

Thermal Equilibrium - when objects have been in contact  
and macroscopic coordinates have stopped changing

↳ involves an exchange of energy between two objects, or an object + its surroundings

↳ volume (constant pressure) → mercury/alcohol

- pressure

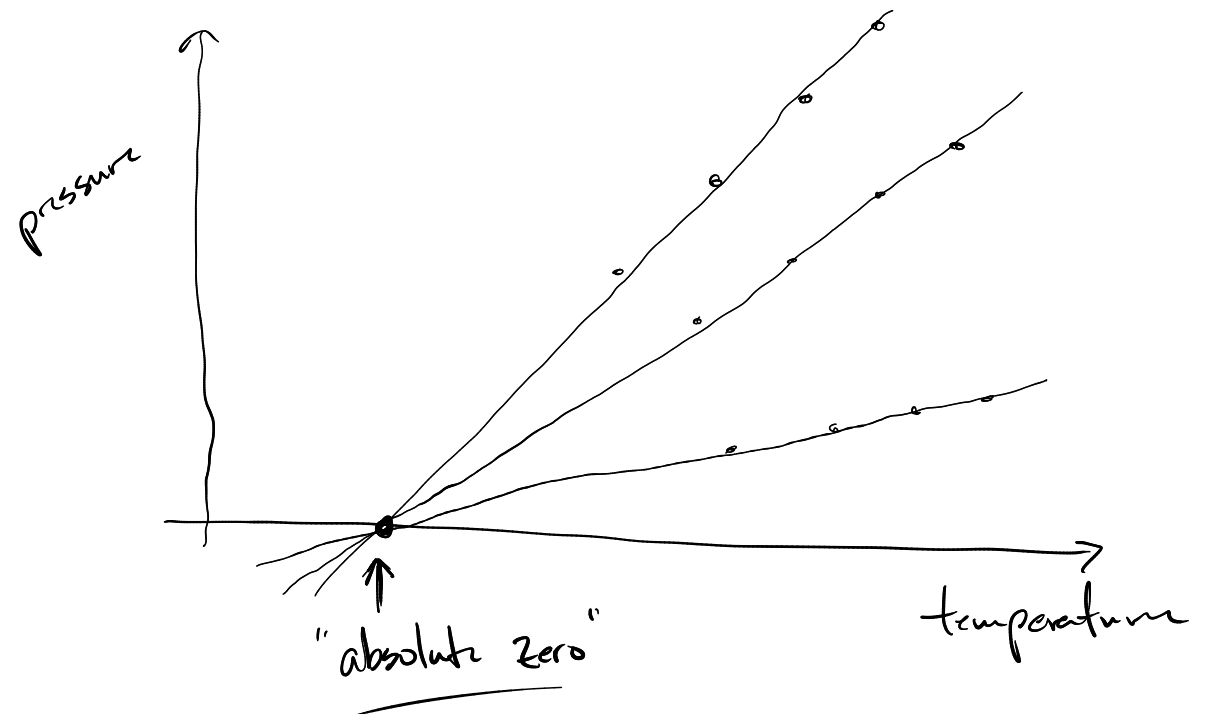
- electrical resistance

- radiation / thermal emf

$$1) \quad C_1 = 0^\circ\text{C} \quad C_2 = 100^\circ\text{C}$$

$$F_1 = 32^\circ\text{F} \quad F_2 = 212^\circ\text{F}$$

$$C = mF + b$$



$$0\text{K} = -273.15^\circ\text{C}$$

1.2 Ideal Gas Law  $\rightarrow$  Equation of State  $\rightarrow$  relates all of the state variables together

$$pV = N k_B T$$

$p = \frac{F}{\text{Area}}$   
 $[Pa] = \left[ \frac{N}{m^2} \right]$

$N$   $\rightarrow$  number of particles  
 $k_B$  Boltzmann's constant  
 $k_B = 1.38 \cdot 10^{-23} \frac{J}{K}$

$$pV = nRT$$

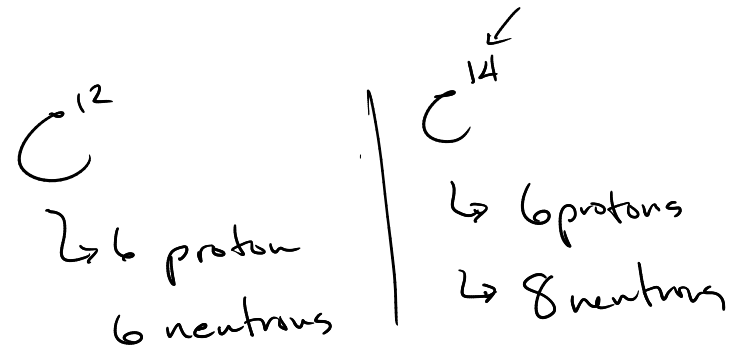
$n$   $\rightarrow$  number of moles  
universal gas constant  
 $R = 8.31 \frac{J}{mol \cdot K} ?$

$$N k_B = nR$$
$$6.022 \cdot 10^{23} \cdot 1.38 \cdot 10^{-23} = R$$
$$8.31 \frac{J}{mol K} = R$$

