

Modeling Active Fluctuations in Living Matter

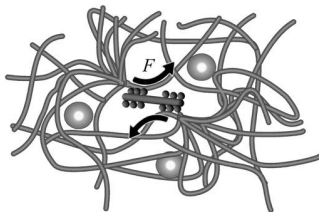
Étienne Fodor¹, Wylie W. Ahmed², Timo Betz², Matthias Bussonnier²,
Nir S. Gov³, Ming Guo⁴, Hisao Hayakawa⁵, Kiyoshi Kanazawa⁵,
Paolo Visco¹, David A. Weitz⁴, Frédéric van Wijland¹



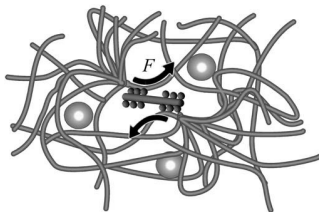
1. Laboratoire Matière et Systèmes Complexes, Université Paris Diderot
2. Laboratoire Physico-Chimie Curie, Institut Curie
3. Department of Chemical Physics, Weizmann Institute of Science
4. School of Engineering and Applied Sciences, Harvard University
5. Yukawa Institute for Theoretical Physics, Kyoto University

Theory Club

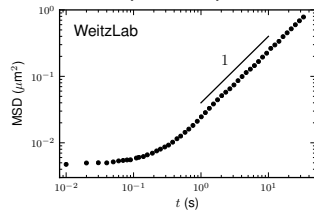
Introduction



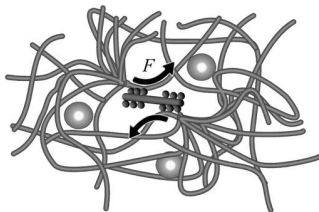
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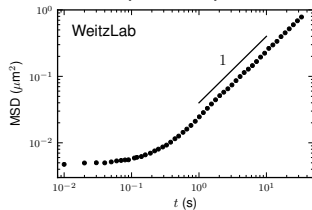
Mean square displacement



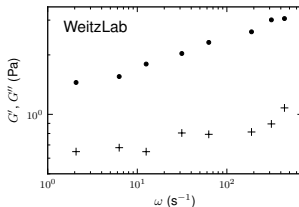
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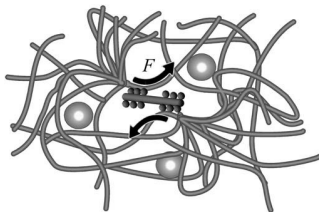
Mean square displacement



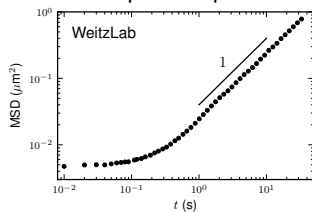
Complex modulus



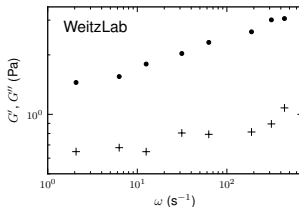
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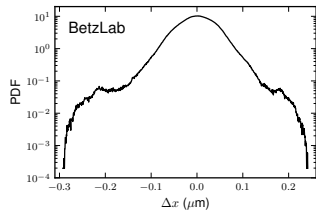
Mean square displacement



