Problem 1

Starting with $\dot{m} = \rho u A$, prove that the mass flowrate through an isentropic choked nozzle can be written in the form:

$$\dot{m} = \frac{p_0 A^*}{\sqrt{T_0}} \sqrt{\frac{\gamma}{R} \left(\frac{2}{\gamma + 1}\right)^{\frac{(\gamma + 1)}{(\gamma - 1)}}}$$

Assumptions:

Isentropic flow through a nozzle with a choked (sonic) throat.

Solution:

The mass flow at a given cross section in a quasi 1-D flow is given by:

$$\dot{m} = \rho u A$$

Choosing the throat of a choked nozzle as the point of interest, we replace the conditions with sonic conditions, denoted by * and indicating the flow property at the location where M=1.

$$\dot{m} = \rho^* u^* A^*$$

In order to cast this in terms of properties that are more easily known ahead of time, we identify relationships involving the total conditions of the flow, beginning with ρ^* . Using the ideal gas law, we can cast the sonic density in terms of pressure and temperature:

$$p^* = \rho^* R T^*$$

$$\rho^* = \frac{p^*}{RT^*}$$

Now, we find relationships for p^* and T^* .

Beginning with the isentropic relationship between total and static pressure:

$$\frac{p_0}{p} = \left(1 + \frac{\gamma - 1}{2}M^2\right)^{\frac{\gamma}{\gamma - 1}}$$

Setting M = 1:

$$\frac{p_0}{p^*} = \left(1 + \frac{\gamma - 1}{2}\right)^{\frac{\gamma}{\gamma - 1}}$$

$$p^* = \frac{p_0}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}}$$

For temperature:

$$\frac{T_0}{T} = \left(1 + \frac{\gamma - 1}{2}M^2\right)$$

Setting M = 1:

$$\frac{T_0}{T^*} = \left(1 + \frac{\gamma - 1}{2}\right)$$

$$T^* = \frac{T_0}{\left(\frac{\gamma+1}{2}\right)}$$

Substituting into the ideal gas equation yields an expression for ρ^* in terms of p_0 , T_0 , R, and γ :

$$\rho^* = \frac{p_0}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}} \frac{\left(\frac{\gamma+1}{2}\right)}{T_0} \frac{1}{R}$$

Next, we examine the sonic velocity term, u^* . Noting that for choked flow M=1, we observe that the flow velocity must be equal to the speed of sound, a.

$$M = \frac{u^*}{a^*} = 1 \to u^* = a^*$$

$$a^* = \sqrt{\gamma R T^*}$$

Substituting our known equation for T^* :

$$a^* = \sqrt{\frac{\gamma R T_0}{\left(\frac{\gamma + 1}{2}\right)}}$$

Substituting everything back into the original mass flow equation:

$$\dot{m} = \frac{p_0}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}} \frac{\left(\frac{\gamma+1}{2}\right)}{T_0} \frac{1}{R} \sqrt{\frac{\gamma R T_0}{\left(\frac{\gamma+1}{2}\right)}} A^*$$

Rearranging:

$$\dot{m} = \frac{p_0 A^*}{R T_0} \frac{\frac{\gamma+1}{2}}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}} \sqrt{\frac{\gamma R T_0}{\left(\frac{\gamma+1}{2}\right)}}$$

$$\dot{m} = p_0 A^* \frac{1}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}} \sqrt{\frac{\gamma R T_0}{R^2 T_0^2} \frac{\left(\frac{\gamma+1}{2}\right)^2}{\left(\frac{\gamma+1}{2}\right)}}$$

$$\dot{m} = \frac{p_0 A^*}{\sqrt{T_0}} \frac{1}{\left(\frac{\gamma+1}{2}\right)^{\frac{\gamma}{\gamma-1}}} \sqrt{\frac{\gamma}{R} \left(\frac{\gamma+1}{2}\right)}$$

$$\dot{m} = \frac{p_0 A^*}{\sqrt{T_0}} \sqrt{\frac{\gamma}{R} \left(\frac{\gamma+1}{2}\right)^{\frac{1-\frac{2\gamma}{\gamma-1}}{\gamma-1}}}$$

$$\dot{m} = \frac{p_0 A^*}{\sqrt{T_0}} \sqrt{\frac{\gamma}{R} \left(\frac{\gamma+1}{2}\right)^{\frac{-\gamma-1}{\gamma-1}}}$$

$$\dot{m} = \frac{p_0 A^*}{\sqrt{T_0}} \sqrt{\frac{\gamma}{R} \left(\frac{2}{\gamma+1}\right)^{\frac{(\gamma+1)}{(\gamma-1)}}}$$