Number, moles, molar, density

$$1 \text{ mole of things} = 6.022 \cdot 10^{23} \text{ things}$$

particle  
atom  
molecule



1 mole is 1 gram of protons + neutrons

Ex: mass of one proton in kg?

mass of one proton  $\times$  number of proton = mass of the collection

$$m \times N = M$$

$$m = \frac{1g}{N_A} = \frac{1}{6.022 \cdot 10^{23}} = 1.7 \cdot 10^{-24} \text{ grams} = \underline{1.7 \cdot 10^{-27} \text{ kg}}$$

What about  $N_2$ ?

$$1 \text{ mole of } N_2 = 2 \left( 14 \frac{g}{\text{mol}} \right) = 28 \frac{g}{\text{mol}}$$

Volume of 1 mole of air at room temp and atmospheric pressure?

$$\hookrightarrow 1 \text{ atm} = \underline{1.013 \cdot 10^5 \text{ Pa}}$$

$$V = \frac{nRT}{P} = \frac{1 \text{ mol} \cdot 8.31 \frac{J}{K \cdot \text{mol}} \cdot 300K}{10^5 \text{ Pa}} = 0.0249 \text{ m}^3$$

$$V_{\text{cube}} = \Delta^3$$

$$\Delta = 0.292 \text{ m} \sim 30 \text{ cm} \sim 1 \text{ ft}$$

1.17]  $PV = nRT \left( 1 + \frac{B(T)}{V/n} \right)$

a)  $PV = nRT \left( 1 + \frac{n}{V} \cdot B(T) \right)$

at atmospheric pressure  $P = 10^5 \text{ Pa}$

→ solve for  $\frac{n}{V}$

$$P = \frac{n}{V} RT \left( 1 + \frac{n}{V} B(T) \right)$$

$$0 = \frac{n}{V} RT + \left( \frac{n}{V} \right)^2 \cdot RT \cdot B(T) - P$$

T	B(T)	$\frac{n}{V}$	$\frac{n}{V} \cdot B(T)$
100	-16.0		
200	-3.5		
300	-4.2		
400	9.0		
500	16.9		
600	21.3		

$$b) \quad PV = nRT \left( 1 + \frac{B(T)}{V/n} \right) \quad \text{if}$$

$$c) \rightarrow PV = nRT \left( 1 + \frac{B(T)}{V/n} + \frac{C(T)}{(V/n)^2} \right) \quad \left. \vphantom{PV = nRT} \right\} \text{ want: } B + C \text{ in terms of } a + b$$

$$\left( P + \frac{an^2}{V^2} \right) \left( V - nb \frac{V}{V} \right) = nRT$$

$$\left( P + \frac{an^2}{V^2} \right) V \left( 1 - \frac{nb}{V} \right) = nRT$$

$$\left( P + \frac{an^2}{V^2} \right) V = nRT \left( 1 - \frac{nb}{V} \right)^{-1}$$

$$PV + \frac{an^2}{V} = nRT \left( 1 - \frac{nb}{V} \right)^{-1}$$

$$PV = nRT \left( 1 - \frac{nb}{V} \right)^{-1} - \frac{an^2}{V}$$

use our approx.

$$(1+x)^p \approx 1 + px + \frac{1}{2}p(p-1)x^2 \quad px \ll 1$$

$$\left( 1 + \left( -\frac{nb}{V} \right) \right)^{-1} \approx 1 + (-1) \left( -\frac{nb}{V} \right) + \frac{1}{2}(-1)(-1-1) \left( -\frac{nb}{V} \right)^2$$

$$\approx 1 + \frac{nb}{V} + \left( \frac{nb}{V} \right)^2$$

$$PV = nRT \left( 1 + \frac{nb}{V} + \left( \frac{nb}{V} \right)^2 \right) - \frac{an^2}{V} \begin{matrix} nRT \\ nRT \end{matrix} \longleftrightarrow PV = nRT \left( 1 + \frac{n}{V} \cdot B(T) + \left( \frac{n}{V} \right)^2 \cdot C(T) \right)$$

