaussian function

$$\text{FWHM} = 2\sigma\sqrt{2\ln 2}$$

$$\text{So, } h_c(x, y; |z|) = \exp \left\{ -4(x^2 + y^2) \ln 2 / R_c^2 (|z|) \right\}$$

↓
depends on source depth

* Additional blurring in scintillator itself
→ called the intrinsic resolution of the Anger camera

Also modeled as a Gaussian:

$$h_I(x, y) = \exp \left\{ -4(x^2 + y^2) \ln 2 / R_I^2 \right\}$$

↑
does not depend on source depth

* So, for a planar source: $A(x, y, z) = A_{z_0}(x, y) \delta(z - z_0)$

$$\phi(x, y) = \frac{A_{z_0}(x, y)}{4\pi z_0^2} \exp \left\{ - \int_{z_0}^0 \mu(x, y, z'; E) dz' \right\} * h_c(x, y; |z_0|) * h_I(x, y)$$

Typically, $R_I \ll R_c$. So, collimator's response dominates.

SNR: $\text{SNR} \propto \sqrt{\bar{N}}$

↑ mean total number of acquired photons