



## Free convection: Applications

- Air conditioning

The diagram shows two cross-sections of a room. In the top section, warm air (orange arrow) rises from a radiator at the bottom, while cool air (blue arrow) enters through a window at the top. In the bottom section, cool air (red arrow) enters through a window at the top, while warm air (orange arrow) rises from a radiator at the bottom. Both diagrams include a watermark for TEC-SCIENCE.COM.

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## Free convection: Applications

- Multi-pane windows



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## Free convection: Applications

- Phase Change Material (PCM)

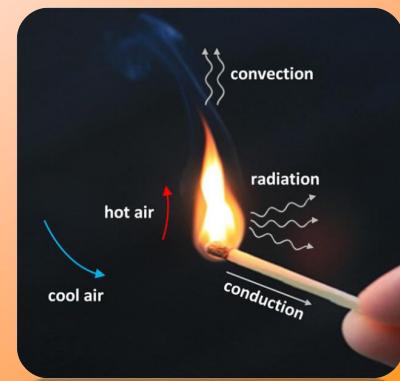


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## Free (or natural) convection

- Fluid motion induced by buoyancy forces
- Buoyancy forces: Density gradients + a body force proportional to density
- In heat transfer: density gradients are due to temperature gradients and the body force is gravitational

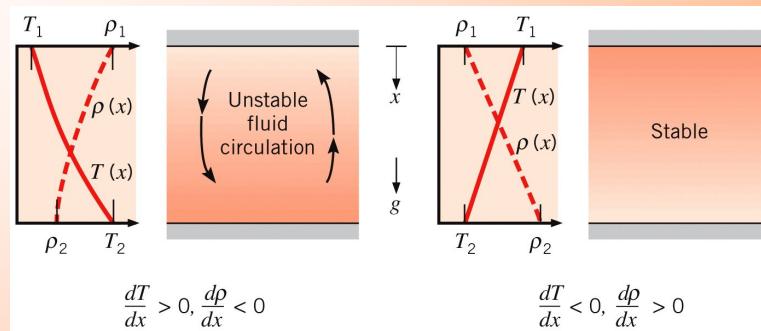


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## Free (or natural) convection

- Stable and Unstable Temperature Gradients



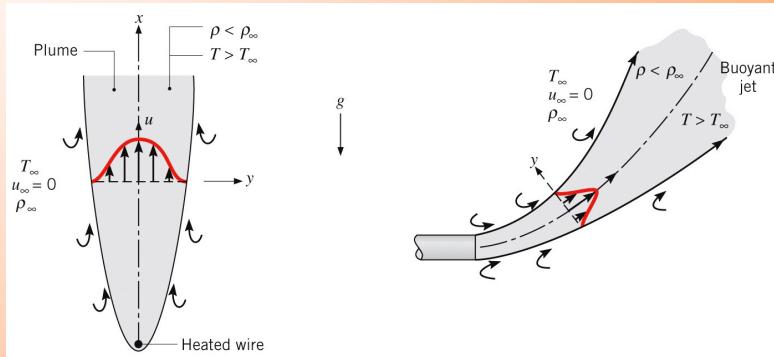
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## Free convection classification

### • Free Boundary Flows

- In an extensive (infinite), quiescent (motionless at locations far from the source of buoyancy) fluid
- Examples: Plumes and Buoyant Jets



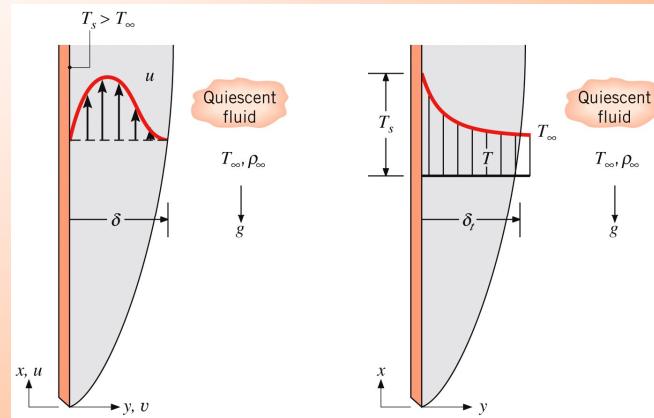
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## Free convection classification

### • Free Convection Boundary Layers

- Boundary layer flow on a hot or cold wall induced by buoyancy



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## Governing equations

- From Heat Transfer I (No phase change and  $\dot{q} = 0$ )