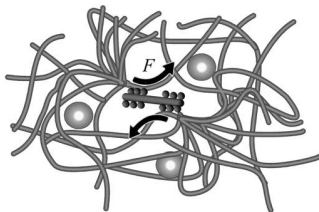
Complex modulus



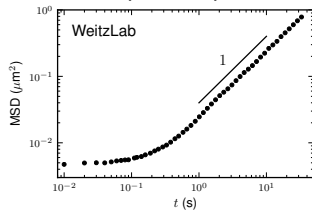
PDF of displacement



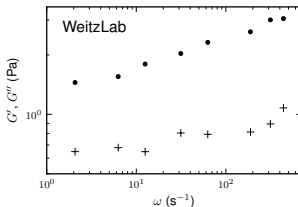
Introduction



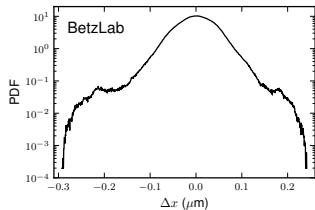
Mean square displacement



Complex modulus



PDF of displacement



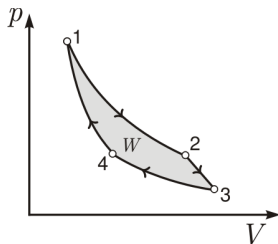
Is it possible to extract information about molecular motor activity?

Propose a model for the tracers' dynamics

- 1 Modeling tracer's dynamics

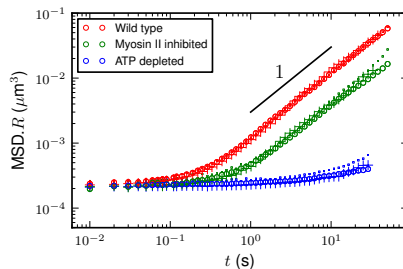
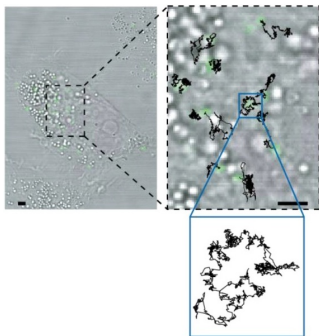
Outline

- 1 Modeling tracer's dynamics
- 2 New kinds of experiments



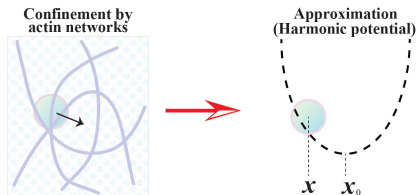
Modeling tracer's dynamics

Experimental results



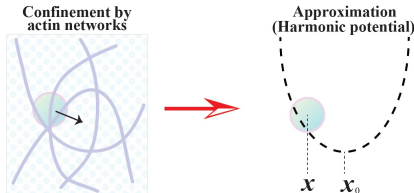
Modeling tracer's dynamics

Equation of motion



Modeling tracer's dynamics

Equation of motion



Tracer's dynamics

$$m \frac{d^2 \mathbf{r}}{dt^2} = -\nabla U + \mathbf{F}_s + \mathbf{F}_{th}$$

Harmonic potential: $U = \frac{k}{2}(\mathbf{r} - \mathbf{r}_0)^2$

Stokes force: $\mathbf{F}_s = -\gamma \frac{d\mathbf{r}}{dt}$

Gaussian white noise: \mathbf{F}_{th}

Modeling tracer's dynamics

Equation of motion

Tracer's dynamics

$$\begin{aligned}\frac{d\mathbf{r}}{dt} &= -\frac{1}{\tau_d}(\mathbf{r} - \mathbf{r}_0) + \sqrt{2D_T}\boldsymbol{\xi} \\ \frac{d\mathbf{r}_0}{dt} &= -\frac{\varepsilon}{\tau_d}(\mathbf{r}_0 - \mathbf{r}) + \sqrt{2\varepsilon D_T}\boldsymbol{\xi}_0\end{aligned}$$

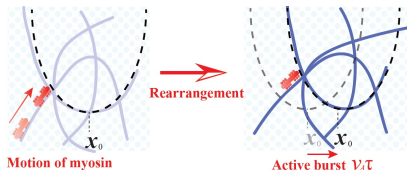
Modeling tracer's dynamics

Equation of motion

Tracer's dynamics

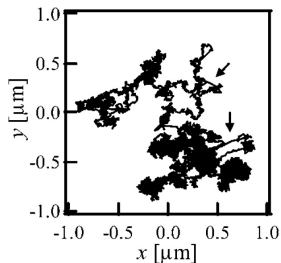
$$\begin{aligned}\frac{d\mathbf{r}}{dt} &= -\frac{1}{\tau_d}(\mathbf{r} - \mathbf{r}_0) + \sqrt{2D_T}\boldsymbol{\xi} \\ \frac{d\mathbf{r}_0}{dt} &= -\frac{\varepsilon}{\tau_d}(\mathbf{r}_0 - \mathbf{r}) + \sqrt{2\varepsilon D_T}\boldsymbol{\xi}_0 + \mathbf{v}_A\end{aligned}$$

