

Chapter 2 - Plane Waves + Index of Refraction

plane wave:

$$\vec{E}(\vec{r}, t) = \vec{E}_0 \cos(\vec{k} \cdot \vec{r} - \omega t + \phi)$$

"wave vector" $\vec{k} = k \hat{u} = \frac{2\pi}{\lambda_{\text{vac}}} \hat{u}$
 "k"
 direction of propagation

angular frequency $\rightarrow \omega = \frac{2\pi c}{\lambda_{\text{vac}}} = 2\pi v$

$k + \omega$ are related just like $\lambda + v$

$$c = \frac{\lambda_{\text{vac}}}{T} = \lambda_{\text{vac}} v = \frac{\omega}{k}$$

\curvearrowleft
 this relationship
 is known as the
 "dispersion relation"

"kappa" \downarrow wavenumber
 $\frac{1}{\lambda_{\text{vac}}} = K = [\text{cm}^{-1}]$

$c = v_{\text{vac}}$ \leftarrow velocity of light in vacuum

v \curvearrowleft "nu"
 \curvearrowright frequency

similarly for magnetic field:

$$\vec{B}(\vec{r}, t) = \vec{B}_0 \cos(\vec{k} \cdot \vec{r} - \omega t + \phi)$$

$$\underbrace{\vec{B}_0 = \frac{\vec{k} \times \vec{E}_0}{\omega}}, \text{ so } \vec{B}_0 \text{ is not independent, but is determined by other parameters.}$$

$$\vec{B}_0 \perp \vec{k} \perp \vec{E}$$

also think about magnitude

$$B_0 = \frac{k E_0}{\omega} = \frac{E_0}{c} \leftarrow \begin{array}{l} \text{since this is so large} \\ \text{we will focus on the } \vec{E} \text{ field.} \end{array}$$

complex plane waves:

$$\vec{E}(\vec{r}, t) = \operatorname{Re} \left\{ \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)} \right\}$$

$$\vec{E}_0 = \vec{E}_0 e^{i\phi} \leftarrow \text{phase shift}$$

$$\vec{E}(\vec{r}, t) = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

Summarize some facts that we know now:

$\lambda \rightarrow$ wavelength

$T \rightarrow$ period

$\gamma \rightarrow$ frequency ($\frac{1}{T}$)

$k \rightarrow$ wave vector ($\frac{2\pi}{\lambda}$)

$\omega \rightarrow$ angular frequency ($\frac{2\pi}{T}$)

$$V = \frac{\omega}{k}$$

$$V = \frac{\lambda}{T} = \lambda \cdot \gamma \leftarrow \text{any wave}$$

$$c = \lambda_{\text{vac}} \cdot \gamma$$

$$\epsilon_0 \mu_0 = \frac{1}{c^2} \Rightarrow c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

permeability
of
free space $\rightarrow \epsilon_0 = 8.85 \cdot 10^{-12} \frac{C^2}{Nm^2}$

permittivity
of
free space $\rightarrow \mu_0 = 4\pi \cdot 10^{-7} \frac{T}{Am^2}$

Speed of light in matter

\hookrightarrow light slows down in materials

$$\frac{c}{v} = n \leftarrow \text{index of refraction}$$

$n=1 \leftarrow$ vacuum

$n=1.0003 \leftarrow$ air

$n=1.33 \leftarrow$ water

$n=1.5 \leftarrow$ glass

$$\nabla^2 \vec{E} = \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\nabla^2 \vec{E} = \frac{1}{V^2} \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\nabla^2 \vec{E} = \frac{k^2}{\omega^2} \frac{\vec{E}}{\partial t^2}$$

Q1: what is light? ✓

Q2: where does light come from? ✓

Q3: can the speed of light change? see Q6

Q4: what affects the brightness of light?