**Continuity:**  $\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0 \quad (1.4)'$

**Momentum:**

$$\rho \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = - \frac{\partial p}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + f_x \quad (2.4)'$$

$$\rho \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = - \frac{\partial p}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + f_y \quad (3.4)'$$

$$\rho \left( \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) = - \frac{\partial p}{\partial z} + \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \tau_{zz}}{\partial z} + f_z \quad (4.4)'$$

**Energy:**

$$\rho \left( \frac{\partial u_s}{\partial t} + u \frac{\partial u_s}{\partial x} + v \frac{\partial u_s}{\partial y} + w \frac{\partial u_s}{\partial z} \right) = \frac{\partial}{\partial x} \left( k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left( k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( k \frac{\partial T}{\partial z} \right) - p \left( \frac{\partial u_x}{\partial x} + \frac{\partial u_y}{\partial y} + \frac{\partial u_z}{\partial z} \right) + \mu \Phi \quad (5.4)'$$

2 state equations, e.g.:

**Chapter 4**  $u_s = \int_{T_{ref}}^T c_v dT \quad (6.4)' \quad \rho = \rho(p, T) \quad (7.4)' \quad \text{By E. Amani}$

## Governing equations

- Boussinesq approximation

- For not too-large temperature gradients
  - Considering a constant density except in the (driving) buoyancy and gravity force terms in the momentum equation
  - Criterion:  $\frac{|\Delta\rho|}{\rho_0} \ll 1$
- Reference temperature  $\rho_0 \equiv \rho(T_0)$
- Reference density  $\Delta\rho \equiv \rho - \rho_0 \quad (1.4)$

Volumetric thermal expansion coefficient ( $K^{-1}$ ):  $\beta \equiv -\frac{1}{\rho} \left( \frac{\partial \rho}{\partial T} \right)_p \approx cte \approx -\frac{1}{\rho_0(1 + \Delta\rho/\rho_0)} \frac{\rho - \rho_0}{T - T_0} \quad | \Delta\rho/\rho_0 \ll 1 |$

Tables A.5 and A.6  $\rho - \rho_0 = \Delta\rho = -\beta\rho_0(T - T_0) \quad (6.4)$

$\beta |T - T_0| \ll 1 \quad (8.4)$

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## Governing equations

### • Boussinesq approximation

#### ➤ Simplifying:

Hydrostatic reference pressure  $\frac{\partial p_0}{\partial x_i} \equiv \rho_0 g_i$  (3.4)

Modified pressure  $P \equiv p - p_0$  (4.4)

$$\rho_0 \left(1 + \frac{\Delta\rho}{\rho_0}\right) \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = -\frac{\partial P}{\partial x} + \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + (\rho_0 + \Delta\rho) g_x \\ |\Delta\rho|/\rho_0 \ll 1$$

$$-\beta \rho_0 (T - T_0)$$

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## Governing equations

### • Boussinesq approximation

**Continuity:**  $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$  (1.4)''

#### Momentum:

$$\rho_0 \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = -\frac{\partial P}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) - \rho_0 g_x \beta (T - T_0) \quad (2.4)''$$

$$\rho_0 \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = -\frac{\partial P}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) - \rho_0 g_y \beta (T - T_0) \quad (3.4)''$$

$$\rho_0 \left( \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) = -\frac{\partial P}{\partial z} + \mu \left( \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} + \frac{\partial^2 w}{\partial z^2} \right) - \rho_0 g_z \beta (T - T_0) \quad (4.4)''$$

#### Energy:

$$\rho_0 \left( \frac{\partial u_s}{\partial t} + u \frac{\partial u_s}{\partial x} + v \frac{\partial u_s}{\partial y} + w \frac{\partial u_s}{\partial z} \right) = \frac{\partial}{\partial x} \left( k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left( k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( k \frac{\partial T}{\partial z} \right) + \mu \Phi \quad (5.4)''$$

**1 state equation**  $u_s = \int_{T_{ref}}^T c_v dT \quad (6.4)''$

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## Governing equations

- **Boussinesq + Boundary Layer (BL) approximations**

➢ 2D + Steady + laminar

➢  $\mu, k, c_v = cte$

➢ **Vertical plane wall ( $g_x = -g$ )**

➢ **BL approximations**

$$\frac{\partial p}{\partial y} \sim 0 \rightarrow p = p(x) = p_\infty(x) \rightarrow \frac{\partial p}{\partial x} = \frac{dp_\infty}{dx}$$

