Linear Algebra Review

Linear Algebra Review

 $A \in \mathbb{C}^{n \times n}$ is Hermitian if $A = A^{\dagger}$ $A \in \mathbb{R}^{n \times n}$ is Hermitian if $A = A^T$

 $O: V \to V$ Hermitian if $\langle Ox, y \rangle = \langle x, Oy \rangle$ for all $x, y \in V$ Hermitian operators have real eigenvalues, orthogonal eigenfunctions and correspond to observable quantities $U \in \mathbb{C}^{n \times n}$ unitary if $UU^{\dagger} = U^{\dagger}U = I \Longrightarrow U^{-1} = U^{\dagger}$ $O \in \mathbb{R}^{n \times n}$ orthogonal if $OO^T = O^TO = I \Longrightarrow O^{-1} = O^T$

Commutator

$$\begin{split} [A,B] &= AB - BA \\ [A,B] &= 0 \Longleftrightarrow AB = BA \Longleftrightarrow B = ABA^{-1} = A^{-1}BA \\ [B,A] &= -[A,B] \qquad [AB,C] = A[B,C] + [A,C]B \\ \{A,B\} &= AB + BA \end{split}$$

Quantum Mechanics in 1D

Some Basic Operators

$$p_x \to -i\hbar \frac{\partial}{\partial x} \quad \vec{p} \to -i\hbar \nabla$$

Uncertainty Principle

$$(\Delta A \Delta B)^2 \ge \left(\frac{1}{2} |\langle [A, B] \rangle|\right)^2$$
$$\Delta x \Delta p \ge \frac{\hbar}{2} \quad \Delta E \Delta t \ge \frac{\hbar}{2}$$

Expectation Values, etc...

$$\begin{split} \langle \mathcal{O} \rangle &= \langle \psi | \, \mathcal{O} \, | \psi \rangle \equiv \langle \psi | \mathcal{O} \psi \rangle \equiv \int \psi^* \mathcal{O} \psi \, \mathrm{d} x \\ \langle \psi | \mathcal{O} \psi \rangle &= \langle \mathcal{O}^* \psi | \psi \rangle \text{ and } \mathcal{O} = \mathcal{O}^* \Longrightarrow \langle \psi | \mathcal{O} \psi \rangle = \langle \mathcal{O} \psi | \psi \rangle \\ \langle \psi_m | \psi_n \rangle &= \int \psi_m^* \psi_n \, \mathrm{d} x \\ \langle \mathcal{O} \psi | \mathcal{O} \psi \rangle &= \| \mathcal{O} \psi \|^2 \\ \langle A^\dagger \rangle &= \langle A \rangle^* \Longrightarrow \langle A \rangle + \langle A^\dagger \rangle = \langle A \rangle + \langle A \rangle^* = 2 \operatorname{Re} \langle A \rangle \end{split}$$

$$\begin{array}{l} \textbf{Uncertainties} \\ \Delta x(t) = \sqrt{\langle x^2, t \rangle - \langle x, t \rangle^2} \quad \Delta p(t) = \sqrt{\langle p^2, t \rangle - \langle p, t \rangle^2} \\ (\Delta A)^2 = \left\langle A^2 \right\rangle - \left\langle A \right\rangle^2 = \left\langle (A - \left\langle A \right\rangle) \right\rangle \end{array}$$

Basis Expansion

Given orthonormal basis
$$\langle \psi_n | \psi_m \rangle = \delta_{nm}$$
:

$$\psi(x,0) = \int c(k) \frac{e^{ikx}}{\sqrt{2\pi}} dk \Longrightarrow c(k) = \int \psi(x,0) \frac{e^{-ikx}}{\sqrt{2\pi}} dx$$
Given orthonormal basis $\langle \psi_n | \psi_m \rangle = \delta_{nm}$:

$$\psi = \sum_n c_n \psi_n \qquad c_n = \int \psi_n^* \psi dx \equiv \langle \psi_n(x) | \psi(x,0) \rangle$$

Free Particle

$$\psi_k(x) = \frac{e^{ikx}}{\sqrt{2\pi}} \qquad E_k = \frac{\hbar^2 k^2}{2m}$$

$$\frac{1}{2\pi} \int e^{i(\tilde{k}-k)x} dx \equiv \delta(\tilde{k}-k)$$

Infinite Potential Well

$$V(|x| < \frac{a}{2}) = 0 \qquad V(|x| > \frac{a}{2}) \to \infty$$

$$\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left[\frac{n\pi}{a}\left(x - x_c + \frac{a}{2}\right)\right] \qquad E_n = \frac{n^2\pi^2\hbar^2}{2ma^2}$$
If $x \in [0, a]$, $\psi_n(x) = \sqrt{\frac{2}{a}} \begin{cases} \sin\frac{n\pi x}{a} & n \text{ even} \\ \cos\frac{n\pi x}{a} & n \text{ odd} \end{cases}$

