Linear Algebra Review

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 $A \in \mathbb{C}^{n \times n}$ is Hermitian if $A = A^{\dagger}$ $A \in \mathbb{R}^{n \times n}$ is Hermitian if $A = A^T$

 $O: V \to V$ Hermitian if $\langle Ox, y \rangle = \langle x, Oy \rangle$ for all $x, y \in V$ Hermitian operators have real eigenvalues, orthogonal eigenfunctions and correspond to observable quantities $U \in \mathbb{C}^{n \times n}$ unitary if $UU^{\dagger} = U^{\dagger}U = I \Longrightarrow U^{-1} = U^{\dagger}$ $O \in \mathbb{R}^{n \times n}$ orthogonal if $OO^T = O^TO = I \Longrightarrow O^{-1} = O^T$

Commutator

$$\begin{split} [A,B] &= AB - BA \\ [A,B] &= 0 \Longleftrightarrow AB = BA \Longleftrightarrow B = ABA^{-1} = A^{-1}BA \\ [B,A] &= -[A,B] \qquad [AB,C] = A[B,C] + [A,C]B \\ \{A,B\} &= AB + BA \end{split}$$

Quantum Mechanics in 1D

Some Basic Operators

$$p_x \to -i\hbar \frac{\partial}{\partial x} \quad \hat{\boldsymbol{p}} \to -i\hbar \nabla$$

Uncertainty Principle

$$(\Delta A \Delta B)^2 \ge \left(\frac{1}{2} |\langle [A, B] \rangle|\right)^2$$
$$\Delta x \Delta p \ge \frac{\hbar}{2} \quad \Delta E \Delta t \ge \frac{\hbar}{2}$$

Expectation Values, etc...

$$\begin{aligned}
\langle \mathcal{O} \rangle &= \langle \psi | \mathcal{O} | \psi \rangle \equiv \langle \psi | \mathcal{O} \psi \rangle \equiv \int \psi^* \mathcal{O} \psi \, \mathrm{d} x \\
\langle \psi | \mathcal{O} \psi \rangle &= \langle \mathcal{O}^* \psi | \psi \rangle \text{ and } \mathcal{O} = \mathcal{O}^* \Longrightarrow \langle \psi | \mathcal{O} \psi \rangle = \langle \mathcal{O} \psi | \psi \rangle \\
\langle \psi_m | \psi_n \rangle &= \int \psi_m^* \psi_n \, \mathrm{d} x \\
\langle \mathcal{O} \psi | \mathcal{O} \psi \rangle &= \| \mathcal{O} \psi \|^2 \\
\langle A^{\dagger} \rangle &= \langle A \rangle^* \Longrightarrow \langle A \rangle + \langle A^{\dagger} \rangle = \langle A \rangle + \langle A \rangle^* = 2 \operatorname{Re} \langle A \rangle
\end{aligned}$$

$$\begin{array}{l} \textbf{Uncertainties} \\ \Delta x(t) = \sqrt{\left\langle x^2, t \right\rangle - \left\langle x, t \right\rangle^2} \quad \Delta p(t) = \sqrt{\left\langle p^2, t \right\rangle - \left\langle p, t \right\rangle^2} \\ (\Delta A)^2 = \left\langle A^2 \right\rangle - \left\langle A \right\rangle^2 = \left\langle (A - \left\langle A \right\rangle) \right\rangle \\ \end{array}$$

Basis Expansion
$$\psi(x,0) = \int c(k) \frac{e^{ikx}}{\sqrt{2\pi}} \, \mathrm{d}k \Longrightarrow c(k) = \int \psi(x,0) \frac{e^{-ikx}}{\sqrt{2\pi}} \, \mathrm{d}x$$
 Given orthonormal basis $\langle \psi_n | \psi_m \rangle = \delta_{nm}$:
$$\psi = \sum_n c_n \psi_n \qquad c_n = \int \psi_n^* \psi \, \mathrm{d}x \equiv \langle \psi_n(x) | \psi(x,0) \rangle$$

Observation

Expand $|\psi\rangle$ in basis of operator being observed Possible outcomes are eigenvalues of each basis state Probability is square of eigenstate coefficient ψ then collapses to eigenfunction of observed eigenvalue

Free Particle

$$\psi_k(x) = \frac{e^{ikx}}{\sqrt{2\pi}} \qquad E_k = \frac{\hbar^2 k^2}{2m}$$
$$\frac{1}{2\pi} \int e^{i(\tilde{k}-k)x} \, \mathrm{d}x \equiv \delta(\tilde{k}-k)$$