Q5: how do human perceive light amount + color

Q6: why does light slow down in materials

Q3

Q7: what is reflection

$\lambda \cdot v = V = \frac{c}{n}$

constant as n increases
 \downarrow
 gets smaller as n increases

$$\lambda_{\text{vac}} \cdot v = c \Leftrightarrow c = n \cdot v$$

\uparrow
 $\lambda \cdot v$

$$\lambda_{\text{vac}} \cdot v = n \cdot \lambda \cdot v$$

\approx wavelength in material

$$\lambda = \frac{\lambda_{\text{vac}}}{n}$$

} λ will be smaller than λ_{vac} since $n > 1$

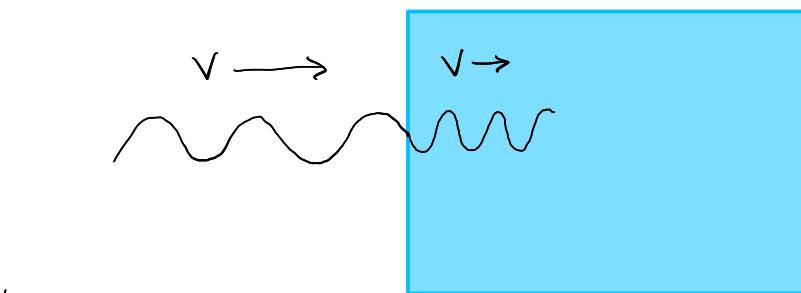
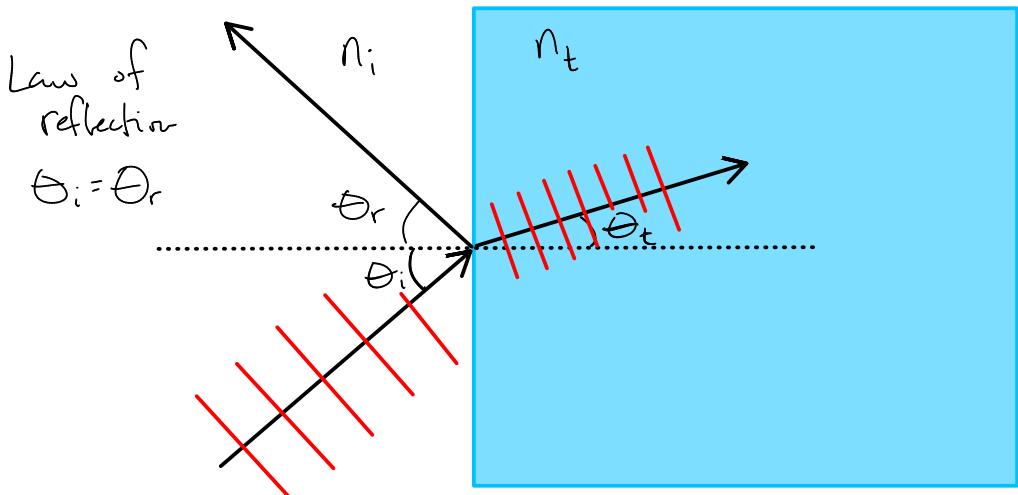


Table 23.1 Indices of Refraction for $\lambda = 589.3 \text{ nm}$ in Vacuum (at 20°C Unless Otherwise Noted)

Material	Index
Solids	
Ice (at 0°C)	1.309
Fluorite	1.434
Fused quartz	1.458
Polystyrene	1.49
Lucite	1.5
Plexiglas	1.51
Crown glass	1.517
Plate glass	1.523
Sodium chloride	1.544
Light flint glass	1.58
Dense flint glass	1.655
Sapphire	1.77
Zircon	1.923
Diamond	2.419
Titanium dioxide	2.9
Gallium phosphide	3.5
Liquids	
Water	1.333
Acetone	1.36
Ethyl alcohol	1.361
Carbon tetrachloride	1.461
Glycerin	1.473
Sugar solution (80%)	1.49
Benzene	1.501
Carbon disulfide	1.628
Methylene iodide	1.74

