

Chapter 2 - 2nd Law of Thermodynamics

↳ heat spontaneously flows from high temp to low temp

Einstein Solid

Ideal Gas

↳ Entropy

→ Anacanda python
install language
Monty Python

Combinatorics

one coin
 $P(\text{heads}) = \frac{1}{2}$

multiple coins

$$P(n) = \frac{n}{N}$$

$$P(3 \text{ heads}) = \frac{\Omega(3)}{\Omega(\text{all})}$$
$$= \frac{\Omega(3)}{\sum_{n=0}^n \Omega(n)}$$

← microstate in the macrostate of 3
← all of the microstates

5 coins

microstates { H H T T H — 3 H } macrostates
 T H H H H — 4 H }

how many microstates are in a macrostate?
↳ multiplicity

$$\Omega(n) = \frac{5!}{n!(5-n)!}$$

↑
number of heads

$$\begin{aligned} \Omega(0) &= 1 \\ \Omega(1) &= 5 \\ \Omega(2) &= 10 \\ \Omega(3) &= 10 \\ \Omega(4) &= 5 \\ \Omega(5) &= 1 \end{aligned}$$


$$\Omega(1) = \frac{5!}{1!(5-1)!} = \frac{5!}{1!4!}$$
$$= \frac{5 \cdot 4 \cdot 3 \cdot 2 \cdot 1}{1 \cdot 4 \cdot 3 \cdot 2 \cdot 1}$$
$$= 5$$
$$\Omega(2) = \frac{5!}{2!3!} = 10$$

$$\Omega(N, n) = \frac{N!}{n!(N-n)!} \leftarrow \text{Notation: } \binom{N}{n}$$

\uparrow
 # of coins

10 atoms each w/ 0 or 1 packets of energy (energy unit)


How many possible ways are there to distribute 4 energy units


 \leftarrow microstate

4 energy packets \leftarrow macrostate

$$\Omega(10, 4) = \frac{10!}{4!6!} = 210$$

What if an atom can have more than one energy packet at a time?

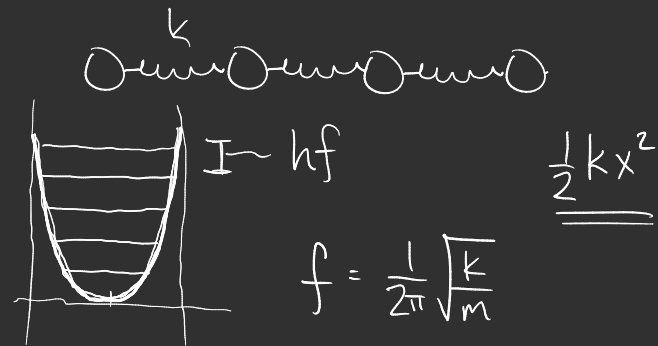

 \leftarrow microstates

4 energy packets \leftarrow macrostate

$$\Omega(N, q) = \frac{(q + N - 1)!}{q!(N-1)!} \quad \binom{q+N-1}{q}$$

\uparrow # of atoms \uparrow # of energy packets
 $q = \text{macrostate}$

This model of a collection of atoms w/ equal size energy quanta distributed among them is the Einstein Solid.