Infinite Potential Well

$$V(|x| < \frac{a}{2}) = 0 \qquad V(|x| > \frac{a}{2}) \to \infty$$

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left[\frac{n\pi}{a}\left(x - x_c + \frac{a}{2}\right)\right] \qquad E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}$$
If $x \in [0, a]$, $\psi_n(x) = \sqrt{\frac{2}{a}} \begin{cases} \sin\frac{n\pi x}{a} & n \text{ even} \\ \cos\frac{n\pi x}{a} & n \text{ odd} \end{cases}$

Finite Potential Well

$$V(|x| < \frac{a}{2}) = 0$$
 $V(|x| > \frac{a}{2}) = V_0$

$$\kappa^{2} = \frac{2mE}{\hbar^{2}} \qquad k^{2} = \frac{2m(V_{0} - E)}{\hbar^{2}}, \quad E > 0$$

$$\text{Sym: } \kappa = k \tan\left(\frac{ka}{2}\right) \qquad \text{Asym: } \kappa = -k \cot\left(\frac{ka}{2}\right)$$

$$u = ka \qquad u_{0}^{2} = \frac{2mV_{0}a^{2}}{\hbar^{2}} \qquad \kappa^{2} = \frac{u_{0}^{2} - u^{2}}{a^{2}}$$

$$\tan\left(\frac{u}{2}\right) = \sqrt{\frac{u_{0}^{2}}{u^{2}} - 1} \qquad \cot\left(\frac{u}{2}\right) = -\sqrt{\frac{u_{0}^{2}}{u^{2}} - 1}$$

Delta Function Well

$$V(x) = -\lambda \delta(x)$$

Boundary Condition:
$$\psi'(0_+) - \psi'(0_-) = -\frac{2m\lambda\psi(0)}{\hbar^2}$$

 $\psi(x) = \sqrt{\kappa}e^{-\kappa|x|}, \quad E_0 = \frac{m\lambda^2}{2\hbar^2}, \quad \kappa = \frac{m\lambda}{\hbar^2}$

Scattering

$$j(x) = \frac{\hbar}{2mi} \left[\psi^* \psi' - {\psi'}^* \psi \right] = \frac{\hbar}{m} \operatorname{Im} \left\{ \psi^* \psi' \right\}$$
Regions 1 and 3 free, region 2 with $E > V = V(x)$

$$\psi_2 = Cf(x) + Dg(x)$$

$$\psi_1 = \frac{A_1}{\sqrt{v_1}} e^{ik_1 x} + \frac{B_1}{\sqrt{v_1}} e^{-ik_1 x} \quad \psi_3 = \frac{A_3}{\sqrt{v_3}} e^{-ik_3 x} + \frac{B_3}{\sqrt{v_3}} e^{ik_3 x}$$

$$\begin{split} v &\equiv \frac{\hbar k}{m}, \qquad k^2 = \frac{2mE}{\hbar^2} \\ j_{1,3} &= |A_{1,3}|^2 - |B_{1,3}|^2 \\ \begin{bmatrix} B_1 \\ B_3 \end{bmatrix} &= \begin{bmatrix} r & \tilde{t} \\ t & \tilde{r} \end{bmatrix} \begin{bmatrix} A_1 \\ A_3 \end{bmatrix} \Longleftrightarrow \boldsymbol{B} = \mathbf{S}\boldsymbol{A} \\ T &= |t|^2 \quad R = |r|^2 \quad \tilde{T} &= |\tilde{t}|^2 \quad \tilde{R} = |\tilde{r}|^2 \\ T + R &= 1 \text{ and } \tilde{T} + \tilde{R} = 1 \\ j_{\text{in}} &= \boldsymbol{A}^{\dagger}\boldsymbol{A} \qquad j_{\text{out}} = \boldsymbol{B}^{\dagger}\boldsymbol{B} \qquad j_{\text{in}} \equiv j_{\text{out}} \\ j_{\text{in}} &= j_{\text{out}} \Longrightarrow (\boldsymbol{A}^{\dagger}\mathbf{S}^{\dagger})(\mathbf{S}\boldsymbol{A}) \Longrightarrow \boldsymbol{S}^{\dagger}\mathbf{S} = \mathbf{I} \end{split}$$