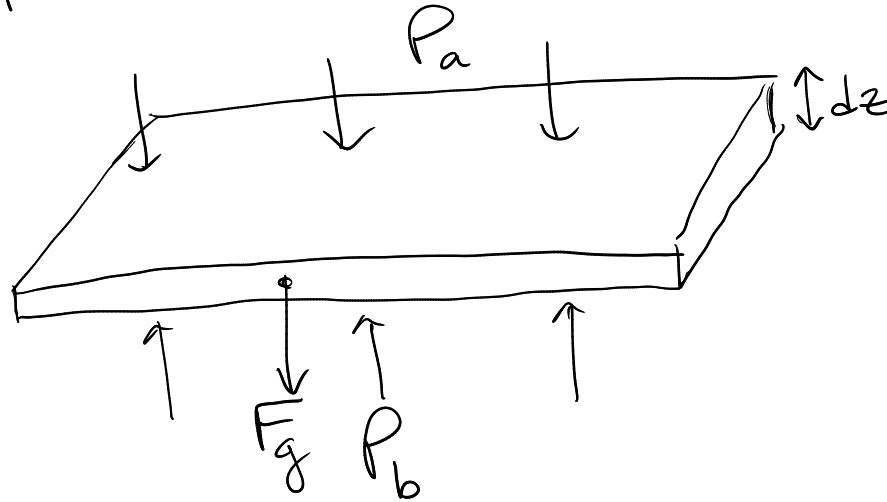
$$PV = nRT \left( 1 + \frac{nb}{V} + \left( \frac{nb}{V} \right)^2 - \frac{an^2}{V \cdot nRT} \right) = nRT \left[ 1 + \underbrace{\frac{n}{V} \left( b - \frac{a}{RT} \right)}_{B(T)} + \underbrace{\left( \frac{n}{V} \right)^2 \cdot b^2}_{C(T)} \right]$$

d) → plot data from that table ←

→ plot  $B(T) = b - \frac{a}{RT}$  choose  $b, a$

1.16 |  $\rho = \frac{M}{V}$

↑  
"rho"  
volumetric  
mass  
density



$$P = \frac{F}{A}$$

$$+ P_b \cdot A - P_a \cdot A - M \cdot g = 0$$

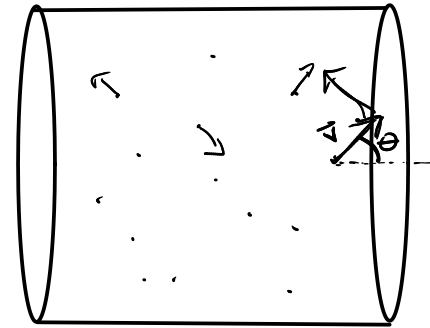
↓

$$\frac{dP}{dz} = \underbrace{\quad}_{\rho}$$

## 1.2b | Kinetic Theory and Equipartition of Energy

pressure  $\longleftrightarrow$  kinetic energy  $\longleftrightarrow$  temperature

$$\text{pressure}_{\text{collision}} = \frac{F}{A} = \frac{\Delta p}{A \Delta t} \quad \leftarrow \underbrace{2mv \cos \theta}$$



$$d(\text{pressure}) = \frac{2mv \cos \theta}{dA \cdot dt} \cdot \left. \begin{array}{l} \text{number of particles hitting area } dA \\ \text{w/ velocity } v \text{ in } dt \text{ amount of time} \end{array} \right\} \begin{array}{l} \text{integrate over all velocities} \\ \text{and over } \theta \end{array}$$

number of atoms  
traveling in a particular  
direction w/ a particular  
speed

• fraction of them  
that are within  
striking distance  
of the surface  $dA$



# Probability

$$P(x) = \frac{\text{desired outcomes}}{\text{total outcomes}}$$

$N(x) \leftarrow$  number of desired outcomes

$\nearrow$   
total outcomes

$$N = \sum_{x=0}^{\infty} N(x)$$

$\nwarrow$   
rearrange

$$P(x) = \frac{N(x)}{N}$$

$$N(x) = P(x) \cdot N$$

normalized probabilities

$$\sum_{x=0}^{\infty} P(x) = 1$$

$$\langle x \rangle = \sum_{x=0}^{\infty} x \cdot P(x)$$

$\uparrow$   
• average value  
• expectation value

$\underbrace{\hspace{10em}}_{\rightarrow \text{weighted average by probability}}$