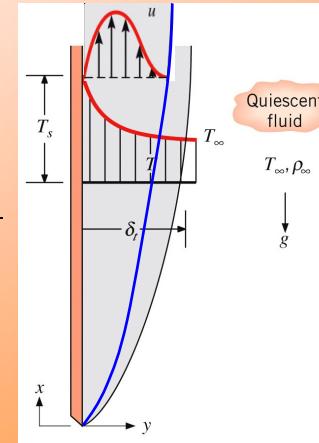
$$u \gg v \quad \frac{\partial u}{\partial y} \gg \frac{\partial u}{\partial x}, \frac{\partial v}{\partial y}, \frac{\partial v}{\partial x} \quad \frac{\partial^2 u}{\partial y^2} \gg \frac{\partial^2 u}{\partial x^2}$$

$$\frac{\partial T}{\partial y} \gg \frac{\partial T}{\partial x} \quad \frac{\partial^2 T}{\partial y^2} \gg \frac{\partial^2 T}{\partial x^2}$$

$$\rho_0 = \rho_\infty \quad T_0 = T_\infty$$

$$(4.4) \quad \frac{\partial P}{\partial x} = \frac{\partial p}{\partial x} - \frac{\partial p_0}{\partial x} = \frac{dp_\infty}{dx} + \rho_0 g = -\rho_\infty g + \rho_\infty g = 0$$

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## Governing equations

- **Boussinesq + Boundary Layer (BL) approximations**

➢ **Equations**

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty) \quad (9.4)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2}$$

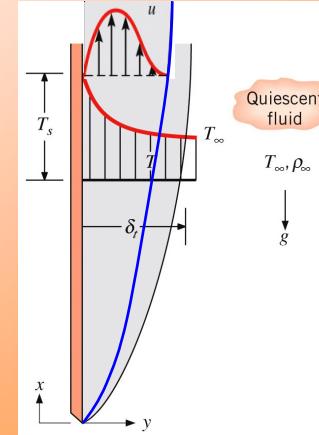
➢ **Boundary conditions:**

$$x = 0; u = 0, T = T_\infty \quad (T - T_\infty = 0)$$

$$y = 0; u = v = 0, T = T_s \quad (T - T_\infty = T_s - T_\infty)$$

$$y \rightarrow \infty; u = 0, T = T_\infty \quad (T - T_\infty = 0) \quad (10.4)$$

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## Dimensional analysis

### • Non-dimensionalization

➢ Therefore,

$$u = u(x, y, g\beta, v, \alpha, T_s - T_\infty, L) \quad \text{For generalization to other geometries}$$

$$T - T_\infty = u(x, y, g\beta, v, \alpha, T_s - T_\infty, L) \quad (11.4)$$

➢ Exercise: Starting from the parametric relations, Eq.

(11.4), show that:

$$\frac{T - T_\infty}{T_s - T_\infty} = f\left(\frac{x^*}{L}, \frac{y^*}{L}, \frac{g\beta(T_s - T_\infty)L^3}{v^2}, \text{Pr}\right) \quad (11.4)'$$

$$\text{Nu}_x = f(x^*, \text{Gr}_L, \text{Pr}) \quad (16.4)'$$

$$\overline{\text{Nu}}_L = f(\text{Gr}_L, \text{Pr}) \quad (16.4)$$

➢ For a flat plate, removing  $L$ , Eq. (16.4)' is simplified to:

$$\text{Nu}_x = f(\text{Gr}_x, \text{Pr}) \quad (16.4)' \quad \text{Gr}_x = \frac{g\beta(T_s - T_\infty)x^3}{v^2}$$

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## Dimensional analysis

### • Non-dimensionalizing equations

➢ Steps:

1. Dependent variables:  $u, v, T$

2. Independent variables:  $x, y$

3. Non-dimensional variables:

$$x^* = \frac{x}{L}, y^* = \frac{y}{L}, u^* = \frac{u}{u_0}, v^* = \frac{v}{u_0}, T^* = \frac{T - T_s}{\theta_0}$$

4. Characteristic parameters:

$$\theta_0 = T_\infty - T_s$$

5. Substituting in the governing equations:

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## Dimensional analysis

### • Non-dimensionalizing equations

➤ Steps:

5. Substituting in the governing equations:

$$\begin{aligned} \frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} &= 0 \\ u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} &= \frac{v}{u_0 L} \frac{\partial^2 u^*}{\partial y^{*2}} + \frac{g\beta(T_s - T_\infty)L}{u_0^2} T^* \\ u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} &= \frac{\alpha}{u_0 L} \frac{\partial^2 T^*}{\partial y^{*2}} \\ x^* = 0; u^* = 0, T^* &= 0 \\ y^* = 0; u^* = v^* = 0, T^* &= 1 \\ y^* \rightarrow \infty; u^* = 0, T^* &= 0 \end{aligned} \tag{12.4}'$$