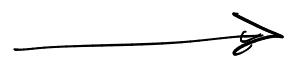


Snell's Law

$$n_i \sin \theta_i = n_t \sin \theta_t$$

static case

$$\vec{E}(\vec{r}) = \frac{kq_0}{r^2}$$



oscillating charge

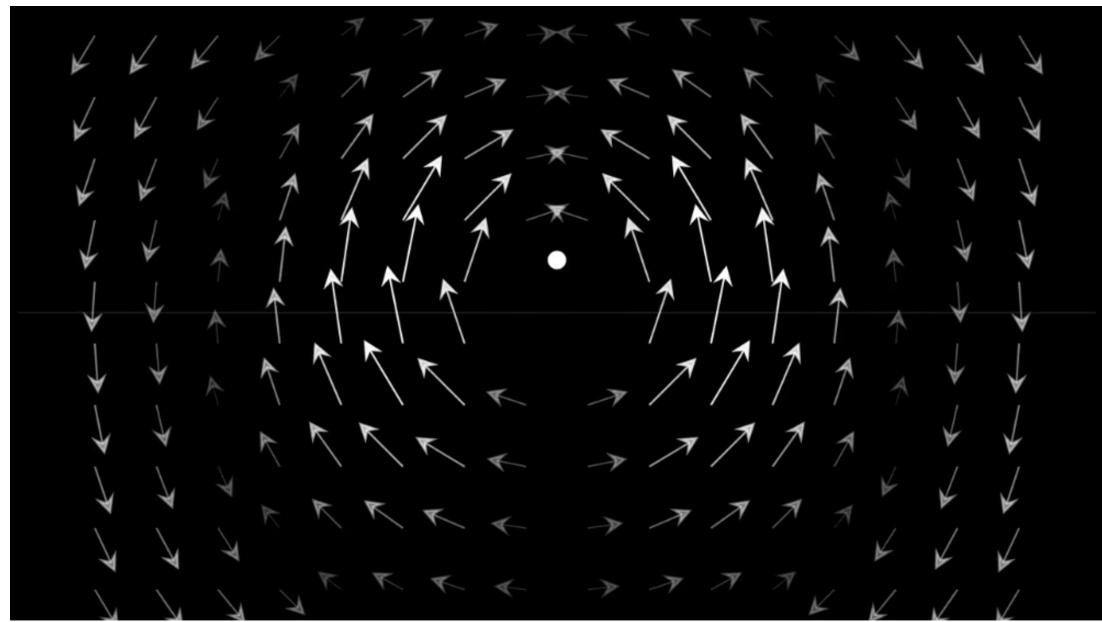
$$\vec{E}(\vec{r}, t) = \vec{E}_{\text{staticish}} + \vec{E}_{\text{rad}}$$

