

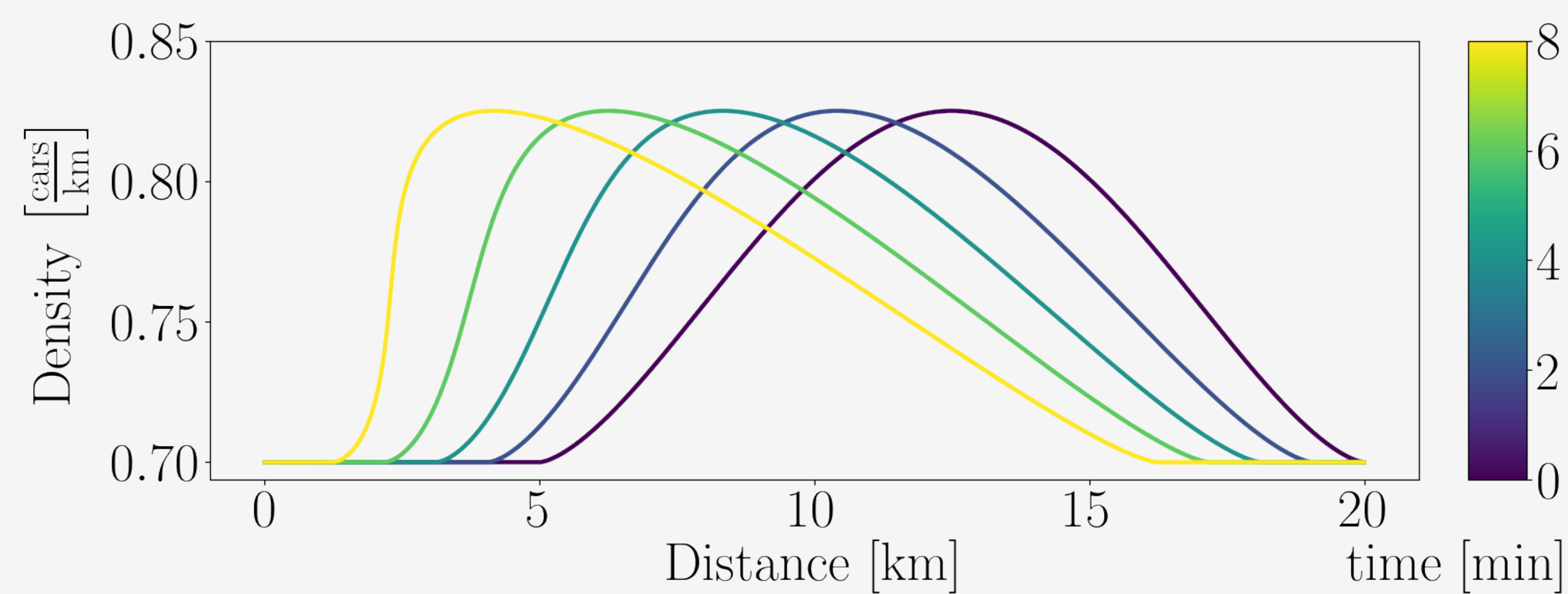
Title

Ingress

A new second order scheme for hyperbolic conversation laws allowing fast and accurate computation of smooth solutions.

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Abstract

The main topic of this poster is to introduce the Lax-Wendroff scheme for solving the hyperbolic systems of conservations laws. In concrete, we will focus on the 1 space dimensional advection equation for different flows to illustrate how our scheme works, and to compare it with other numerical methods. This will lead us to conclude that the Lax-Wendroff scheme works, in general terms, better than the other ones.

Introduction

Hyperbolic systems of conservation laws are time-dependent systems of partial differential equations that describes the conservation of some quantities. In the 1D case the equation is:

$$u_t + f(u)_x = 0$$

Where $u(x,t)$ is a vector of conserved quantities (could be mass, momentum, heat,...) and $f(u)$ is the flux function. This equation must be augmented with initial and boundary conditions. The easiest problem is the Cauchy problem where we need to give initial conditions:

$$u(x, 0) = u_0(x)$$

Flux functions are commonly nonlinear functions, leading to nonlinear systems of PDEs that, in general, the exact solution is unknown. Hence numerical methods are used, to compute an approximation.

