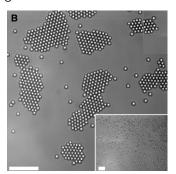
### Self-propelled particles as an active matter system

Étienne Fodor,<sup>1,2</sup> Cesare Nardini,<sup>3</sup> Paolo Visco,<sup>1</sup>
Julien Tailleur,<sup>1</sup> Frédéric van Wijland<sup>1,2</sup>

- 1. Laboratoire Matière et Systèmes Complexes, Université Paris Diderot
  - 2. Yukawa Institute of Theoretical Physics, University of Kyoto
- 3. Department of Applied Mathematics and Theoretical Physics, University of Cambridge

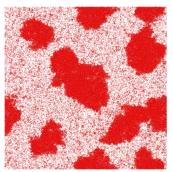
Yukawa Institute of Theoretical Physics University of Kyoto

#### Light-induced clusters of colloids



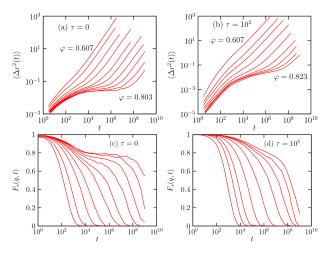
J. Palacci, S. Sacanna, A. P. Steinberg, D. J. Pine, and P. M. Chaikin, Science **339**, 936 (2013)

#### Simulated interacting active colloids



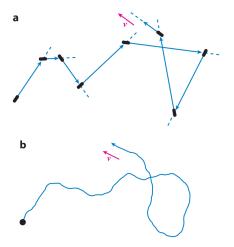
G. S. Redner, M. F. Hagan, and A. Baskaran, Phys. Rev. Lett. 110, 055701 (2013)

#### Dense suspension of active colloids: glassy dynamics



L. Berthier, Phys. Rev. Lett. 112, 220602 (2014)

Run-and-tumble and active Brownian particles



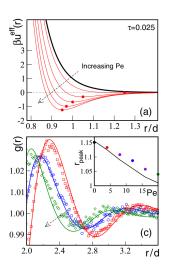
M. E . Cates and J. Tailleur, Ann. Rev. CMP 6, 219 (2015)

#### Competition self-propulsion/interaction

Péclet number

$$\mbox{Pe} = \frac{\mbox{persistence length}}{\mbox{interaction range}} \label{eq:persistence}$$

$$Pe = \sqrt{\frac{persistence time}{relaxation time}}$$



T. F. F. Farage, P. Krinninger, and J. M. Brader, Phys. Rev. E 91, 042310 (2015)

#### Brownian particles

$$\gamma\dot{\mathbf{r}}_{i}=-\sum_{j}
abla_{i}\phi\left(\mathbf{r}_{i}-\mathbf{r}_{j}
ight)+oldsymbol{\xi}_{i}$$
 Mobility  $\mu=\gamma^{-1}$ 

#### Equilibrium fluctuations

$$\langle \boldsymbol{\xi}_i(t) \boldsymbol{\xi}_j(0) \rangle = 2 \gamma T \delta_{ij} \delta(t)$$

$$P_{\mathrm{S}} = \exp \left[ -rac{1}{2T} \sum_{i,j} \phi \left( \mathbf{r}_i - \mathbf{r}_j 
ight) 
ight]$$

#### Self-propelled particles

$$\dot{\mathbf{r}}_{i} = -\mu \sum_{j} \nabla_{i} \phi \left( \mathbf{r}_{i} - \mathbf{r}_{j} \right) + \mathbf{v}_{i}$$

Nonequilibrium fluctuations v

- Random direction with fixed norm
- Memory and non-Gaussian

#### Main questions of interest

- Equation of state: temperature, pressure
- Phase diagram: stationary measure
- Structural relaxation: time scales

Searching for a minimal model of self-propulsion

1 Searching for a minimal model of self-propulsion

Quantifying nonequilibrium

1 Searching for a minimal model of self-propulsion

Quantifying nonequilibrium

1 Searching for a minimal model of self-propulsion

- Quantifying nonequilibrium
- 3 Effective equilibrium regime

Overdamped dynamics

$$\dot{\mathbf{r}}_i = -\mu \nabla_i U + \mathbf{v}_i$$

External potential + pair-wise interactions

$$U = \sum_{i} V_{\text{ext}}(\mathbf{r}_{i}) + \frac{1}{2} \sum_{i,j} \phi(\mathbf{r}_{i} - \mathbf{r}_{j})$$