Active motion of local minimum



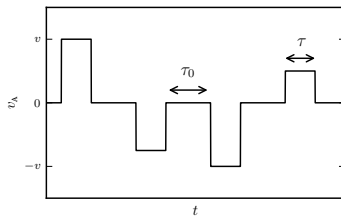
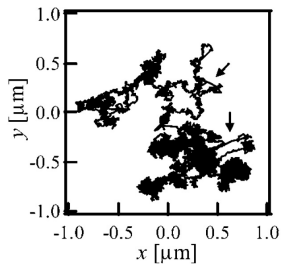
Modeling tracer's dynamics

Active burst's statistics



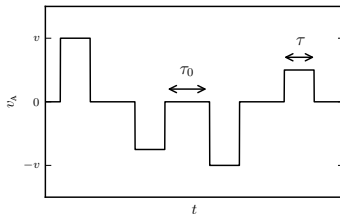
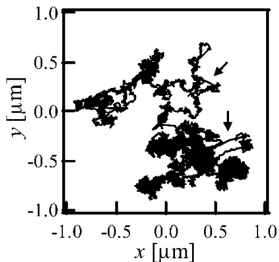
Modeling tracer's dynamics

Active burst's statistics



Modeling tracer's dynamics

Active burst's statistics

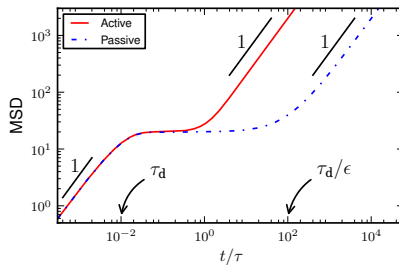


2-time correlation function

$$\langle v_A(t) v_A(0) \rangle = \frac{D_A}{\tau} e^{-|t|/\tau}, \quad D_A = \frac{(v\tau)^2}{3(\tau + \tau_0)}$$

Modeling tracer's dynamics

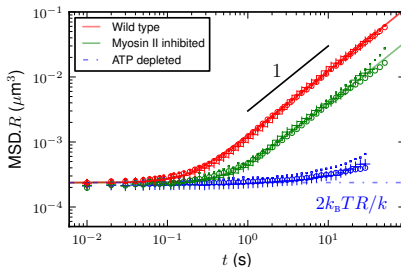
Tracer's statistics



Modeling tracer's dynamics

Fitting experimental results

Short time Confinement
Large time Free diffusion
 $\text{MSD} \sim 2D_A t$

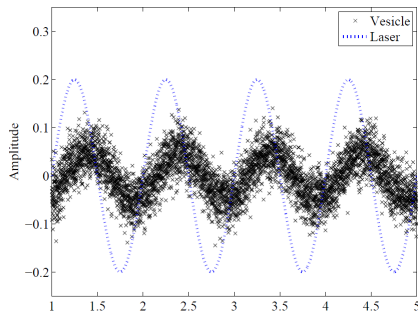


Microscopic features

- Typical time of activity: $\tau_{\text{wt}} \simeq 0.15 \text{ s}$, $\tau_b \simeq 0.41 \text{ s}$
- Amplitude of active fluctuations:
 $D_{A,\text{wt}} \simeq 5.2 \cdot 10^{-3} D_T$, $D_{A,b} \simeq 1.6 \cdot 10^{-3} D_T$