Delta Function Scattering

$$H = \frac{p^2}{2m} + \lambda \delta(x), \quad E > 0$$

$$r = \tilde{r} = -\frac{i\kappa}{k+i\kappa}, \quad t = \tilde{t} = \frac{k}{k+i\kappa}$$

$$k^2 = \frac{2mE}{\hbar^2} \qquad \kappa = \frac{m\lambda}{\hbar^2}$$

$$\mathbf{S} = \begin{bmatrix} r & t \\ t & r \end{bmatrix} = \frac{1}{k+i\kappa} \begin{bmatrix} -i\kappa & k \\ k & -i\kappa \end{bmatrix}$$

Harmonic Oscillator

$$H = \frac{p^2}{2m} + \frac{kx^2}{2} = \frac{p^2}{2m} + \frac{m\omega}{2}x^2, \qquad \omega = \sqrt{\frac{k}{m}}$$

$$H | n \rangle = \hbar \omega \left(n + \frac{1}{2} \right) | n \rangle, \quad n = 0, 1, 2, \dots$$

$$a = \frac{1}{\sqrt{2}} \left(\frac{x}{x_0} + i \frac{p}{p_0} \right) \qquad a^\dagger = \frac{1}{\sqrt{2}} \left(\frac{x}{x_0} - i \frac{p}{p_0} \right)$$

$$[a, a^\dagger] = 1 \qquad \qquad [a^\dagger, a] = -1$$

$$a | n \rangle = \sqrt{n} | n - 1 \rangle \qquad a^\dagger | n \rangle = \sqrt{n+1} | n+1 \rangle$$

$$x_0 = \sqrt{\frac{\hbar}{m\omega}} \qquad p_0 = \frac{\hbar}{x_0}$$

$$x = \frac{x_0}{\sqrt{2}} (a + a^\dagger) \qquad p = \frac{p_0}{\sqrt{2}i} (a - a^\dagger) \qquad H = \hbar \omega \left(a^\dagger a + \frac{1}{2} \right)$$

$$a^\dagger a | n \rangle = n | n \rangle \qquad | n \rangle = \frac{(a^\dagger)^n}{\sqrt{n!}} | 0 \rangle$$
Ehrenfest:
$$\frac{d}{dt} \langle x, t \rangle = \frac{\langle p, t \rangle}{m}, \frac{d}{dt} \langle p, t \rangle = \langle -\frac{dV}{dx} \rangle = -k \langle x, t \rangle$$

Coherent States

$$\begin{aligned} a & |z\rangle = z & |z\rangle, \ z \in \mathbb{C} \qquad \sum_{n} c_{n} a & |n\rangle = \sum_{n} c_{n} z & |n\rangle \\ c_{n+1}\sqrt{n+1} = c_{n}z \Longrightarrow c_{n} = \frac{z^{n}}{\sqrt{n!}}c_{0} \end{aligned}$$
$$|z\rangle = \sum_{n} c_{n} |n\rangle = c_{0} \sum_{n} \frac{z^{n}}{\sqrt{n!}} |n\rangle, \quad c_{0} = e^{-\frac{|z|^{2}}{2}}$$
$$|z\rangle = e^{-\frac{|z|^{2}}{2}} e^{za^{\dagger}} |0\rangle$$
$$|z,t\rangle = e^{-i\frac{\omega}{2}t} |ze^{-i\omega t}\rangle \qquad z(t) = ze^{-i\omega t}$$
$$\langle z| (a^{\dagger})^{n} a^{m} |z\rangle = (z^{*})^{n} z^{m}$$
$$\langle z| (a^{\dagger})^{n} a^{m} |z\rangle = (z^{*})^{n} z^{m}$$
$$\langle x\rangle = \sqrt{2}x_{0} \operatorname{Re} z \qquad \langle x^{2}\rangle = x_{0}^{2} \left(2 \operatorname{Re}^{2} z + \frac{1}{2}\right) \\\langle p\rangle = \sqrt{2}p_{0} \operatorname{Im} z \qquad \langle p^{2}\rangle = p_{0}^{2} \left(2 \operatorname{Im}^{2} z + \frac{1}{2}\right) \\\langle H\rangle = \frac{\hbar \omega}{2} + \hbar \omega |z|^{2} \qquad \langle H^{2}\rangle = \hbar^{2} \omega^{2} (|z|^{4} + 2|z|^{2} + \frac{1}{4}) \end{aligned}$$