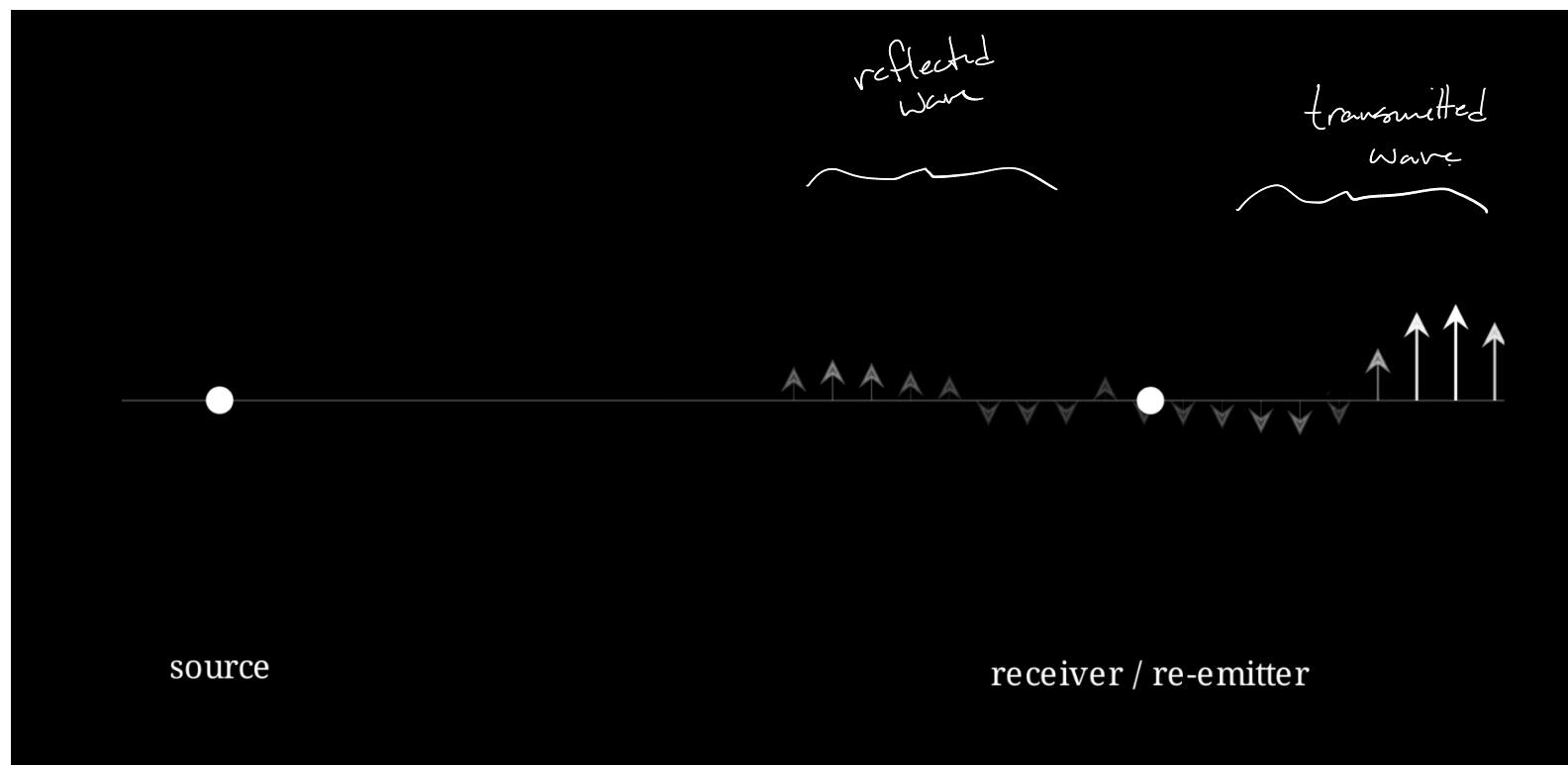