The Lax-Wendroff Method

For linear hyperbolic systems $u_t + au_x = 0$ the Lax-Wendroff methods is given as:

$$U_j^{n+1} = U_j^n - \frac{k}{2h}a(U_{j+1}^n - U_{j-1}^n) + \frac{k^2}{2h^2}a^2(U_{j+1}^n - 2U_j^n + U_{j-1}^n)$$

The idea is to take the first three terms of the Taylor series expansion and use centered difference approximations for the derivatives appearing there. The stepsizes k in time and h in space are chosen such that $|a|\frac{k}{h} \leq 1$ to fulfil the CFL condition. For non-linear equations the scheme can be generalized to:

$$U_j^{n+1} = U_j^n - \frac{k}{h} \left(f \left(U_{j+\frac{1}{2}}^{n+\frac{1}{2}} \right) - f \left(U_{j-\frac{1}{2}}^{n+\frac{1}{2}} \right) \right) \text{ where } U_{j+\frac{1}{2}}^{n+\frac{1}{2}} = \frac{1}{2}(U_j^n + U_{j+1}^n) - \frac{k}{2h}[f(U_{j+1}^n) - f(U_j^n)]$$

while the step sizes are in the following considered to be controlled by $\frac{k}{h} = \text{const}$. The computation of the next time step can be written as a function of the solution of the current timestep by $U^{n+1} = \mathcal{H}(U^n)$. This notation can be used to define the local truncation error of a method by

$$L_k(x, t) = \frac{u(x, t+k) - \mathcal{H}(u(\cdot, t); x)}{k}$$

contingent on the used step size in time k . This error will now be considered in the linear advection equation with $a > 0$. By Taylor expansion around $u(x, t)$ and some simplifications this yields to:

$$L_k(x, t) = \frac{a}{6}(a^2k^2 - h^2)u_{xxx} + \mathcal{O}(h^3)$$

As $L_k(x, t) = \mathcal{O}(h^2)$ the method is consistent and has order 2. Additional for stability the condition $|\nu| < 1$ where $\nu = a\frac{k}{h}$ is the Courant number must be fulfilled. In the case of an linear advection equation the convergence of a method follows from consistency and stability by the Lax equivalence theorem. In the nonlinear case Lax and Wendroff proved that a consistent numerical method with bounded solutions always converges to a weak solution of the equation. The local truncation error also leads to the modified equation which is given as:

$$u_t + au_x = \frac{a}{6}(a^2k^2 - h^2)u_{xxx} := \mu u_{xxx}$$

This is also a dispersive equation. Regarding the theory of dispersive waves a solution $u(x, t)$ can be represented in Fourier space. By isolating each wavenumber ξ to apply solutions of the form $u(x, t) = e^{i(\xi x - c(\xi)t)}$ to the linear advection equation yields to the dispersion relation:

$$c(\xi) = a\xi + \mu\xi^3$$

Based on this relation its possible to calculate the group velocity $c'(\xi)$ for wavenumber ξ which describes in which speed a wave peak travel. It is given by:

$$c'(\xi) = a + 3\mu\xi^2$$

As $\mu = \frac{a}{6}h^2(\nu^2 - 1)$ and since $a > 0$ and for stability $|\nu| < 1$ the group velocity $c'(\xi)$ is smaller than a for all ξ , but tens to a as $h \rightarrow 0$. It follows that the scheme leads to an oscillatory wave train lagging behind the discontinuity, which is traveling with speed a . This can also later be seen later in the numerical experiments.