$$\langle x^2 \rangle = \sum_{x=0}^{\infty} x^2 \cdot P(x)$$

$$\langle f(x) \rangle = \sum_{x=0}^{\infty} f(x) \cdot P(x)$$

Example (Blundell 3.3)

a.  $P(x) = \frac{e^{-m} m^x}{x!}$ ; show

$$\sum_{x=0}^{\infty} P(x) = 1$$

Looked up Taylor Series:

$$e^x = \sum_{n=0}^{\infty} \frac{x^n}{n!}$$

$$\sum_{x=0}^{\infty} \frac{e^{-m} m^x}{x!} \stackrel{?}{=} 1$$

$$e^{-m} \cdot \sum_{x=0}^{\infty} \frac{m^x}{x!} \stackrel{?}{=} 1$$

$$\sum_{x=0}^{\infty} \frac{m^x}{x!} \stackrel{?}{=} e^m \checkmark \checkmark$$

b. Show  $\langle x \rangle = \sum_{x=0}^{\infty} x \cdot P(x) \stackrel{?}{=} m$

$$= \sum_{x=0}^{\infty} x \cdot \frac{e^{-m} m^x}{x!} \stackrel{?}{=} m$$

$$= e^{-m} \sum_{x=0}^{\infty} \frac{x \cdot m^x}{x!} \stackrel{?}{=} m$$

$$\sum_{x=0}^{\infty} \frac{x \cdot m^x}{x!} \stackrel{?}{=} m e^m$$

$$\underbrace{\quad}_{x \cdot (x-1) \cdot (x-2) \cdot (x-3) \cdots 1}$$

$$\cancel{\frac{0 \cdot m}{0!}} + \frac{1 \cdot m^1}{1!} + \frac{2 m^2}{2!} + \dots$$

$$\sum_{x=1}^{\infty} \frac{x m^x}{x!} = \sum_{x=1}^{\infty} \frac{m^x}{(x-1)!} = \sum_{x=1}^{\infty} \frac{m^{x+1-1}}{(x-1)!} = \sum_{x=1}^{\infty} \frac{m \cdot m^{x-1}}{(x-1)!}$$

$$x \cdot (x-1) \cdot (x-2) \cdots$$

$$\begin{aligned} x' &= x-1 \\ \hookrightarrow x &= x'+1 \end{aligned}$$

$$m \cdot \sum_{x'=0}^{\infty} \frac{m^{x'}}{x'!}$$

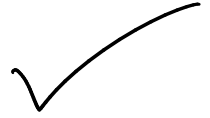
drop the primes!

C.