6. Choosing  $u_0$  to minimize the governing parameters

$$u_0 = [g\beta(T_s - T_\infty)L]^{1/2}$$

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## Dimensional analysis

### • Non-dimensionalizing equations

$$\frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} = 0$$

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{Gr_L^{1/2}} \frac{\partial^2 u^*}{\partial y^{*2}} + T^* \tag{12.4}$$

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{1}{Gr_L^{1/2} Pr} \frac{\partial^2 T^*}{\partial y^{*2}} \quad Gr_L = \frac{g\beta(T_s - T_\infty)L^3}{v^2} \tag{14.4}$$

$$x^* = 0; u^* = 0, T^* = 0$$

$$y^* = 0; u^* = v^* = 0, T^* = 1 \tag{13.4}$$

$$y^* \rightarrow \infty; u^* = 0, T^* = 0$$

➤ Therefore,

$$u^*, T^* = f(x^*, y^*, Gr_L, Pr) \quad Nu_x = f(x^*, Gr_L, Pr) \quad \overline{Nu}_L = f(Gr_L, Pr)$$

➤ What is the ratio of buoyancy to viscous force?

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## Dimensional analysis

### • Non-dimensionalizing equations

➤ Corollary:

$$\overline{\text{Nu}}_L = f(\text{Ra}_L, \text{Pr}) \quad (17.4)$$

**Rayleigh number**  $\text{Ra}_L = \text{Gr}_L \text{Pr} = \frac{g\beta(T_s - T_\infty)L^3}{\nu\alpha} \quad (18.4)$

➤ For liquid metals:

$$\text{Gr}_L^{1/2} \gg 1 \xrightarrow{(12.4)} \overline{\text{Nu}}_L = f(\text{Gr}_L^{1/2} \text{Pr}) = f(\text{Ra}_L^{1/2} \text{Pr}^{1/2}) \longrightarrow$$

$$\overline{\text{Nu}}_L = f(\text{Ra}_L \text{Pr}) \quad (19.4)$$

$$\frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} = 0 \quad u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{\text{Gr}_L^{1/2}} \frac{\partial^2 u^*}{\partial y^{*2}} + T^* \quad u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{1}{\text{Gr}_L^{1/2} \text{Pr}} \frac{\partial^2 T^*}{\partial y^{*2}}$$

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## Dimensional analysis

### • Mixed (or combined free and force) convection

➤ Boussinesq + Boundary Layer (BL) approximations

➤ In Eq. (2.4)'',  $\frac{\partial P}{\partial x} \neq 0$ .

➤ Show that Eqs. (12.4) and (13.4) are extended to:

$$\frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} = 0 \quad u_0 = u_\infty, \quad p^* = \frac{P}{\rho_0 u_\infty^2}$$

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{\text{Re}_L} \frac{\partial^2 u^*}{\partial y^{*2}} + \frac{\text{Gr}_L}{\text{Re}_L^2} T^* - \frac{dp^*}{dx^*} \quad (21.4)$$

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{1}{\text{Re}_L \text{Pr}} \frac{\partial^2 T^*}{\partial y^{*2}}$$

$$x^* = 0; u^* = 1, T^* = 0$$

$$y^* = 0; u^* = v^* = 0, T^* = 1 \quad (22.4)$$

$$y^* \rightarrow \infty; u^* = 1, T^* = 0$$

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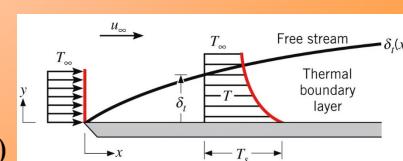


Figure 6.2  
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## Dimensional analysis

- **Mixed (or combined free and force) convection**

➤ Therefore,

$$u^*, T^* = f(x^*, y^*, \text{Gr}_L, \text{Re}_L, \text{Pr}) \quad (23.4)$$

$$\text{Nu}_x = f(x^*, \text{Gr}_L, \text{Re}_L, \text{Pr}) \quad \overline{\text{Nu}}_L = f(\text{Gr}_L, \text{Re}_L, \text{Pr}) \quad (24.4)$$

➤ The ratio of buoyancy to (forced convection) inertial force:  $\text{Gr}_L/\text{Re}_L^2 \ll 1$

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{\text{Re}_L} \frac{\partial^2 u^*}{\partial y^{*2}} + \frac{\text{Gr}_L}{\text{Re}_L^2} T^* - \frac{dp^*}{dx^*}$$