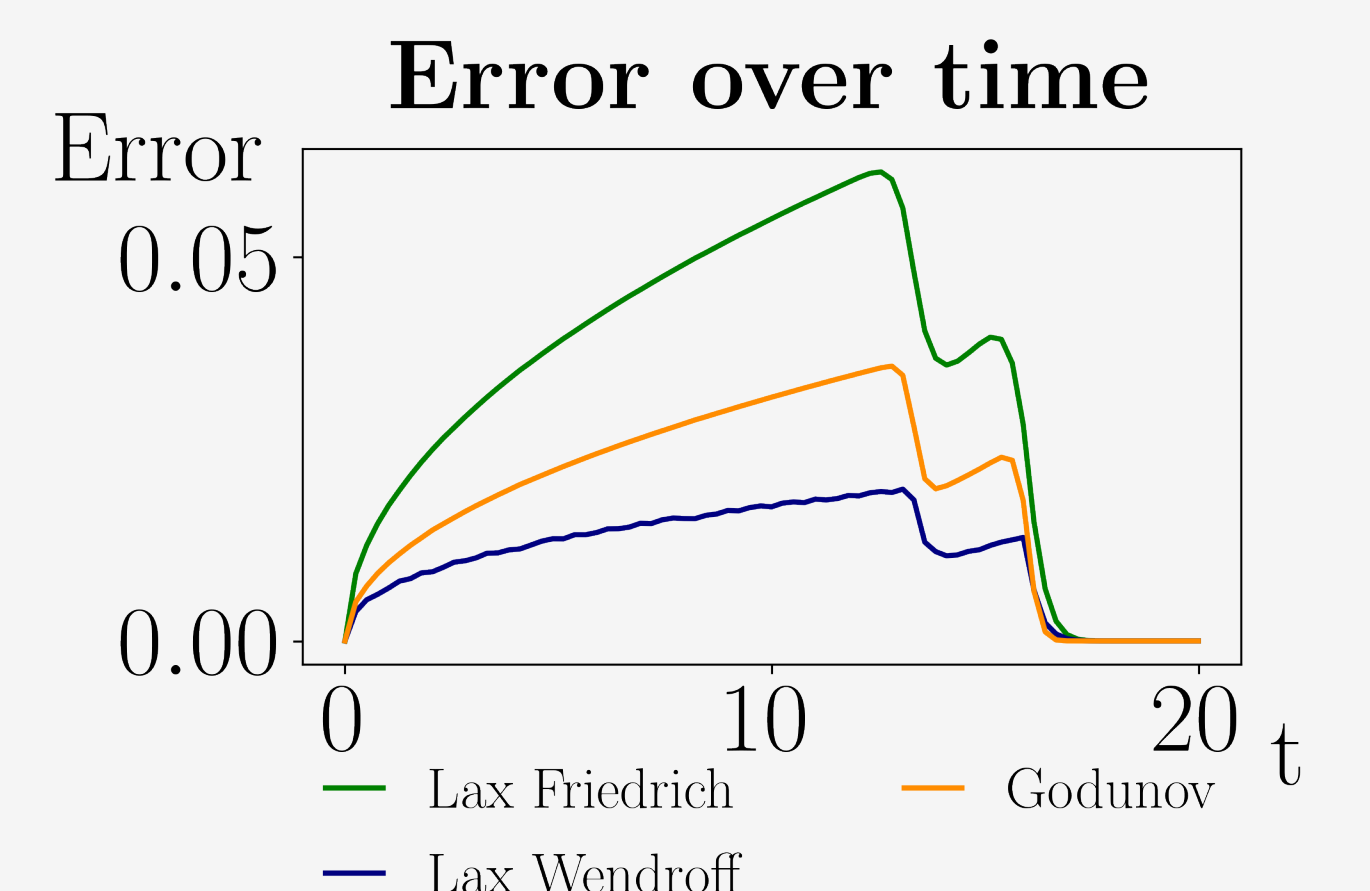
Numerical Analysis

In the following the approximations of the Lax-Wendroff method will be presented in comprehension to the Lax-Friedrich and Godunov method.

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REFERENCES

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