Consider a gas at temperature T. Let N be the number of molecules in the gas and let  $N_n$  be the number of molecules with internal energy  $\varepsilon_n$ . By the Maxwell-Boltzmann distribution we have

$$\frac{N_n}{N} = p_n \exp\left(-\frac{\varepsilon_n}{kT}\right)$$

where k is Boltzmann's constant. The coefficient  $p_n$  is a statistical weighting factor that does not depend on T.

Let us now consider the processes by which a molecule transitions between energy levels. The processes are absorption, induced emission, and spontaneous emission. Let  $\varepsilon_m$  be an energy level such that  $\varepsilon_m > \varepsilon_n$ . Let  $B_n$ ,  $B_m$ , and  $A_m$  be coefficients of transition rates such that

$$\frac{dN_n}{dt} = B_n \rho(\nu) N_n \qquad \frac{dN_m}{dt} = B_m \rho(\nu) N_m \qquad \frac{dN_m}{dt} = A_m N_m$$
absorption induced emission spontaneous emission  $\varepsilon_m \to \varepsilon_m$ 

Absorption and induced emission are proportional to  $\rho(\nu)$  which is the radiant energy density of the gas as a function of radiant frequency  $\nu$ . The A and B coefficients are presumed to not depend on temperature T.

At equilibrium, transition rates between  $\varepsilon_m$  and  $\varepsilon_n$  are equal.

$$B_n \rho(\nu) N_n = B_m \rho(\nu) N_m + A_m N_m$$
absorption induced emission spontaneous emission
$$\varepsilon_n \to \varepsilon_m \qquad \text{induced emission} \qquad \varepsilon_m \to \varepsilon_n$$

Divide through by N to obtain

$$B_n \rho(\nu) p_n \exp\left(-\frac{\varepsilon_n}{kT}\right) = B_m \rho(\nu) p_m \exp\left(-\frac{\varepsilon_m}{kT}\right) + A_m p_m \exp\left(-\frac{\varepsilon_m}{kT}\right)$$
absorption
$$\lim_{\varepsilon_n \to \varepsilon_n} \text{ induced emission spontaneous emission}$$

$$\lim_{\varepsilon_m \to \varepsilon_n} \text{ spontaneous emission }$$

Note that for increasing T we have

$$\lim_{T \to \infty} \exp\left(-\frac{\varepsilon_n}{kT}\right) = 1$$

It follows that for  $T \to \infty$  the equilibrium formula is

$$B_n \rho(\nu) p_n = B_m \rho(\nu) p_m + A_m p_m$$

Divide through by  $\rho(\nu)$ .

$$B_n p_n = B_m p_m + A_m p_m / \rho(\nu)$$

Energy density  $\rho(\nu)$  increases with temperature T hence  $A_m p_m/\rho(\nu)$  vanishes for  $T \to \infty$  leaving

$$B_n p_n = B_m p_m$$

Einstein reasoned that the above relation is true in general based on the assumption that the factors involved do not depend on T. By substitution in the absorption term we now have

$$B_{m}\rho(\nu)p_{m}\exp\left(-\frac{\varepsilon_{n}}{kT}\right) = B_{m}\rho(\nu)p_{m}\exp\left(-\frac{\varepsilon_{m}}{kT}\right) + A_{m}p_{m}\exp\left(-\frac{\varepsilon_{m}}{kT}\right)$$
absorption
$$\sum_{n \to \varepsilon_{m}} \text{induced emission} \text{spontaneous emission}$$

$$\sum_{m \to \varepsilon_{m}} \text{spontaneous emission}$$

Divide through by  $B_m p_m$  and rearrange terms.

$$\rho(\nu) \exp\left(-\frac{\varepsilon_n}{kT}\right) - \rho(\nu) \exp\left(-\frac{\varepsilon_m}{kT}\right) = (A_m/B_m) \exp\left(-\frac{\varepsilon_m}{kT}\right)$$
absorption induced emission spontaneous emission
$$\varepsilon_m \to \varepsilon_m \qquad \qquad \varepsilon_m \to \varepsilon_n$$

Solve for energy density  $\rho(\nu)$ .

$$\rho(\nu) = \frac{(A_m/B_m) \exp\left(-\frac{\varepsilon_m}{kT}\right)}{\exp\left(-\frac{\varepsilon_n}{kT}\right) - \exp\left(-\frac{\varepsilon_m}{kT}\right)} = \frac{A_m/B_m}{\exp\left(\frac{\varepsilon_m - \varepsilon_n}{kT}\right) - 1}$$

From Wien's law  $\rho(\nu) = \alpha \nu^3 \exp(-h\nu/kT)$  which is accurate for large  $\nu$ , we have

$$A_m/B_m = \alpha \nu^3$$

and

$$\varepsilon_m - \varepsilon_n = h\nu$$

Planck's law follows directly.

$$\rho(\nu) = \frac{\alpha \nu^3}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$