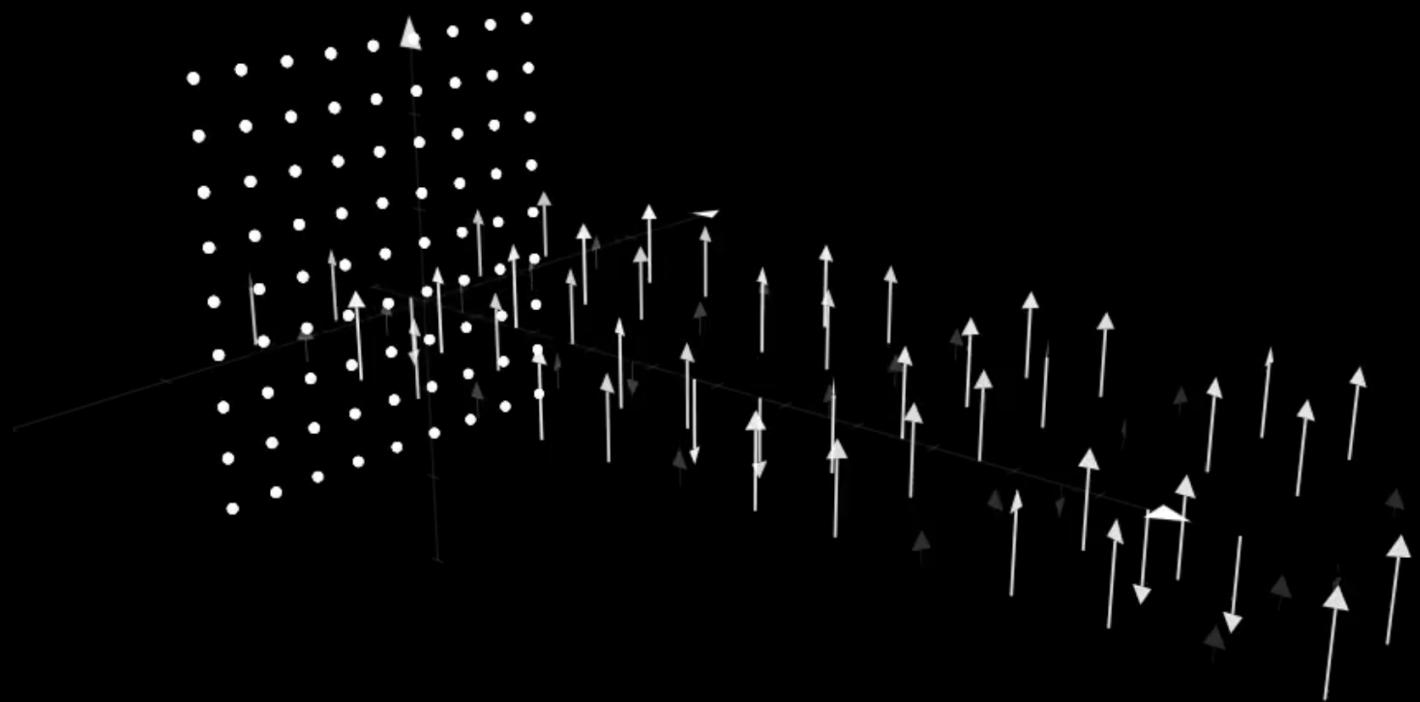
\sim

