

Ultrasound wave travels at the speed of sound

$$c = \sqrt{\frac{1}{k\rho}}$$

k : compressibility
 ρ : density $(\text{m.s}^2/\text{kg})$ } properties of material
 kg/m^3

$c = 330 \text{ m/s}$ in air

$c = 1480 \text{ m/s}$ in water

$c = 4080 \text{ m/s}$ in bone

$c = 1540 \text{ m/s}$ in tissue (average value)

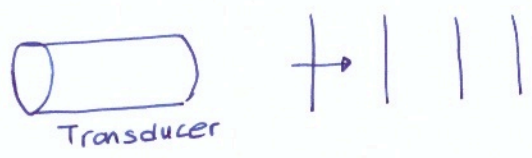
* $\lambda = \frac{c}{f}$: wavelength

Plane waves

an acoustic wave that varies only in one spatial direction and time.

$p(x,y,z,t) = p(z,t)$
 \uparrow
 acoustic pressure

plane wave moving in $+z$ or $-z$ direction



This wave has to satisfy the wave equation:

$$\nabla^2 p = \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} \quad (\text{in 3D}) \quad \left(\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)$$

Laplacian

(reduce to 1D for $p(z,t)$)

$$\frac{\partial^2 p}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2}$$

General solution is:

$$p(z,t) = \phi_f(t - z/c) + \phi_b(t + z/c)$$

\uparrow
 forward-traveling wave

\nwarrow backward-traveling wave

we are interested in forward traveling wave

e.g., $p(z,t) = \cos(k(z - ct))$

\hookrightarrow for constant z , pressure varies sinusoidally with $\omega = kc$
 called "wave number"

$$f = \frac{\omega}{2\pi} = \frac{kc}{2\pi} \quad , \quad \lambda = \frac{c}{f} = \frac{2\pi}{k}$$

Example: $f = 2 \text{ MHz}$ sinusoidal wave. what is λ ?

average $c = 1540 \text{ m/s}$ in tissue

$$\lambda = \frac{c}{f} = \frac{1540 \text{ m/s}}{2 \cdot 10^6 \text{ 1/s}} = 0.77 \text{ mm}$$

Example: An ultrasound transducer is pointing down the $+z$ axis. At time $t=0$, it starts to generate an acoustic pulse with form

$$\phi(t) = (1 - e^{-t/z_1}) e^{-t/z_2}$$

* what is the forward-traveling wave? At what time does the leading edge of the impulse hit an interface 10 cm away from the transducer?

$$\phi_f(z, t) = (1 - e^{-(t-z/c)/z_1}) e^{-(t-z/c)/z_2}$$

it will take $\Delta t = \frac{10 \text{ cm}}{1540 \text{ m/s}} = 64.9 \mu\text{s}$

Acoustic intensity: $I = \frac{p^2}{Z}$

↑ intensity
(or acoustic energy flux)

← pressure

← characteristic impedance of material ($Z = \rho \cdot c$)

↑ density

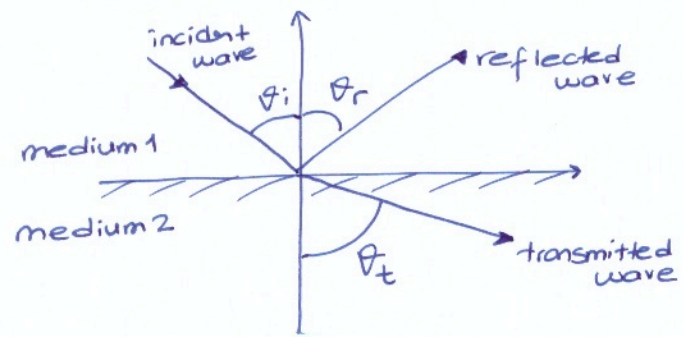
Analogy to electrical circuits:

pressure \longleftrightarrow voltage

Z \longleftrightarrow resistance

I \longleftrightarrow power

Reflection and Refraction at Plane Interfaces



$\theta_i = \theta_r$ (reflection)

