

Aliasing error: Eq. (7.21)

To prove this behavior, we derive the relation between \hat{f}_k and \tilde{f}_k :

Eq. (7.15) $\hat{f}_k = \frac{1}{N} \sum_{j=0}^{N-1} f(x_j) e^{-ikx_j} = \frac{1}{N} \sum_{j=0}^{N-1} \left[\sum_{K=-\infty}^{+\infty} \tilde{f}_K e^{iKx_j} \right] e^{-ikx_j} = \sum_{K=-\infty}^{+\infty} \tilde{f}_K \frac{1}{N} \sum_{j=0}^{N-1} e^{i(K-k)x_j}$

Eq. (7.14) $\frac{1}{N} \sum_{j=0}^{N-1} e^{i(K-k)x_j} = \begin{cases} 1 & \text{if } K=k \\ 0 & \text{otherwise} \end{cases}$

Eq. (7.20) $\hat{f}_k = \sum_{m=-\infty}^{+\infty} \tilde{f}_{k+mN} = \tilde{f}_k + \sum_{m \neq 0} \tilde{f}_{k+mN}$

Note: $I_N f(x) = \sum_{k=-N/2}^{N/2-1} \hat{f}_k e^{ikx} = \mathcal{P}_N f(x) + \sum_{k=-N/2}^{N/2-1} \left(\sum_{m \neq 0} \tilde{f}_{k+mN} \right) e^{ikx}$ (7.21)

\mathcal{P}_N (see Eq. (7.18))

Convolution sum: Eq. (7.22)

We can avoid the aliasing error resulted from the product operation. For this purpose we begin with:

$\hat{H}_m = (\hat{f} \hat{g})_m = \frac{1}{N} \sum_{j=0}^{N-1} f_j g_j e^{-imx_j} = \frac{1}{N} \sum_{j=0}^{N-1} \sum_{K=-N/2}^{N/2-1} \sum_{K'=-N/2}^{N/2-1} \hat{f}_K \hat{g}_{K'} e^{i(K+K'-m)x_j}$

$= \sum_{K=-N/2}^{N/2-1} \sum_{K'=-N/2}^{N/2-1} \hat{f}_K \hat{g}_{K'} \frac{1}{N} \sum_{j=0}^{N-1} e^{i(K+K'-m)x_j}$

(I)

According to Eq. (7.14), the term (I) is non-zero for $K = K' - m = pN$, $p = 0, \pm 1, \pm 2, \dots$. As K, K' , and m are all integers in the range $(-N/2, N/2-1)$ the only possible cases are $K = -N, 0, +N$. As you know the maximum presentable wavenumber in the grid is $|k|_{\max} = N/2$, therefore the waves $K = \pm N$ which was generated by product operation alias the other modes. The remedy is simply to discard these modes, $K = \pm N$, and only keep the mode $K = 0$ (or $K = m - K'$) in the summation. Hence

$\hat{H}_m = \sum_{K=-N/2}^{N/2-1} \hat{f}_K \hat{g}_{m-K} \quad m = -N/2, \dots, N/2-1$ (7.22)

This is called the convolution sum of the Fourier coefficients of f and g .

Note that only \hat{g}_K ; $-N/2 \leq K \leq N/2-1$ are available and other \hat{g}_K in Eq. (7.22) are considered to be zero.

Energy spectrum function

proof (7.73)

$$(7.86) \rightarrow \phi_{ij}(\vec{k}) = \langle \hat{u}_i^*(\vec{k}) \hat{u}_j(\vec{k}) \rangle = \sum_{\vec{k}'} \sum_{\vec{k}''} \underbrace{\delta(\vec{k}-\vec{k}') \delta(\vec{k}-\vec{k}'')}_{\text{Kronecker delta}} \langle \hat{u}_i^*(\vec{k}') \hat{u}_j(\vec{k}'') \rangle$$

① $\delta(\vec{k}-\vec{k}') \delta(\vec{k}-\vec{k}'') = \delta(\vec{k}-\vec{k}') \delta(\vec{k}'-\vec{k}'')$

$$= \sum_{\vec{k}'} \left[\left(\sum_{\vec{k}''} \delta(\vec{k}-\vec{k}') \delta(\vec{k}'-\vec{k}'') \right) \langle \hat{u}_i^*(\vec{k}') \hat{u}_j(\vec{k}'') \rangle \right]$$

② $\delta(\vec{k}-\vec{k}') \delta(\vec{k}'-\vec{k}'') = \delta(\vec{k}-\vec{k}') \delta(\vec{k}'-\vec{k})$

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proof (7.82)

$$\int_0^\infty E(k) dk = \int_0^\infty \sum_{\vec{k}'} \delta(k'-k) \hat{E}(\vec{k}') dk = \sum_{\vec{k}'} \hat{E}(\vec{k}') \int_0^\infty \delta(k'-k) dk = \sum_{\vec{k}'} \hat{E}(\vec{k}') = k$$

proof (7.83)

$$\int_0^\infty 2\nu k^2 E(k) dk = \int_0^\infty \sum_{\vec{k}'} \delta(k'-k) \hat{E}(\vec{k}') dk 2\nu k^2 dk = \sum_{\vec{k}'} \hat{E}(\vec{k}') \int_0^\infty 2\nu k^2 \delta(k'-k) dk$$

① $\int_0^\infty 2\nu k^2 \delta(k'-k) dk = 2\nu k'^2$

$$= \sum_{\vec{k}'} 2\nu k'^2 \hat{E}(\vec{k}') = \epsilon$$

proof (7.81)

$$\iint \frac{1}{2} \phi_{ii}(\vec{k}') \delta(\vec{k}-\vec{k}') d\vec{k}' = \iint \frac{1}{2} \sum_{\vec{k}''} \delta(\vec{k}-\vec{k}'') \hat{E}(\vec{k}'') \delta(\vec{k}'-\vec{k}'') d\vec{k}' = \sum_{\vec{k}''} \hat{E}(\vec{k}'') \iint \delta(\vec{k}-\vec{k}'') \delta(\vec{k}'-\vec{k}'') d\vec{k}'$$

① $\delta(\vec{k}-\vec{k}'') \delta(\vec{k}'-\vec{k}'') = \delta(\vec{k}-\vec{k}'') \delta(\vec{k}'-\vec{k}'')$

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Math note: shifting property of delta function, see [7] appendix C

① $\iint g(\vec{x}) \delta(\vec{x}-\vec{y}) d\vec{x} = g(\vec{y})$ (7.84)