wave behavior

compare
to
 $\vec{r} - vt$

$$\vec{E}_{\text{rad}} \propto \frac{\vec{a}_\perp(t - r/c)}{r}$$





Dielectric Media

↳ insulator (like glass, water, air)

Assumptions: ① isotropic \rightarrow same in all directions

② homogeneous \rightarrow same over distance

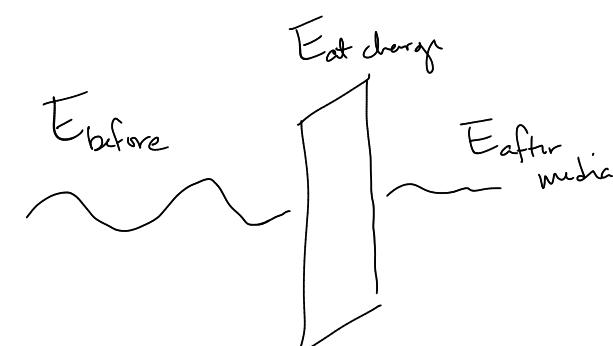
③ non-conducting (insulator/dielectric) $\vec{J}_{\text{free}} = 0$

④ very thin material
(air or some other
thin gas)

} we will
relax this
later on

Light in a thin gas \rightarrow Feynman, vol 1, Ch 31

Light in matter = incident field + field reradiated
by driven charges



$$\vec{E}_{\text{after}} = \underbrace{\vec{E}_s}_{\text{plate}} + \sum_{\text{all other charges}} \vec{E}_{\text{each charge}}$$

$$\vec{E}_s = \vec{E}_o \cos(kx - \omega t)$$

$$\vec{E}_s = \vec{E}_o \cos\left(k\left(x - \frac{\omega}{k}t\right)\right)$$

$$\vec{E} = \vec{E}_o \cos\left(\omega\left(\frac{k}{\omega}x - t\right)\right)$$

$$\vec{E} = \vec{E}_o \cos\left(-\omega\left(t - \frac{k}{\omega}x\right)\right)$$

$$\vec{E} = \vec{E}_0 \cos\left(\omega(t - \frac{x}{c})\right)$$

$$\vec{E} = \vec{E}_0 e^{i\omega(t - \frac{x}{c})}$$

So now what about the field from other charges?

$$\text{plate thickness} = \Delta x$$

$$t_{\text{vac}} = \frac{\Delta x}{c}$$

$$t_{\text{plate}} = \frac{n \Delta x}{c}$$

$$\Delta t = t_{\text{plate}} - t_{\text{vac}}$$

$$= (n-1) \frac{\Delta x}{c}$$

$$\begin{aligned} E_{\text{after}} &= E_0 e^{i\omega(t - \Delta t - \frac{x}{c})} \\ \text{plate} &= E_0 e^{i\omega(t - (n-1)\frac{\Delta x}{c} - \frac{x}{c})} \\ &= E_0 e^{-i\omega(n-1)\frac{\Delta x}{c}} E_0 e^{i\omega(t - \frac{x}{c})} \\ &\quad \underbrace{e^{-i\omega(n-1)\frac{\Delta x}{c}}}_{\text{w/ a phase shift}} \underbrace{E_0 e^{i\omega(t - \frac{x}{c})}}_{\substack{\text{just like the} \\ \text{original wave} \\ \text{BUT w/ a}}} \end{aligned}$$

$$e^{-i\omega(n-1)\frac{\Delta x}{c}}, \Delta x \text{ is very small}$$

$$\left\{ e^S \approx 1 + S, \text{ for small } S \right.$$

$$e^{-i\omega(n-1)\frac{\Delta x}{c}} \approx 1 - i\omega(n-1)\frac{\Delta x}{c}$$

$$E_{\text{after plate}} = E_0 e^{i\omega(t-\frac{x}{c})} - \frac{i\omega(n-1)\Delta x}{c} \cdot E_0 e^{i\omega(t-\frac{x}{c})}$$

$$E_s \qquad \qquad \qquad E_a$$



$$ma = \sum F = -kz$$

$$m \frac{d^2 z}{dt^2} = -kz$$

$$\frac{d^2 z}{dt^2} = -\frac{k}{m} z$$

~~$$f(A\omega^2 \cos(\omega t)) = f\frac{k}{m} \cdot A \cos(\omega t)$$~~

$$\omega^2 = \frac{k}{m} \Rightarrow \omega = \sqrt{\frac{k}{m}}$$

one of our constants

$$z = A \cos(\omega t)$$

$$\frac{dz}{dt} = -A\omega \sin(\omega t)$$

$$\frac{d^2 z}{dt^2} = -A\omega^2 \cos(\omega t)$$

$$z = A \cos(\sqrt{\frac{k}{m}} t)$$

$$z = A e^{\omega t}$$

$$\frac{dz}{dt} = A\omega e^{\omega t}$$

$$\frac{d^2 z}{dt^2} = A\omega^2 e^{\omega t}$$

$$A\omega^2 e^{\omega t} = -\frac{k}{m} \cdot A e^{\omega t}$$

$$\omega^2 = -\frac{k}{m}$$

Given by initial conditions
at $t=0, z=1 \Rightarrow A=1$

$$z=1\cdot \cos\left(\sqrt{\frac{k}{m}}\cdot t\right)$$