$$M \cdot \sum_{x=0}^{\infty} \frac{m^x}{x!}$$

$\underbrace{\hspace{10em}}_{=e^m}$

$$M e^{m?} = m e^m$$



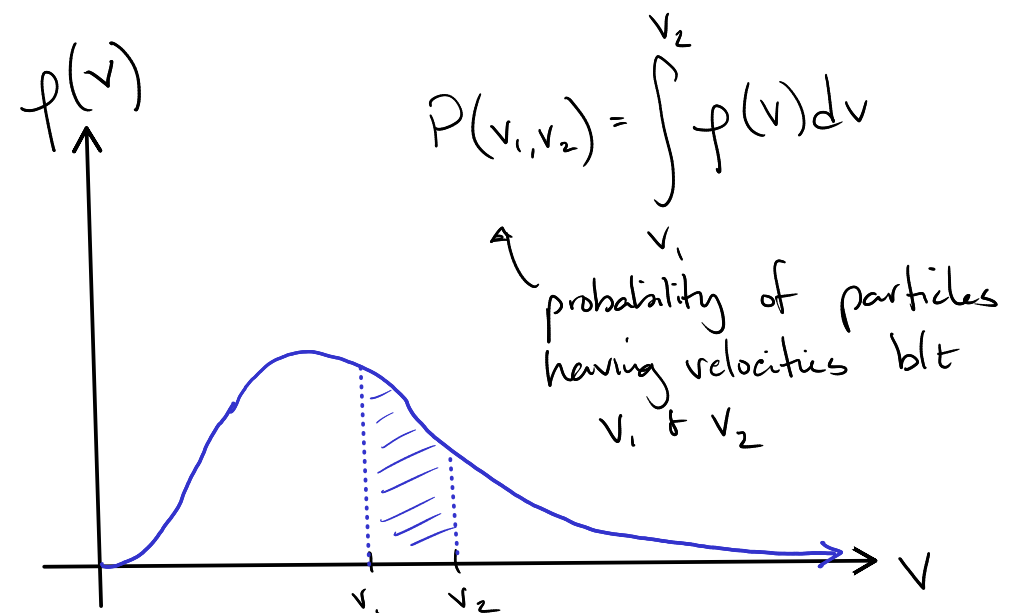
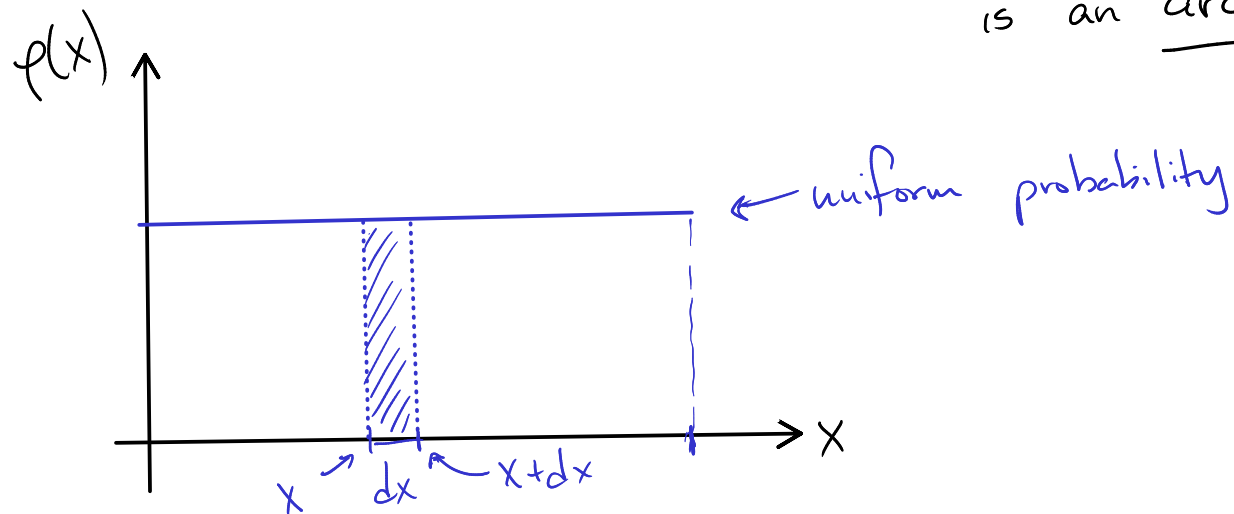
$$\langle x \rangle = m$$

The continuous variable case follows by analogy but w/ some clarifications

$$\left\{ \begin{array}{l} \text{probability of choosing} \\ \text{a value at random} \\ \text{between } x \text{ and } x+dx \end{array} \right\} = p(x) \cdot dx$$

$\rightarrow$  distribution function

probability itself  
is an area.



Normalized

$$\int_{-\infty}^{\infty} p(v) dv = 1$$

Since we will be dealing w/ such large numbers

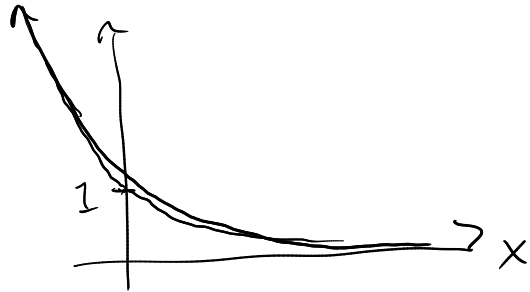
$$\left\{ \begin{array}{l} \text{probability that a} \\ \text{particle has a} \\ \text{velocity between} \\ v \text{ and } v+dv \end{array} \right\} = \left\{ \begin{array}{l} \text{fraction of particles} \\ \text{having a velocity} \\ \text{between } v \text{ + } v+dv \end{array} \right\}$$

Example: (Blundell 3.4)

$$p(x)dx = A e^{-x/\lambda} dx$$

a) Find A so that

$$\int_0^{\infty} p(x)dx = 1$$



Also mean

$$\langle v \rangle = \int_{-\infty}^{\infty} v \cdot p(v) dv$$

or

$$\langle v^2 \rangle = \int_{-\infty}^{\infty} v^2 \cdot p(v) dv$$

or

$$\langle f(v) \rangle = \int_{-\infty}^{\infty} f(v) \cdot p(v) dv$$