

# PHYS 311 Lecture 17

#physics

#phys311

#lecturenotes

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## Finite Potential Well (Finite Square Well)

More can be found in the book.

### Finite Attractive Potential Well

Finite (attractive) potential well.

Bound stationary states  $E = -|E|$

$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi_{(L)} = -|E| \psi_{(L)}$$
$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi_M = (-|E| + V_0) \psi_M$$

$> 0$

$$\psi_L = A_L e^{Kx}$$
$$\psi_R = A_R e^{-Kx}$$
$$\frac{\hbar^2 K^2}{2m} = |E|$$

### Bound Stationary States

Should be:  $E = -|E|$

$$\frac{-\hbar}{2m} \frac{\partial^2}{\partial x^2} \psi_{LR} = -|E| \psi_{LR}$$

$$\frac{-\hbar}{2m} \frac{\partial^2}{\partial x^2} \psi_M = (-|E| + V_0) \psi_M$$

We see that  $-|E| + V_0 > 0$

$$\begin{aligned}\psi_L &= A_L e^{\kappa x} \\ \psi_R &= A_R e^{-\kappa x} \\ \frac{\hbar \kappa}{2m} &= |E|\end{aligned}$$

We see that wave function and its derivative should be continuous.

$$\psi_M = \alpha \sin(kx) + \beta \cos(kx), \quad \frac{\hbar^2 k^2}{2m} = V_0 - |E|$$

Since this is an **even potential** we can write this to apply the boundary condition easier.

$$\begin{aligned}\psi^{even}(x) &= \begin{cases} A e^{\kappa x} \\ \beta \cos(kx) \\ A e^{-\kappa x} \end{cases} \\ \psi^{odd}(x) &= \begin{cases} \tilde{A} e^{\kappa x} \\ \alpha \sin(kx) \\ -\tilde{A} e^{-\kappa x} \end{cases}\end{aligned}$$

Thus

$$\frac{\hbar k^2}{2m} + \frac{\hbar \kappa^2}{2m} = V_0$$

**even:**

$$\begin{aligned}A e^{\kappa L/2} &= \beta \cos(k \frac{L}{2}) \\ \kappa A e^{-\kappa L/2} &= -k \beta \sin(k(-\frac{L}{2})) = k \beta \sin(k \frac{L}{2})\end{aligned}$$

even

$$\tan(\frac{kL}{2}) \frac{kL}{2} = \frac{\kappa L}{2}$$

See  $A/\beta$ :

$$= e^{\kappa L/2} \cos(kL/2) = \frac{k}{\kappa} e^{\kappa L/2} \sin(k \frac{L}{2})$$

**odd:**

$$\begin{aligned}\tilde{A}e^{kL/2} &= -\alpha \sin(kL/2) \\ \kappa \tilde{A}e^{kL/2} &= \alpha k \cos\left(\frac{kL}{2}\right) \\ \frac{\tilde{A}}{\alpha} &= -e^{-\kappa L/2} \sin\left(\frac{kL}{2}\right) \\ \frac{\tilde{A}}{\alpha} &= \frac{k}{\kappa} \cos\left(\frac{kL}{2}\right) e^{-\kappa L/2}\end{aligned}$$

## Summarizing all the equations

$$\begin{aligned}\text{even: } \tan\left(\frac{kL}{2}\right) \frac{kL}{2} &= \frac{\kappa L}{2} \\ \text{odd: } -\cot\left(\frac{kL}{2}\right) \frac{kL}{2} &= \frac{\kappa L}{2}\end{aligned}$$

and for all

$$\frac{\hbar^2 k^2 L^2}{4 \cdot 2m} + \frac{\hbar^2 \kappa^2 L^2}{2m \cdot 4} = \frac{V_0 L^2}{4}$$

$$\begin{aligned}\tan(x)x &= y \text{ even} \\ -\cot(x)x &= y \text{ odd}\end{aligned}$$

