

NOTES:

→ Notion of system and state don't really change

But, in quantum mechanics, you can't directly observe the state itself.

- Schrödinger's Eqn. is NOT deterministic.

Remember Schrödinger's Equation:

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi$$

(linear equation → leads to entanglement)  
(non-relativistic)

if  $\psi(\vec{r}, 0)$  is known,

one knows  $\psi(\vec{r}, t)$

(Penrose U process)

"diffusion equation"  
 $\frac{\partial P}{\partial t} = D \nabla^2 P$

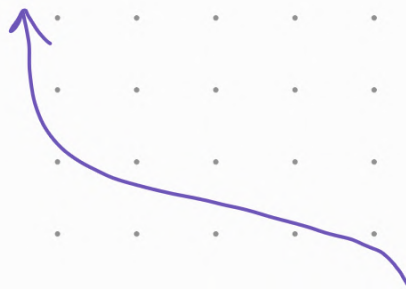
↳ Looks similar mathematically, but they aren't related or smth.

•  $\psi$  is called the state of the system, generally complex function. **AND ISN'T DIRECTLY OBSERVABLE.**

- An observable operator (differential) acting on  $\psi$ .

↳ Must have "real" eigenvalues. ↳ must have an associated eigenvector.  
(hermitian operators)

⊗ Only one of those eigenvalues can be measured in one measurement.



What happened after the measurement?

→ State changes somehow.

### Collapse of The Wave Function (Penrose Rprocess)

- ⊗ The state right after this measurement is the corresponding eigenvector.

(correlation, Bell experiments)

- The outcome of many repetitions of a measurement will give a distribution of values.

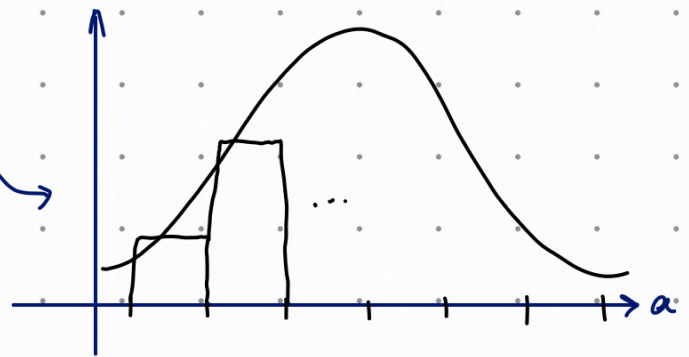
observable is  $A$

$\langle A \rangle$  expectation value.

$\langle A^2 \rangle$  ...

$\langle A^3 \rangle$  ...

$\langle A^n \rangle$  ...



Calculating  $\langle A \rangle$ :

$$\langle A \rangle_{\psi} \equiv \int d^3 \vec{r} \psi^*(\vec{r}, t) (A \psi(\vec{r}, t))$$

- ⊗ Time is not an observable! You measure time with motion. We assume it's observable externally.

- Operators you can measure simultaneously:

→ Simultaneously diagonalizable

- Commuting matrices (Common eigenvalues)

Complete Set of commuting operators.

- Those who don't commute, are bound on Heisenberg type uncertainty.

If  $V$  is not a function of  $t$ .

- one can ask for stationary states of definite energy.

$$\Psi(\vec{r}, t) = e^{-i\frac{Et}{\hbar}} \Psi_E(\vec{r}) \longrightarrow \left( -\frac{\hbar^2}{2m} \nabla^2 + V \right) \Psi_E = E \Psi_E$$

$$\hookrightarrow \Psi(\vec{r}, t) = \sum_E C_E e^{-i\frac{Et}{\hbar}} \Psi_E(\vec{r})$$

expansion  
coefficients

• Take the norm:

$$|\Psi_E|^2 = |\Psi(\vec{r})|^2$$

### Notes!

- Hermitian Operators: • Eigenvectors are orthogonal (linearly independent).

- You can have both discrete and continuous eigenvalues  
( $E$  in Bohr's model when bound  $\rightarrow$  discrete)

### Particle in a Box.

(infinite box, Energies are discrete, particle stays in the box.)

$$|\Psi(x, t)|^2 = 0 \text{ whenever particle is out } x \leq 0 \text{ and } x \geq L$$

$$\Psi(x, t) = 0$$

for  $x \leq 0$  and  $x \geq L$ .

$$\Psi(x=0, t) = \Psi(x=L, t) = 0$$

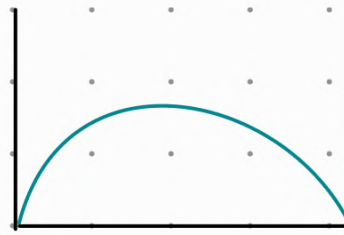
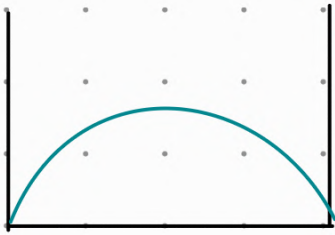
$$-\frac{\hbar^2}{2m} \frac{d^2 \Psi_E(x)}{dx^2} = E \Psi_E(x) \longrightarrow \Psi_n = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$
$$E_n = \frac{n^2 \pi^2 \hbar^2}{2m L^2}$$

$$\Psi(x, t) = \sum_{n=1}^{\infty} C_n e^{-i\frac{E_n t}{\hbar}} \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

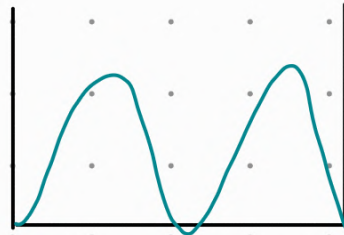
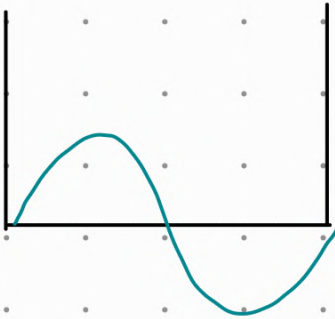
are fixed by  $\Psi(x, 0)$



$n=1$

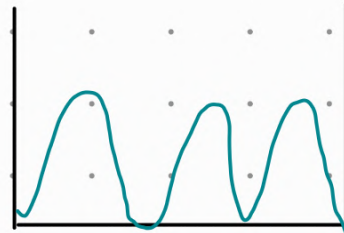
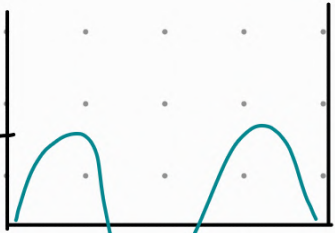


$n=2$



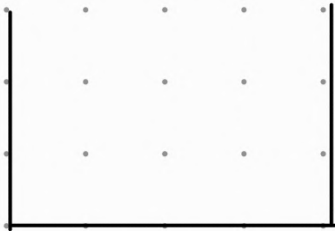
$n=3$

$\sqrt{\frac{2}{L}}$

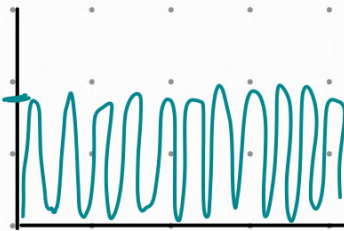


$-\sqrt{\frac{2}{L}}$

$n \rightarrow \infty$



$\frac{2}{L}$



It oscillates so much, you expect it to converge to



$\frac{1}{L}$