2D Isotropic Oscillator

$$k_x = k_y \equiv k$$
 $\omega_x = \omega_y \equiv \omega = \sqrt{\frac{k}{m}}$
 $V(\mathbf{r}) = \frac{1}{2}k(x^2 + y^2) = \frac{1}{2}kr^2$
 $H |n_x n_y\rangle = \hbar\omega(n_x + n_y + 1) |n_x n_y\rangle$

Gaussian Wave Packet
$$\psi(x) = \frac{1}{\sqrt[4]{2\pi\sigma^2}} \exp\left(\frac{(x - \langle x \rangle)^2}{4\sigma^2}\right) e^{-i\frac{\langle p \rangle}{\hbar}x}$$
$$\Delta x(0) = \sigma \Longrightarrow \langle x^2, 0 \rangle = \sigma^2 + \langle x, 0 \rangle^2$$
$$\Delta p(0) = \frac{\hbar}{2\Delta x(0)} = \frac{\hbar}{2\sigma} \Longrightarrow \langle p^2, 0 \rangle = \frac{\hbar^2}{4\sigma^2} + \langle p, 0 \rangle^2$$

Time Evolution

Expand
$$\psi(x,0)$$
 over eigenfunctions $\psi_n(x)$ solving $H\psi_n(x) = E_n\psi_n(x)$ to get $\psi(x,0) = \sum_n c_n\psi_n(x)$.
Find $\psi(x,t)$ with $\psi(x,t) = \sum_n c_n e^{-\frac{iE_n}{\hbar}t}\psi_n(x)$.
 $|\psi,t\rangle \equiv \sum_n c_n e^{-i\frac{E_n}{\hbar}t} |n\rangle = e^{-i\frac{H}{\hbar}t} |\psi,0\rangle$

Time-Dependent Expectation Value

Heisenberg approach: $\langle \mathcal{O}, t \rangle = \langle \psi, 0 | \mathcal{O}(t) | \psi, 0 \rangle$ Schrödinger approach: $\langle \mathcal{O}, t \rangle = \langle \psi, t | \mathcal{O} | \psi, t \rangle$

Time-Dependent Operator

Time-Dependent Operator
$$\mathcal{O}(t) = e^{i\frac{H}{\hbar}t}\mathcal{O}e^{-i\frac{H}{\hbar}t} \quad \frac{\mathrm{d}}{\mathrm{d}t}\mathcal{O}(t) = \frac{i}{\hbar}[H,\mathcal{O}](t)$$

$$\mathcal{O}(t)\big|_{t=0} = e^{i\frac{H}{\hbar}\cdot 0}\mathcal{O}e^{-i\frac{H}{\hbar}\cdot 0} = \mathcal{O}$$

$$(A\cdot B)(t) = A(t)B(t)$$

Angular Momentum

$$L_{z} = xp_{y} - yp_{x} \rightarrow -\hbar \frac{\partial}{\partial \phi}$$

$$L_{z} |lm\rangle = m\hbar |lm\rangle \qquad L^{2} |lm\rangle = l(l+1)\hbar^{2} |lm\rangle$$

$$m = -l, -l+1, \dots, l-1, l$$

$$\langle \mathbf{r}|lm\rangle = Y_{l}^{m}(\vartheta, \phi) = (-1)^{m} \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_{lm}(\cos \vartheta) e^{im\phi}$$

$$L_{\pm} |lm\rangle = \hbar \sqrt{l(l+1-m(m\pm 1))} |l,m\pm 1\rangle$$

$$\begin{split} \boldsymbol{L} \cdot \hat{\boldsymbol{e}} & |\psi\rangle = \hbar l \, |\psi\rangle \Longrightarrow \boldsymbol{L} \parallel \hat{\boldsymbol{e}} \\ L_{+} \equiv L_{x} + iL_{y} & L_{-} \equiv L_{+}^{\dagger} = L_{x} - iL_{y} \\ L_{+} & |lm\rangle = \hbar \sqrt{l(l+1) - m(m+1)} \, |l, m+1\rangle \\ L_{-} & |lm\rangle = \hbar \sqrt{l(l+1) - m(m-1)} \, |l, m-1\rangle \\ L_{x} = \frac{L_{+} + L_{-}}{2} & L_{y} = \frac{L_{+} - L_{-}}{2i} \\ \langle L_{x}, t\rangle = \operatorname{Re} \langle L_{+}, t\rangle & \langle L_{y}, t\rangle = \operatorname{Im} \langle L_{+}, t\rangle \, \boldsymbol{\mu} = \gamma \boldsymbol{L} \end{split}$$