Snell's Law:

$$\frac{\sin \theta_i}{\sin \theta_t} = \frac{c_1}{c_2}$$

Example: medium 1: fat
medium 2: liver

, $\theta_i = 45^\circ$, $\theta_r = ?$, $\theta_t = ?$

* $\theta_r = \theta_i = 45^\circ$

$c_1 = 920 \text{ m/s}$ in fat
 $c_2 = 1060 \text{ m/s}$ in liver

(Table 10.1)

$$\left. \begin{array}{l} \sin \theta_t = \sin \theta_i \cdot \frac{c_2}{c_1} = \frac{1}{\sqrt{2}} \cdot \frac{1060}{920} \\ = 0.8147 \\ \theta_t = 54.6^\circ \end{array} \right\}$$

Reflectivity:

$R_I = \frac{I_r}{I_i} = \left(\frac{Z_2 \cos \theta_i - Z_1 \cos \theta_t}{Z_2 \cos \theta_i + Z_1 \cos \theta_t} \right)^2$, intensity reflectivity

$T_I = \frac{I_t}{I_i} = \frac{4 Z_1 Z_2 \cos^2 \theta_i}{(Z_2 \cos \theta_i + Z_1 \cos \theta_t)^2}$, intensity transmittivity

Example: what fraction of acoustic intensity is reflected back from a fat/liver interface at normal incidence?
(error in book) in results

* $\theta_i = 0 \Rightarrow R_I = \left(\frac{Z_2 - Z_1}{Z_2 + Z_1} \right)^2$

if $Z_1 = Z_2$, $R_I = 0$ no reflection!
if $Z_1 \gg Z_2$, $R_I \approx 1$ all reflected
some for $Z_2 \gg Z_1$

Using Table 10.1, $Z_{\text{liver}} = 1.66 \times 10^{-6} \text{ kg/m}^2/\text{s}$
 $Z_{\text{fat}} = 1.35 \times 10^{-6} \text{ kg/m}^2/\text{s}$

$R_I = \left(\frac{1.66 - 1.35}{1.66 + 1.35} \right)^2 = 0.0106$

$\rightarrow \sim 1\%$ reflected
99% transmitted through

* How about muscle/bone interface?

$Z_{\text{bone}} \sim 6 \times 10^{-6} \text{ kg/m}^2/\text{s}$

$Z_{\text{muscle}} \sim 1.7 \times 10^{-6} \text{ kg/m}^2/\text{s}$

$R_I = \left(\frac{6 - 1.7}{6 + 1.7} \right)^2 = 0.32$

$\rightarrow 32\%$ reflected
 \rightarrow more difficult to see behind bone

Attenuation: amplitude of the wave decreases as wave propagates

due to absorption
(conversion to thermal energy),
scattering,
mode conversion

$p(z, t) = A_0 \cdot e^{-\mu_a z} \cdot f(t - z/c)$

μ_a : amplitude attenuation coefficient (cm^{-1})

$\frac{A_z}{A_0} = e^{-\mu_a z}$

$\rightarrow 20 \log_{10} \left(\frac{A_z}{A_0} \right) = -\mu_a z \cdot 20 \log_{10} e$

(attenuation in dB) = $-z \cdot \underbrace{20 (\log_{10} e)}_{\alpha} \mu_a$

phenomenological
(agrees with practice, but not easily supported by theory)

$$\alpha = 8.7 \mu\text{a}$$

↳ attenuation coefficient in dB/cm.
(or absorption coefficient)

← very large for air & bones !!

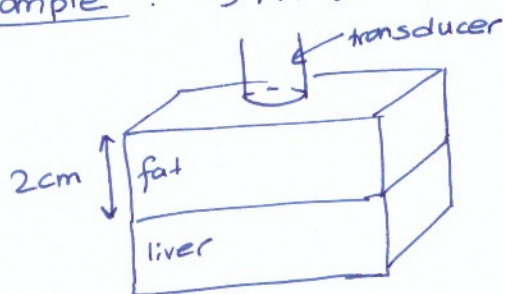
* α depends on frequency

$$\alpha \approx a.f \quad (\text{approximately})$$

↳ dB/cm/MHz

* more attenuation at higher frequencies !!