➤ The free convection is omitted when:  $\boxed{\text{Gr}_L/\text{Re}_L^2 \ll 1}$

➤ The forced convection is omitted when:  $\boxed{\text{Gr}_L/\text{Re}_L^2 \gg 1}$

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## Laminar free convection on a vertical plate

- **Semi-analytical solution: The similarity solution**

- **Details of derivation: Appendix A**

➤ With a change of variables:

$$\eta = \frac{y}{x} \left( \frac{\text{Gr}_x}{4} \right)^{1/4} \quad (32.4) \quad \text{Gr}_x = \frac{g\beta(T_s - T_\infty)x^3}{v^2} \quad (31.4)$$

$$f'(\eta) = \frac{u}{2v} \frac{x}{\sqrt{\text{Gr}_x}} \quad \text{or} \quad u = \frac{2v}{x} \sqrt{\text{Gr}_x} f'(\eta) \quad (33.4)$$

$$\frac{T - T_\infty}{T_s - T_\infty} = T^*(\eta) \quad (34.4)$$

➤ Eqs. (12.4) and (13.4) are simplified to

$$\begin{aligned} f''' + 3ff'' - 2f'^2 + T^* &= 0 \\ T^{***} + 3\text{Pr}fT^{**} &= 0 \end{aligned} \quad (27.4)$$

$$\begin{aligned} \eta &= 0; f = f' = 0, T^* = 1 \\ \eta &\rightarrow \infty; f' = 0, T^* = 0 \end{aligned} \quad (30.4)$$

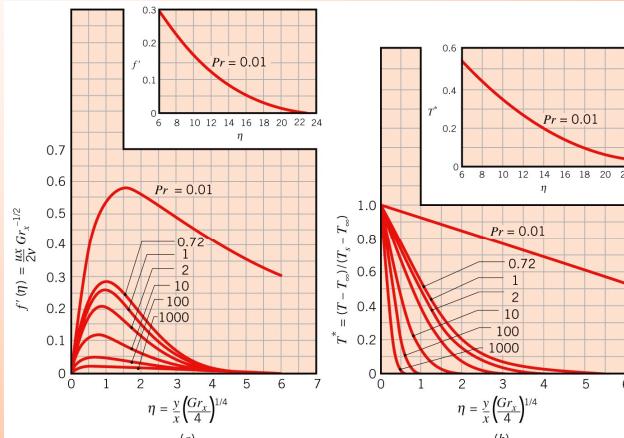
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## Laminar free convection on a vertical plate

- **Semi-analytical solution: The similarity solution**

➤ Numerical solution of Eqs. (27.4) and (30.4):



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Figure 9.4  
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## Laminar free convection on a vertical plate

- **Semi-analytical solution: The similarity solution**

➤ Show that:

$$\text{Nu}_x \equiv \frac{hx}{k} = - \left. \frac{dT^*}{d\eta} \right|_{\eta=0} \left( \frac{d\eta}{dy} \right)_x = \left( \frac{\text{Gr}_x}{4} \right)^{\frac{1}{4}} \underbrace{g(\text{Pr})}_{T^{*'}(\eta=0)}$$

$$\text{Nu}_x = \left( \frac{\text{Gr}_x}{4} \right)^{\frac{1}{4}} g(\text{Pr}) \quad (35.4)$$

➤ By curve fitting to the numerical solution:

$$g(\text{Pr}) = \frac{0.75 \text{Pr}^{1/2}}{(0.609 + 1.221 \text{Pr}^{1/2} + 1.238 \text{Pr})^{1/4}} \quad (0 < \text{Pr} < \infty) \quad (36.4)$$

➤ For a plate of height  $L$ :  $(35.4)$

$$\bar{h} = \frac{1}{L} \int_0^L h dx = \int_0^1 h(x^*) dx^* = \frac{k}{L} \int_0^1 \text{Nu}_x(x^*) dx^* = \frac{4}{3} \left( \frac{\text{Gr}_L}{4} \right)^{\frac{1}{4}} g(\text{Pr}) \rightarrow$$

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$$\overline{\text{Nu}}_L = \frac{4}{3} \text{Nu}_x(x = L) = \frac{4}{3} \text{Nu}_L \quad (37.4)$$

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## Laminar free convection on a vertical plate

