

Eulerian-Lagrangian SS

- Ensemble averaging Eqs. of DNS-approach 2, dropping the index k assuming one continuous phase, for the derivation see references [1, 2]
- Fluctuation definition $(Q'' \equiv Q \tilde{Q})$
- Averaging non-linear terms produces additional unclosed terms (closure problem)

Example ▼

 $\frac{\partial}{\partial t} \overline{\langle \rho_k \chi_k \rangle} + \frac{\partial}{\partial x_i} \overline{\langle \rho_k U_{k,j} \chi_k \rangle} = \left\langle S_m^{(I_k)} \right\rangle$

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The Eulerian continuous-phase equations

• The continuous phase continuity:

$$\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial}{\partial x_j} \left(\bar{\rho} \widetilde{U}_j \right) = \langle S_m^{(I)} \rangle$$

• The continuous phase momentum:

$$\frac{\partial}{\partial t} (\bar{\rho} \widetilde{U}_{i}) + \frac{\partial}{\partial x_{j}} (\bar{\rho} \widetilde{U}_{j} \widetilde{U}_{i})$$

$$= -\frac{\partial \bar{p}}{\partial x_{j}} + \frac{\partial \bar{\tau}_{ij}}{\partial x_{j}} + \bar{\rho} g_{i} - \frac{\partial}{\partial x_{j}} (\bar{\rho} \widetilde{u_{i}''} u_{j}'') + \langle S_{u_{i}}^{(I)} \rangle + \langle U_{I,i} S_{m}^{(I)} \rangle$$

• The continuous phase scalars:

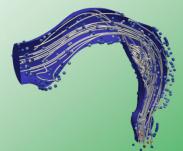
$$\frac{\partial}{\partial t} (\bar{\rho}\tilde{Q}) + \frac{\partial}{\partial x_{j}} (\bar{\rho}\tilde{U}_{j}\tilde{Q})$$

$$= -\frac{\partial \bar{J}_{Q,j}}{\partial x_{j}} - \frac{\partial}{\partial x_{j}} (\bar{\rho}\tilde{u}_{j}^{"Q"}) + \bar{\rho}\tilde{S}_{Q} + \langle S_{Q}^{(I)} \rangle + \langle Q_{I}S_{m}^{(I)} \rangle$$

The Eulerian continuous-phase equations

• Volume fraction by sampling particles:

$$\alpha_c = 1 - \alpha_d$$



• For the summary of the Eulerian continuousphase equations, see file "Chap8-ELSSContinuousPhaseEquations.pdf".

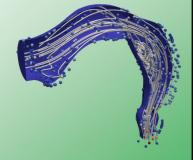
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The Eulerian continuous-phase equations

- Closures:
 - + closures for interface transfer (two-way coupling) source terms, $\langle S_m^{(I)} \rangle$, $\langle S_{U_i}^{(I)} \rangle$, $\langle U_{I,i} S_m^{(I)} \rangle$ etc. satisfying averaged jump condition constraints
 - + closures for Reynolds stresses and fluxes, $\widetilde{u_i''u_j''}$ and $\widetilde{u_i''Q''}$



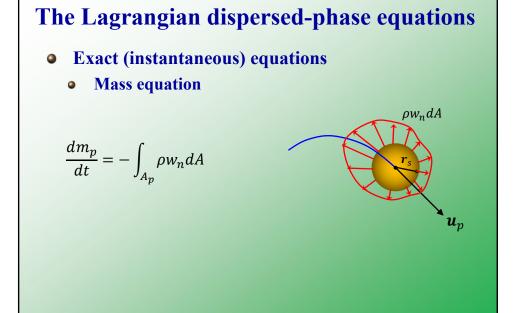


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3



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The Lagrangian dispersed-phase equations