$$x^2 + y^2 = \frac{2mV_0L^2}{???}$$

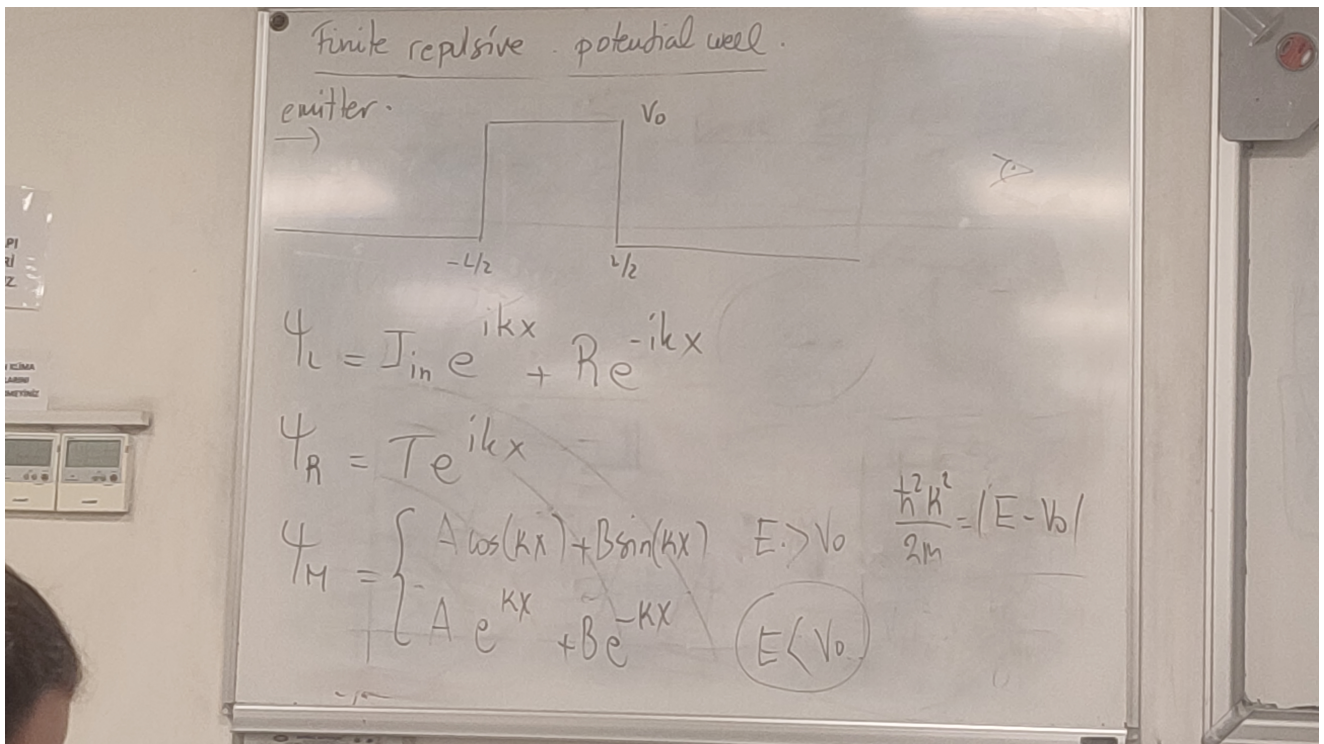
{PUT GRAPH HERE}

???

"Bir yere kadar odd diye bir sey yok!"

"Sonsuz tane bound state olamaz!" Hidrojen atomu buna tam uymuyor.  $r=0$ 'dan itibaren sonsuz potansiyel var. Potansiyel sonlu kaliyorsa bound state'leri de sayilir olmak zorunda

## Finite Repulsive Potential Well



$$\psi_L = I_{in} e^{ikx} + R e^{-ikx}$$

$$\psi_R = T e^{ikx}$$

$$\psi_M = \begin{cases} A \cos(\kappa x) + B \sin(\kappa x), & E > V_0 \\ A e^{\kappa x} + B e^{-\kappa x}, & E < V_0 \end{cases}$$

Remember:  $\frac{\hbar^2 \kappa^2}{2m} = |E - V_0|$

( $I_{in}$  is on me  $\rightarrow$  4 equations, 4 unknowns.)

Then:

At ( $x = -L/2$ ):

$$I_{in} e^{-ikL/2} + R e^{ikL/2} = A e^{kL/2} + B e^{-kL/2}$$

$$k \left( I_{in} e^{-ikL/2} - R e^{ikL/2} \right) = k \left( A e^{kL/2} - B e^{-kL/2} \right)$$

At ( $x = L/2$ ):

$$T e^{ikL/2} = A e^{kL/2} + B e^{-kL/2}$$

$$k T e^{ikL/2} = k \left( A e^{kL/2} - B e^{-kL/2} \right)$$

Then we did a shit ton of calculations for too long...

Final: 6 soru cevap

Next Time: [Spin](#) and [Angular Momentum](#)