- **Semi-analytical solution: The similarity solution**

➤ Therefore,

$$\overline{Nu}_L = \frac{4}{3} \left( \frac{\text{Gr}_L}{4} \right)^{\frac{1}{4}} \frac{0.75 \text{Pr}^{1/2}}{(0.609 + 1.221\text{Pr}^{1/2} + 1.238\text{Pr})^{1/4}} = \frac{0.707 \text{Ra}_L^{1/4}}{(0.609\text{Pr}^{-1} + 1.221\text{Pr}^{-1/2} + 1.238)^{1/4}} \quad (37.4)$$

➤ Alternatively, a purely empirical correlation suggests that:

$$\overline{Nu}_L = 0.68 + \frac{0.670 \text{Ra}_L^{1/4}}{\left[ 1 + (0.492 / \text{Pr})^{9/16} \right]^{4/9}} \quad (9.27)[1]$$

➤ **Exercise:** Are Eq. (37.4) and (9.27)[1] consistent for  $\text{Ra}_L \rightarrow 0$  and  $\text{Ra}_L \rightarrow \infty$ ? Why? Hint: Which condition leads to the violation of BL assumptions, i.e., the thin BL?

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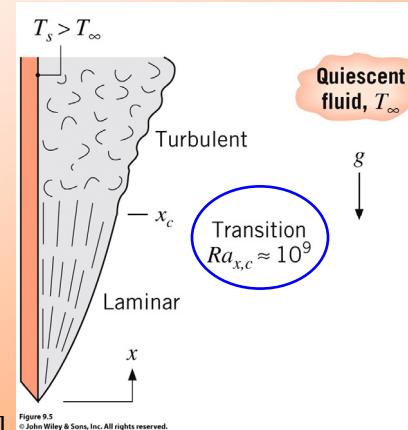
## Turbulent free convection on a vertical plate

- **Transition to turbulence**

$$\text{Ra}_x = \text{Gr}_x \text{Pr} = \frac{g\beta(T_s - T_\infty)x^3}{\nu\alpha} \quad (38.4)$$

➤ Empirical correlation:

$$\overline{Nu}_L = \left\{ 0.825 + \frac{0.387 \text{Ra}_L^{1/6}}{\left[ 1 + (0.492 / \text{Pr})^{9/16} \right]^{8/27}} \right\}^2 \quad (9.26)[1]$$



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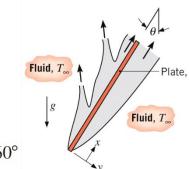
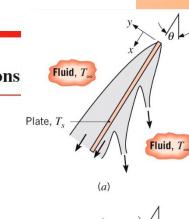
## Empirical correlations

- External free convection: Table 1 in “chap4-summary.pdf”

**TABLE 9.2** Summary of free convection empirical correlations for immersed geometries

Geometry	Recommended Correlation	Restrictions
1. Vertical plates <sup>a</sup>	Equation 9.26	None
2. Inclined plates Cold surface up or hot surface down	Equation 9.26 $g \rightarrow g \cos \theta$	$0 \leq \theta \leq 60^\circ$

<sup>a</sup> The correlation may be applied to a vertical cylinder if  $(D/L) \geq (35/Gr_L^{1/4})$ .



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## Empirical correlations

- External free convection: Table 1 in “chap4-summary.pdf”

**TABLE 9.2** Summary of free convection empirical correlations for immersed geometries

Geometry	Recommended Correlation	Restrictions
3. Horizontal plates (a) Hot surface up or cold surface down	Equation 9.30 Equation 9.31	$10^4 \leq Ra_L \leq 10^7, Pr \geq 0.7$ $10^7 \leq Ra_L \leq 10^{11}$
(b) Cold surface up or hot surface down	Equation 9.32	$10^4 \leq Ra_L \leq 10^9, Pr \geq 0.7$

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## Empirical correlations

- External free convection: Table 1 in “chap4-summary.pdf”

**TABLE 9.2** Summary of free convection empirical correlations for immersed geometries

Geometry	Recommended Correlation	Restrictions
4. Horizontal cylinder	Equation 9.34	$Ra_D \leq 10^{12}$
5. Sphere	Equation 9.35	$Ra_D \leq 10^{11}$ $Pr \geq 0.7$

- Notes:

- Film temperature:  $T_f = \frac{T_s + T_\infty}{2}$
- For  $q''_s = cte$ , use  $T_s = T_s(L/2)$      $q''_s = \bar{h} \left( T_s \left( \frac{L}{2} \right) - T_\infty \right)$

Chapter 4

By E. Amani

## Empirical correlations

- Internal free convection: Table 2 in “chap4-summary.pdf”