Exact (instantaneous) equations

Momentum equation

Center of mass position

ter of mass position
$$\frac{dx_p}{dt} = u_p$$
 Center of mass velocity Fluid-particle interaction force $F = m_p \frac{du_p}{dt}$; $F = F_{\rm FPI} + F_{\rm B} + F_{\rm col}$ Body force Collision force $F_{\rm FPI} = \int_{A_p} \sigma . \, n dA = \int_{A_p} (-p n + \tau . \, n) dA$

$$\mathbf{F} = m_p \frac{d\mathbf{u}_p}{dt}; \quad \mathbf{F} = \mathbf{F}_{\text{FPI}} + \mathbf{F}_{\text{B}} + \mathbf{F}_{\text{col}}$$

$$\mathbf{F}_{\text{FPI}} = \int_{A_p} \boldsymbol{\sigma} \cdot \boldsymbol{n} dA = \int_{A_p} (-p\boldsymbol{n} + \boldsymbol{\tau} \cdot \boldsymbol{n}) dA$$

- Assumptions:
 - Neglecting internal flow effect
 - · Neglecting tangential outflow effect
 - Assuming uniform efflux (zero-thrust)

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Assuming spherical particle

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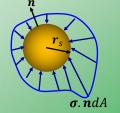
The Lagrangian dispersed-phase equations

- **Exact (instantaneous) equations**
 - **Angular momentum equation**

(mass) moment

of inertia
$$oldsymbol{T} = I_p rac{doldsymbol{\omega}_p}{dt}; oldsymbol{T} = oldsymbol{T}_{\mathrm{FPI}} + oldsymbol{T}_{\mathrm{col}}$$

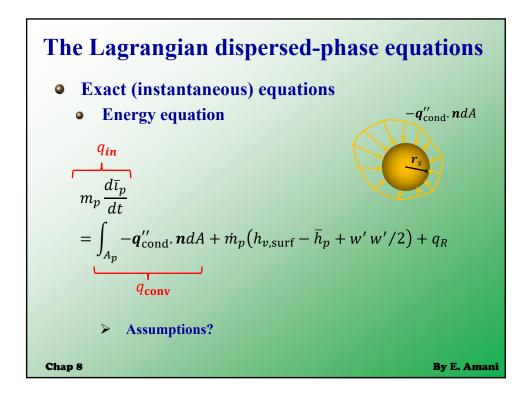
$$T_{\text{FPI}} = \int_{A_p} r_{\scriptscriptstyle S} \times (\boldsymbol{n}.\,\boldsymbol{\tau}) dA$$

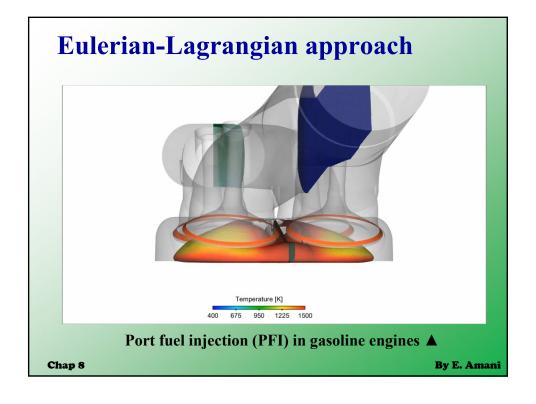


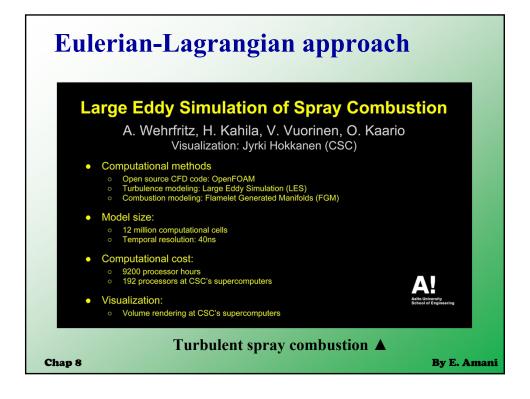
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Chap 8

Assuming spherical particle







References

- Michaelides, E., C.T. Crowe, and J.D. Schwarzkopf, Multiphase flow handbook. 2016: CRC Press.
- 2. Naud, B., PDF modeling of turbulent sprays and flames using a particle stochastic approach. 2003.
- 3. Kataoka, I., *Local instant formulation of two-phase flow.* International Journal of Multiphase Flow, 1986. 12(5): p. 745-758.
- 4. Tryggvason, G., R. Scardovelli, and S. Zaleski, *Direct Numerical Simulations of Gas-Liquid Multiphase Flows*. 2011: Cambridge University Press.

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