Finite Potential Well

$$\begin{split} &V(|x|<\frac{a}{2})=0 \qquad V(|x|>\frac{a}{2})=V_0\\ &\kappa^2=\frac{2mE}{\hbar^2} \qquad k^2=\frac{2m(V_0-E)}{\hbar^2}, \quad E>0\\ &\mathrm{Sym:} \ \kappa=k\tan\left(\frac{ka}{2}\right) \quad \mathrm{Asym:} \ \kappa=-k\cot\left(\frac{ka}{2}\right)\\ &u=ka \qquad u_0^2=\frac{2mV_0a^2}{\hbar^2} \qquad \kappa^2=\frac{u_0^2-u^2}{a^2}\\ &\tan\left(\frac{u}{2}\right)=\sqrt{\frac{u_0^2}{u^2}-1} \qquad \cot\left(\frac{u}{2}\right)=-\sqrt{\frac{u_0^2}{u^2}-1} \end{split}$$

Delta Function Well

$$V(x) = -\lambda \delta(x)$$

Boundary Condition:
$$\psi'(0_+) - \psi'(0_-) = -\frac{2m\lambda\psi(0)}{\hbar^2}$$

 $\psi(x) = \sqrt{\kappa}e^{-\kappa|x|}, \quad E_0 = \frac{m\lambda^2}{2\hbar^2}, \quad \kappa = \frac{m\lambda}{\hbar^2}$

Harmonic Oscillator

$$\begin{split} H &= \frac{p^2}{2m} + \frac{kx^2}{2} = \frac{p^2}{2m} + \frac{m\omega}{2}x^2, \qquad \omega = \sqrt{\frac{k}{m}} \\ H &|n\rangle = \hbar\omega \left(n + \frac{1}{2}\right)|n\rangle \;, \quad n = 0, 1, 2, \dots \\ a &= \frac{1}{\sqrt{2}}\left(\frac{x}{x_0} + i\frac{p}{p_0}\right) \qquad a^\dagger = \frac{1}{\sqrt{2}}\left(\frac{x}{x_0} - i\frac{p}{p_0}\right) \\ \left[a, a^\dagger\right] &= 1 \qquad \qquad \left[a^\dagger, a\right] = -1 \\ a &|n\rangle &= \sqrt{n}\,|n-1\rangle \qquad a^\dagger &|n\rangle = \sqrt{n+1}\,|n+1\rangle \\ x_0 &= \sqrt{\frac{\hbar}{m\omega}} \qquad p_0 = \frac{\hbar}{x_0} \\ x &= \frac{x_0}{\sqrt{2}}(a+a^\dagger) \qquad p = \frac{p_0}{\sqrt{2}i}(a-a^\dagger) \qquad H = \hbar\omega \left(a^\dagger a + \frac{1}{2}\right) \\ a^\dagger a &|n\rangle = n\,|n\rangle \qquad |n\rangle = \frac{(a^\dagger)^n}{\sqrt{n!}}\,|0\rangle \\ \text{Ehrenfest: } \frac{\mathrm{d}}{\mathrm{d}t}\,\langle x, t\rangle &= \frac{\langle p, t\rangle}{m}, \, \frac{\mathrm{d}}{\mathrm{d}t}\,\langle p, t\rangle = \langle -\frac{\mathrm{d}V}{\mathrm{d}x}\rangle = -k\,\langle x, t\rangle \end{split}$$

Coherent States

$$\begin{array}{l} a\mid z\rangle = z\mid z\rangle,\,z\in\mathbb{C} \qquad \sum_{n}c_{n}a\mid n\rangle = \sum_{n}c_{n}z\mid n\rangle \\ c_{n+1}\sqrt{n+1} = c_{n}z\Longrightarrow c_{n} = \frac{z^{n}}{\sqrt{n!}}c_{0} \\ \mid z\rangle = \sum_{n}c_{n}\mid n\rangle = c_{0}\sum_{n}\frac{z^{n}}{\sqrt{n!}}\mid n\rangle, \quad c_{0} = e^{-\frac{\mid z\mid^{2}}{2}} \\ \mid z\rangle = e^{-\frac{\mid z\mid^{2}}{2}}e^{za^{\dagger}}\mid 0\rangle \\ \mid z,t\rangle = e^{-i\frac{\omega}{2}t}\mid ze^{-i\omega t}\rangle \qquad z(t) = ze^{-i\omega t} \\ \langle z\mid (a^{\dagger})^{n}a^{m}\mid z\rangle = (z^{*})^{n}z^{m} \\ \langle x\rangle = \sqrt{2}x_{0}\operatorname{Re}z \qquad \langle x^{2}\rangle = x_{0}^{2}\left(2\operatorname{Re}^{2}z+\frac{1}{2}\right) \\ \langle p\rangle = \sqrt{2}p_{0}\operatorname{Im}z \qquad \langle p^{2}\rangle = p_{0}^{2}\left(2\operatorname{Im}^{2}z+\frac{1}{2}\right) \\ \langle H\rangle = \frac{\hbar\omega}{2}+\hbar\omega|z|^{2} \qquad \langle H^{2}\rangle = \hbar^{2}\omega^{2}\left(|z|^{4}+2|z|^{2}+\frac{1}{4}\right) \end{array}$$