Table 2: Internal free convection

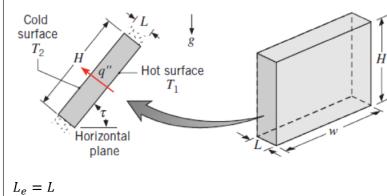
Geometry	Recommended Correlation	Restrictions
1-parallel plate channels <sup>a</sup>	Vertical (9.45),(9.37),(9.38) Table 9.3 1 <sup>st</sup> and 3 <sup>rd</sup> rows  (9.46),(9.41),(9.42) Table 9.3 2 <sup>nd</sup> and 4 <sup>th</sup> rows  <b>Inclined</b> (9.47)	Constant Temperature  Constant flux  Constant Temperature $0 \leq \theta \leq 45^\circ$ $Ra_S(S/L) > 200$ fluid: water
2-Rectangular cavities ( $W/L \gg 1$ ) <sup>b</sup>	(9.48)	

Chapter 4

By E. Amani

## Empirical correlations

- Internal free convection: Table 2 in “chap4-summary.pdf”

2-Rectangular cavities ( $W/L \gg 1$ ) <sup>b</sup>	(9.48)	
	<b>Horizontal (<math>\tau = 0</math>)</b> (9.49) $(Ra_{L,cr} = 3 \times 10^5, Ra_{L,cr2} = 1708)$ $3 \times 10^5 \leq Ra_L \leq 7 \times 10^9$	
	<b>Vertical (<math>\tau = 90</math>)</b> (9.50)-(9.53) $(Ra_{L,cr2} = 1000)$ See each correlation	
	<b>Inclined</b> $\tau^* \rightarrow$ Table 9.4 (9.54),(9.55) (9.56),(9.57) $(Ra_{L,cr2} = 1708/\cos \tau)$ $\tau \leq \tau^*$ $\tau > \tau^*$	
3-Concentric cylinders and spheres <sup>b</sup>		$(Ra_{L,cr2} = 100)$

a  $T_f = (T_s(x = L) + T_f)/2$

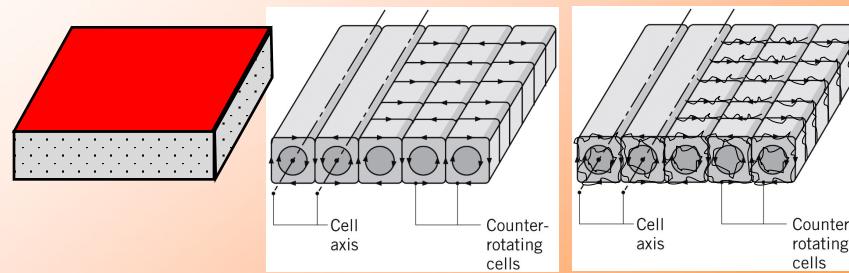
b  $T_f = \frac{T_1+T_2}{2}$ ,  $Ra_L = \frac{g\beta(T_1-T_2)L^3}{\alpha v}$ , and  $Ra_{L,cr2}$  indicates the transition from conduction to free convection.

Chapter 4

By E. Amani

## Empirical correlations

- Internal free convection: Table 2 in “chap4-summary.pdf”



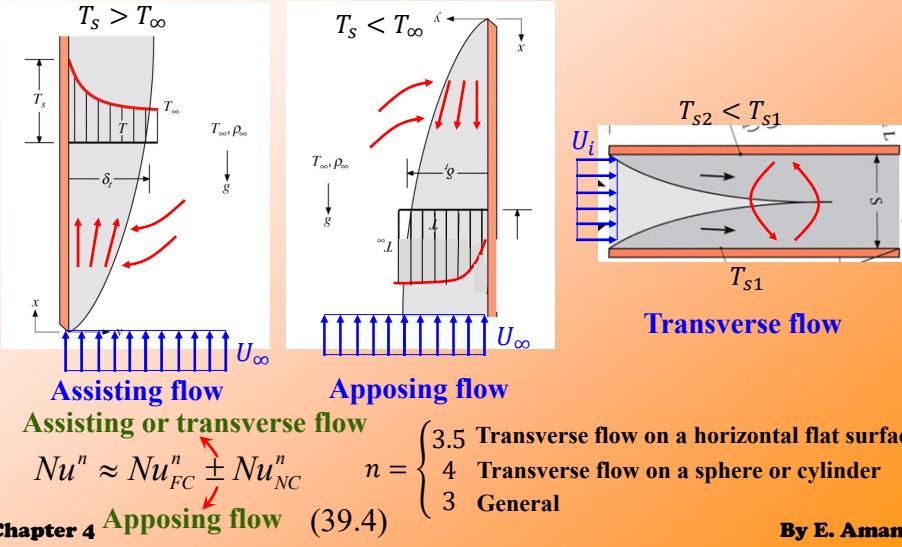
$$\begin{aligned} Ra_L < Ra_{L,cr2} = 1708 & \quad Ra_{L,cr2} < Ra_L < Ra_{L,cr} \\ \overline{Na}_L = 1 & \end{aligned}$$