Example : 5 MHz acoustic pulse



When does the "echo" reflected from the fat/liver interface arrive back at the transducer?

$$c_{\text{fat}} = 1450 \text{ m/s}$$

$$t = \frac{2 \times \overset{\text{round-trip}}{2 \text{ cm}}}{1450 \text{ m/s}} = 27.6 \mu\text{s}$$

* What is the amplitude loss of the returning wave?

from before, $R_I = 0.103$ for fat/liver interface

↑ intensity reflectivity. remember $I = \frac{p^2}{z}$

$$\text{So, amplitude reflectivity} = \sqrt{0.0106} = 0.103$$

$$\text{↳ convert to dB: } 20 \log_{10}(0.103) = -19.7 \text{ dB}$$

Also, attenuation:

$$a_{\text{fat}} = 0.63 \text{ dB/cm/MHz}$$

$$\alpha = 0.63 \times 5 = 3.15 \text{ dB/cm}$$

$$\text{(can also do } 10 \log_{10}(0.0106) = -19.7 \text{ dB)}$$

* So, for 4 cm round-trip : $4 \times 3.15 = 12.6 \text{ dB}$ attenuation

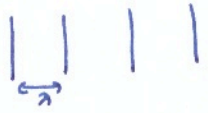
* In total

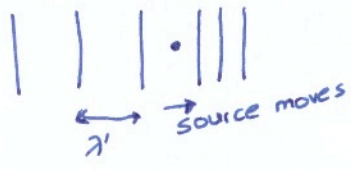
$$\text{dB loss} = 20 \log_{10} \frac{A_z}{A_0} = -19.7 - 12.6 = -32.3 \text{ dB}$$

Doppler Effect

change in frequency of sound due to relative motion of the source and receiver

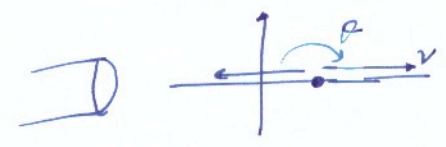
(ambulance siren : approaching vs. moving away)
 ↑ higher pitch ↑ lower pitch

* Distance between wave peaks: $\lambda = \frac{c}{f_0}$ 

* if source moves: $\lambda' = \frac{c}{f_0} + \frac{v}{f_0}$ ^{speed of source} $\Rightarrow f' = \frac{c}{c+v} f_0$


more generally: $f' = \frac{c}{c - v \cos \theta} f_0$

$f_D = f' - f_0 = \left(\frac{v \cos \theta}{c - v \cos \theta} \right) f_0$
 (Doppler freq)



* for $v \ll c$,

$f_D \approx \frac{v \cos \theta}{c} \cdot f_0$

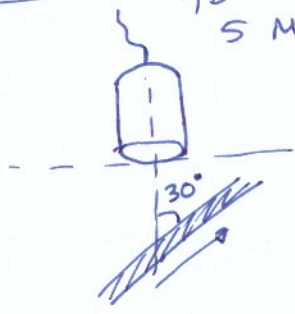
* In ultrasound this is double due to round-trip from source to object

$f_D = \frac{2v \cdot \cos \theta}{c} f_0$

if object moving away, $\theta = 180^\circ \rightarrow f_D$ is negative (lower pitch)

if " " closer, $\theta = 0 \rightarrow f_D$ positive (higher pitch)

Example: $f_D = +500 \text{ Hz}$ measured
 5 MHz transducer, makes 30° angle with direction of blood motion in a vessel



$f_D > 0 \rightarrow$ moving toward transducer

$$v = \frac{c f_D}{2 \cdot \cos \theta \cdot f_0} = \frac{500 \text{ Hz} \cdot 1540 \text{ m/s}}{2 \cdot \sqrt{3}/2 \cdot 5 \cdot 10^6 \text{ Hz}}$$

$$= 0.0889 \text{ m/s}$$

Note: need to know angle. if we thought $\theta = 0^\circ$, then

$$v' = \frac{c f_D}{2 f_0} = 0.077 \text{ m/s}$$

CHAPTER 11 : Ultrasound Imaging Systems

(9)