→ Debye Model

• Large Number → addition of small numbers is not important

$$10^{23} + 23 = 10^{23}$$

• Very Large Number

$$10^{10^{23}} \times 10^{23} = 10^{10^{23} + 23} \approx 10^{10^{23}}$$

Stirling's Approximation

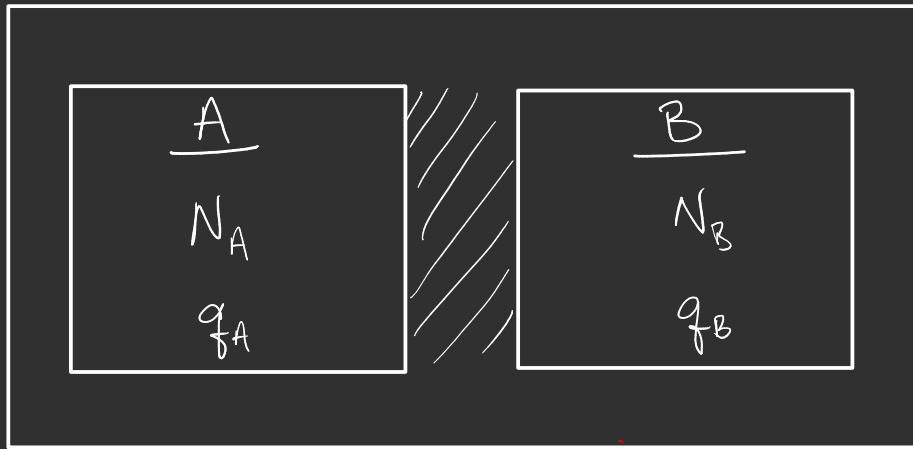
$$N! \approx N^N e^{-N} \sqrt{2\pi N} = \frac{N^N}{e^N} \underbrace{\sqrt{2\pi N}}_{\text{very large numbers we leave this off}}$$

↓

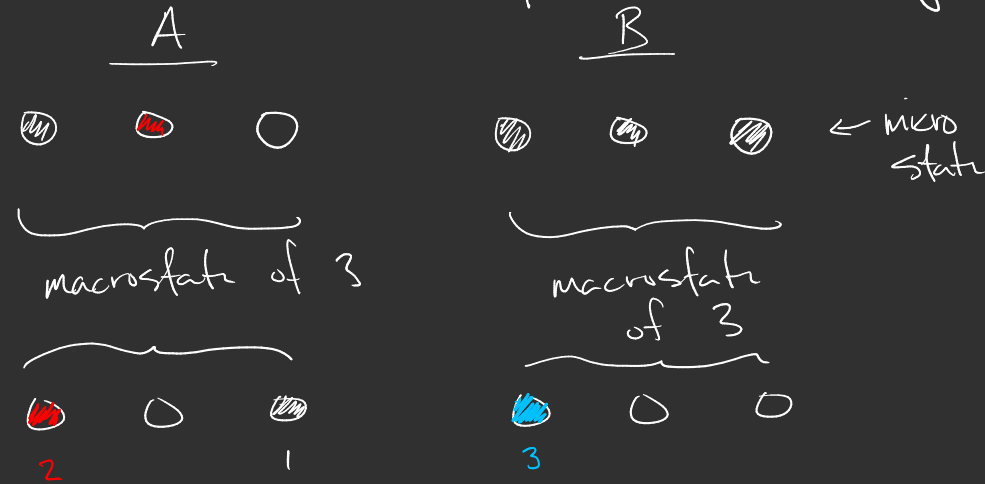
$$\ln N! \approx N \ln N - N$$

very large numbers
we leave this off

Two Systems



Ex. Each system has 3 particles and 6 total quanta of energy

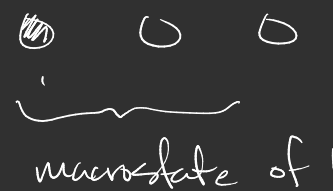


q_A	mult_A	q_B	mult_B	mult_total
0	1	6	28	28
1	3	5	21	63
2	6	4	15	90
3	10	3	10	100
4	15	2	6	90
5	21	1	3	63
6	28	0	1	28

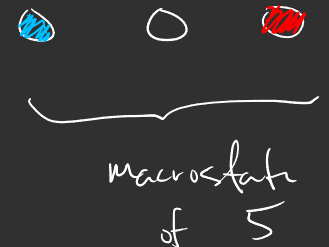
macrostate w/ the most microstates

total number of microstates

(462)



$$\Omega(3,1) = 3$$



$$\Omega(3,5)$$

$$= \frac{7!}{5!2!} = 21$$

Fundamental Assumption of Stat Mech:

all microstates are possible and equally probable

But that does not mean that every microstate will occur.