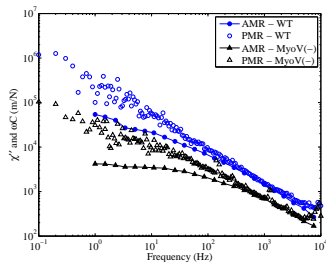
New kinds of experiments

Response function



New kinds of experiments

Fluctuation dissipation

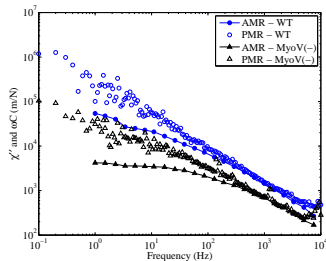


Fluctuation dissipation

$$\chi''(\omega) = \frac{\omega C(\omega)}{2k_B T}$$

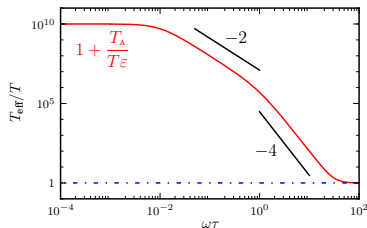
New kinds of experiments

Fluctuation dissipation



Fluctuation dissipation

$$\chi''(\omega) = \frac{\omega C(\omega)}{2k_B T}$$

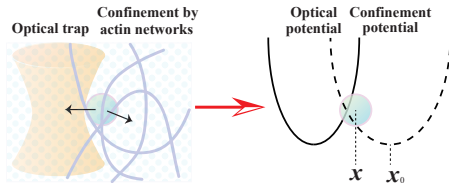


Effective temperature

$$k_B T_{\text{eff}}(\omega) = \frac{\omega C(\omega)}{2\chi''(\omega)}$$

New kinds of experiments

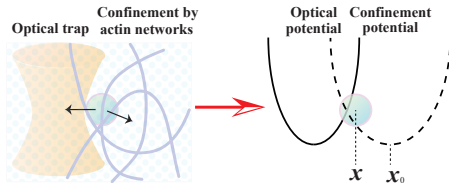
Applying external quadratic potential



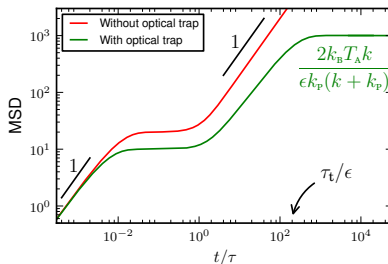
$$U_P = \frac{k_P}{2} x^2$$

New kinds of experiments

Applying external quadratic potential



$$U_P = \frac{k_P}{2} x^2$$



New kinds of experiments

Applying external quadratic potential

$$U_P = U_P(x, \{a_i(t)\}) \quad \rightarrow \quad \mathcal{W} = \sum_i \int da_i \frac{\partial U_P}{\partial a_i}$$

New kinds of experiments

Applying external quadratic potential

$$U_P = U_P(x, \{a_i(t)\}) \quad \rightarrow \quad \mathcal{W} = \sum_i \int da_i \frac{\partial U_P}{\partial a_i}$$

Quasistatic protocol

$$U_P = \frac{k_P}{2} x^2, \quad k_P : k_i \longrightarrow k_f$$

$$W_H = \frac{1}{2} \int_{k_i}^{k_f} dk_P \langle x^2 \rangle_{ss}$$

New kinds of experiments

Applying external quadratic potential

$$U_P = U_P(x, \{a_i(t)\}) \quad \rightarrow \quad \mathcal{W} = \sum_i \int da_i \frac{\partial U_P}{\partial a_i}$$