Received Signal amplitude goes down exponentially as time passes

Gain compensation: $g(t) = \frac{(t)^2 e^{\mu_{act} t}}{K}$

compensate for amplitude loss of signal

Problem : Depth of penetration

if wave travels too deep, the received signal level can go below the detection threshold of the system.

$$d_p = \frac{L}{2af}$$

, where L : loss the system can handle (~ 80 dB)
depth of penetration (we cannot receive signal beyond this depth)

* assuming $a = 1 \text{ dB} \cdot \text{cm}^{-1} \cdot \text{MHz}^{-1}$, $L = 80 \text{ dB}$

freq (MHz)	Depth of Penetration (cm)
1	40
2	20
3	13
5	8
10	4
20	2

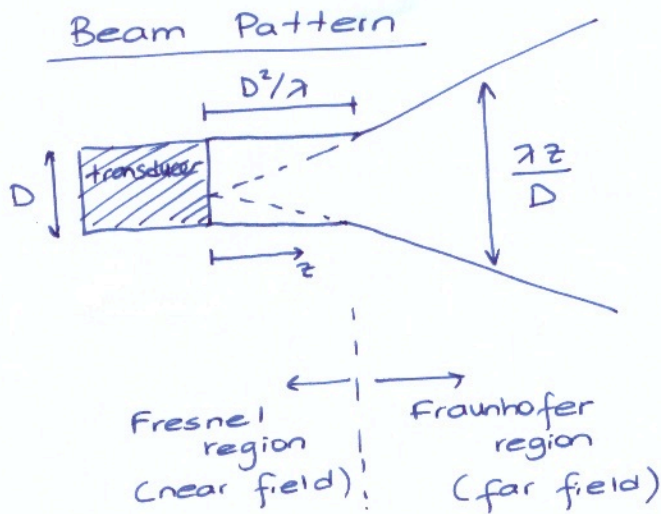
} high frequencies can only image superficially

Pulse Repetition Rate (TR)

Send a short pulse, wait until all echoes die out
Round trip distance $\sim 2d_p$
time it takes $\sim \frac{2d_p}{c}$

So, $TR \geq \frac{2d_p}{c}$

Example: $d_p = 2 \text{ cm}$ @ $20 \text{ MHz} \Rightarrow TR = \frac{2 \times 2 \times 10^{-2}}{1540} \approx 26 \mu\text{s}$
 $d_p = 40 \text{ cm}$ @ $1 \text{ MHz} \Rightarrow TR = \frac{2 \times 40 \times 10^{-2}}{1540} \approx 519 \mu\text{s}$



beamwidth, $w(z) = \begin{cases} D, & z \leq D^2/\lambda \\ \frac{\lambda z}{D}, & z > D^2/\lambda \end{cases}$

lateral
← resolution
worsens
deep in the body

Also; λ larger for lower freq ($\lambda = \frac{c}{f}$)
→ worse resolution

"Optimum" transducer size

$$\frac{D_{opt}^2}{\lambda} = z_{max} \quad (\text{remain in Fresnel region})$$

$$\Rightarrow D_{opt} = \sqrt{\lambda \cdot z_{max}}$$

→ choose as small as possible, limited by attenuation.
(small $\lambda \rightarrow$ high $f \rightarrow$ more att.)

1) $z_{max} = 20 \text{ cm}$ (deep body imaging)

from table → $f = 2 \text{ MHz}$, $\lambda = \frac{1540 \text{ m/s}}{2 \cdot 10^6 \text{ 1/s}} = 0.77 \text{ mm}$

$$D_{opt} = \sqrt{0.77 \text{ mm} \times 20 \text{ cm}} \approx 1.2 \text{ cm}$$

2) $z_{max} = 4 \text{ cm}$ (thyroid imaging)

$$f = 10 \text{ MHz}, \quad \lambda = \frac{1540}{10 \cdot 10^6} = 0.15 \text{ mm}$$

$$D_{opt} = \sqrt{0.15 \text{ mm} \times 4 \text{ cm}} = 2.4 \text{ mm}$$