Spin

$$S^{2} = |sm_{s}\rangle = \hbar^{2}s(s+1) |sm_{s}\rangle$$

$$S_{z} |sm_{s}\rangle = \hbar m_{s} |sm_{s}\rangle$$

$$\left|\frac{1}{2}\frac{1}{2}\rangle \equiv |\uparrow\rangle \text{ and } \left|\frac{1}{2} - \frac{1}{2}\rangle \equiv |\downarrow\rangle$$

$$\begin{split} S_{+} &= S_{x} + iS_{y} & S_{-} = S_{x} - iS_{y} \\ S_{x} &= \frac{S_{+} + S_{-}}{2} & S_{y} = \frac{S_{+} - S_{-}}{2i} \\ S_{+} \left| sm_{s} \right\rangle &= \hbar \sqrt{s(s+1) - m_{s}(m_{s}+1)} \left| s, m_{s} + 1 \right\rangle \\ S_{-} \left| sm_{s} \right\rangle &= \hbar \sqrt{s(s+1) - m_{s}(m_{s}-1)} \left| s, m_{s} + 1 \right\rangle \\ S \cdot \hat{e} \left| \psi_{s} \right\rangle &= \frac{\hbar}{2} \left| \psi_{s} \right\rangle \text{ implies} \\ S \parallel \hat{e} &= (\cos \varphi \sin \theta, \sin \varphi \sin \theta, \cos \theta) \\ \left| \psi_{s} \right\rangle &= \cos \frac{\theta}{2} \left| \uparrow \right\rangle + \sin \frac{\theta}{2} e^{i\varphi} \left| \downarrow \right\rangle \\ \text{Useful: } \mathbf{k} &= (k_{x}, k_{y}) \text{ to } k_{x} + ik_{y} = ke^{i\phi}, \ k_{x} - ik_{y} = ke^{-i\phi} \end{split}$$

Spin Matrices

$$\vec{S} = \frac{\hbar}{2} \boldsymbol{\sigma} \qquad \boldsymbol{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$$

$$\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \qquad \sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \qquad \sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -i \end{bmatrix}$$

$$\{\sigma_a, \sigma_b\} = 2\delta_{ab}I$$

Two-Particle Spin System

Product basis: $|s_1m_1s_2m_2\rangle$.

Total angular momentum basis: $|sm_ss_1s_2\rangle$

$$|s_1 - s_2| \le s \le s_1 + s_2$$
 $m_s = -s, -s + 1, \dots, s - 1, s.$

Time Reversal

TODO

For spin s = 1/2 particles $T = i\sigma_{\eta}K$; $K : \psi \mapsto \psi^*$

Nondegenerate Perturbations

$$H = H_0 + H' \qquad H_0 |n_0\rangle = E_n^{(0)} |n_0\rangle$$

$$E_n = E_n^{(0)} + \langle n_0 | H' | n_0 \rangle + \sum_{m \neq n} \frac{|\langle m_0 | H' | n_0 \rangle|^2}{E_n^{(0)} - E_m^{(0)}},$$

$$|n\rangle = |n_0\rangle + \sum_{m \neq n} \frac{|\langle m_0 | H' | n_0 \rangle|^2}{E_n^{(0)} - E_m^{(0)}} |m_0\rangle$$

Degenerate Perturbations

Theory: Degenerate Perturbation Theory

$$H = H_0 + H' \qquad H_0 |\psi\rangle = E^{(0)} |\psi\rangle$$

$$E^{(0)} \text{ is } N\text{-times degenerate...}$$

$$\dots \text{ and } |\psi\rangle = \sum_{i}^{N} c_i |i\rangle \qquad H |i\rangle = E^{(0)} |i\rangle \ \forall \ i$$

$$\text{Create matrix } \mathbf{P} = \begin{pmatrix} \langle 1|H'|1\rangle & \cdots & \langle 1|H'|N\rangle \\ \langle 2|H'|1\rangle & \cdots & \langle 2|H'|N\rangle \\ \vdots & \ddots & \vdots \\ \langle N|H'|1\rangle & \cdots & \langle N|H'|N\rangle \end{pmatrix}$$

Diagonalize **P**; find eigenvalues λ_j and eigenvectors \boldsymbol{u}_j . $E_j = E^{(0)} + \lambda_j$ and $|\psi_j\rangle = \sum_{k=1}^N u_{j_k} |k\rangle$