Quasistatic protocol

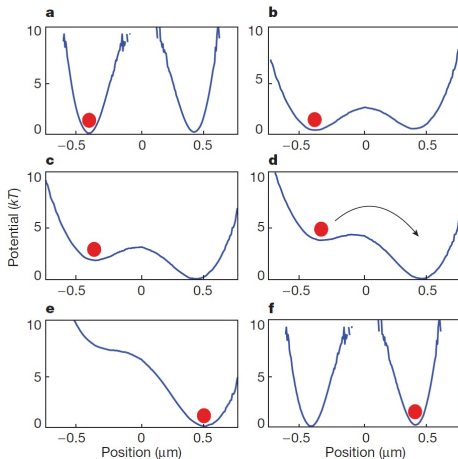
$$U_P = \frac{k_P}{2} x^2, \quad k_P : k_i \longrightarrow k_f$$

$$\mathcal{W}_H = \frac{1}{2} \int_{k_i}^{k_f} dk_P \langle x^2 \rangle_{ss}$$

$$\mathcal{W}_H = \frac{T_A}{2\varepsilon} \ln \left[\frac{k_f(k + k_i)}{k_i(k + k_f)} \right] + \mathcal{O}(\varepsilon^0)$$

New kinds of experiments

Thermodynamic cycle with quartic potential



$$U_P = \frac{k_P}{2}x^2 + \frac{b_P}{4}x^4$$

Naert *et al.*, Nature 2012

New kinds of experiments

Thermodynamic cycle with quartic potential

Quasistatic protocol

$$U_P = \frac{k_P}{2}x^2 + \frac{b_P}{4}x^4, \quad k_P : k_i \longrightarrow k_f$$

$$W_Q = \frac{1}{2} \int_{k_i}^{k_f} dk_P \langle x^2 \rangle_{ss}$$

New kinds of experiments

Thermodynamic cycle with quartic potential

Quasistatic protocol

$$U_P = \frac{k_P}{2}x^2 + \frac{b_P}{4}x^4, \quad k_P : k_i \longrightarrow k_f$$

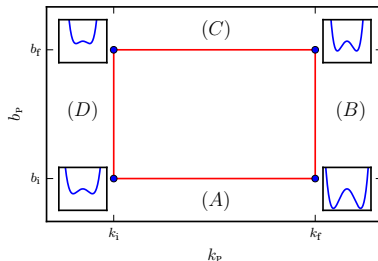
$$W_Q = \frac{1}{2} \int_{k_i}^{k_f} dk_P \langle x^2 \rangle_{ss}$$

$$W_Q = W_H + W_P + \mathcal{O}(b_P^2)$$

$$\frac{W_P}{b_P} = \left(\frac{T_A}{\varepsilon} \right)^2 f(k, k_i, k_f, \tau, \tau_d) + \mathcal{O}(1/\varepsilon)$$

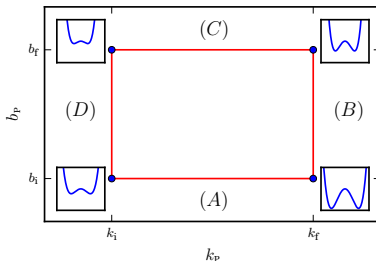
New kinds of experiments

Thermodynamic cycle with quartic potential



New kinds of experiments

Thermodynamic cycle with quartic potential

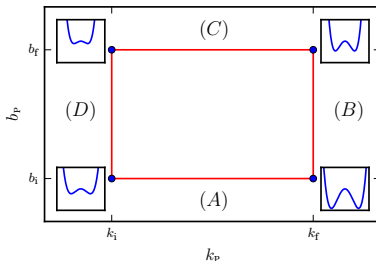


Quasistatic protocol

$$W_C = \frac{1}{2} \oint_C dk_P \langle x^2 \rangle_{ss} + \frac{1}{4} \oint_C db_P \langle x^4 \rangle_{ss}$$

New kinds of experiments

Thermodynamic cycle with quartic potential



Quasistatic protocol

$$W_C = \frac{1}{2} \oint_C dk_P \langle x^2 \rangle_{ss} + \frac{1}{4} \oint_C db_P \langle x^4 \rangle_{ss}$$

$$W_C = \left(\frac{T_A}{\varepsilon} \right)^2 f_C(k, k_i, k_f, b_i, b_f, \tau, \tau_d) + \mathcal{O}(1/\varepsilon)$$