Not all macrostates are equally probable.

$$P_A(q) = \frac{\Omega_{\text{total}}(q_A)}{\sum_{q_A} \Omega_{\text{total}}(q_A)}$$

We could find the total # of microstates:

$$\Omega(6, 6) = \frac{(6+6-1)!}{6!(6-1)!} = 462$$

$\uparrow \quad \uparrow$
 $N \quad q$

So let's apply Stirling's Approx to Multiplicity $[\ln N! = N \ln N - N]$

$$\Omega(N, q) = \frac{(q+N-1)!}{q!(N-1)!} \approx \frac{(q+N)!}{q!N!}$$

very large number

$$\ln \Omega = \ln(q+N)! - \ln q! - \ln N! \rightarrow N \ln N - N$$

$\underbrace{\ln(q+N)!}_{\ln(q+N)! = (q+N) \ln(q+N) - (q+N)}$ $\underbrace{\ln q!}_{q \ln q - q}$ $\underbrace{\ln N!}_{N \ln N - N}$

$$\ln \Omega = (q+N) \ln(q+N) - \cancel{q} - \cancel{N} - q \ln q + \cancel{q} - N \ln N + \cancel{N}$$

$\ln \Omega = (q+N) \ln(q+N) - q \ln q - N \ln N$

high temperature limit $\rightarrow q \gg N$

$$\begin{aligned}
 \ln \Omega &= (q+N) \ln(q+N) - q \ln q - N \ln N \\
 &= \underbrace{q \ln(q+N)}_{\downarrow} + \underbrace{N \ln(q+N)}_{\downarrow} - q \ln q - N \ln N \\
 &= \ln \left[q \cdot \left(1 + \frac{N}{q} \right) \right] \\
 &= \ln q + \ln \left(1 + \frac{N}{q} \right) \\
 &\quad \downarrow \qquad \qquad \downarrow \begin{array}{l} \ln(1+x) \approx x \text{ for small } x \\ \rightarrow \frac{N}{q} \end{array} \\
 &= \ln q + \frac{N}{q} \approx \ln(q+N)
 \end{aligned}$$

$$\begin{aligned}
 \ln \Omega &= \cancel{q \ln q} + N + \underbrace{N \ln q}_{\text{small}} + \frac{N^2}{q} - \cancel{q \ln q} - \underbrace{N \ln N}_{\text{small}} \\
 &= N \ln \left(\frac{q}{N} \right) + N + \frac{N^2}{q}
 \end{aligned}$$

$$\ln \Omega(q \gg N) \approx N \ln \left(\frac{q}{N} \right) + N$$

$$\Omega(q \gg N) = e^{N \ln \left(\frac{q}{N} \right) + N} = e^{N \ln \left(\frac{q}{N} \right)} \cdot e^N = \underbrace{\left(e^{\ln \left(\frac{q}{N} \right)} \right)^N} \cdot e^N$$

$$\Omega(q \gg N) = \left(\frac{q}{N} \right)^N e^N = \left| \left(\frac{eq}{N} \right)^N \right|$$