Overdamped dynamics

$$\dot{\mathbf{r}}_i = -\mu \nabla_i U + \mathbf{v}_i$$

Gaussian fluctuations

$$\langle \mathbf{v}_i(t)\mathbf{v}_j(0)
angle = \delta_{ij}rac{\mu\,T}{ au}\mathrm{e}^{-|t|/ au}$$

Equilibrium reference

$$\langle \mathbf{v}_i(t)\mathbf{v}_j(0)\rangle \underset{\tau\to 0}{\longrightarrow} 2\mu T\delta_{ij}\delta(t)$$

No interaction: harmonic oscillator

$$\dot{r} = -\kappa \mu r + v$$

$$\langle v(t)v(0)
angle = rac{\mu T}{ au} \mathrm{e}^{-|t|/ au}$$

Stationary distribution

$$P_{\mathsf{S}} = \exp\left[-\left(1 + \tau\kappa\mu\right)rac{\kappa r^2}{2T}
ight]$$

No interaction: harmonic oscillator

$$\dot{r} = -\kappa \mu r + v$$

$$\langle v(t)v(0)
angle = rac{\mu T}{ au} \mathrm{e}^{-|t|/ au}$$

Stationary distribution

$$P_{\mathsf{S}} = \exp\left[-\left(1 + \mathsf{Pe}^2\right) rac{\kappa r^2}{2T}
ight]$$

No interaction: harmonic oscillator

$$\dot{r} = -\kappa \mu r + v$$

$$\dot{v} = -\frac{v}{\tau} + \xi$$

$$\langle \xi(t)\xi(0) \rangle = \frac{2\mu T}{\tau^2} \delta(t)$$

Equilibrium underdamped dynamics

$$\dot{r}=p$$
 
$$\dot{p}=-\frac{p}{\tau}-\frac{\mu\kappa}{\tau}\left(1+\frac{\kappa\tau}{\gamma}\right)r+\xi$$

#### Interacting particles

$$\dot{r}_{i} = -\mu \nabla_{i} U + v_{i}$$

$$\dot{v}_{i} = -\frac{v_{i}}{\tau} + \xi_{i}$$

$$\langle \xi_{i}(t)\xi_{j}(0) \rangle = \frac{2\mu T}{\tau^{2}} \delta_{ij} \delta(t)$$

Nonequilibrium underdamped dynamics

$$\dot{p}_i = p_i$$
 
$$\dot{p}_i = -\frac{p_i}{\tau} - \mu p_k \partial_{ik}^2 U - \frac{\mu}{\tau} \nabla_i U + \xi_i$$

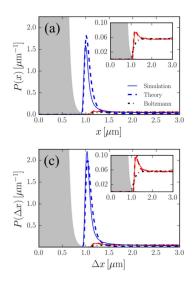
Previous attempt: neglecting inertia

$$\dot{r}_i = p_i$$

$$\not > = -\frac{p_i}{\tau} - \mu p_k \partial_{ik}^2 U - \frac{\mu}{\tau} \nabla_i U + \xi_i$$

Stationary measure

$$P_{\mathsf{S}} = \exp\left[-rac{U}{T} - rac{ au^2}{2T} \left(\partial_i U
ight)^2
ight] |\mathsf{det}\,\mathbb{M}|$$
  $\mathbb{M}_{ij} = \delta_{ij} + au\mu\partial_{ij}^2 U$ 



C. Maggi, U. Marini Bettolo Marconi, N. Gnan, and R. Di Leonardo, Sci. Rep. 5, 10724 (2015)

Previous attempt: neglecting inertia

$$\dot{r}_i = p_i$$

$$\not p_i = -\frac{p_i}{\tau} - \mu p_k \partial^2_{ik} U - \frac{\mu}{\tau} \nabla_i U + \xi_i$$

Mapping to equilibrium dynamics

1 Searching for a minimal model of self-propulsion

Quantifying nonequilibrium

Entropy production rate

$$\sigma = \lim_{t \to \infty} \frac{1}{t} \ln \frac{\mathcal{P}}{\mathcal{P}^{\mathsf{R}}}$$