2D Isotropic Oscillator

$$k_x = k_y \equiv k$$
 $\omega_x = \omega_y \equiv \omega = \sqrt{\frac{k}{m}}$
 $V(\mathbf{r}) = \frac{1}{2}k(x^2 + y^2) = \frac{1}{2}kr^2$
 $H |n_x n_y\rangle = \hbar\omega(n_x + n_y + 1) |n_x n_y\rangle$

Gaussian Wave Packet
$$\psi(x) = \frac{1}{\sqrt[4]{2\pi\sigma^2}} \exp\left(\frac{(x-\langle x\rangle)^2}{4\sigma^2}\right) e^{-i\frac{\langle p\rangle}{\hbar}x}$$

$$\Delta x(0) = \sigma \Longrightarrow \langle x^2, 0\rangle = \sigma^2 + \langle x, 0\rangle^2$$

$$\Delta p(0) = \frac{\hbar}{2\Delta x(0)} = \frac{\hbar}{2\sigma} \Longrightarrow \langle p^2, 0\rangle = \frac{\hbar^2}{4\sigma^2} + \langle p, 0\rangle^2$$

Time Evolution

Expand $\psi(x,0)$ over eigenfunctions $\psi_n(x)$ solving $H\psi_n(x) = E_n\psi_n(x)$ to get $\psi(x,0) = \sum_n c_n\psi_n(x)$. Find $\psi(x,t)$ with $\psi(x,t) = \sum_{n} c_n e^{-\frac{iE_n}{\hbar}t} \psi_n(x)$. $|\psi,t\rangle \equiv \sum_{n} c_n e^{-i\frac{E_n}{\hbar}t} |n\rangle = e^{-i\frac{H}{\hbar}t} |\psi,0\rangle$

Time-Dependent Expectation Value

Heisenberg approach: $\langle \mathcal{O}, t \rangle = \langle \psi, 0 | \mathcal{O}(t) | \psi, 0 \rangle$ Schrödinger approach: $\langle \mathcal{O}, t \rangle = \langle \psi, t | \mathcal{O} | \psi, t \rangle$

Time-Dependent Operator

Time-Dependent Operator
$$\mathcal{O}(t) = e^{i\frac{H}{\hbar}t}\mathcal{O}e^{-i\frac{H}{\hbar}t} \qquad \frac{\mathrm{d}}{\mathrm{d}t}\mathcal{O}(t) = \frac{i}{\hbar}[H,\mathcal{O}](t)$$

$$\mathcal{O}(t)\big|_{t=0} = e^{i\frac{H}{\hbar}\cdot 0}\mathcal{O}e^{-i\frac{H}{\hbar}\cdot 0} = \mathcal{O}$$

$$(A\cdot B)(t) = A(t)B(t)$$

Angular Momentum

$$\begin{split} L_z &= x p_y - y p_x \rightarrow -\hbar \frac{\partial}{\partial \phi} \\ L_z &| lm \rangle = m \hbar \, | lm \rangle \qquad L^2 \, | lm \rangle = l(l+1) \hbar^2 \, | lm \rangle \\ m &= -l, -l+1, \dots, l-1, l \end{split}$$

$$\langle \mathbf{r} | lm \rangle = Y_l^m(\vartheta, \phi) = (-1)^m \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_{lm}(\cos \vartheta) e^{im\phi} \quad \begin{array}{l} L_x = \frac{L_+ + L_-}{2} \\ \langle L_x, t \rangle = \operatorname{Re} \langle L_+, t \rangle \end{array} \quad \begin{array}{l} L_y = \frac{L_+ - L_-}{2i} \\ \langle L_y, t \rangle = \operatorname{Im} \langle L_+, t \rangle \end{array}$$

$$L_{\pm} | lm \rangle = \hbar \sqrt{l(l+1-m(m\pm 1))} | l, m\pm 1 \rangle \qquad \qquad \mu = \gamma \mathbf{L}$$

$$\begin{split} \boldsymbol{L} \cdot \hat{\boldsymbol{e}} & |\psi\rangle = \hbar l \, |\psi\rangle \Longrightarrow \boldsymbol{L} \parallel \hat{\boldsymbol{e}} \\ L_{+} &\equiv L_{x} + iL_{y} \qquad L_{-} \equiv L_{+}^{\dagger} = L_{x} - iL_{y} \\ L_{+} & |lm\rangle = \hbar \sqrt{l(l+1) - m(m+1)} \, |l,m+1\rangle \\ L_{-} & |lm\rangle = \hbar \sqrt{l(l+1) - m(m-1)} \, |l,m-1\rangle \end{split}$$

$$L_{x} = \frac{L_{+} + L_{-}}{2}$$

$$\langle L_{x}, t \rangle = \operatorname{Re} \langle L_{+}, t \rangle$$

$$L_{y} = \frac{L_{+} - L_{-}}{2i}$$

$$\langle L_{y}, t \rangle = \operatorname{Im} \langle L_{+}, t \rangle$$

$$\mu = \gamma L$$

Observation

Expand $|\psi\rangle$ in basis of operator being observed Possible outcomes are eigenvalues of each basis state Probability is square of eigenstate coefficient ψ then collapses to eigenfunction of observed eigenvalue