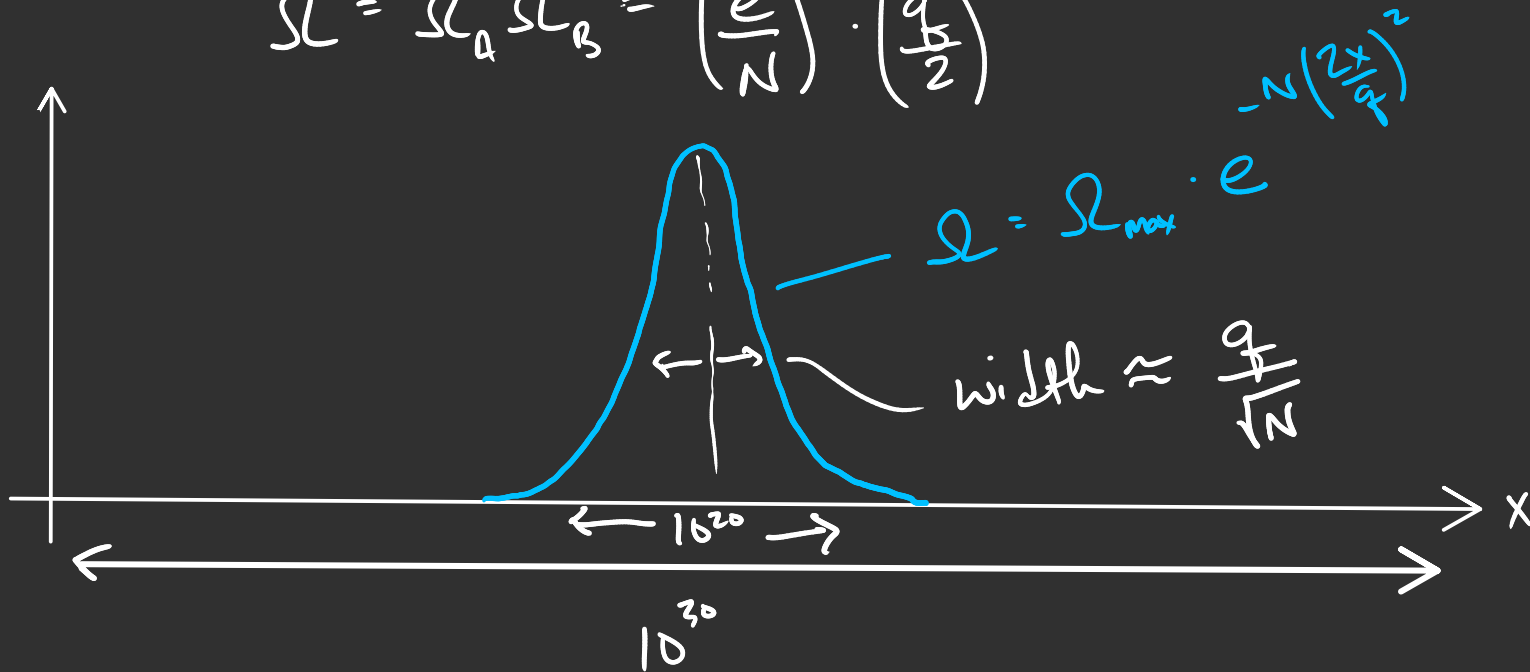
- 2 Einstein solids (high temperature limit) ($q \gg N$)

$$\Omega_A = \left(\frac{e q_A}{N_A} \right)^{N_A} \quad \Omega_B = \left(\frac{e q_B}{N_B} \right)^{N_B}$$

$$q = q_A + q_B$$

if $N_A = N_B$, multiplicity max will occur $\frac{q}{2} = q_A = q_B$

$$\Omega = \Omega_A \Omega_B = \left(\frac{e}{N} \right)^{2N} \cdot \left(\frac{q}{2} \right)^{2N}$$



$$q = 10^{30}$$

$$N = 10^{20}$$

$$\frac{q}{\sqrt{N}} = 10^{20}$$

Entropy + the 2nd Law:

Any large system in equilibrium will be found in the macrostate with the greatest multiplicity

2nd Law of Thermodynamics

Multiplicity tends to increase.

Multiplicities are very large! Take the natural log of them.

$$\text{entropy} \rightarrow S = k_B \ln \Omega$$

$$\frac{S}{k_B} = \ln \Omega$$

$$\Omega = \left(\frac{eq}{N}\right)^N \quad q \gg N \text{ Einstein solid}$$

$$N = 10^{23} \quad q = 10^{25}$$

$$S = k_B \ln \left(\left(\frac{eq}{N} \right)^N \right) = N \cdot k_B \ln \left(\frac{eq}{N} \right) = N k_B \left(1 + \ln \left(\frac{q}{N} \right) \right)$$

\uparrow
 $10^{23} \cdot 1.38 \cdot 10^{-23}$

$$= 1.38 \left(1 + \underbrace{\ln(10^2)}_{4.6} \right)$$

$$S = 7.7 \text{ J/K}$$