Brownian particles + non-conservative force

$$\gamma \dot{r}_i = F_i - \partial_i U + \xi_i$$
 $\langle \xi_i(t) \xi_j(0) \rangle = 2 \gamma T \delta_{ij} \delta(t)$ 
 $\sigma = \frac{1}{T} \langle \dot{r}_i F_i \rangle$ 

#### Self-propelled particles

$$\dot{r}_i = -\mu \partial_i U + v_i$$

$$\langle v_i(t)v_j(0)\rangle = \delta_{ij} \frac{\mu T}{\tau} \mathrm{e}^{-|t|/ au}$$

#### Entropy production

$$\sigma = \frac{\tau^2}{2T} \sum_{ijk} \left\langle \dot{r}_i \dot{r}_j \dot{r}_k \partial_{ijk}^3 U \right\rangle$$

#### Péclet number

$$\mathrm{Pe} = \frac{\mathrm{persistence\ length}}{\mathrm{interaction\ range}} = \sqrt{\frac{\tau}{\tau_{\mathrm{R}}}}$$

Relaxation time

$$\tau_{\rm R} = \frac{\left( \text{interaction range} \right)^2}{\mu T}$$

Self-propelled particles: scaled dynamics

- space unit = interaction range
- time unit =  $\tau/Pe$
- energy unit = T

Self-propelled particles: scaled dynamics

$$\dot{p}_i = p_i$$
  $\dot{p}_i = -rac{1}{\mathsf{Pe}} \left( p_i + \mathsf{Pe}^2 p_j \partial_{ij}^2 U 
ight) - \partial_i U + \xi_i$   $\langle \xi_i(t) \xi_i(0) 
angle = rac{2}{\mathsf{Pe}} \delta_{ij} \delta(t)$ 

Self-propelled particles: scaled dynamics

$$\dot{p}_i = p_i$$
  $\dot{p}_i = -rac{1}{\mathsf{Pe}} \left( p_i + \mathsf{Pe}^2 p_j \partial_{ij}^2 U 
ight) - \partial_i U + \xi_i$   $\langle \xi_i(t) \xi_i(0) 
angle = rac{2}{\mathsf{Pe}} \delta_{ij} \delta(t)$ 

Maggi *et al.*: Pe finite and  $Pe \rightarrow 0$ 

Entropy production

$$\sigma = \frac{\mathsf{Pe}^2}{2} \sum_{ijk} \left\langle p_i p_j p_k \partial_{ijk}^3 U \right\rangle$$

Small Pe expansion

$$\sum_{ijk}\left\langle p_{i}p_{j}p_{k}\partial_{ijk}^{3}U
ight
angle =c_{1}\mathsf{Pe}+c_{2}\mathsf{Pe}^{2}+c_{3}\mathsf{Pe}^{3}+\mathcal{O}\left(\mathsf{Pe}^{4}
ight)$$

#### Stationary measure

$$P_{S} = e^{-U - \frac{\rho_{i}^{2}}{2}} \left\{ 1 - \frac{\mathsf{Pe}^{2}}{2} \left[ (\partial_{i} U)^{2} + \mathcal{L} U \right] + \frac{\mathsf{Pe}^{3}}{6} p_{i} \partial_{i} \mathcal{L} U + \mathcal{O} \left( \mathsf{Pe}^{4} \right) \right\}$$

$$\mathcal{L} U = (p_{j} p_{k} - 3\delta_{jk}) \partial_{jk}^{2} U$$

- Position-velocity coupling
- Term Pe<sup>3</sup> odd in velocity

#### Entropy production

$$\sigma = \frac{\mathsf{Pe}^5}{2\mathcal{Z}} \int_r \left( \partial_{ijk}^3 U \right)^2 \mathsf{e}^{-U} + \mathcal{O}\left( \mathsf{Pe}^6 \right)$$

No contribution from the Pe<sup>2</sup> term of the measure

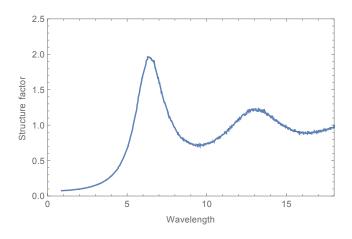
Searching for a minimal model of self-propulsion