Chapter 4

By E. Amani

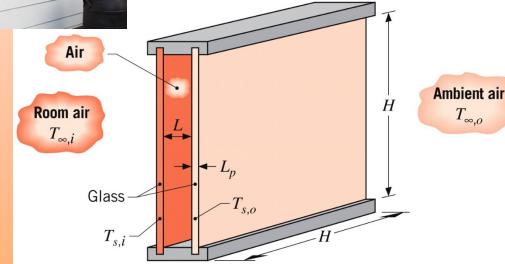
## Empirical correlations

### • Mixed convection



## Sample problems

### Two-pane window



Chapter 4

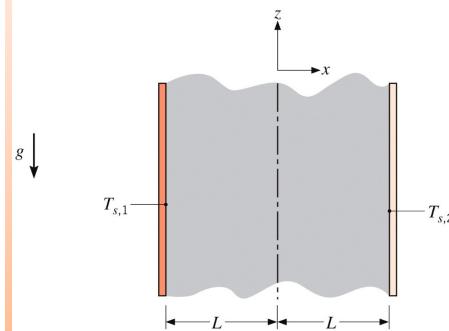
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## Sample problems

### Two-pane window

**9.78** Consider two long vertical plates maintained at uniform temperatures  $T_{s,1} > T_{s,2}$ . The plates are open at their ends and are separated by the distance  $2L$ .

- Sketch the velocity distribution in the space between the plates.
- Write appropriate forms of the continuity, momentum, and energy equations for laminar flow between the plates.
- Evaluate the temperature distribution, and express your result in terms of the mean temperature,  $T_m = (T_{s,1} + T_{s,2})/2$ .
- Estimate the vertical pressure gradient by assuming the density to be a constant  $\rho_m$  corresponding to  $T_m$ . Substituting from the Boussinesq approximation, obtain the resulting form of the momentum equation.
- Determine the velocity distribution.



→ **Lecture Notes: IV.10.1**

**Chapter 4**

**By E. Amani**

## Sample problems

### Two-pane window

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad (1.4)''$$

$$\rho_0 \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = - \frac{\partial P}{\partial x} + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) - \rho_0 g_x \beta (T - T_0) \quad (2.4)''$$

$$\rho_0 \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = - \frac{\partial P}{\partial y} + \mu \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} + \frac{\partial^2 v}{\partial z^2} \right) - \rho_0 g_y \beta (T - T_0) \quad (3.4)''$$

$$\rho_0 \left( \frac{\partial u_s}{\partial t} + u \frac{\partial u_s}{\partial x} + v \frac{\partial u_s}{\partial y} + w \frac{\partial u_s}{\partial z} \right) = \frac{\partial}{\partial x} \left( k \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left( k \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( k \frac{\partial T}{\partial z} \right) + \mu \Phi \quad (5.4)''$$

**Chapter 4**

**By E. Amani**

## Sample problems

### Automobile rear window heater

**9.27** The vertical rear window of an automobile is of thickness  $L = 8 \text{ mm}$  and height  $H = 0.5 \text{ m}$  and contains fine-meshed heating wires that can induce nearly uniform heat flux on the internal surface of the window.

- (a) Consider steady-state conditions for which the interior surface of the window is exposed to quiescent air at  $10^\circ\text{C}$ , while the exterior surface is exposed to air at  $-10^\circ\text{C}$  moving in parallel flow over the surface with a velocity of  $20 \text{ m/s}$ . Determine the heat flux needed to maintain the interior window surface at  $T_{s,i} = 15^\circ\text{C}$ .



→ **Lecture Notes: IV.10.2**

**Chapter 4**

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## Sample problems

### Automobile rear window heater

$$\overline{Nu}_L = \left\{ 0.825 + \frac{0.387 Ra_L^{1/6}}{[1 + (0.492/Pr)^{9/16}]^{8/27}} \right\}^2 \quad (9.26)$$

$$\overline{Nu}_L = 0.68 + \frac{0.670 Ra_L^{1/4}}{[1 + (0.492/Pr)^{9/16}]^{4/9}} \quad Ra_L \lesssim 10^9 \quad (9.27)$$

**Chapter 4**

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## **The end of chapter 4**

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