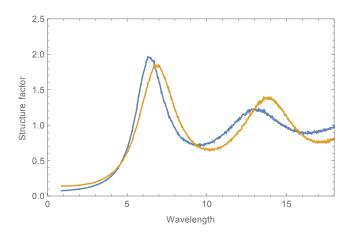
- Quantifying nonequilibrium
- 3 Effective equilibrium regime

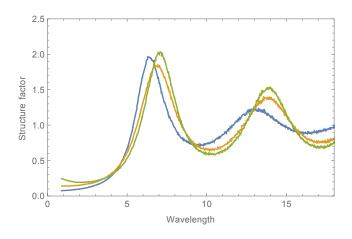
#### Stationary distribution

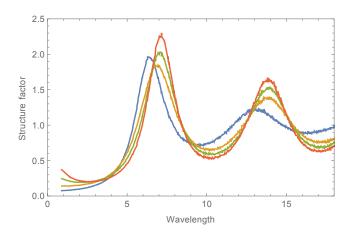
$$P_{\text{S}} = \text{exp}\left\{-U - \text{Pe}^2\left[\frac{\left(\partial_i U\right)^2}{2} - \partial_{ii}^2 U\right] + \mathcal{O}\left(\text{Pe}^4\right)\right\}$$

- Three-body interactions
- Attractive effects from bare repulsion









Brownian particles: external perturbation

$$\gamma \dot{r}_i = f_i - \partial_i U + \xi_i$$

Response function

$$R(t) = \left. \frac{\delta \left\langle r_i(t) \right\rangle}{\delta f_i(0)} \right|_{t=0}$$

Fluctuation-dissipation theorem

$$R(t) = -\frac{1}{T} \frac{\mathsf{d}}{\mathsf{d}t} \langle r_i(t) r_i(0) \rangle$$

Self-propelled particles: external perturbation

$$\dot{r}_i = \mu \left( f_i - \partial_i U \right) + v_i$$

Scaled dynamics

$$\dot{p}_{i} = -rac{1}{\mathsf{Pe}}\left(p_{i} + \mathsf{Pe}^{2}p_{j}\partial_{ij}^{2}U\right) + f_{i} + \mathsf{Pe}\dot{f}_{i} - \partial_{i}U + \xi_{i}$$

$$\langle \xi_{i}(t)\xi_{i}(0)\rangle = rac{2}{\mathsf{Pe}}\delta_{ij}\delta(t)$$

Fluctuation-dissipation relation with Pe = 0

$$R(t) = -\frac{d}{dt} \underbrace{\left[ \langle r_i(t) r_i(0) \rangle + \frac{\text{Pe}^2}{\langle \dot{r}_i(t) \dot{r}_i(0) \rangle} \right]}_{C(t)}$$

Susceptibility 
$$\chi(t) = \int_0^t R(s) ds$$

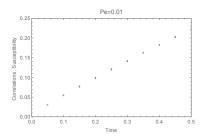
Correlation 
$$\Delta C(t) = C(0) - C(t)$$

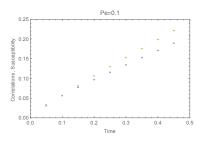
Fluctuation-dissipation relation when  $\text{Pe} \to 0$ 

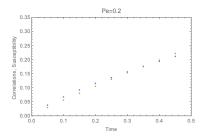
$$R(t) = -\frac{d}{dt} \underbrace{\left[ \langle r_i(t) r_i(0) \rangle + \text{Pe}^2 \langle \dot{r}_i(t) \dot{r}_i(0) \rangle \right]}_{C(t)} + \mathcal{O}\left(\text{Pe}^4\right)$$

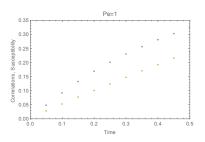
Susceptibility 
$$\chi(t) = \int_0^t R(s) ds$$

Correlation 
$$\Delta C(t) = C(0) - C(t)$$









#### Conclusion

# Minimal model of self-propulsion Gaussian exponentially correlated fluctuations

- Stationary measure in position-velocity
- Entropy production
- Fluctuation dissipation relation

#### Outlook

#### Future directions

- Collective modes
   Hydrodynamic equations, tracer dynamics
- Non-Gaussian white fluctuations | H. Hayakawa
   Static structure, phase separation